PREDICTIVE MONTE CARLO TOOLS FOR COLLIDERS

FABIO MALTONI
CENTRE FOR COSMOLOGY, PARTICLE PHYSICS AND PHENOMENOLOGY (CP3), BELGIUM

LECTURE I

TEST: HOW MUCH DO I KNOW ABOUT MC'S?

Statements		TRUE	FALSE	IT DEPENDS	I have no clue
0	MC's are black boxes, I don't need to know the details as long as there are no bugs.				
1	A MC generator produces "unweighted" events, i.e., events distributed as in Nature.				
2	MC's are based on a classical approximation (Markov Chain), QM effects are not included.				
3	The "Sudakov form factor" directly quantifies how likely it is for a parton to undergo branching.				
4	A calculation/code at NLO for a process provides NLO predictions for any IR safe observable.				
5	Tree-level based MC's are less accurate than those at NLO.				

TEST: HOW MUCH DO I KNOW ABOUT MC'S?

	Statements	TRUE	FALSE	IT DEPENDS	I have no clue
0	MC's are black boxes, I don't need to know the details as long as there are no bugs.		✓		
1	A MC generator produces "unweighted" events, i.e., events distributed as in Nature.	✓			
2	MC's are based on a classical approximation (Markov Chain), QM effects are not included.		✓		
3	The "Sudakov form factor" directly quantifies how likely it is for a parton to undergo branching.		✓		
4	A calculation/code at NLO for a process provides NLO predictions for any IR safe observable.		✓		
5	Tree-level based MC's are less accurate than those at NLO.			✓	

TEST: HOW MUCH DO I KNOW ABOUT MC'S?

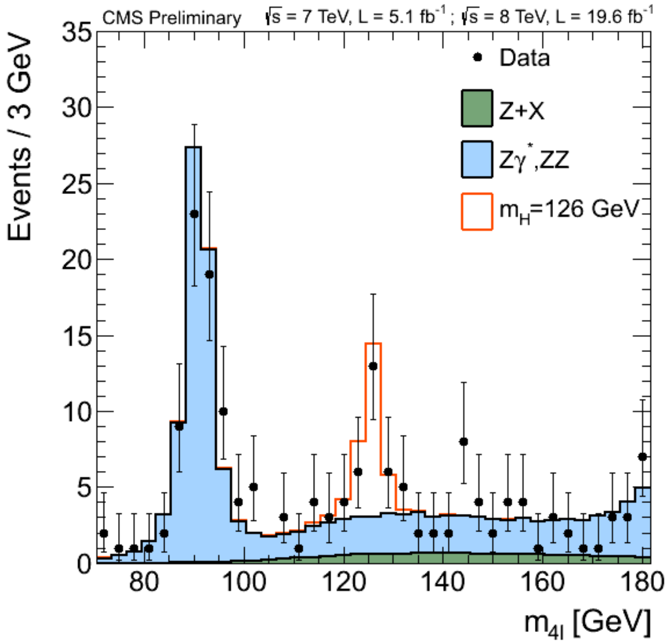
Score	Result	Comment
≥5	Addict	Always keep in mind that there are also other interesting activities in the field.
4	Excellent	No problem in following these lectures.
3	Fair	Check out carefully the missed topics.
≤2	Room for improvement	Enroll in a MC crash course at your home institution.
6 x no clue	No clue	

DISCOVERIES AT HADRON COLLIDERS

DISCOVERIES AT HADRON COLLIDERS

peak

$pp \to H \to 4l$



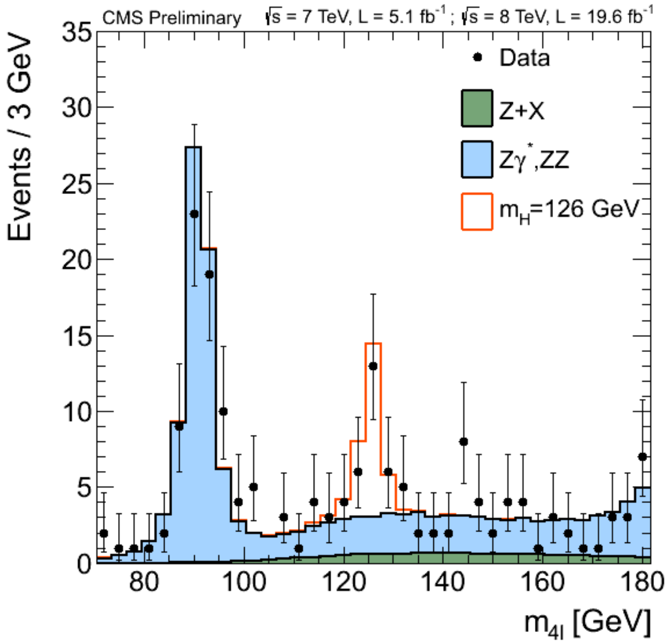
“easy”

Background directly measured from data. TH needed only for parameter extraction (Normalization, acceptance,...)

Discoveries at hadron colliders

peak

$pp \to H \to 4l$

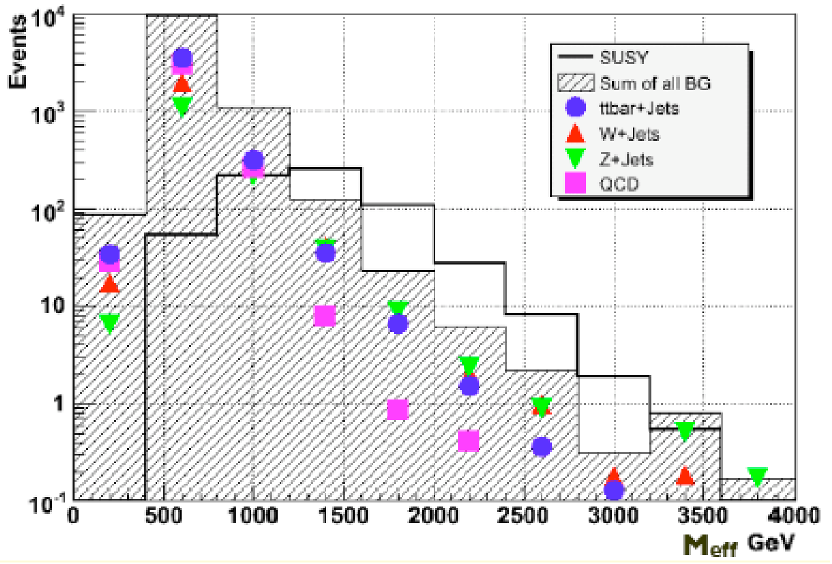


"easy"

Background directly measured from data. TH needed only for parameter extraction (Normalization, acceptance,...)

shape

$pp \to \tilde{g}\tilde{g}, \tilde{g}\tilde{q}, \tilde{q}\tilde{q} \to \text{jets} + \not{E}_T$



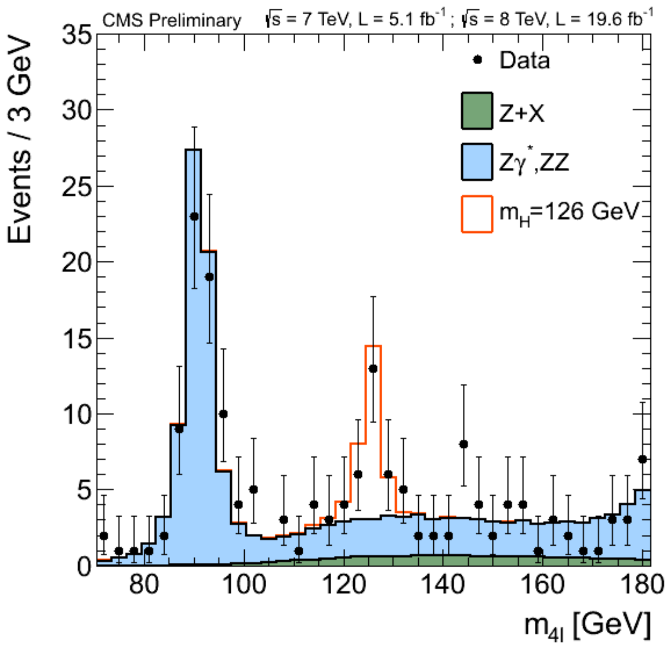
hard

Background shapes needed. Flexible MC for both signal and background tuned and validated with data.

DISCOVERIES AT HADRON COLLIDERS

peak

$pp \to H \to 4l$

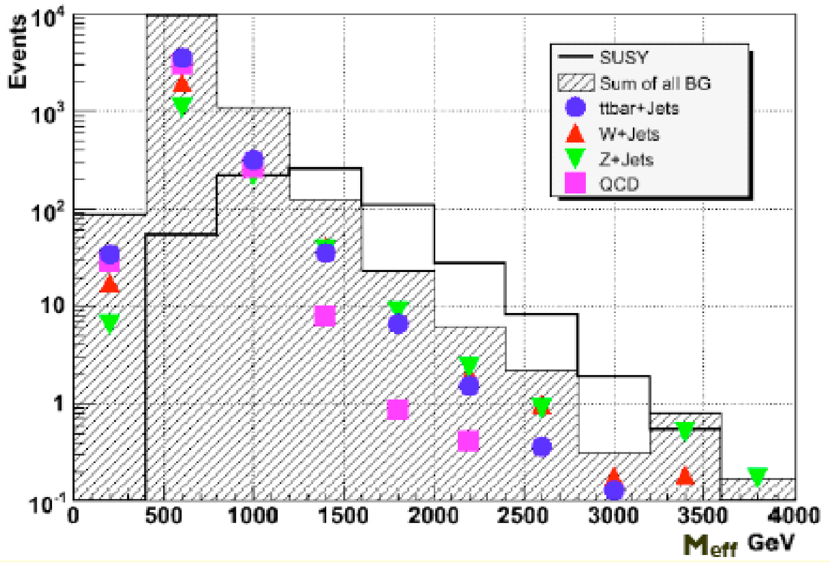


“easy”

Background directly measured from data. TH needed only for parameter extraction (Normalization, acceptance,...)

shape

$pp \to \tilde{g}\tilde{g}, \tilde{g}\tilde{q}, \tilde{q}\tilde{q} \to \text{jets} + \not{E}_T$

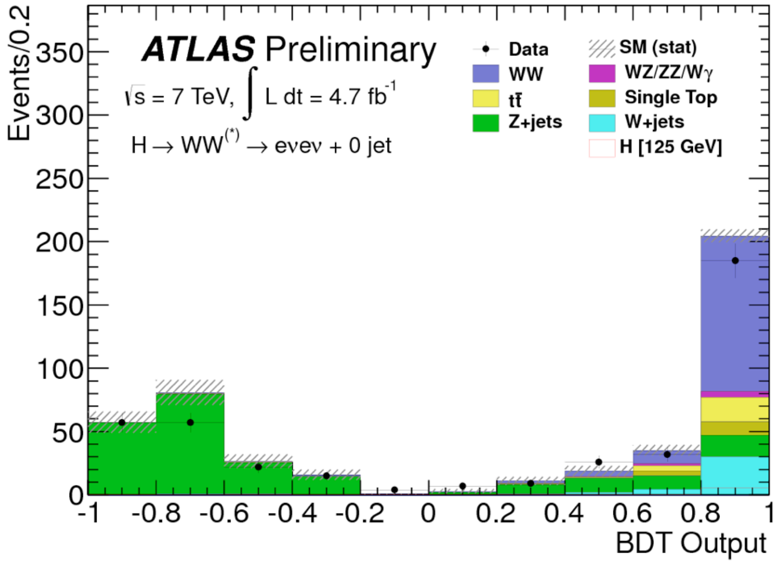


hard

Background shapes needed. Flexible MC for both signal and background tuned and validated with data.

discriminant

$pp \to H \to W^+W^-$



very hard

Background normalization and shapes known very well. Interplay with the best theoretical predictions (via MC) and data.

NO SIGN OF NEW PHYSICS (SO FAR)!

MC developer

WHY HAPPY?

WHY HAPPY?

- Optimism: New Physics could be hiding there already, just need to dig it out.

WHY HAPPY?

- Optimism: New Physics could be hiding there already, just need to dig it out.

- Democratization: No evidence of most beaten BSM proposals, means more and more room for diversification. Possibility for small teams to make a big discovery.

WHY HAPPY?

- Optimism: New Physics could be hiding there already, just need to dig it out.
- Democratization: No evidence of most beaten BSM proposals, means more and more room for diversification. Possibility for small teams to make a big discovery.
- Ingenuity/Creativity: From new signatures to smart and new analysis techniques (MVA), and combination with non-collider searches (DM, Flavor...).

WHY HAPPY?

- Optimism: New Physics could be hiding there already, just need to dig it out.
- Democratization: No evidence of most beaten BSM proposals, means more and more room for diversification. Possibility for small teams to make a big discovery.
- Ingenuity/Creativity: From new signatures to smart and new analysis techniques (MVA), and combination with non-collider searches (DM, Flavor...).
- Massification (the practice of making luxury products available to the mass market) : MC's in the hands of every th/exp might turn out to be the best overall strategy for discovering the Unexpected.

WHY HAPPY?

- Optimism: New Physics could be hiding there already, just need to dig it out.

- Democratization: No evidence of most beaten BSM proposals, means more and more room for diversification. Possibility for small teams to make a big discovery.

- Ingenuity/Creativity: From new signatures to smart and new analysis techniques (MVA), and combination with non-collider searches (DM, Flavor...).

- Massification (the practice of making luxury products available to the mass market) : MC's in the hands of every th/exp might turn out to be the best overall strategy for discovering the Unexpected.

- Flexibility: We need MC that are able to predict the pheno of the Unexpected.

WHY HAPPY?

- Optimism: New Physics could be hiding there already, just need to dig it out.

- Democratization: No evidence of most beaten BSM proposals, means more and more room for diversification. Possibility for small teams to make a big discovery.

- Ingenuity/Creativity: From new signatures to smart and new analysis techniques (MVA), and combination with non-collider searches (DM, Flavor...).

- Massification (the practice of making luxury products available to the mass market) : MC's in the hands of every th/exp might turn out to be the best overall strategy for discovering the Unexpected.

- Flexibility: We need MC that are able to predict the pheno of the Unexpected.

- Accuracy: accurate simulations for both SM and BSM are a must.

CHALLENGES FOR LHC PHYSICISTS

CHALLENGES FOR LHC PHYSICISTS

- Accurate and experimental friendly predictions for collider physics range from being very useful to strictly necessary.

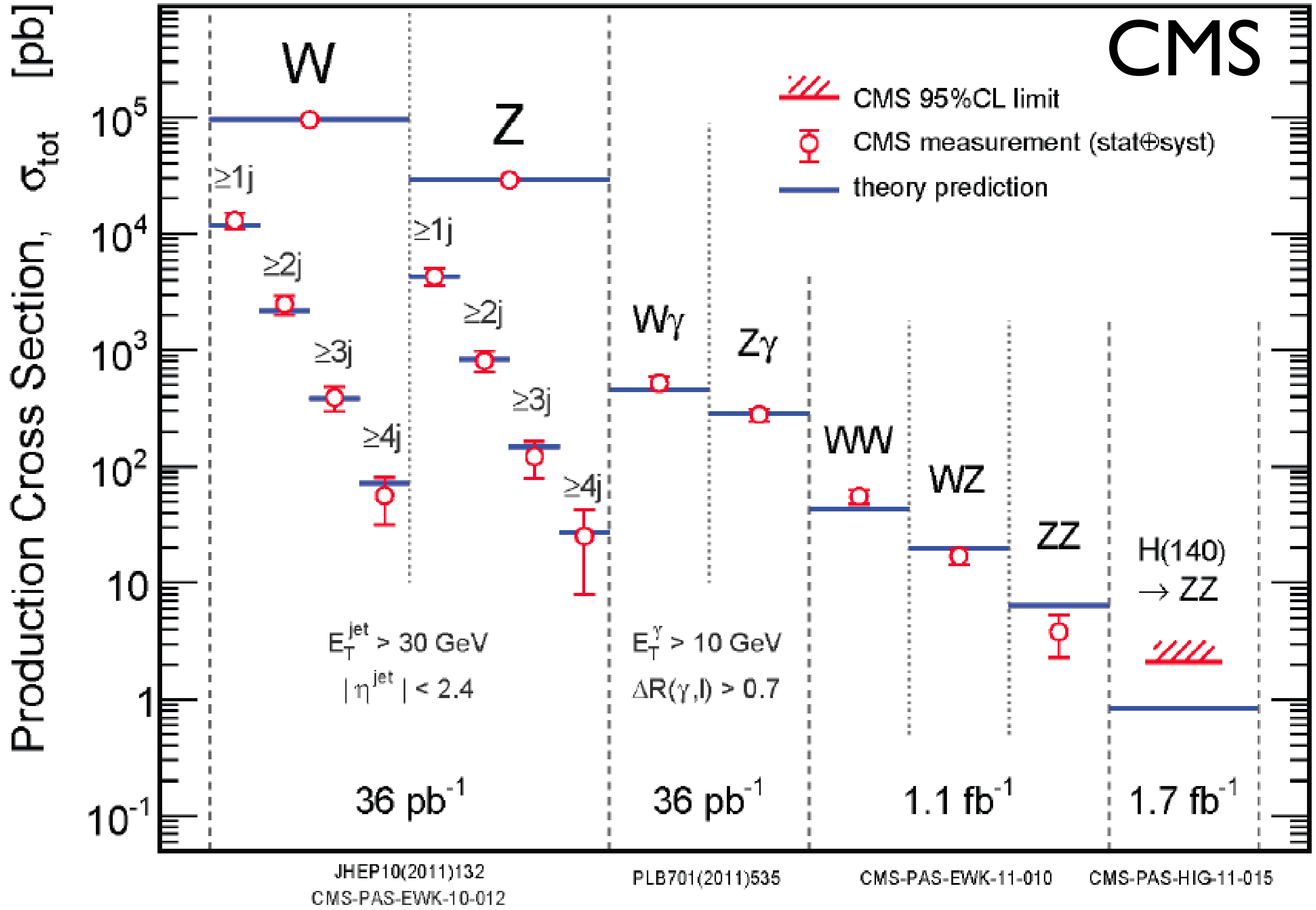
CHALLENGES FOR LHC PHYSICISTS

- Accurate and experimental friendly predictions for collider physics range from being very useful to strictly necessary.
- Confidence on possible excesses, evidences and eventually discoveries builds upon an intense (and often non-linear) process of description/ prediction of data via MC's.

CHALLENGES FOR LHC PHYSICISTS

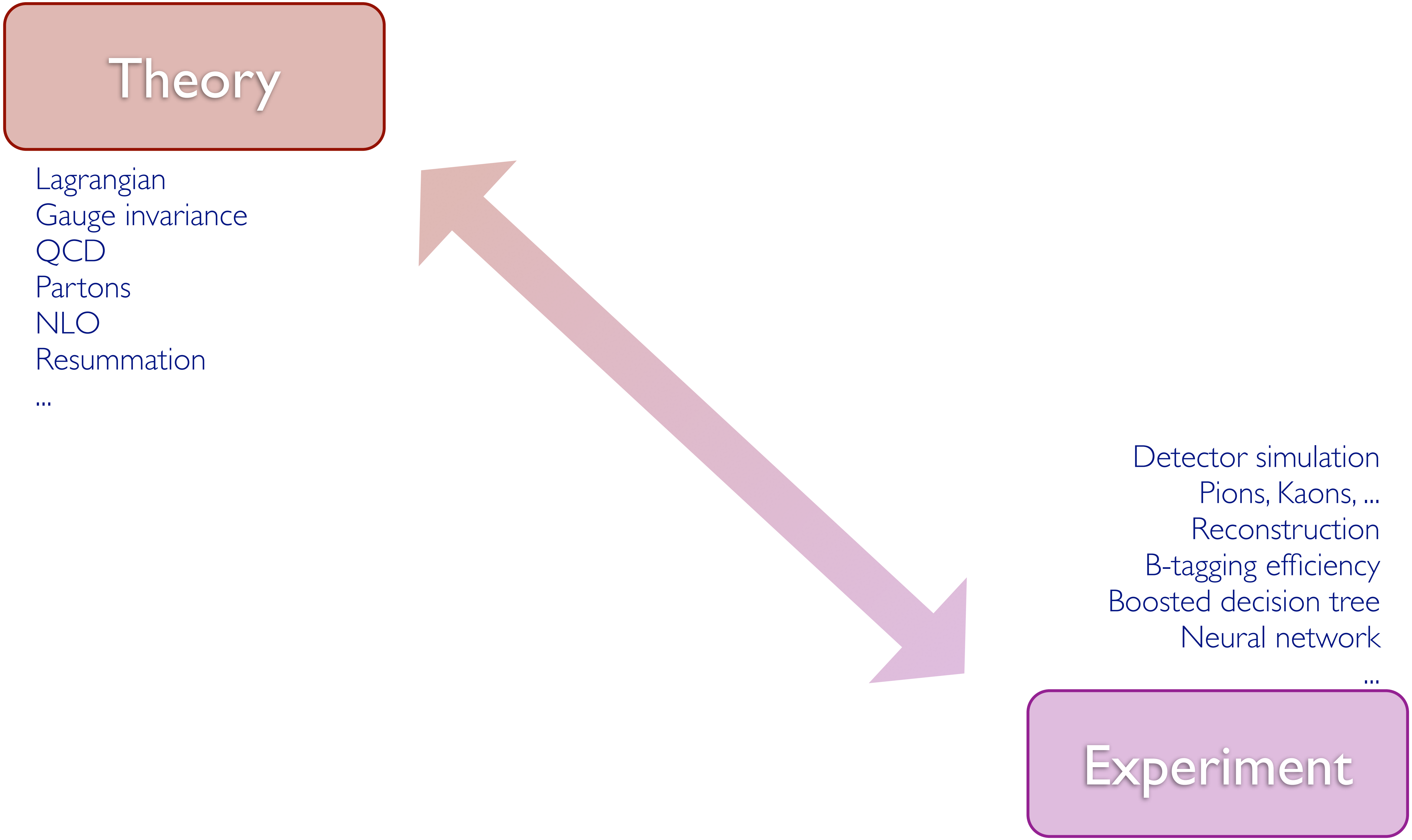
- Accurate and experimental friendly predictions for collider physics range from being very useful to strictly necessary.
- Confidence on possible excesses, evidences and eventually discoveries builds upon an intense (and often non-linear) process of description/prediction of data via MC's.
- Both **measurements** and **exclusions** rely on accurate predictions.

CHALLENGES FOR LHC PHYSICISTS



Even this plot actually needs theory input (and the total quoted uncertainty in the measurements does have a contribution from theory)!!!

NEW GENERATION (LHC) OF MC TOOLS



Theory

Lagrangian
Gauge invariance
QCD
Partons
NLO
Resummation
...

Detector simulation
Pions, Kaons, ...
Reconstruction
B-tagging efficiency
Boosted decision tree
Neural network
...

Experiment

NEW GENERATION (LHC) OF MC TOOLS

Theory

Lagrangian
Gauge invariance
QCD
Partons
NLO
Resummation
...

Detector simulation
Pions, Kaons, ...
Reconstruction
B-tagging efficiency
Boosted decision tree
Neural network
...

Experiment

AIMS FOR THESE LECTURES

- Recall the basics of the necessary QCD concepts to understand what is going on in a pp event at the TeV scale.
- Critically revisit the "old" ways of making predictions for hadron colliders: either via fixed-order predictions or parton showers.
- Present the new *predictive* techniques which allow to:
  - Merge tree-level calculations with parton showers (CKKW/MLM).
  - Match NLO calculations with parton showers (MC@NLO and POWHEG) automatically.

WIKI

PLAN

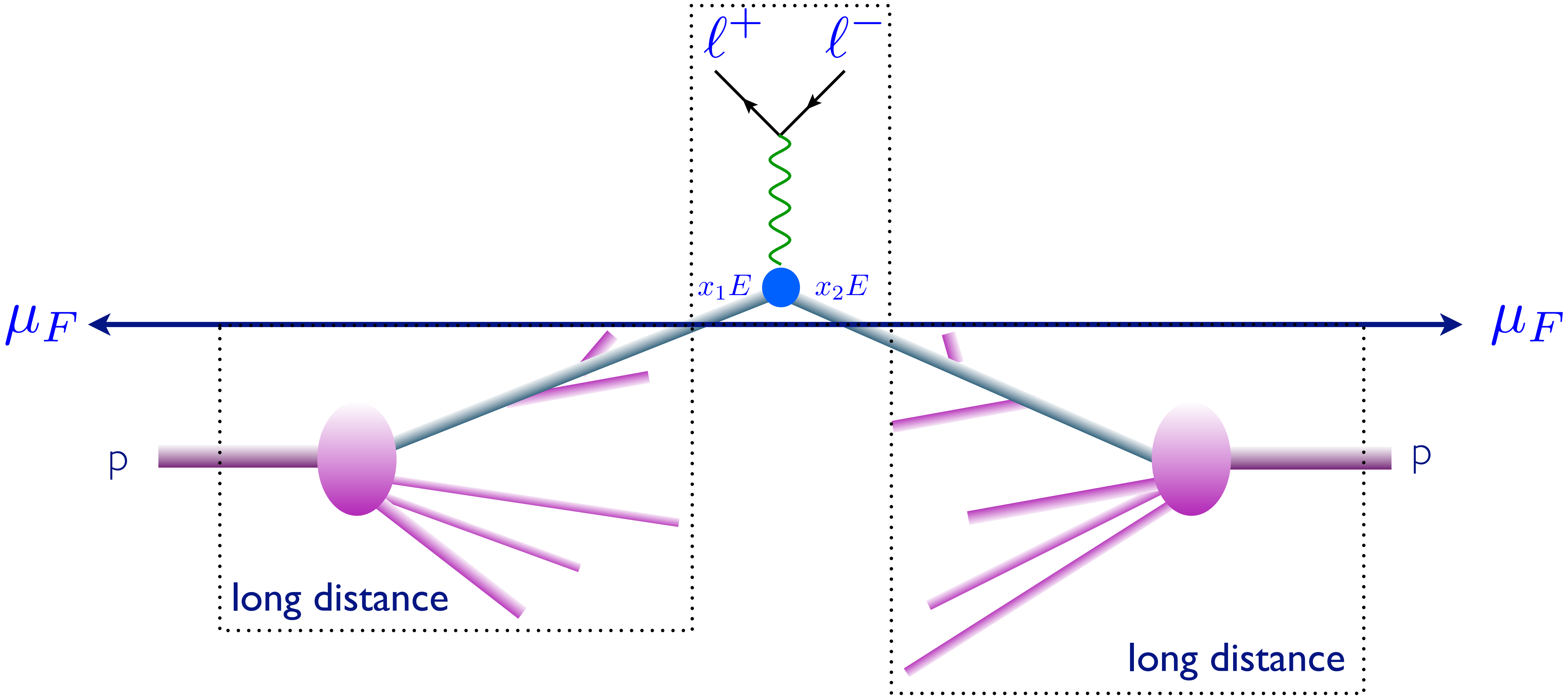
- Basics : LO predictions and event generation
- Fixed-order calculations : from NLO to NNLO
- Exclusive predictions : Parton Shower
- Merging ME+PS
- Matching NLO with PS

PLAN

- Basics : LO predictions and event generation
- Fixed-order calculations : from NLO to NNLO
- Exclusive predictions : Parton Shower
- Merging ME+PS
- Matching NLO with PS

Today

MASTER FORMULA FOR THE LHC



$$\sum_{a,b} \int dx_1 dx_2 d\Phi_{\rm FS}\, f_a(x_1,\mu_F) f_b(x_2,\mu_F)\, \hat\sigma_{ab\to X}(\hat s,\mu_F,\mu_R)$$

Phase-space integral

Parton density functions

Parton-level cross section

MASTER FORMULA FOR THE LHC

$$\sum_{a,b} \int dx_1 dx_2 d\Phi_{\rm FS}\, f_a(x_1,\mu_F) f_b(x_2,\mu_F)\, \hat{\sigma}_{ab\to X}(\hat{s},\mu_F,\mu_R)$$

Phase-space integral

Parton density functions

Parton-level cross section

Two ingredients necessary:

1. Parton distribution functions : non perturbative (fit from experiments, but evolution from theory)

2. Parton-level cross section: short distance coefficients as an expansion in α_S (from theory)

PERTURBATIVE EXPANSION

$\hat{\sigma}_{ab\to X}(\hat{s},\mu_F,\mu_R)$ Parton-level cross section

PERTURBATIVE EXPANSION

$\hat{\sigma}_{ab\to X}(\hat{s}, \mu_F, \mu_R)$ Parton-level cross section

- The parton-level cross section can be computed as a series in perturbation theory, using the coupling constant as an expansion parameter

PERTURBATIVE EXPANSION

$\hat{\sigma}_{ab\to X}(\hat{s}, \mu_F, \mu_R)$ **Parton-level cross section**

- The parton-level cross section can be computed as a series in perturbation theory, using the coupling constant as an expansion parameter

$$\hat{\sigma} = \sigma^{\mathrm{Born}}\left(1 + \frac{\alpha_s}{2\pi}\sigma^{(1)} + \left(\frac{\alpha_s}{2\pi}\right)^2 \sigma^{(2)} + \left(\frac{\alpha_s}{2\pi}\right)^3 \sigma^{(3)} + \ldots\right)$$

PERTURBATIVE EXPANSION

$\hat{\sigma}_{ab \to X}(\hat{s}, \mu_F, \mu_R)$ Parton-level cross section

- The parton-level cross section can be computed as a series in perturbation theory, using the coupling constant as an expansion parameter

$$\hat{\sigma} = \sigma^{\text{Born}} \left(1 + \frac{\alpha_s}{2\pi} \sigma^{(1)} + \left(\frac{\alpha_s}{2\pi} \right)^2 \sigma^{(2)} + \left(\frac{\alpha_s}{2\pi} \right)^3 \sigma^{(3)} + \dots \right)$$

LO predictions

Perturbative expansion

$\hat{\sigma}_{ab\to X}(\hat{s},\mu_F,\mu_R)$ **Parton-level cross section**

- The parton-level cross section can be computed as a series in perturbation theory, using the coupling constant as an expansion parameter

$$\hat{\sigma} = \sigma^{\rm Born}\left(1+\frac{\alpha_s}{2\pi}\sigma^{(1)}+\left(\frac{\alpha_s}{2\pi}\right)^2\sigma^{(2)}+\left(\frac{\alpha_s}{2\pi}\right)^3\sigma^{(3)}+\ldots\right)$$

LO predictions

NLO corrections

PERTURBATIVE EXPANSION

$\hat{\sigma}_{ab\to X}(\hat{s},\mu_F,\mu_R)$ **Parton-level cross section**

- The parton-level cross section can be computed as a series in perturbation theory, using the coupling constant as an expansion parameter

$$\hat{\sigma} = \sigma^{\mathrm{Born}}\left(1+\frac{\alpha_s}{2\pi}\sigma^{(1)}+\left(\frac{\alpha_s}{2\pi}\right)^2\sigma^{(2)}+\left(\frac{\alpha_s}{2\pi}\right)^3\sigma^{(3)}+\dots\right)$$

LO predictions

NLO corrections

NNLO corrections

Perturbative expansion

$\hat{\sigma}_{ab\to X}(\hat{s},\mu_F,\mu_R)$ **Parton-level cross section**

- The parton-level cross section can be computed as a series in perturbation theory, using the coupling constant as an expansion parameter

$$\hat{\sigma} = \sigma^{\rm Born}\left(1+\frac{\alpha_s}{2\pi}\sigma^{(1)}+\left(\frac{\alpha_s}{2\pi}\right)^2\sigma^{(2)}+\left(\frac{\alpha_s}{2\pi}\right)^3\sigma^{(3)}+\ldots\right)$$

LO predictions

NLO corrections

NNLO corrections

NNNLO corrections

PERTURBATIVE EXPANSION

$\hat{\sigma}_{ab\to X}(\hat{s},\mu_F,\mu_R)$ **Parton-level cross section**

- The parton-level cross section can be computed as a series in perturbation theory, using the coupling constant as an expansion parameter

$$\hat{\sigma} = \sigma^{\text{Born}}\left(1+\frac{\alpha_s}{2\pi}\sigma^{(1)}+\left(\frac{\alpha_s}{2\pi}\right)^2\sigma^{(2)}+\left(\frac{\alpha_s}{2\pi}\right)^3\sigma^{(3)}+\ldots\right)$$

LO predictions

NLO corrections

NNLO corrections

NNNLO corrections

- Including higher corrections improves predictions and reduces theoretical uncertainties

PREDICTIONS AT LO

How do we calculate a LO cross section for 3 jets at the LHC?

PREDICTIONS AT LO

How do we calculate a LO cross section for 3 jets at the LHC?

1. Identify all subprocesses (gg→ggg, qg→qgg....) in:

$$\sigma(pp \to 3j) = \sum_{ijk} \int f_i(x_1) f_j(x_2) \hat{\sigma}(ij \to k_1 k_2 k_3)$$

PREDICTIONS AT LO

How do we calculate a LO cross section for 3 jets at the LHC?

I. Identify all subprocesses (gg→ggg, qg→qgg....) in:

$$\sigma(pp \to 3j) = \sum_{ijk} \int f_i(x_1) f_j(x_2) \hat{\sigma}(ij \to k_1 k_2 k_3)$$

II. For each one, calculate the amplitude:

$$\mathcal{A}(\{p\},\{h\},\{c\}) = \sum_i D_i$$

PREDICTIONS AT LO

How do we calculate a LO cross section for 3 jets at the LHC?

I. Identify all subprocesses (gg→ggg, qg→qgg....) in:

$$\sigma(pp \to 3j) = \sum_{ijk} \int f_i(x_1) f_j(x_2) \hat{\sigma}(ij \to k_1 k_2 k_3)$$

II. For each one, calculate the amplitude:

$$\mathcal{A}(\{p\},\{h\},\{c\}) = \sum_i D_i$$

III. Square the amplitude, sum over spins & color, integrate over the phase space (D ~ 3n)

$$\hat{\sigma} = \frac{1}{2\hat{s}} \int d\Phi_p \sum_{h,c} |\mathcal{A}|^2$$

PREDICTIONS AT LO

How do we calculate a LO cross section for 3 jets at the LHC?

I. Identify all subprocesses (gg→ggg, qg→qgg....) in:

$$\sigma(pp \to 3j) = \sum_{ijk} \int f_i(x_1) f_j(x_2) \hat{\sigma}(ij \to k_1 k_2 k_3)$$

easy

II. For each one, calculate the amplitude:

$$\mathcal{A}(\{p\},\{h\},\{c\}) = \sum_i D_i$$

III. Square the amplitude, sum over spins & color, integrate over the phase space (D ~ 3n)

$$\hat{\sigma} = \frac{1}{2\hat{s}} \int d\Phi_p \sum_{h,c} |\mathcal{A}|^2$$

PREDICTIONS AT LO

How do we calculate a LO cross section for 3 jets at the LHC?

I. Identify all subprocesses (gg→ggg, qg→qgg....) in:

$$\sigma(pp \to 3j) = \sum_{ijk} \int f_i(x_1) f_j(x_2) \hat{\sigma}(ij \to k_1 k_2 k_3)$$

easy

II. For each one, calculate the amplitude:

$$\mathcal{A}(\{p\},\{h\},\{c\}) = \sum_i D_i$$

difficult

III. Square the amplitude, sum over spins & color, integrate over the phase space (D ~ 3n)

$$\hat{\sigma} = \frac{1}{2\hat{s}} \int d\Phi_p \sum_{h,c} |\mathcal{A}|^2$$

PREDICTIONS AT LO

How do we calculate a LO cross section for 3 jets at the LHC?

I. Identify all subprocesses (gg→ggg, qg→qgg....) in:

$$\sigma(pp \to 3j) = \sum_{ijk} \int f_i(x_1) f_j(x_2) \hat{\sigma}(ij \to k_1 k_2 k_3)$$

easy

II. For each one, calculate the amplitude:

$$\mathcal{A}(\{p\},\{h\},\{c\}) = \sum_i D_i$$

difficult

III. Square the amplitude, sum over spins & color, integrate over the phase space (D ~ 3n)

$$\hat{\sigma} = \frac{1}{2\hat{s}} \int d\Phi_p \sum_{h,c} |\mathcal{A}|^2$$

quite hard

PHASE-SPACE INTEGRAL

- Calculations of cross section or decay widths involve integrations over phase space of very complex functions

PHASE-SPACE INTEGRAL

- Calculations of cross section or decay widths involve integrations over phase space of very complex functions

$$\sigma = \frac{1}{2s}\int |\mathcal{M}|^2 d\Phi(n)$$

PHASE-SPACE INTEGRAL

- Calculations of cross section or decay widths involve integrations over phase space of very complex functions

$$Dim[\Phi(n)] \sim 3n$$

$$\sigma = \frac{1}{2s} \int |\mathcal{M}|^2 d\Phi(n)$$

PHASE-SPACE INTEGRAL

- Calculations of cross section or decay widths involve integrations over phase space of very complex functions

$$Dim[\Phi(n)] \sim 3n$$

$$\sigma = \frac{1}{2s}\int |\mathcal{M}|^2 d\Phi(n)$$

General and flexible method is needed:
Numerical (Monte Carlo) integration

PHASE-SPACE

PHASE-SPACE

$$d\Phi_n = \left[\Pi_{i=1}^{n} \frac{d^3 p_i}{(2\pi)^3 (2E_i)}\right] (2\pi)^4 \delta^{(4)}(p_0 - \sum_{i=1}^{n} p_i)$$

PHASE-SPACE

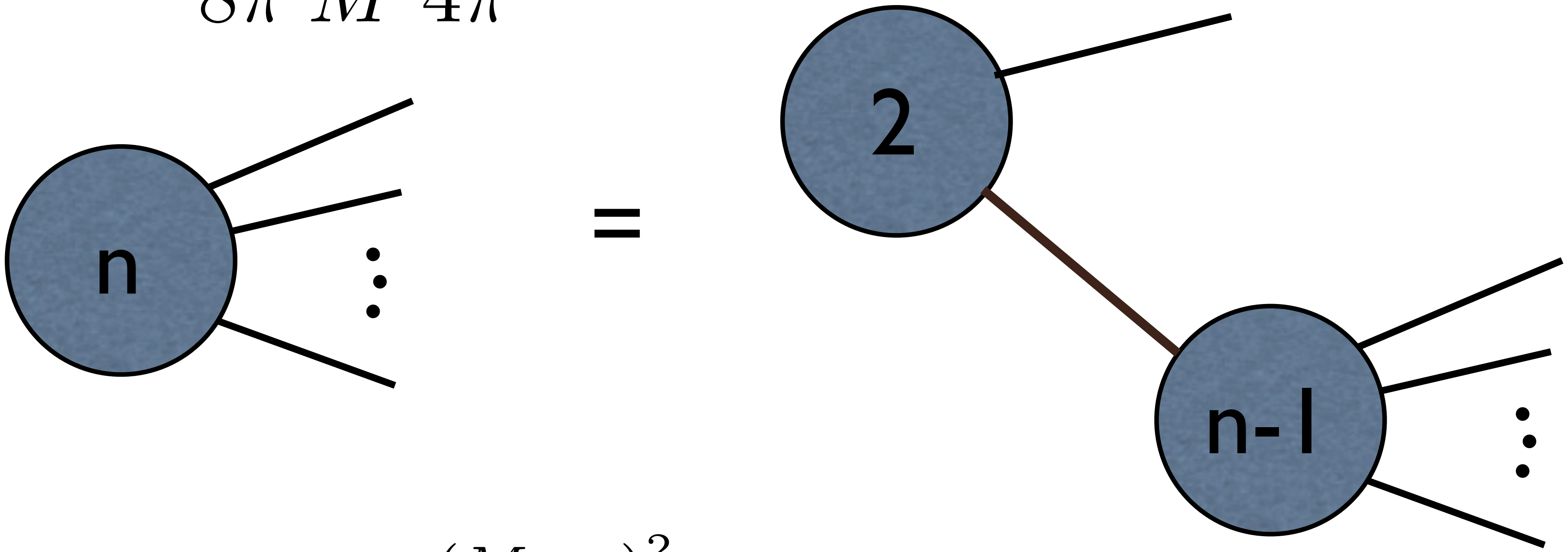
$$d\Phi_n = \left[\Pi_{i=1}^{n} \frac{d^3 p_i}{(2\pi)^3 (2E_i)} \right] (2\pi)^4 \delta^{(4)} (p_0 - \sum_{i=1}^{n} p_i)$$

$$d\Phi_2(M) = \frac{1}{8\pi} \frac{2p}{M} \frac{d\Omega}{4\pi}$$

PHASE-SPACE

$$d\Phi_n = \left[\Pi_{i=1}^n \frac{d^3 p_i}{(2\pi)^3 (2E_i)}\right] (2\pi)^4 \delta^{(4)}(p_0 - \sum_{i=1}^n p_i)$$

$$d\Phi_2(M) = \frac{1}{8\pi} \frac{2p}{M} \frac{d\Omega}{4\pi}$$



$$d\Phi_n(M) = \frac{1}{2\pi} \int_0^{(M-\mu)^2} d\mu^2 d\Phi_2(M) d\Phi_{n-1}(\mu)$$

INTEGRALS AS AVERAGES

$$I = \int_{x_1}^{x_2} f(x) dx \quad \Rightarrow \quad I_N = (x_2 - x_1) \frac{1}{N} \sum_{i=1}^{N} f(x)$$

$$V = (x_2 - x_1) \int_{x_1}^{x_2} [f(x)]^2 dx - I^2 \quad \Rightarrow \quad V_N = (x_2 - x_1)^2 \frac{1}{N} \sum_{i=1}^{N} [f(x)]^2 - I_N^2$$

INTEGRALS AS AVERAGES

$$I = \int_{x_1}^{x_2} f(x)dx \Longrightarrow I_N = (x_2 - x_1)\frac{1}{N}\sum_{i=1}^{N} f(x)$$

$$V = (x_2 - x_1)\int_{x_1}^{x_2} [f(x)]^2 dx - I^2 \Longrightarrow V_N = (x_2 - x_1)^2 \frac{1}{N}\sum_{i=1}^{N} [f(x)]^2 - I_N^2$$

$$I = I_N \pm \sqrt{V_N/N}$$

INTEGRALS AS AVERAGES

$$I = \int_{x_1}^{x_2} f(x)dx \quad \Longrightarrow \quad I_N = (x_2 - x_1)\frac{1}{N}\sum_{i=1}^{N} f(x)$$

$$V = (x_2 - x_1)\int_{x_1}^{x_2} [f(x)]^2 dx - I^2 \quad \Longrightarrow \quad V_N = (x_2 - x_1)^2 \frac{1}{N}\sum_{i=1}^{N} [f(x)]^2 - I_N^2$$

$$I = I_N \pm \sqrt{\enclose{circle}{V_N}/N}$$

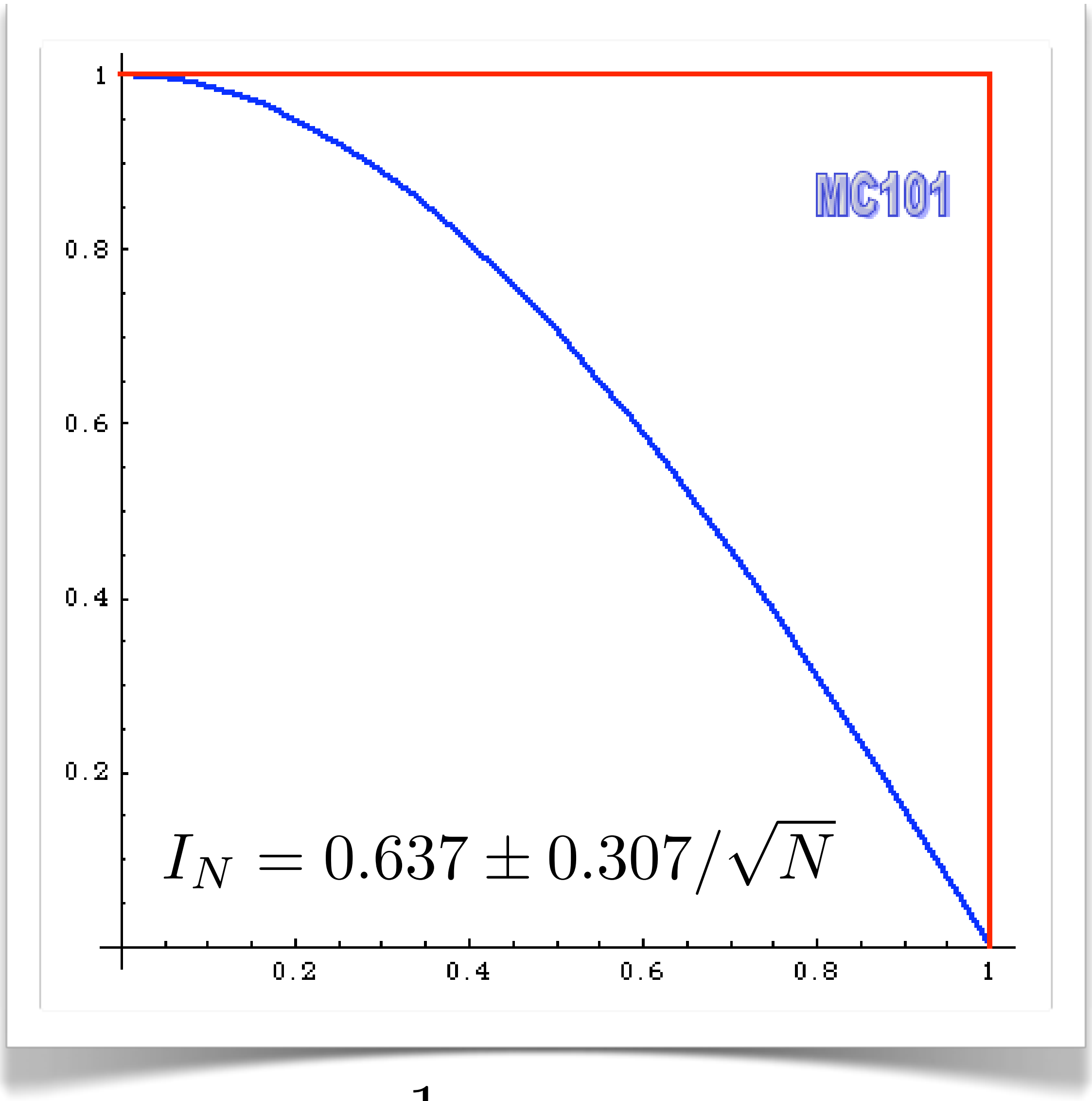
☞ Convergence is slow but it can be estimated easily

☞ Error does not depend on # of dimensions!

☞ Improvement by minimizing V_N

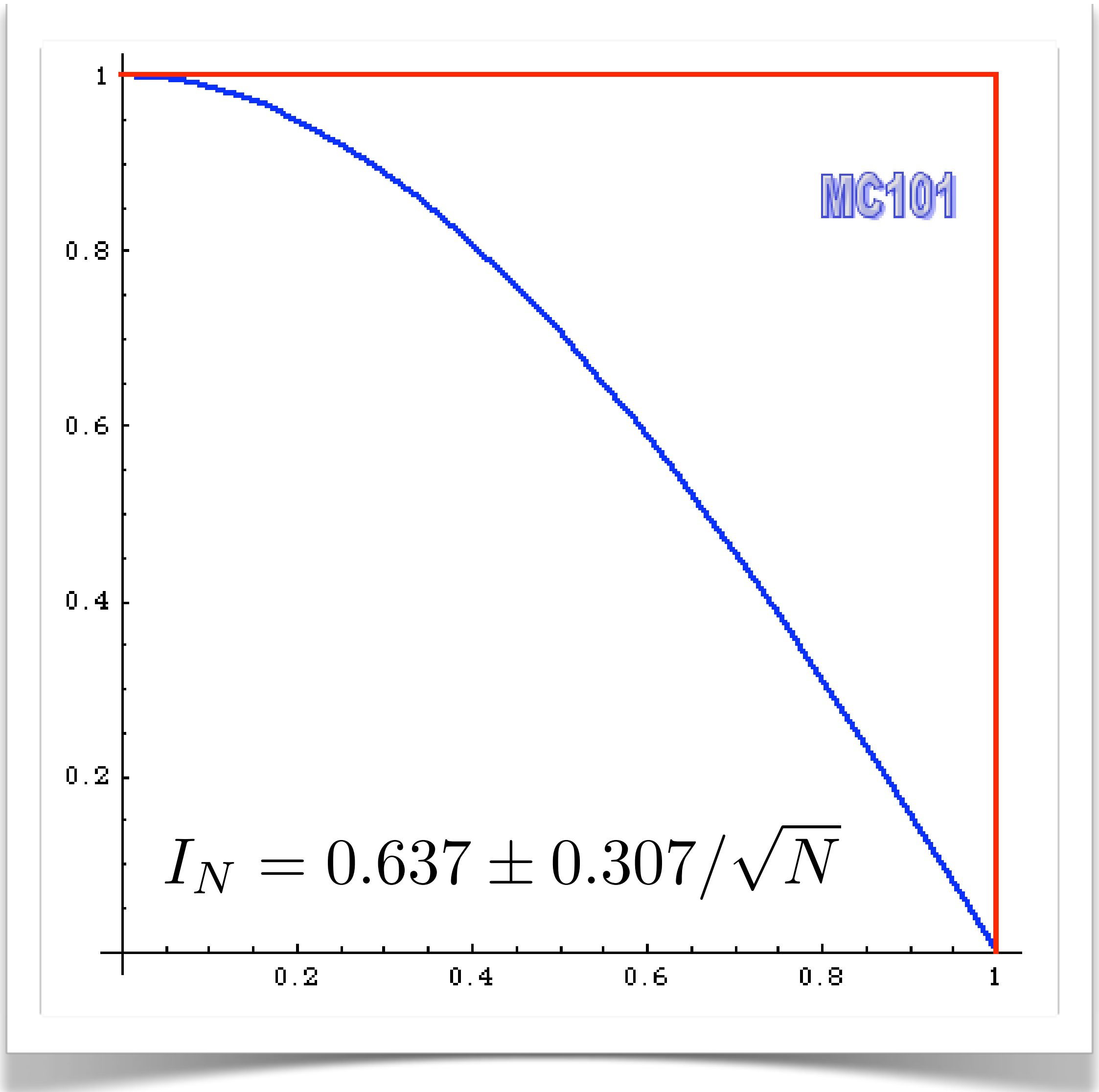
☞ Optimal/Ideal case: f(x) = Constant ⇒ $V_N = 0$

IMPORTANCE SAMPLING

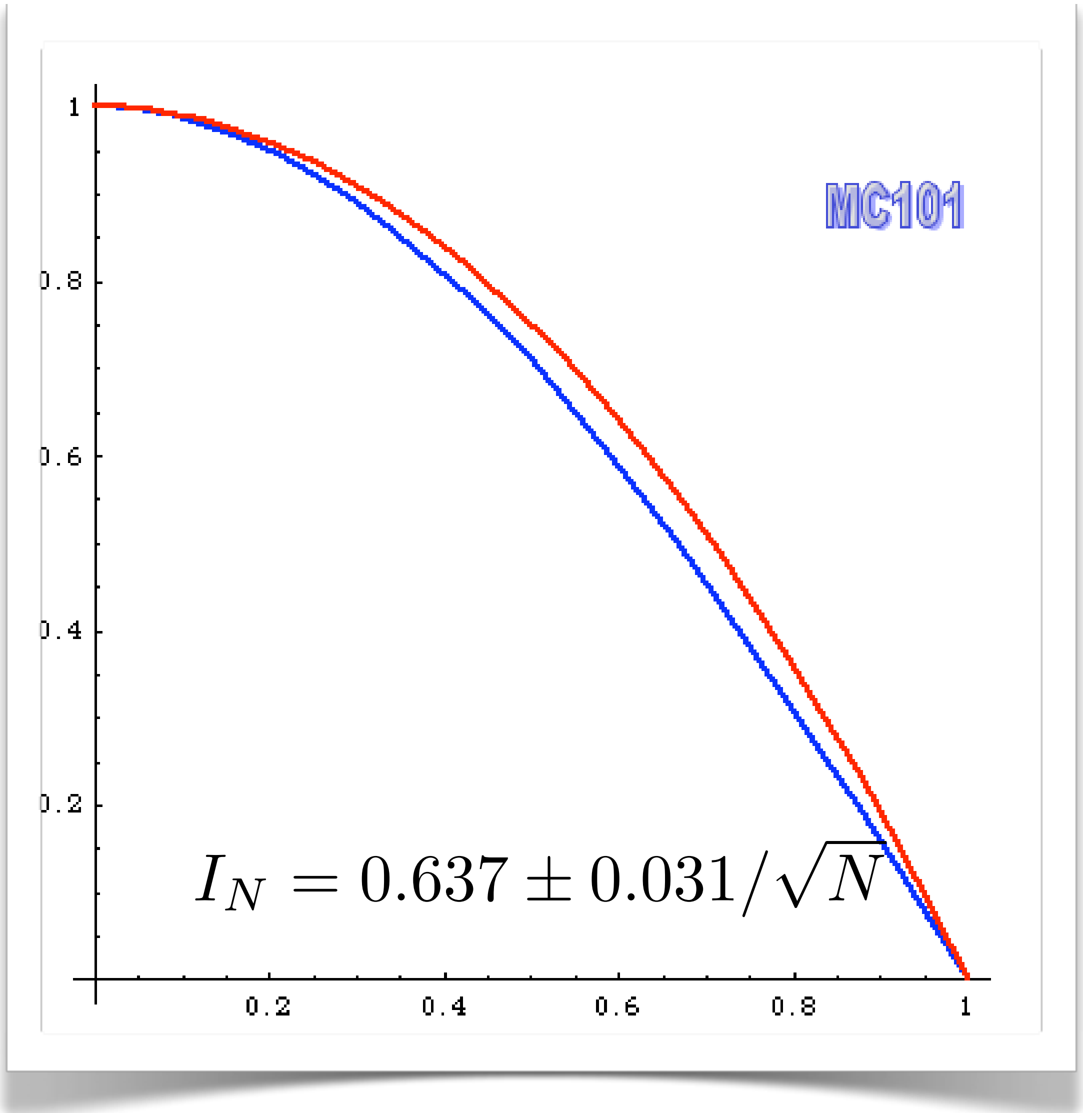


$$I = \int_0^1 dx \cos \frac{\pi}{2} x$$

IMPORTANCE SAMPLING

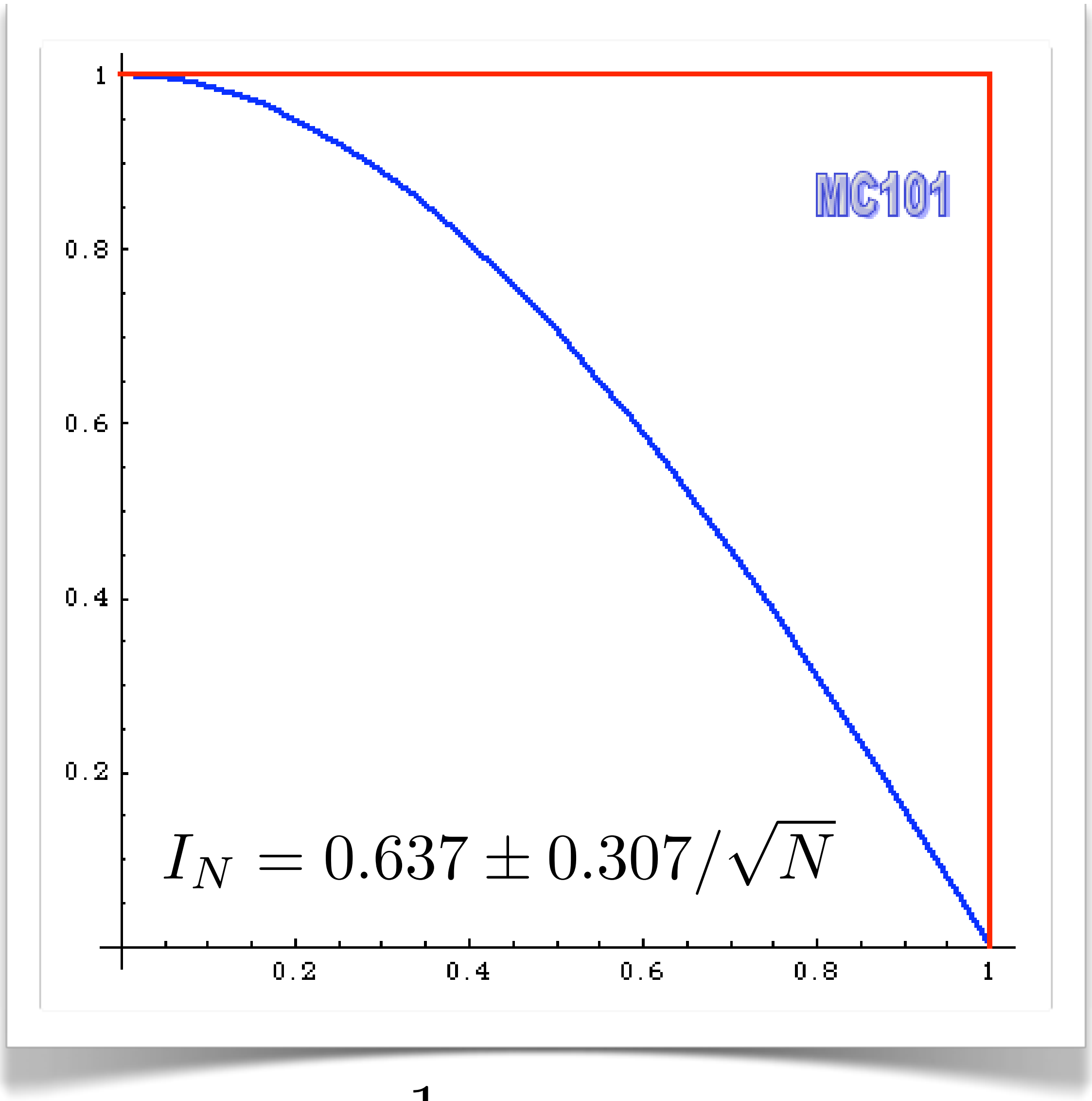


$$I = \int_0^1 dx \cos\frac{\pi}{2}x$$

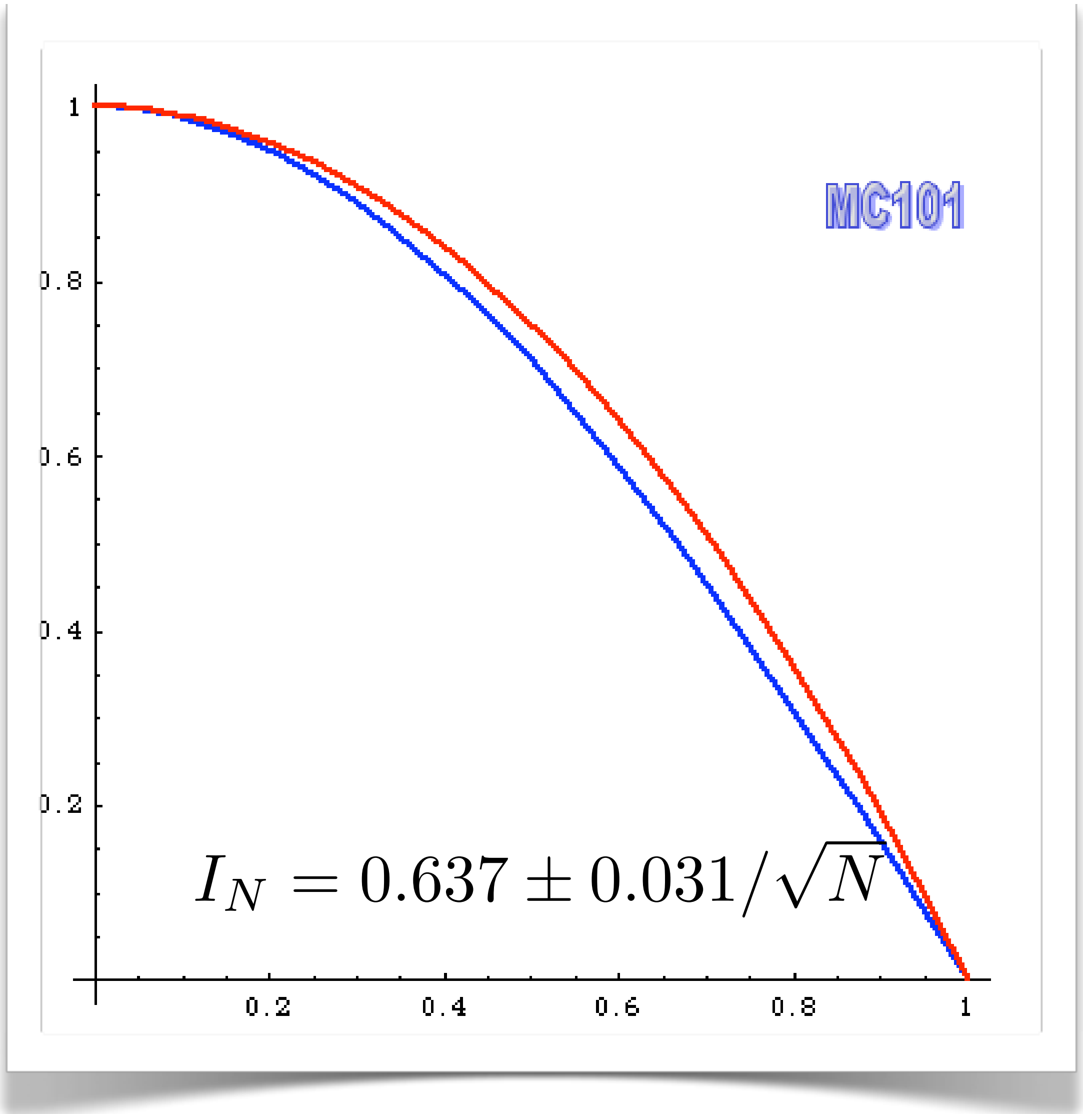


$$I = \int_0^1 dx(1-x^2)\frac{\cos\frac{\pi}{2}x}{1-x^2}$$

IMPORTANCE SAMPLING



$I_N = 0.637 \pm 0.307/\sqrt{N}$



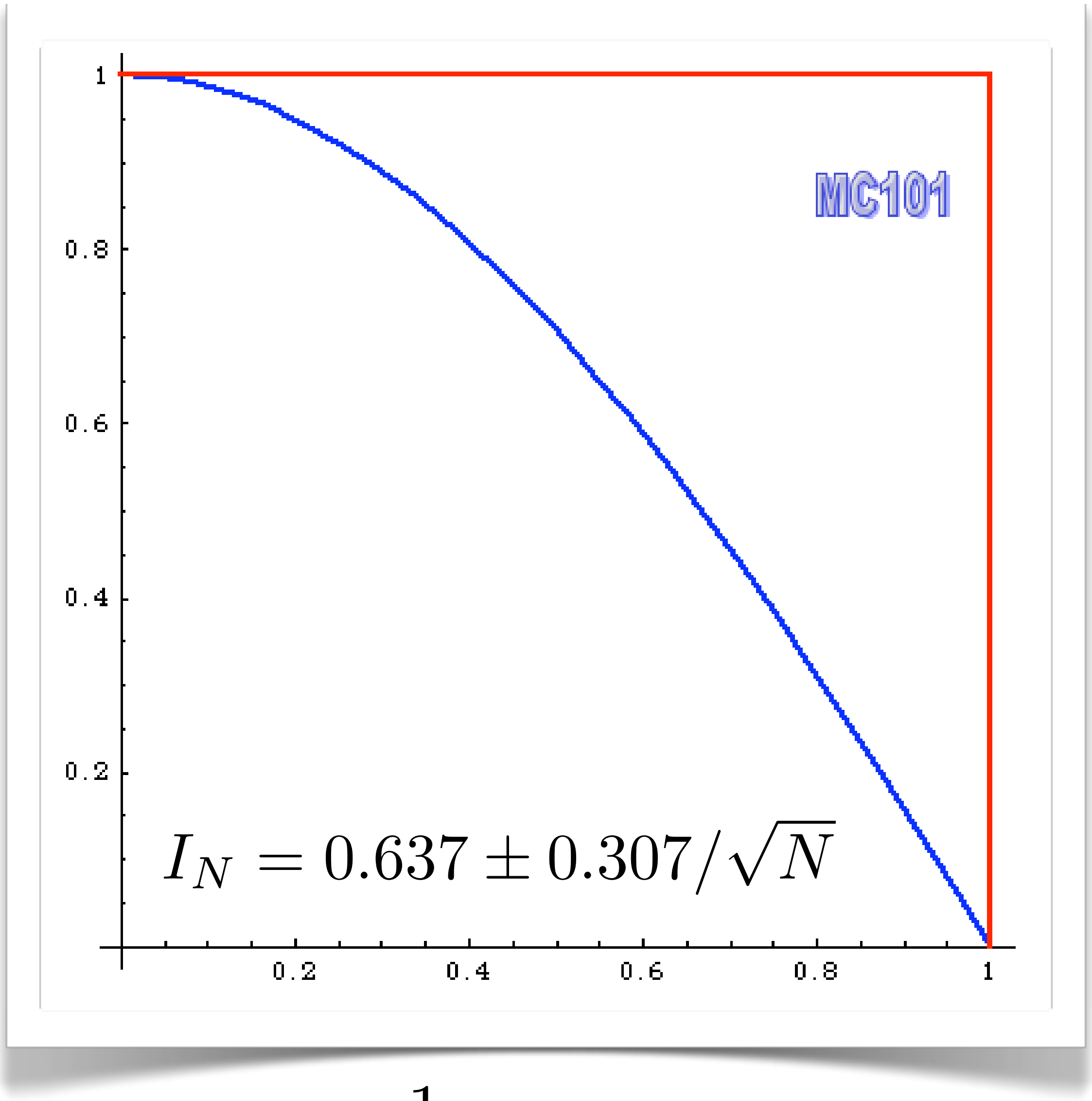
$I_N = 0.637 \pm 0.031/\sqrt{N}$

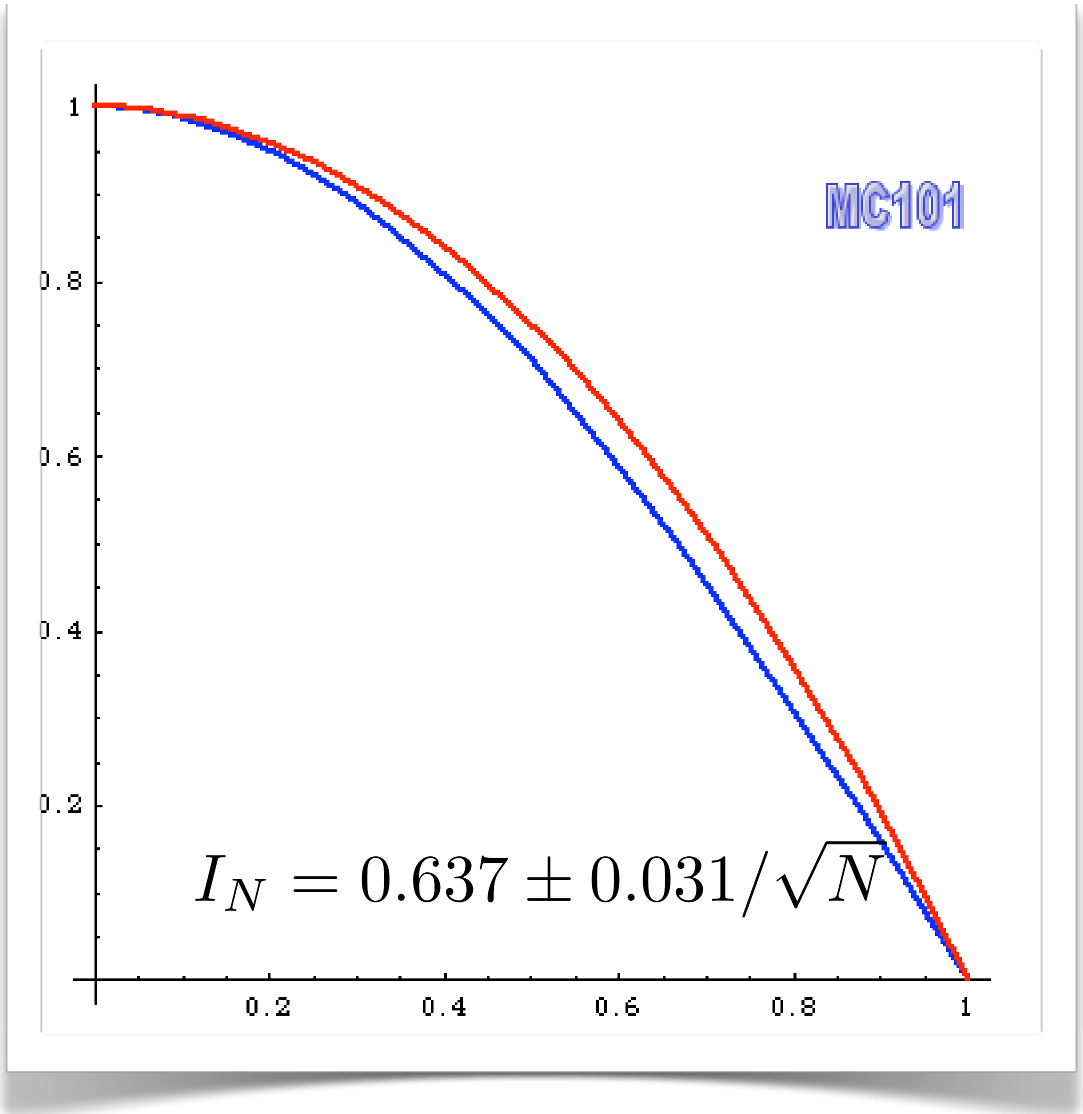
$$I = \int_0^1 dx \cos\frac{\pi}{2}x$$

$$I = \int_0^1 dx(1-x^2)\frac{\cos\frac{\pi}{2}x}{1-x^2}$$

$$= \int_{\xi_1}^{\xi_2} d\xi \frac{\cos\frac{\pi}{2}x[\xi]}{1-x[\xi]^2}$$

IMPORTANCE SAMPLING



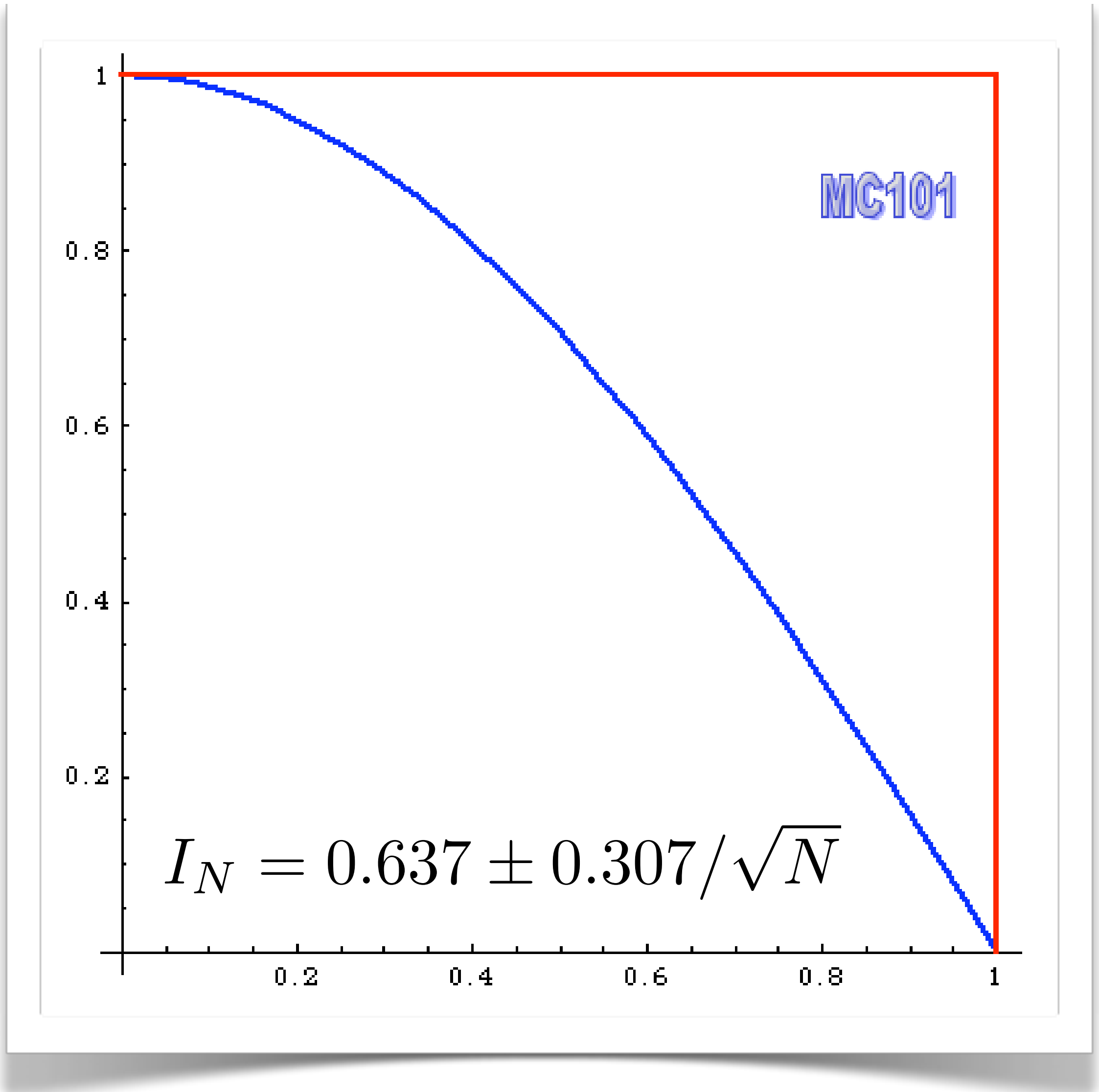


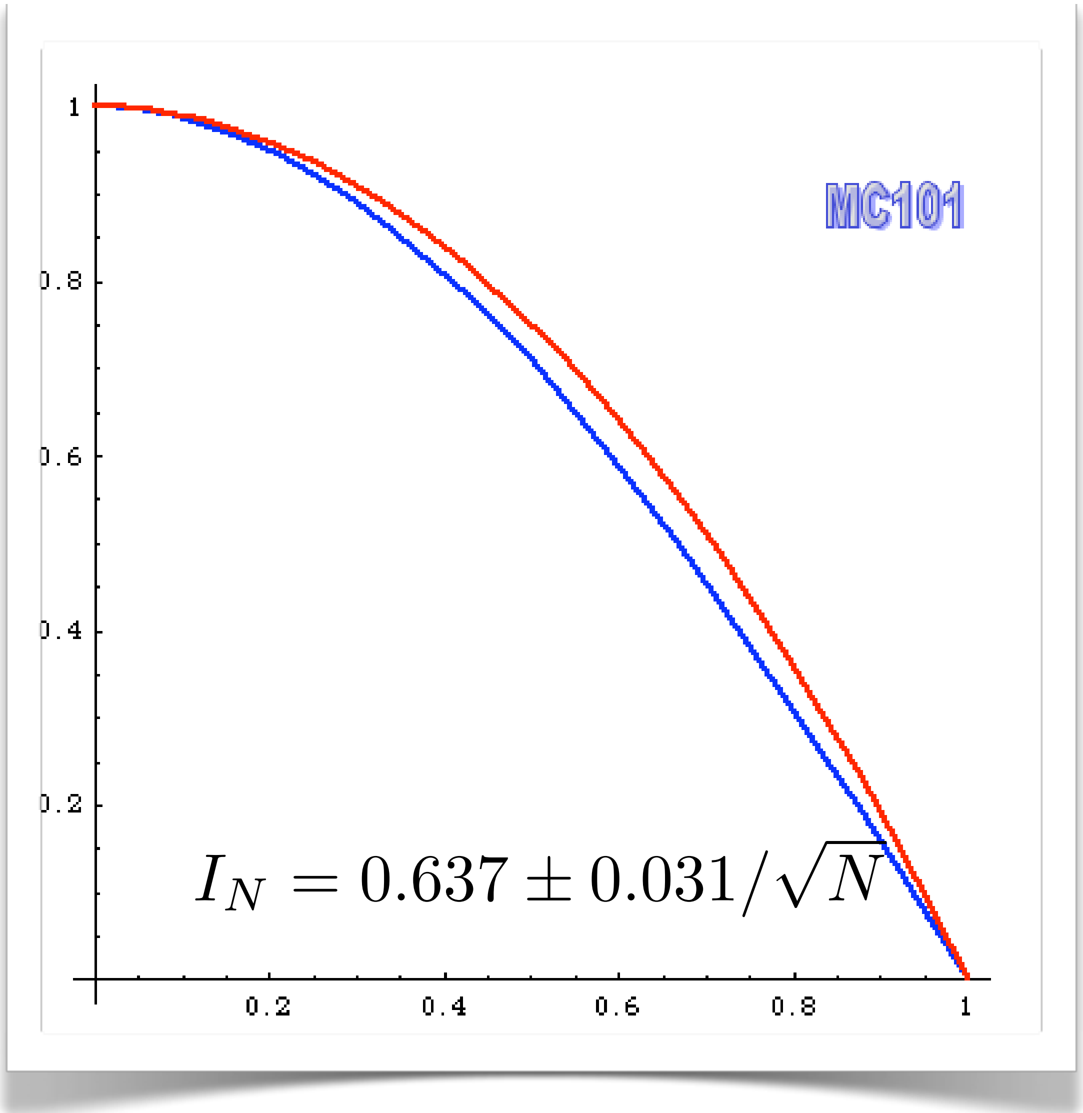
$$I = \int_0^1 dx \cos\frac{\pi}{2}x$$

$$I = \int_0^1 dx(1 - x^2)\frac{\cos\frac{\pi}{2}x}{1 - x^2}$$

$$= \int_{\xi_1}^{\xi_2} d\xi \frac{\cos\frac{\pi}{2}x[\xi]}{1-x[\xi]^2}$$

IMPORTANCE SAMPLING





$$I = \int_0^1 dx \cos\frac{\pi}{2}x$$

$$I = \int_0^1 dx(1-x^2)\frac{\cos\frac{\pi}{2}x}{1-x^2}$$

$$= \int_{\xi_1}^{\xi_2} d\xi \frac{\cos\frac{\pi}{2}x[\xi]}{1-x[\xi]^2} \Rightarrow \simeq 1$$

IMPORTANCE SAMPLING

IMPORTANCE SAMPLING

But... you need to know too much about f(x)!

IMPORTANCE SAMPLING

But... you need to know too much about f(x)!

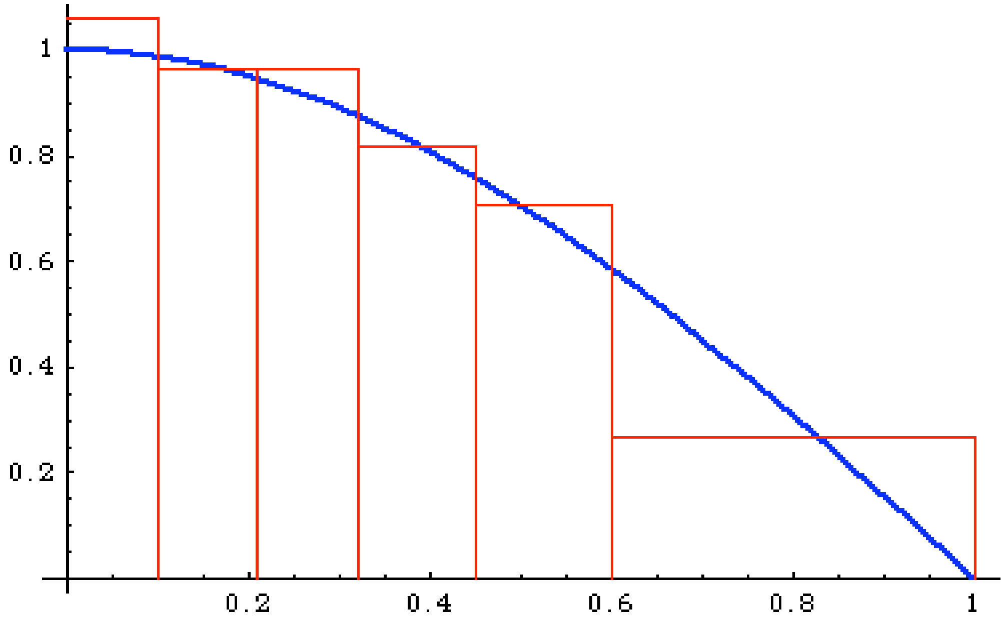
Idea: learn during the run and build a step-function approximation p(x) of f(x) → VEGAS

IMPORTANCE SAMPLING

But... you need to know too much about f(x)!

Idea: learn during the run and build a step-function approximation p(x) of f(x) ⇨ VEGAS

MC101

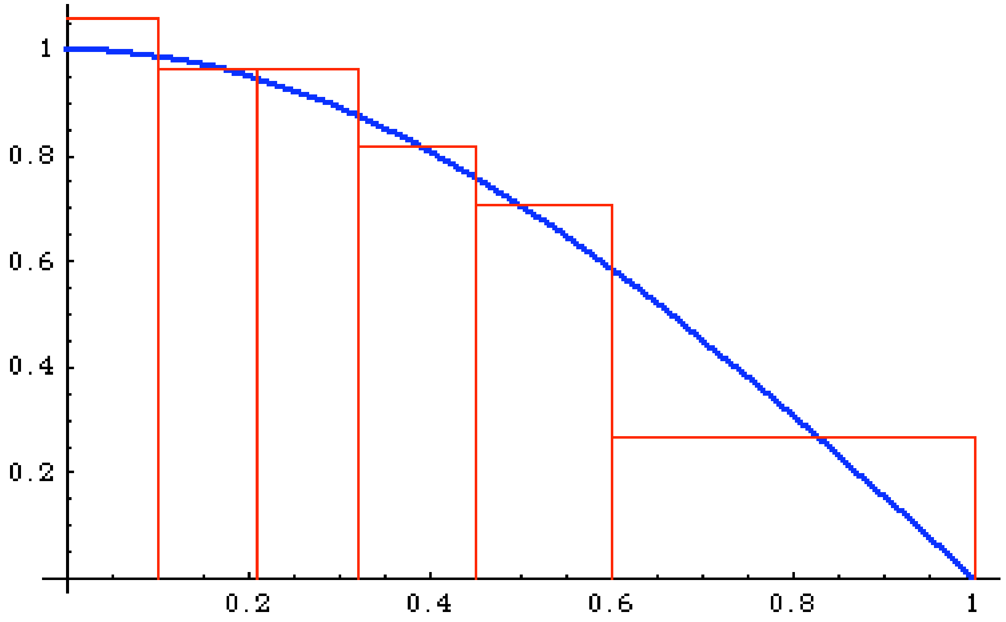


IMPORTANCE SAMPLING

But... you need to know too much about f(x)!

Idea: learn during the run and build a step-function approximation p(x) of f(x) ⇨ VEGAS

MC101



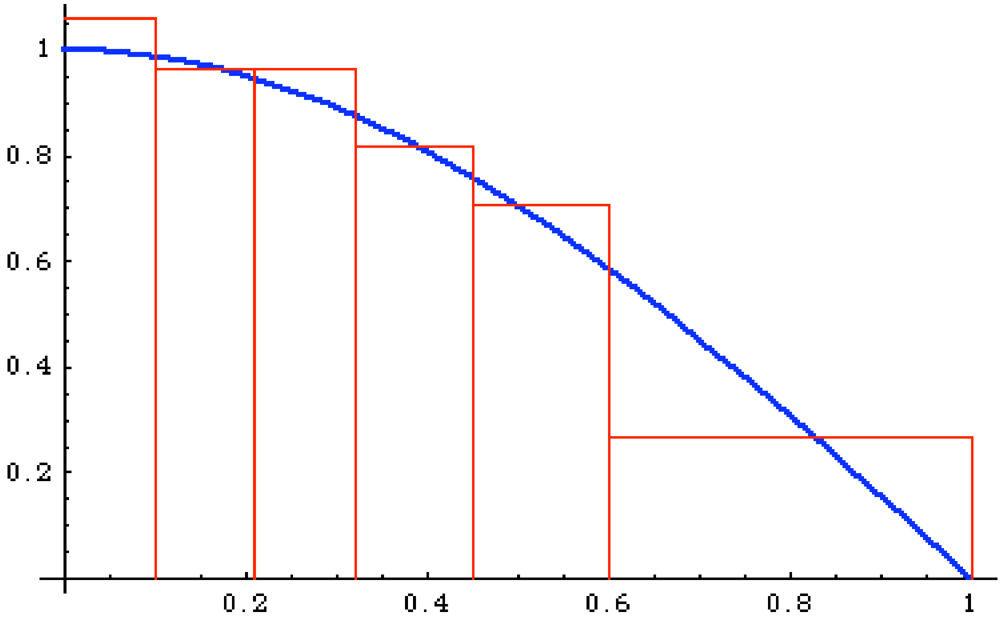
more bins where f(x) is large

IMPORTANCE SAMPLING

But... you need to know too much about f(x)!

Idea: learn during the run and build a step-function approximation p(x) of f(x) ⇒ VEGAS

MC101



more bins where f(x) is large

$$p(x) = \frac{1}{N_b \Delta x_i}, \quad x_i - \Delta x_i < x < x_i$$

IMPORTANCE SAMPLING

IMPORTANCE SAMPLING

can be generalized to n dimensions:

$$p(\vec{x}) = p(x) \bullet p(y) \bullet p(z) \ldots$$

IMPORTANCE SAMPLING

can be generalized to n dimensions:

$$p(\vec{x}) = p(x) \bullet p(y) \bullet p(z) \ldots$$

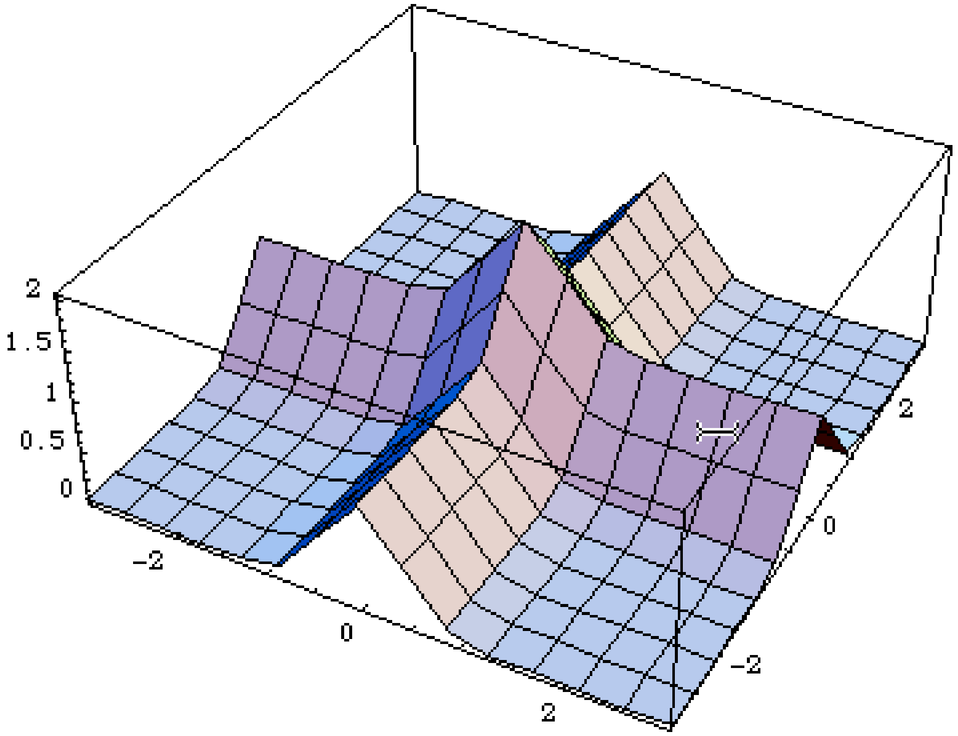
but the peaks of $f(\vec{x})$ need to be "aligned" to the axis!

IMPORTANCE SAMPLING

can be generalized to n dimensions:

$$p(\vec{x}) = p(x) \bullet p(y) \bullet p(z) \ldots$$

but the peaks of $f(\vec{x})$ need to be "aligned" to the axis!



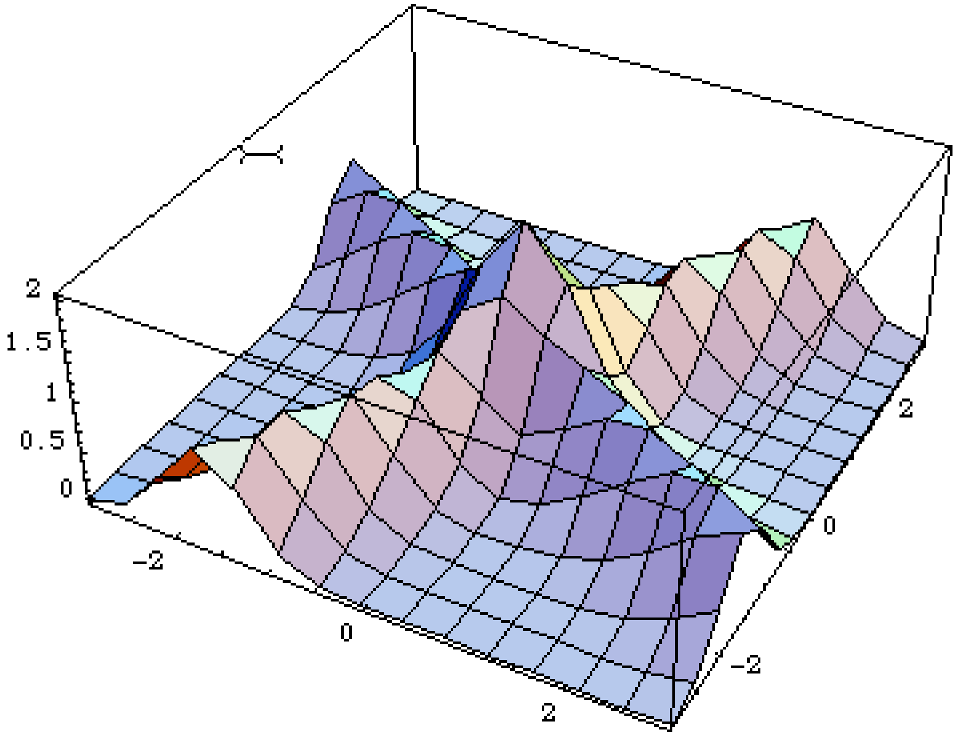
This is ok...

IMPORTANCE SAMPLING

can be generalized to n dimensions:

$$p(\vec{x}) = p(x) \bullet p(y) \bullet p(z) \ldots$$

but the peaks of $f(\vec{x})$ need to be "aligned" to the axis!



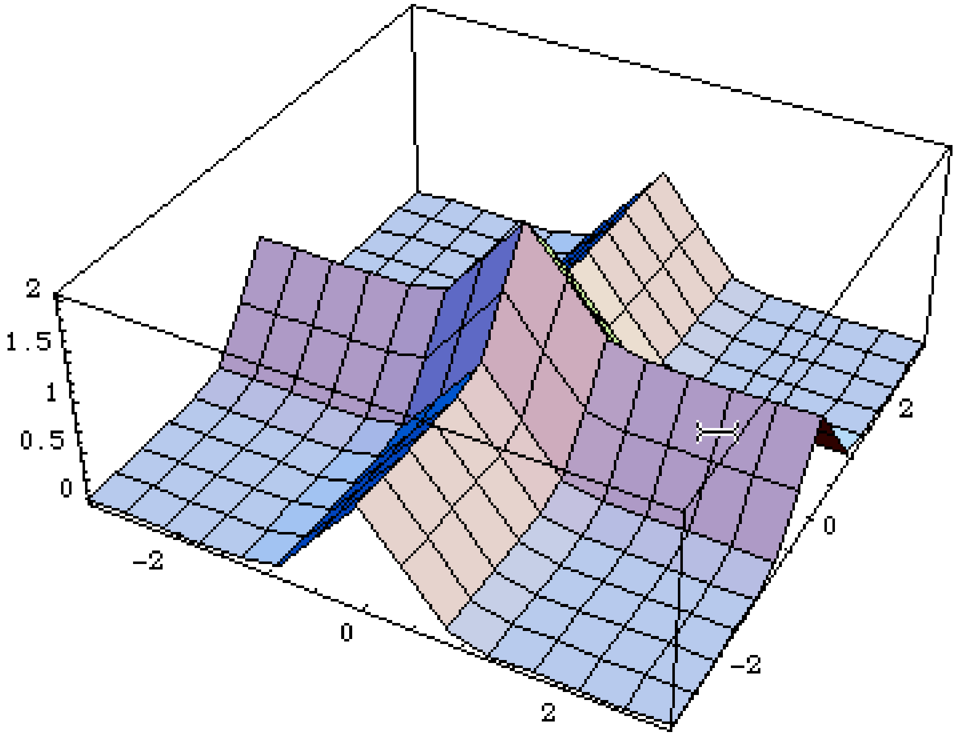
This is not ok...

IMPORTANCE SAMPLING

can be generalized to n dimensions:

$$p(\vec{x}) = p(x)\bullet p(y)\bullet p(z)\ldots$$

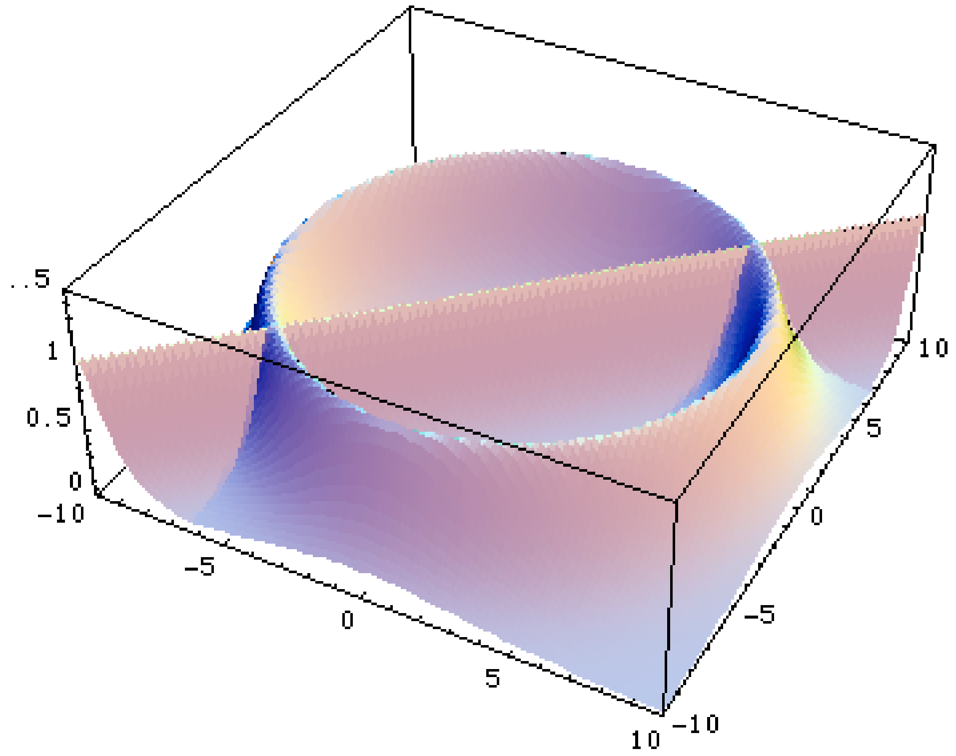
but the peaks of $f(\vec{x})$ need to be "aligned" to the axis!



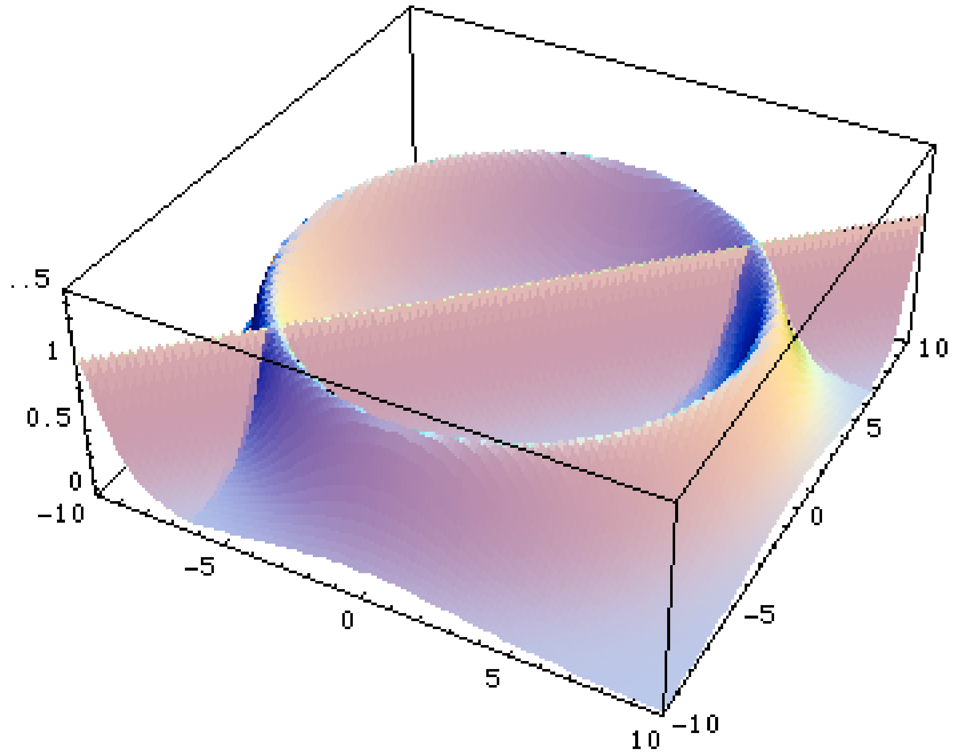
but it is sufficient to make
a change of variables!

MULTI-CHANNEL

MULTI-CHANNEL

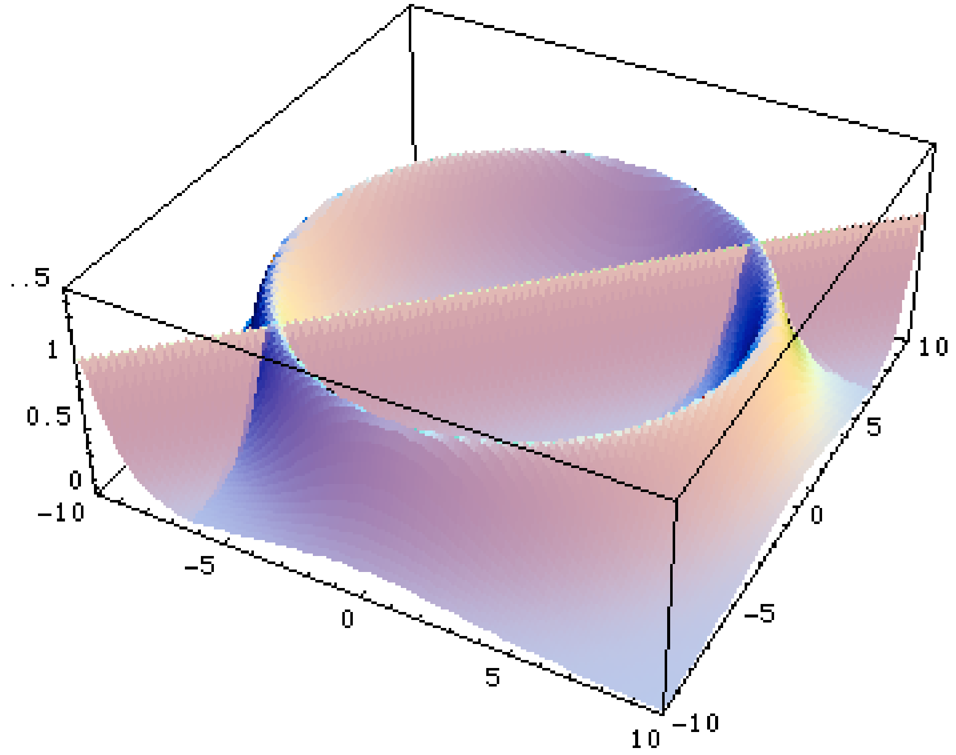


MULTI-CHANNEL



In this case there is no unique tranformation: Vegas is bound to fail!

MULTI-CHANNEL



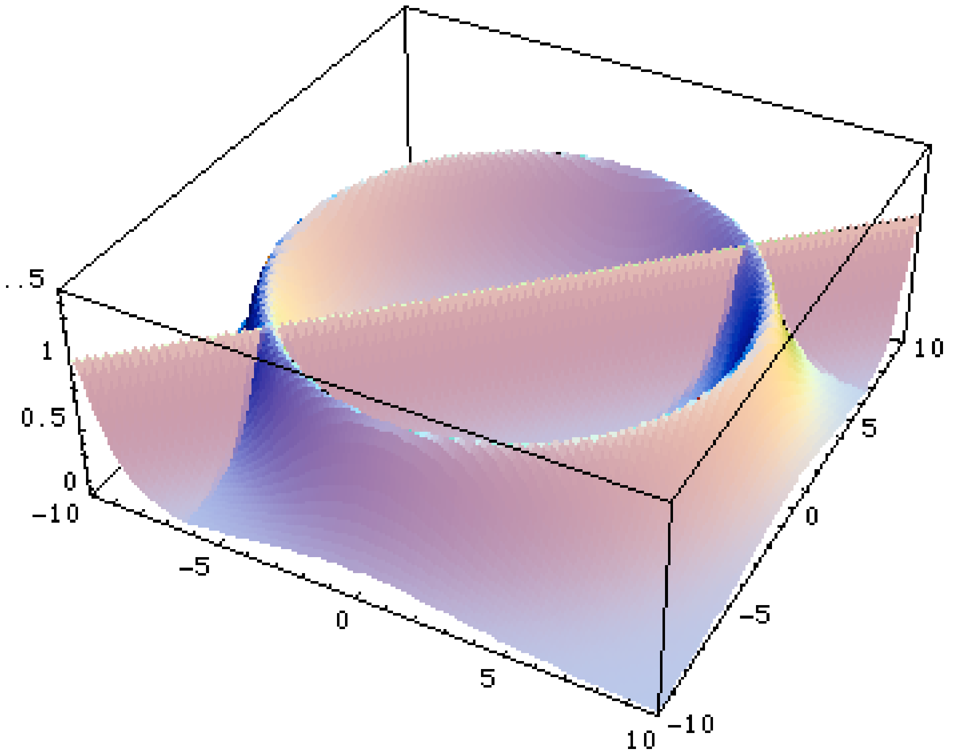
In this case there is no unique tranformation: Vegas is bound to fail!

Solution: use different transformations= channels

$$p(x) = \sum_{i=1}^{n} \alpha_i p_i(x) \qquad \text{with} \qquad \sum_{i=1}^{n} \alpha_i = 1$$

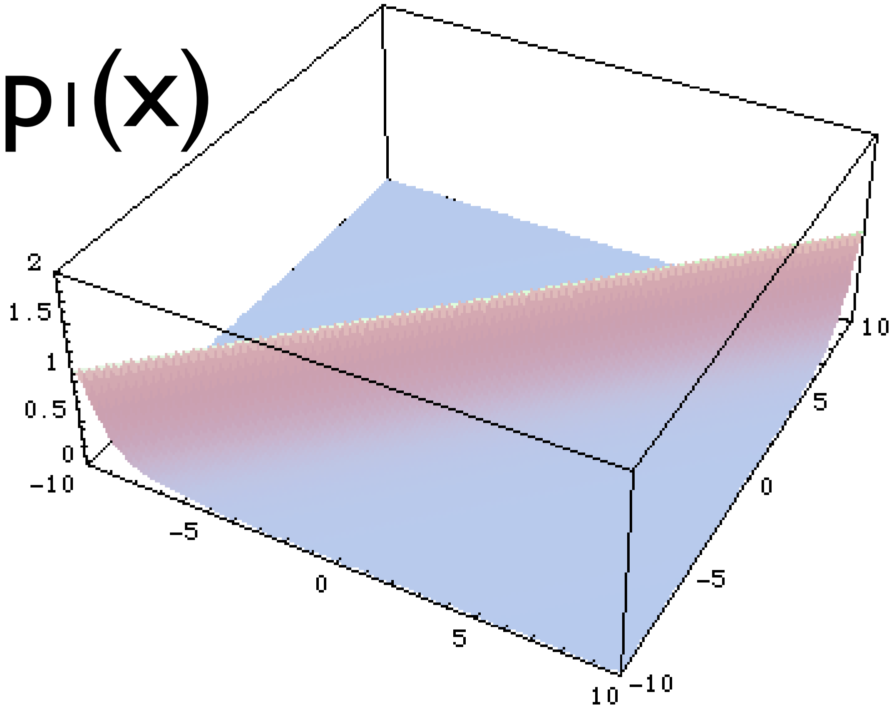
with each $p_i(x)$ taking care of one "peak" at the time

MULTI-CHANNEL

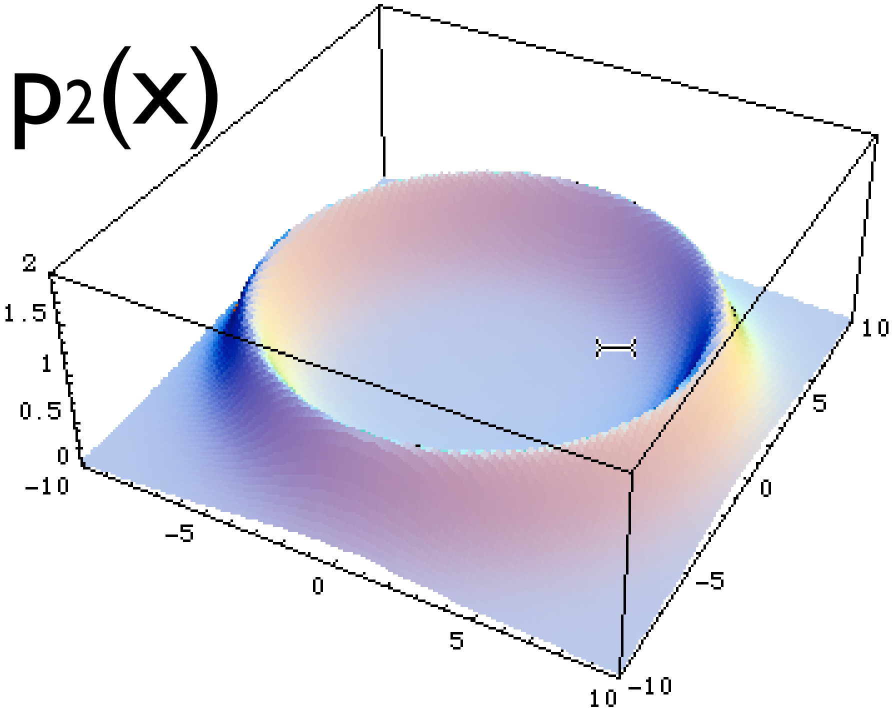


In this case there is no unique tranformation: Vegas is bound to fail!

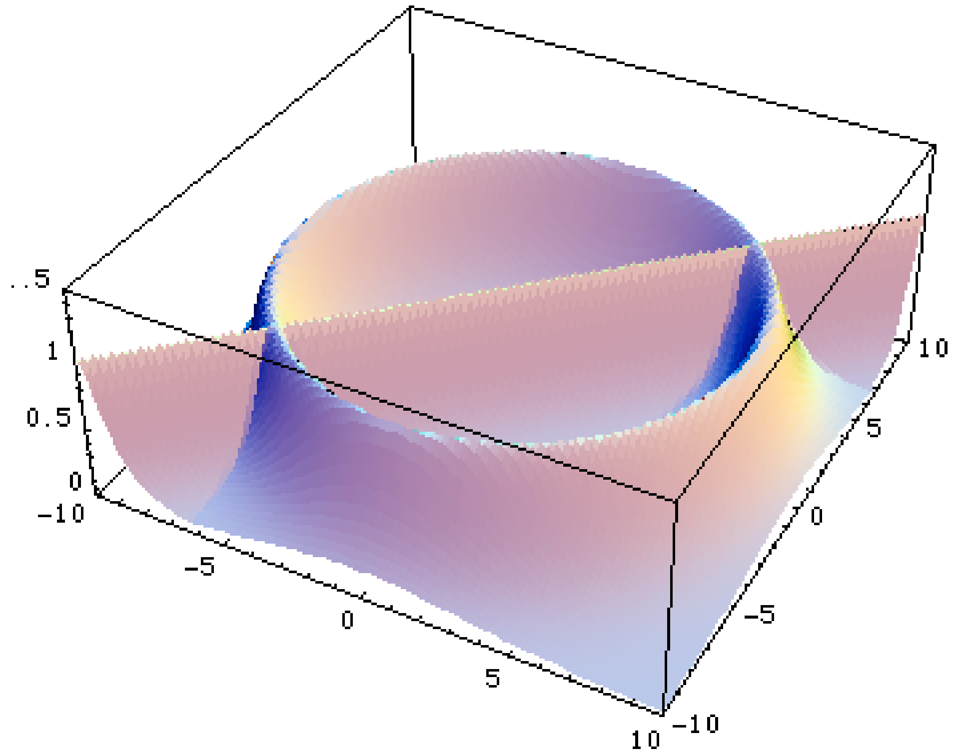
$p_1(x)$



$p_2(x)$



MULTI-CHANNEL



In this case there is no unique tranformation: Vegas is bound to fail!

But if you know where the peaks are (=in which variables) we can use different transformations= channels:

$$p(x) = \sum_{i=1}^{n} \alpha_i p_i(x) \qquad \textbf{with} \qquad \sum_{i=1}^{n} \alpha_i = 1$$

$$I = \int f(x)dx = \sum_{i=1}^{n} \alpha_i \int \frac{f(x)}{p(x)} p_i(x)dx$$

MULTI-CHANNEL

- Advantages
  - The integral does not depend on the α_i but the variance does and can be minimised by a careful choice
- Drawbacks
  - Need to calculate all g_i values for each point
  - Each phase space channel must be invertible
  - N coupled equations for α_i so it might only work for small number of channels

MULTI-CHANNEL

- Advantages
  - The integral does not depend on the α_i but the variance does and can be minimised by a careful choice
- Drawbacks
  - Need to calculate all g_i values for each point
  - Each phase space channel must be invertible
  - N coupled equations for α_i so it might only work for small number of channels

Very popular method!

MULTI-CHANNEL BASED ON SINGLE DIAGRAMS

Consider the integration of an amplitude |M|^2 at treel level which lots of diagrams contribute to. If there were a basis of functions,

$$f = \sum_{i=1}^{n} f_i \quad \text{with} \quad f_i \geq 0\,, \quad \forall\ i\,,$$

such that:

1. we know how to integrate each one of them,
2. they describe all possible peaks,

then the problem would be solved:

$$I = \int d\vec{\Phi} f(\vec{\Phi}) = \sum_{i=1}^{n} \int d\vec{\Phi}\, g_i(\vec{\Phi})\, \frac{f_i(\vec{\Phi})}{g_i(\vec{\Phi})} = \sum_{i=1}^{n} I_i\,,$$

MULTI-CHANNEL BASED ON SINGLE DIAGRAMS

Consider the integration of an amplitude |M|^2 at treel level which lots of diagrams contribute to. If there were a basis of functions,

$$f = \sum_{i=1}^{n} f_i \quad \text{with} \quad f_i \geq 0\,, \quad \forall\ i\,,$$

such that:

1. we know how to integrate each one of them,
2. they describe all possible peaks,

then the problem would be solved:

$$I = \int d\vec{\Phi} f(\vec{\Phi}) = \sum_{i=1}^{n} \int d\vec{\Phi}\, g_i(\vec{\Phi})\, \frac{f_i(\vec{\Phi})}{g_i(\vec{\Phi})} = \sum_{i=1}^{n} I_i\,,$$

Does such a basis exist?

MULTI-CHANNEL BASED ON SINGLE DIAGRAMS

Consider the integration of an amplitude |M|^2 at treel level which lots of diagrams contribute to. If there were a basis of functions,

$$f = \sum_{i=1}^{n} f_i \quad \text{with} \quad f_i \geq 0\,, \quad \forall \;\; i\,,$$

such that:

1. we know how to integrate each one of them,
2. they describe all possible peaks,

then the problem would be solved:

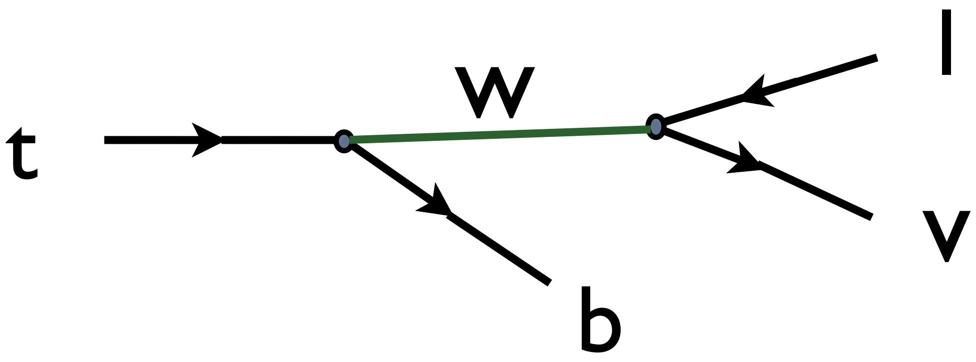
$$I = \int d\vec{\Phi} f(\vec{\Phi}) = \sum_{i=1}^{n} \int d\vec{\Phi}\, g_i(\vec{\Phi})\, \frac{f_i(\vec{\Phi})}{g_i(\vec{\Phi})} = \sum_{i=1}^{n} I_i\,,$$

Does such a basis exist? YES! $f_i = \frac{|A_i|^2}{\sum_i |A_i|^2} |A_{\rm tot}|^2$

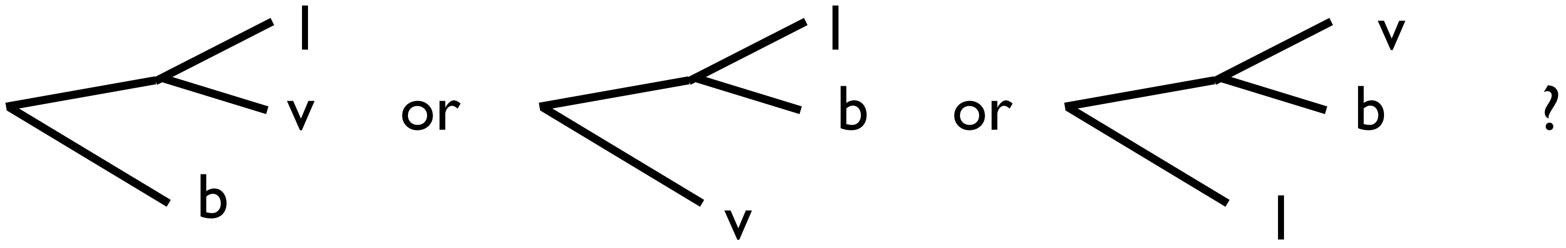
MULTI-CHANNEL : MADGRAPH

- Key Idea
  - Any single diagram is "easy" to integrate
  - Divide integration into pieces, based on diagrams
- Get N independent integrals
  - Errors add in quadrature so no extra cost
  - No need to calculate "weight" function from other channels.
  - Can optimize # of points for each one independently
  - Parallel in nature
- What about interference?
  - Never creates "new" peaks, so we're OK!

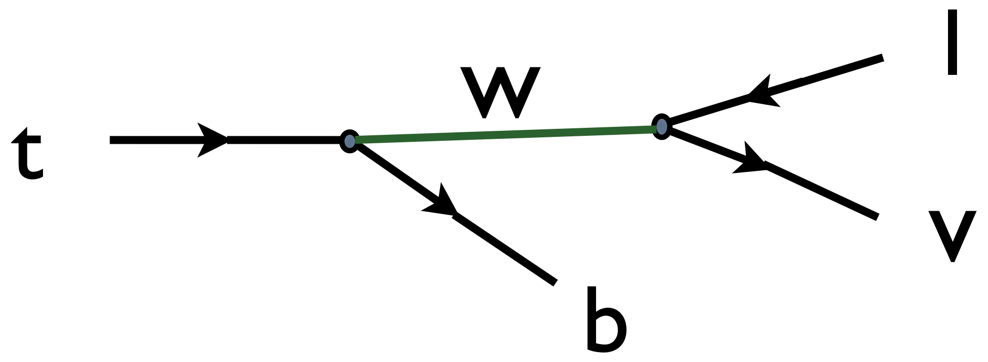
EXERCISE: TOP DECAY

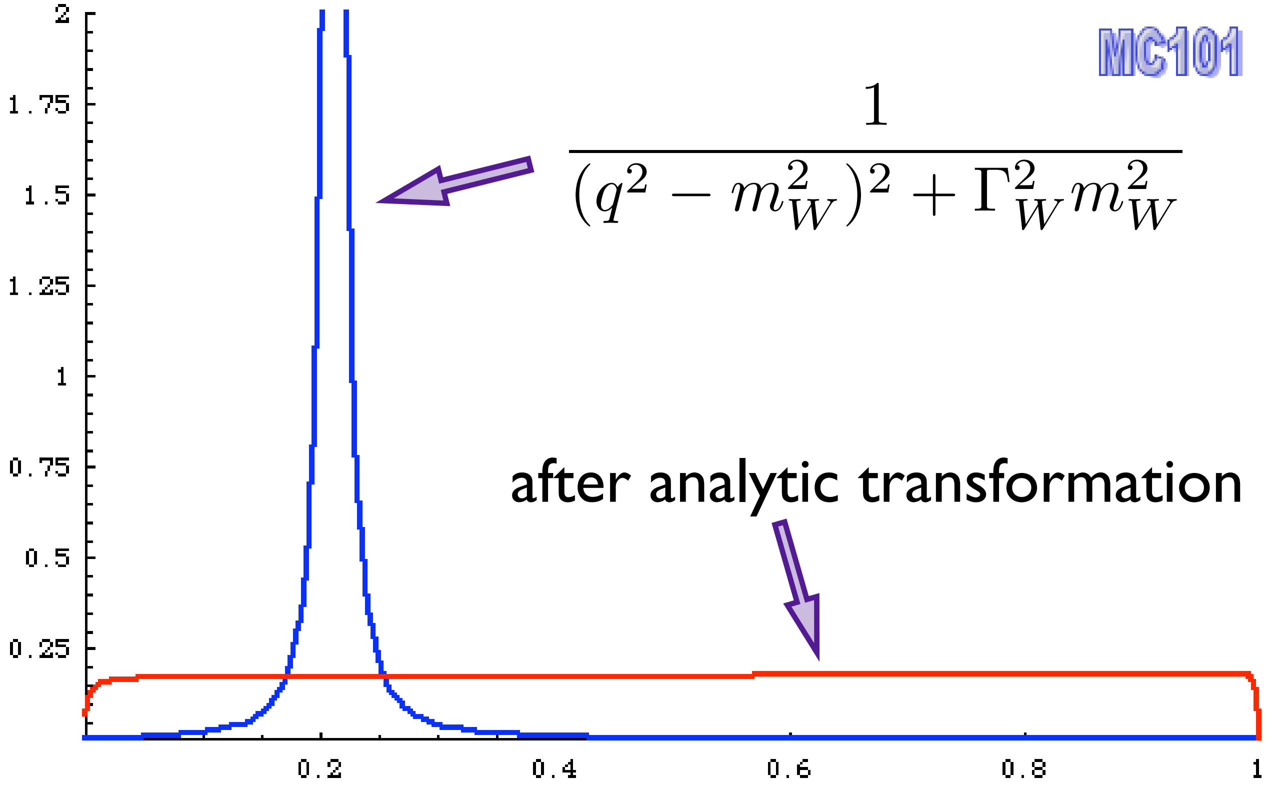


- Easy but non-trivial
- Breit-Wigner peak $\frac{1}{(q^2 - m_W^2)^2 + \Gamma_W^2 m_W^2}$ to be "flattened":
- Choose the right "channel" for the phase space:



Exercise: Top decay





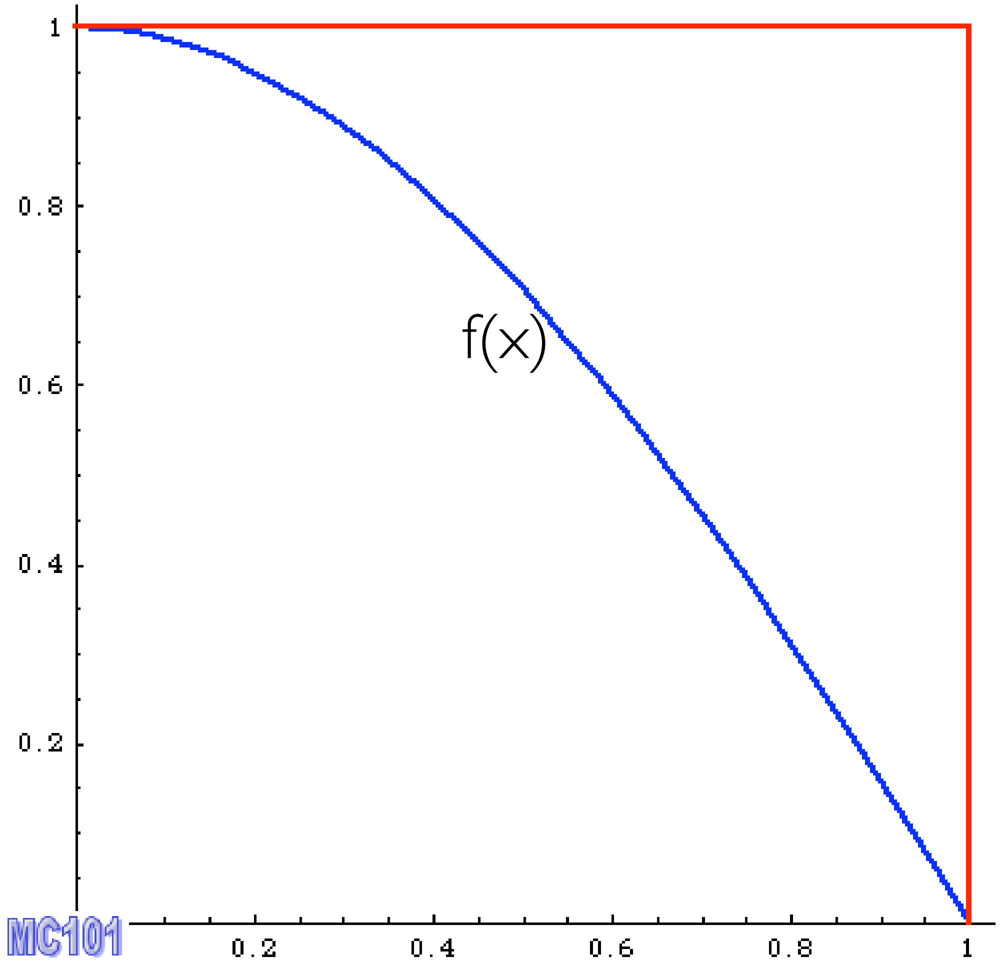
$$\frac{1}{(q^2 - m_W^2)^2 + \Gamma_W^2 m_W^2}$$

after analytic transformation

EVENT GENERATION

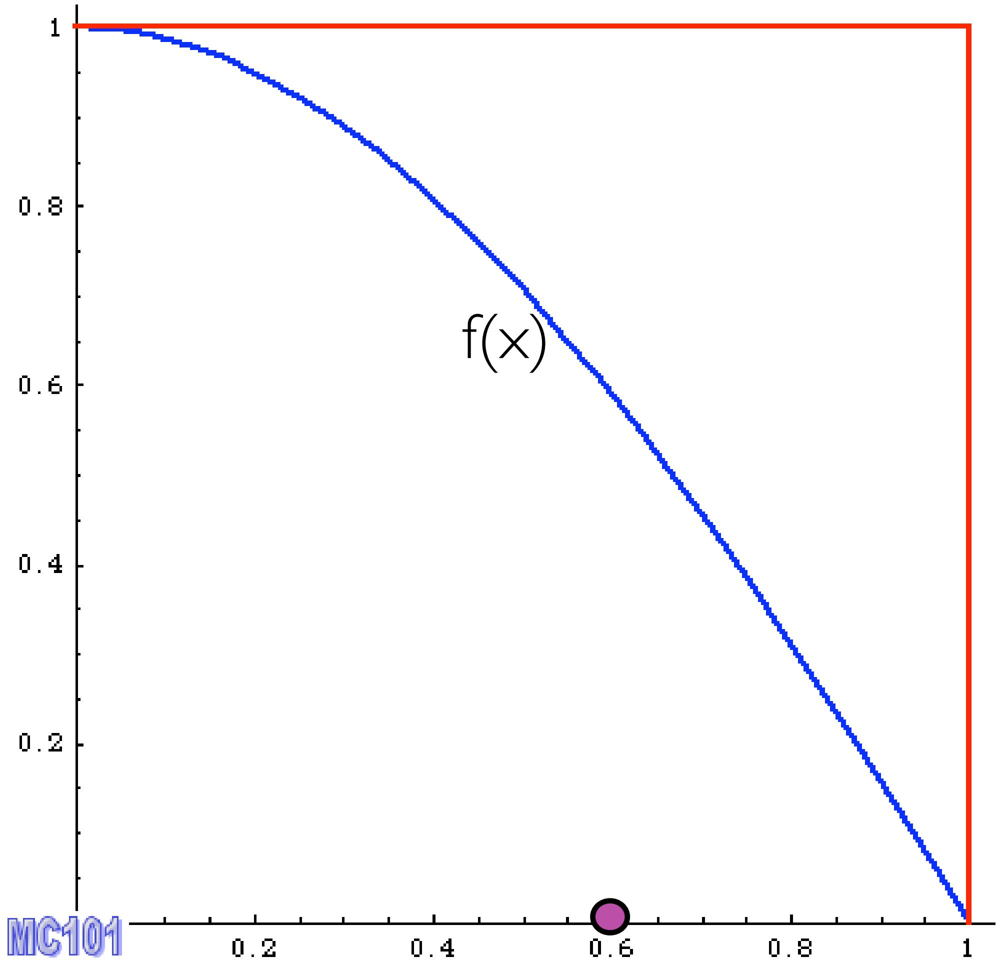
- Every phase-space point computed in this way, can be seen as an event (=collision) in a detector
- However, they still carry the "weight" of the matrix elements:
  - ▻ events with large weights where the cross section is large
  - ▻ events with small weights where the cross section is small
- In nature, the events don't carry a weight:
  - ▻ more events where the cross section is large
  - ▻ less events where the cross section is small
- How to go from weighted events to unweighted events?

EVENT GENERATION



Alternative way

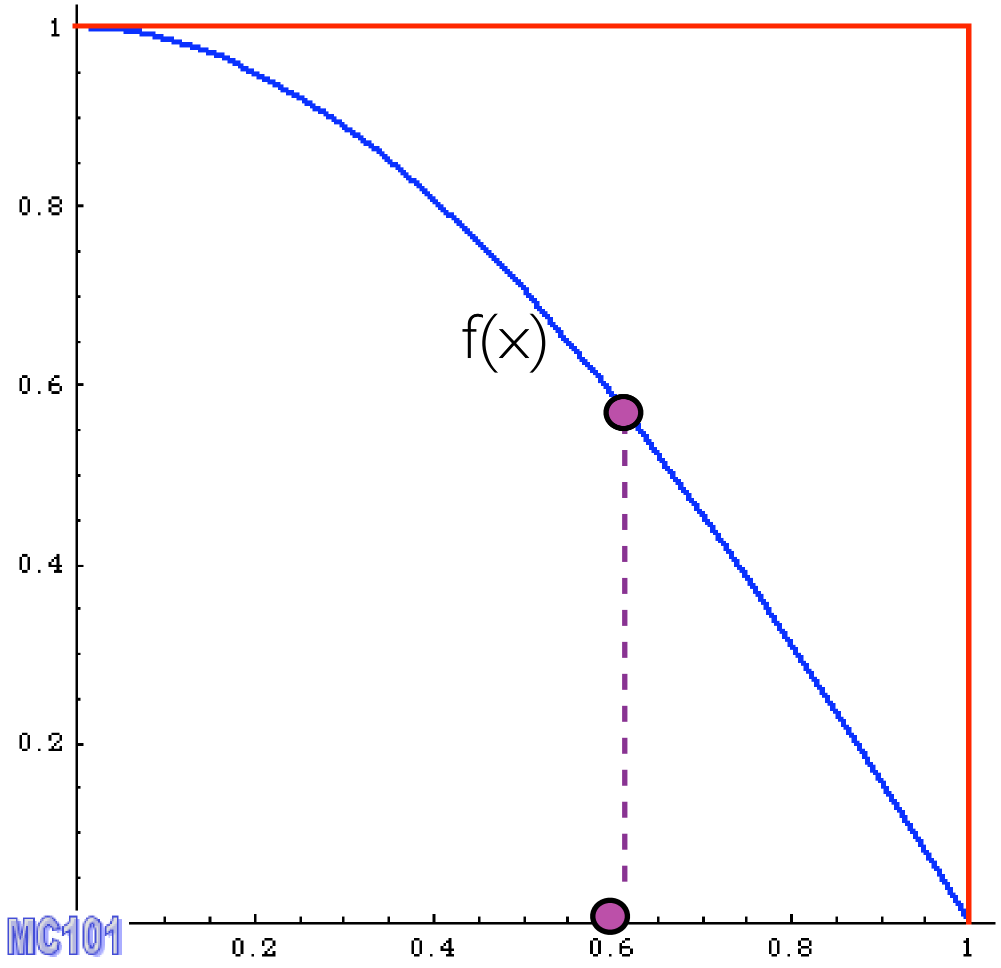
EVENT GENERATION



Alternative way

1. (randomly) pick x

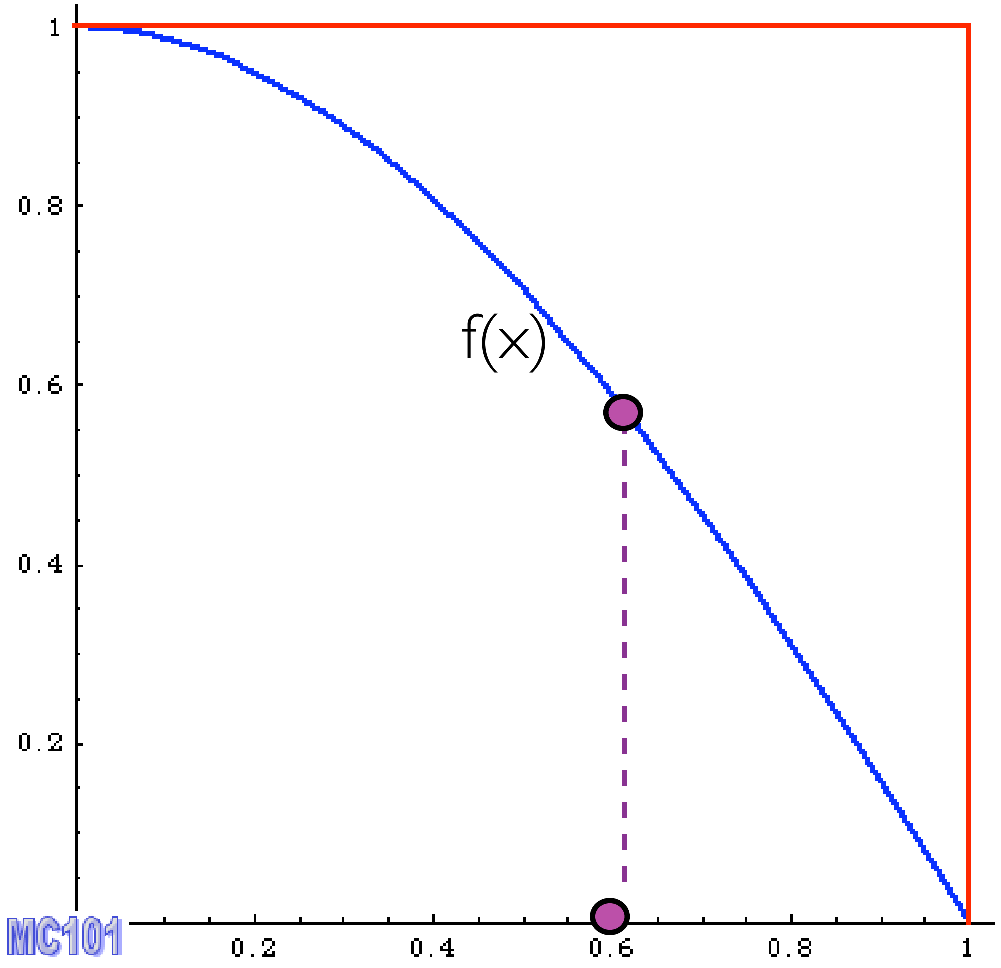
EVENT GENERATION



Alternative way

1. (randomly) pick x
2. calculate f(x)

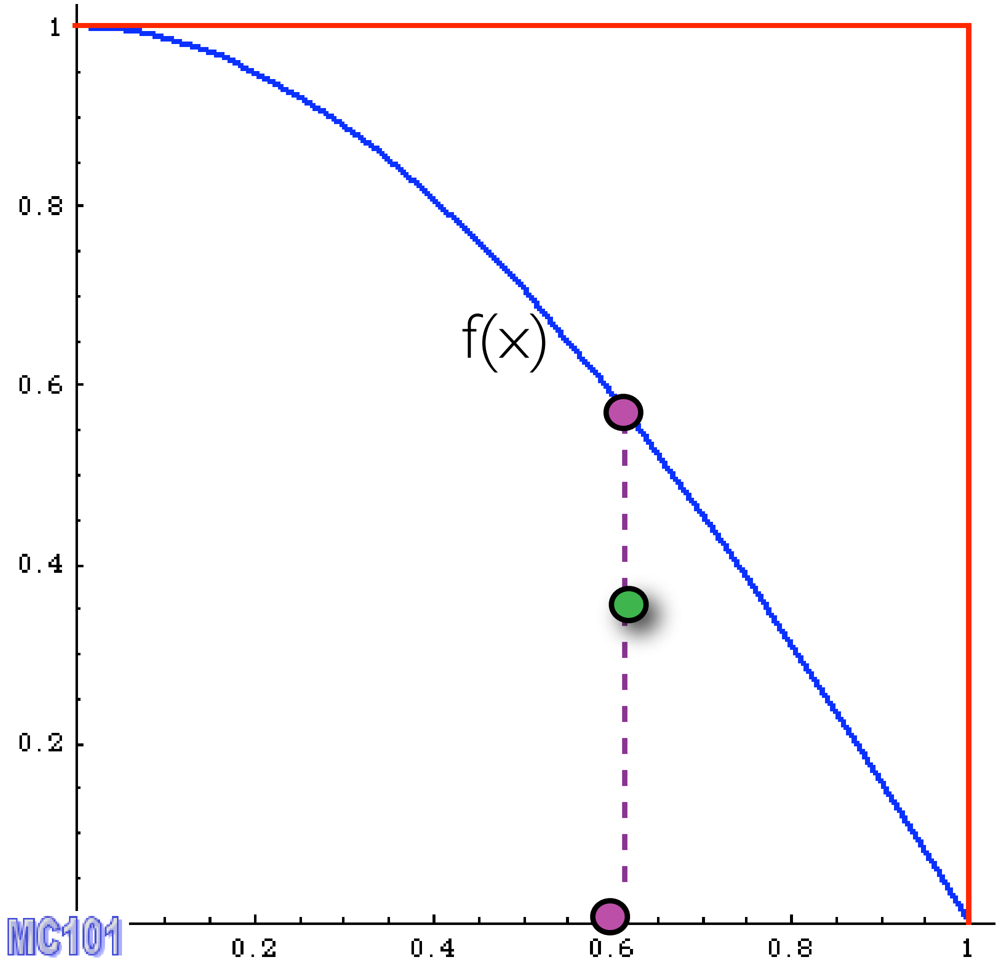
EVENT GENERATION



Alternative way

1. (randomly) pick x
2. calculate f(x)
3. (randomly) pick 0<y<fmax

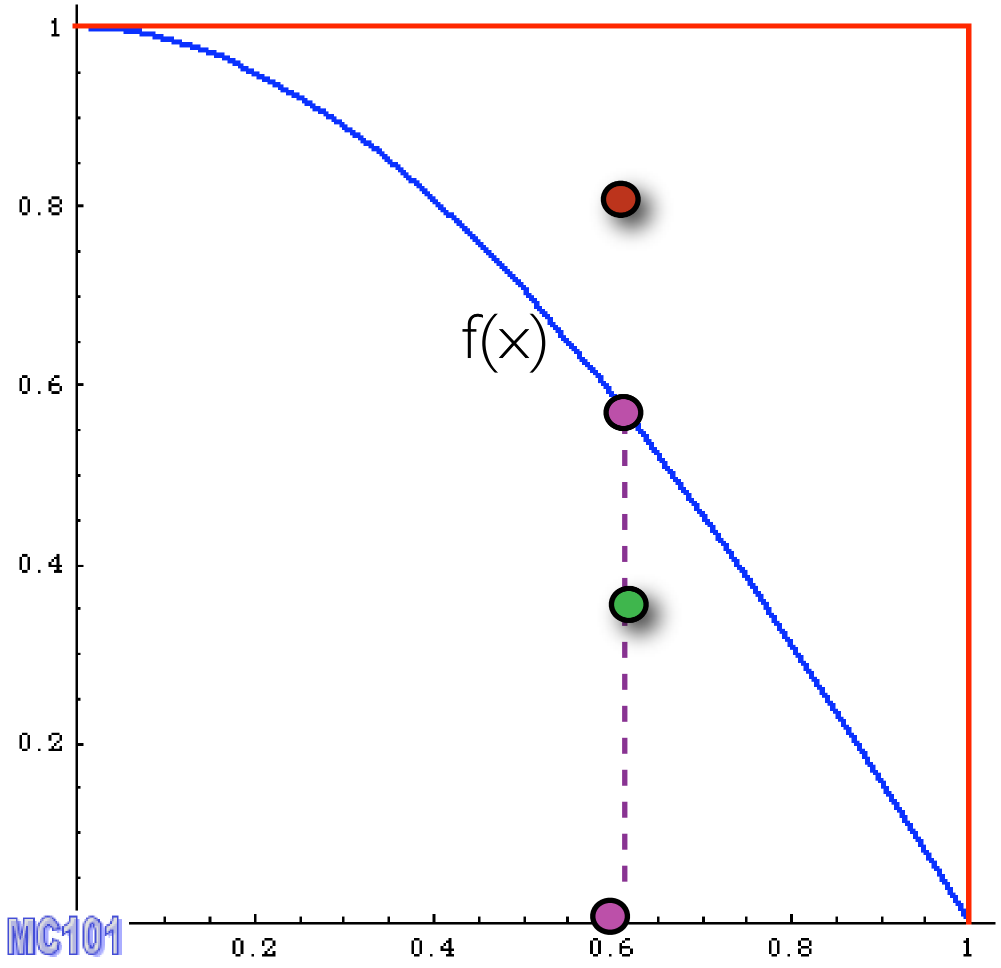
EVENT GENERATION



Alternative way

1. (randomly) pick x
2. calculate f(x)
3. (randomly) pick 0<y<fmax
4. Compare:
   if f(x)>y accept event,

EVENT GENERATION

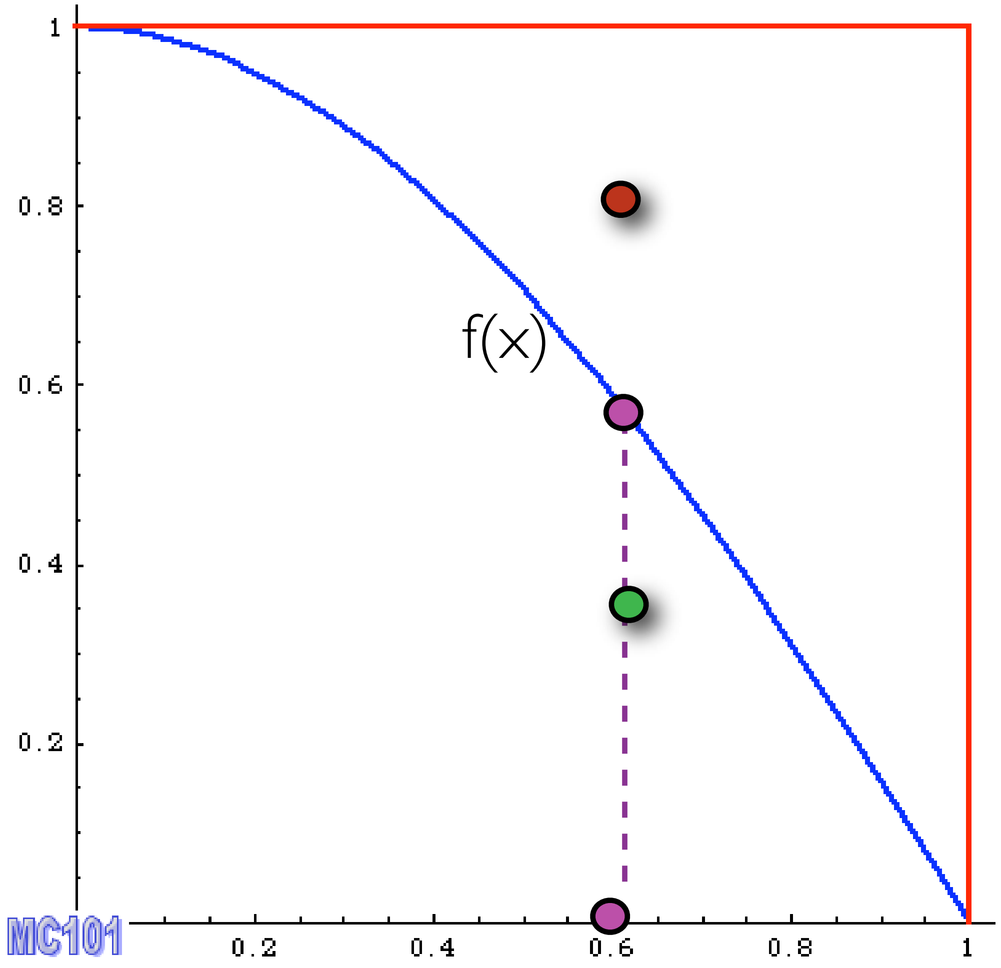


Alternative way

1. (randomly) pick x
2. calculate f(x)
3. (randomly) pick 0<y<fmax
4. Compare:
   if f(x)>y accept event,

   else reject it.

EVENT GENERATION



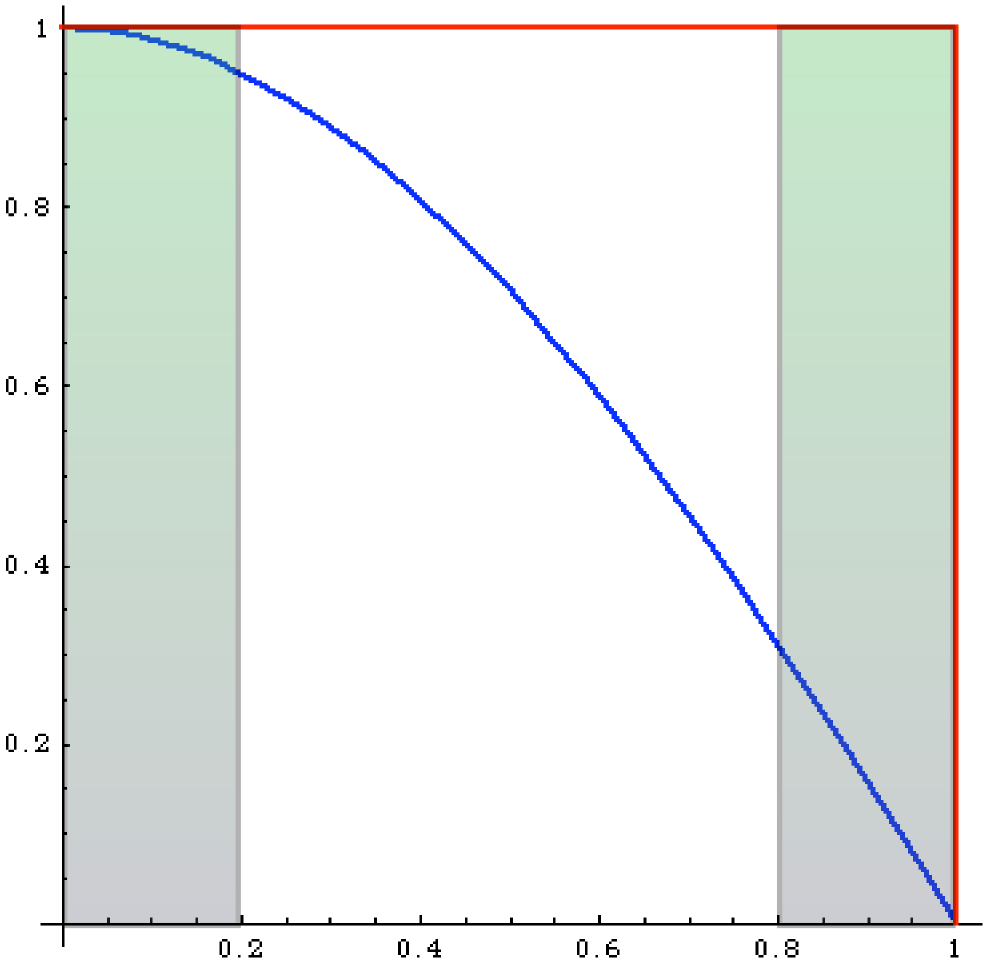
Alternative way

1. (randomly) pick x
2. calculate f(x)
3. (randomly) pick 0<y<fmax
4. Compare:
   if f(x)>y accept event,

   else reject it.

$$\text{Integral} = \frac{\text{accepted}}{\text{total tries}} = \text{efficiency}$$

EVENT GENERATION



What's the difference?

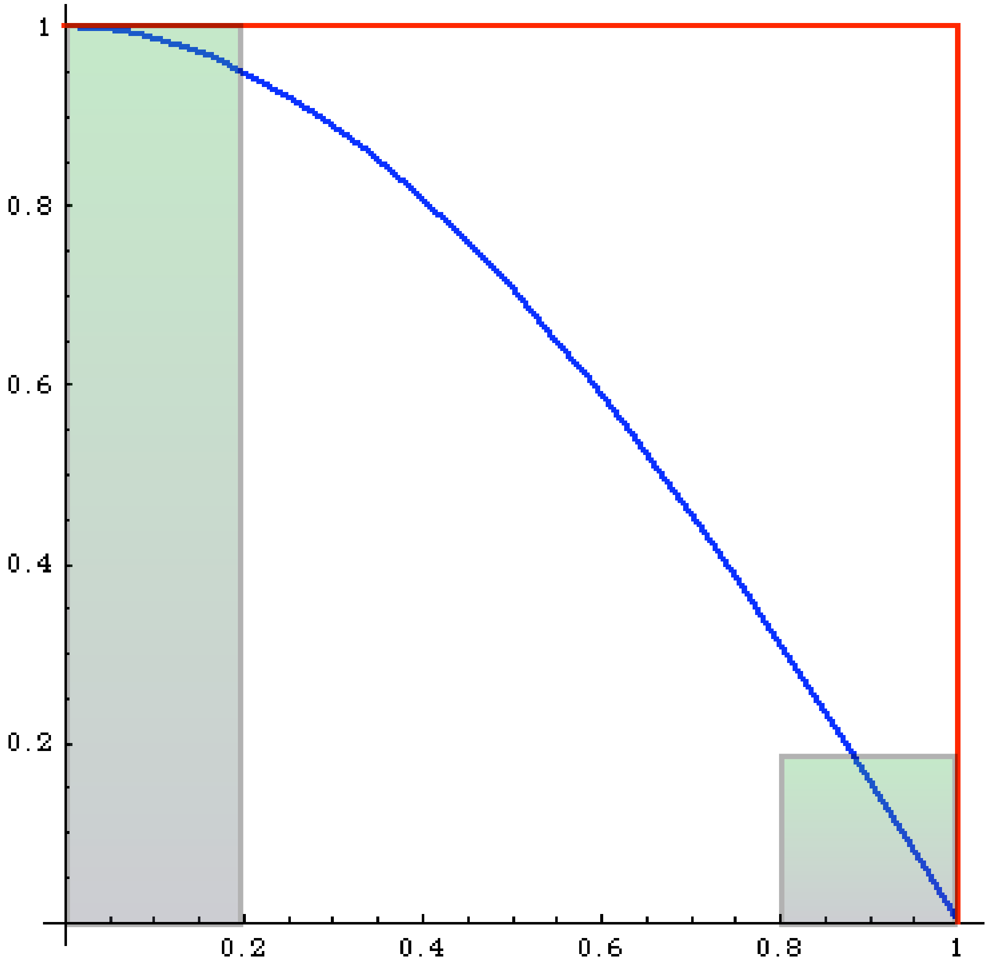
before:

Same # of events in areas of phase space with very different probabilities:

Events must have different weights:

$$w_i = p(x_i)$$

EVENT GENERATION



What's the difference?

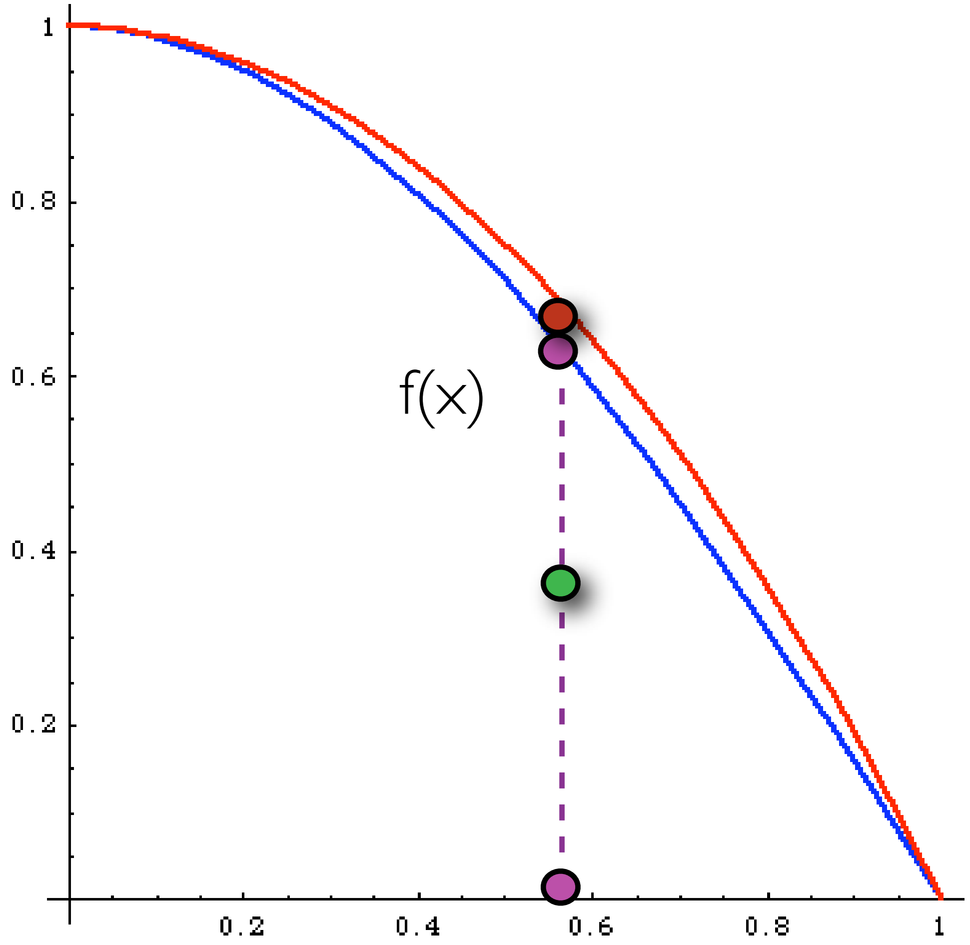
after:

events is proportional to the probability of areas of phase space:

Events have all the same weight ("unweighted")

Events distributed as in Nature

EVENT GENERATION



Improved

1. pick x distributed as p(x)

2. calculate f(x) and p(x)

3. pick 0<y<1

4. Compare:
   if f(x)>y p(x) accept event,

else reject it.

much better efficiency!!!

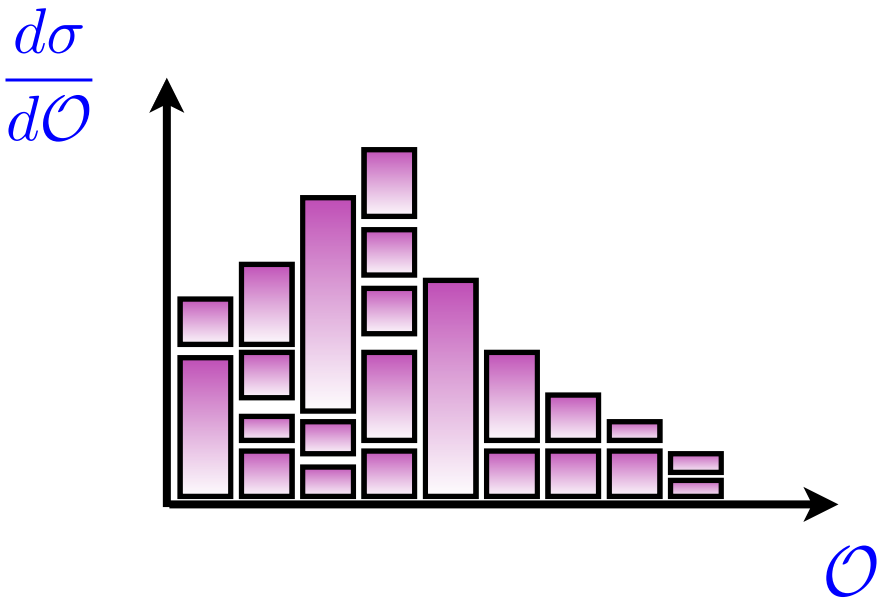
EVENT GENERATION

EVENT GENERATION

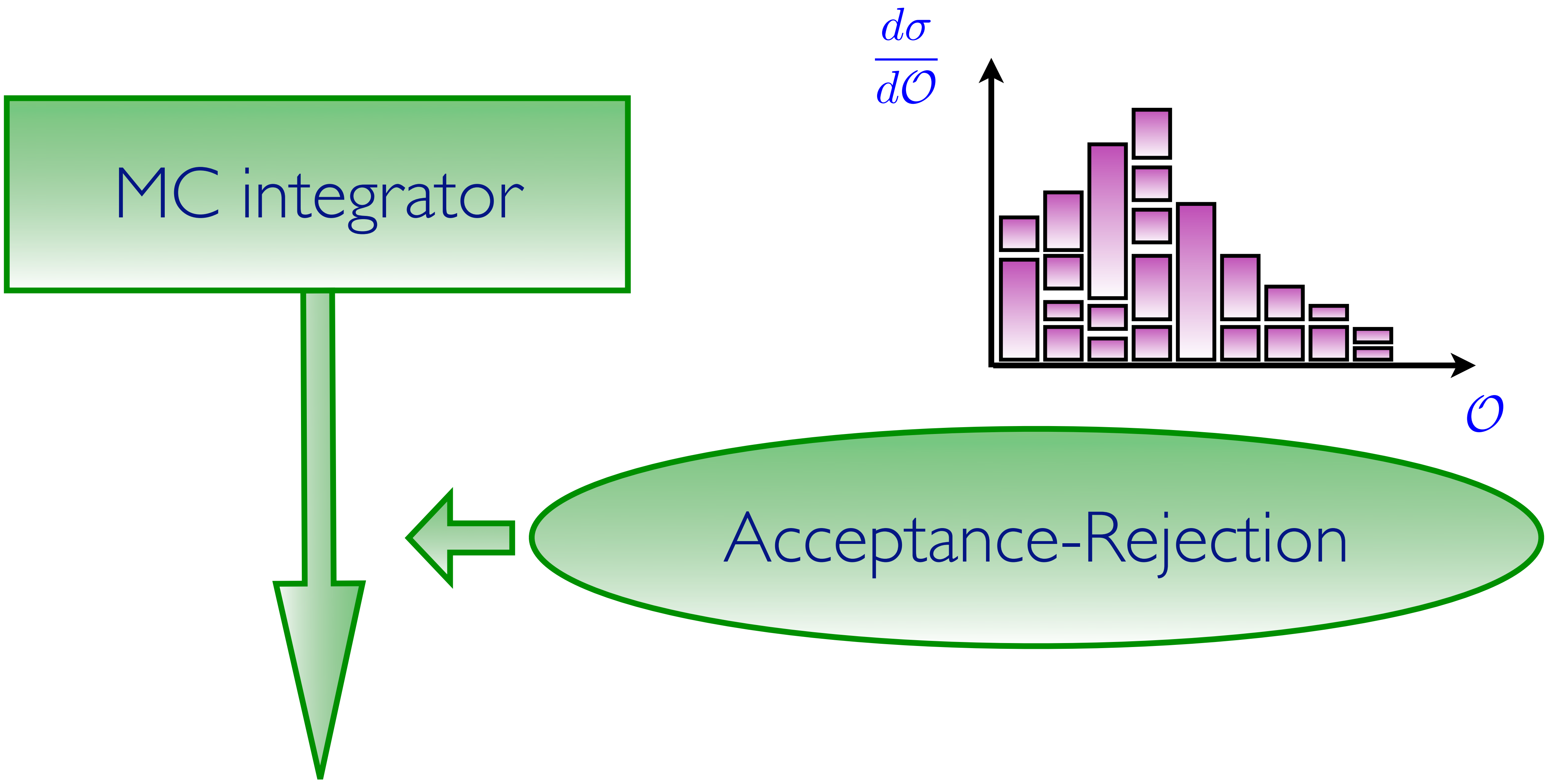
MC integrator

EVENT GENERATION

MC integrator



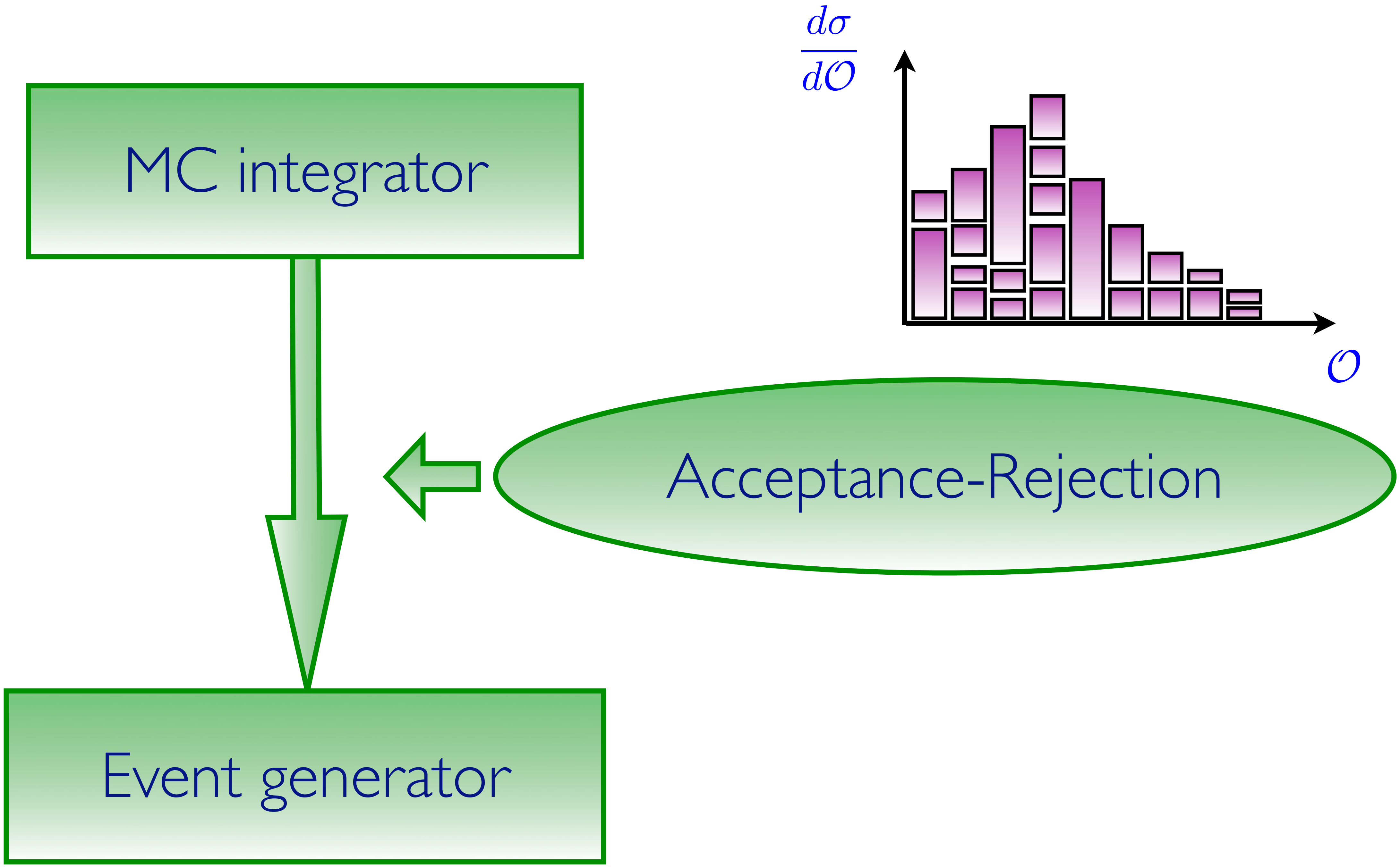
EVENT GENERATION



MC integrator

Acceptance-Rejection

EVENT GENERATION

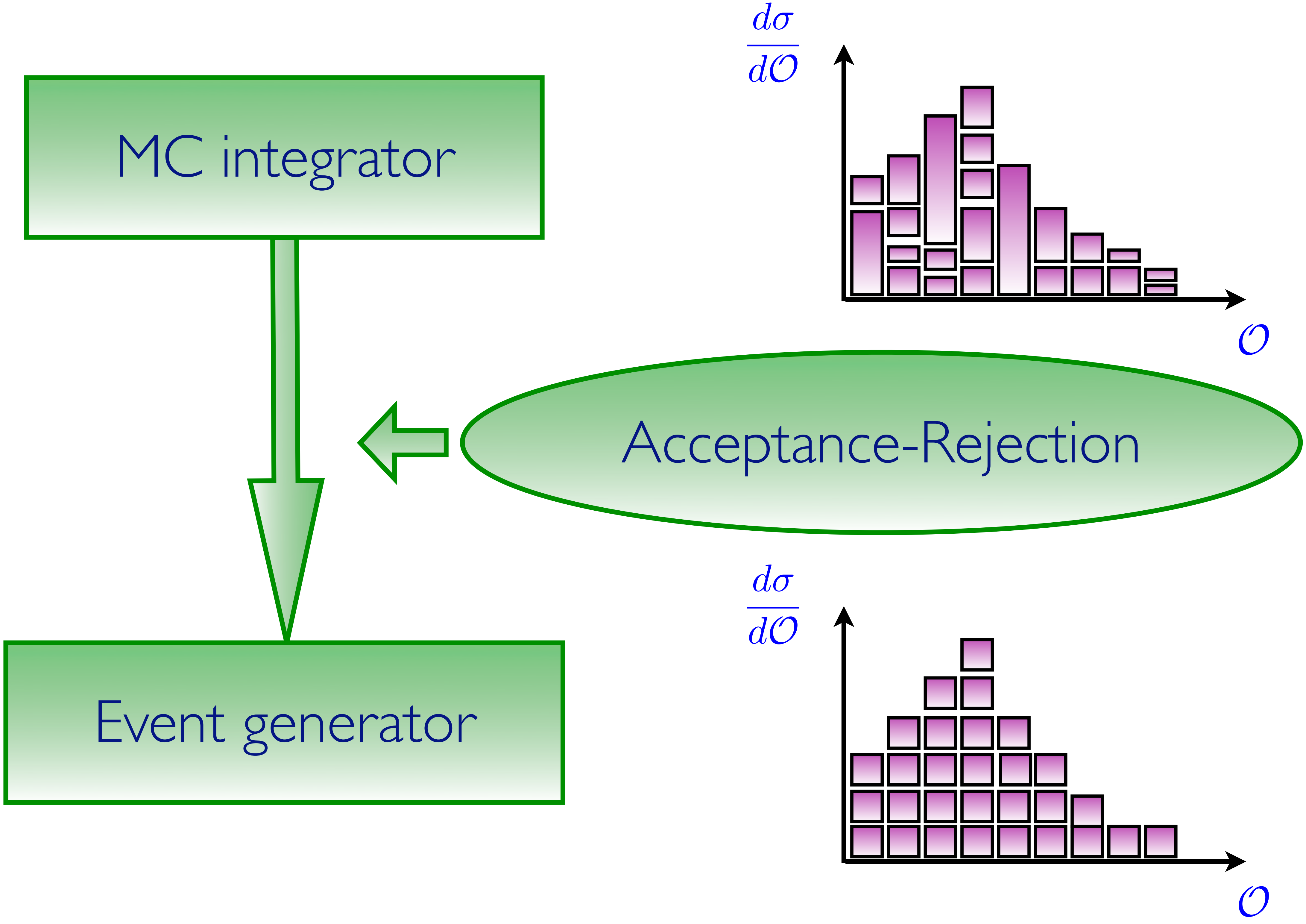


MC integrator

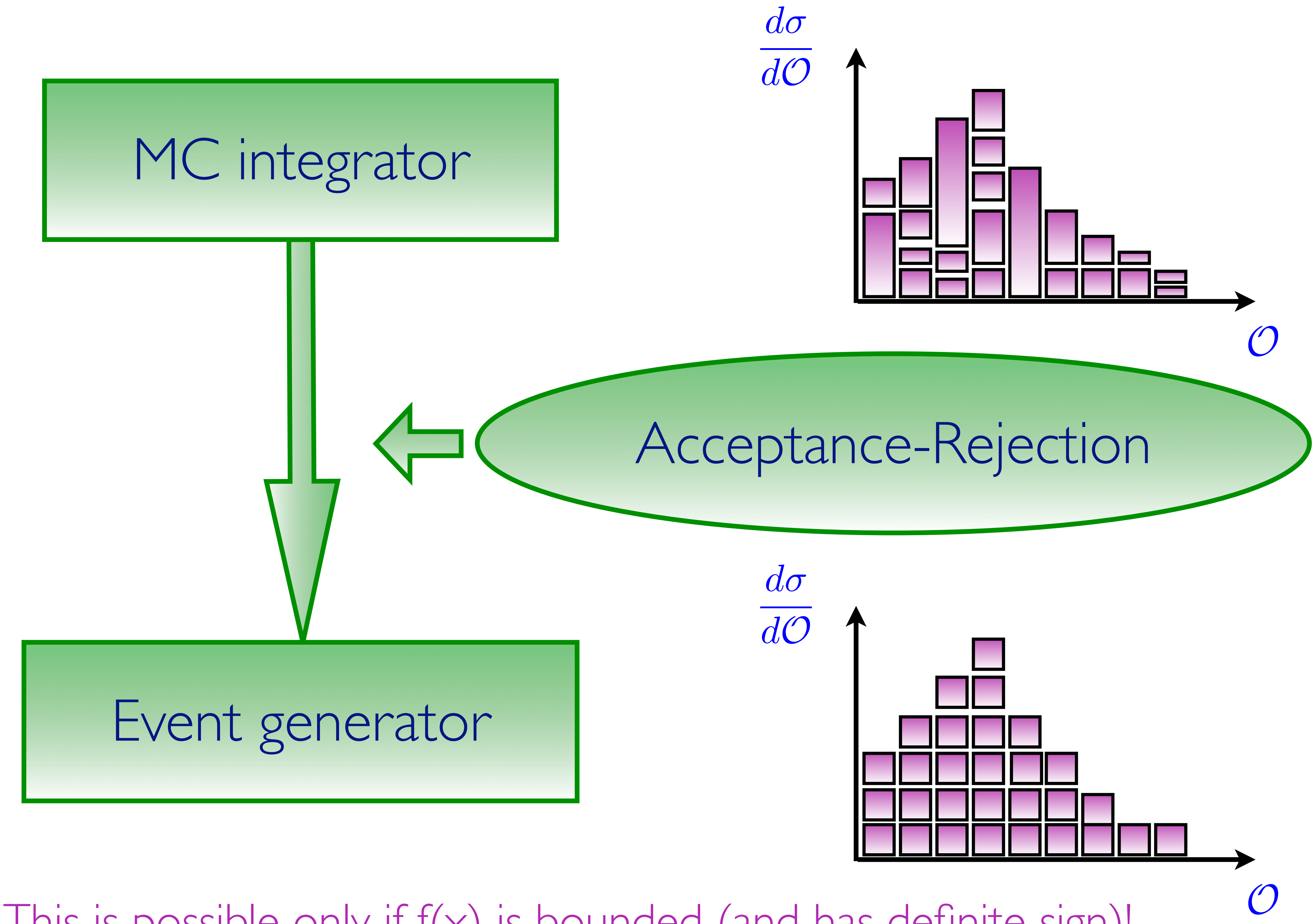
Acceptance-Rejection

Event generator

EVENT GENERATION



EVENT GENERATION



MC integrator

Acceptance-Rejection

Event generator

☞ This is possible only if f(x) is bounded (and has definite sign)!

MC EVENT GENERATOR: DEFINITION

At the most basic level a Monte Carlo event generator is a program which produces particle physics events with the same probability as they occur in nature (virtual collider).

In practice it performs (a possibly large) number of (sometimes very difficult) integrals and then unweights to give the four momenta of the particles that interact with the detector (simulation).

Note that, at least among theorists, the definition of a "Monte Carlo program" also includes codes which don't provide a fully exclusive information on the final state but only cross sections or distributions at the parton level, even when no unweighting can be performed (typically at NLO).

I will refer to these kind of codes as "MC integrators".

GENERAL STRUCTURE

SM

subprocs handler

Includes all possible subprocess leading to a given multi-jet final state automatically or manually (done once for all)

d~ d -> a d d~ u u~ g
d~ d -> a d d~ c c~ g
s~ s -> a d d~ u u~ g
s~ s -> a d d~ c c~ g
...

ME calculator

"Automatically" generates a code to calculate |M|2 for arbitrary processes with many partons in the final state. Use Feynman diagrams with tricks to reduce the factorial growth, others have recursive relations to reduce the complexity to exponential. ☺

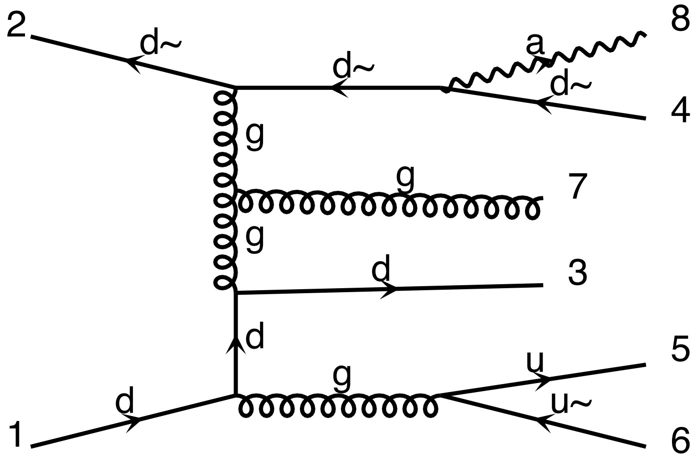
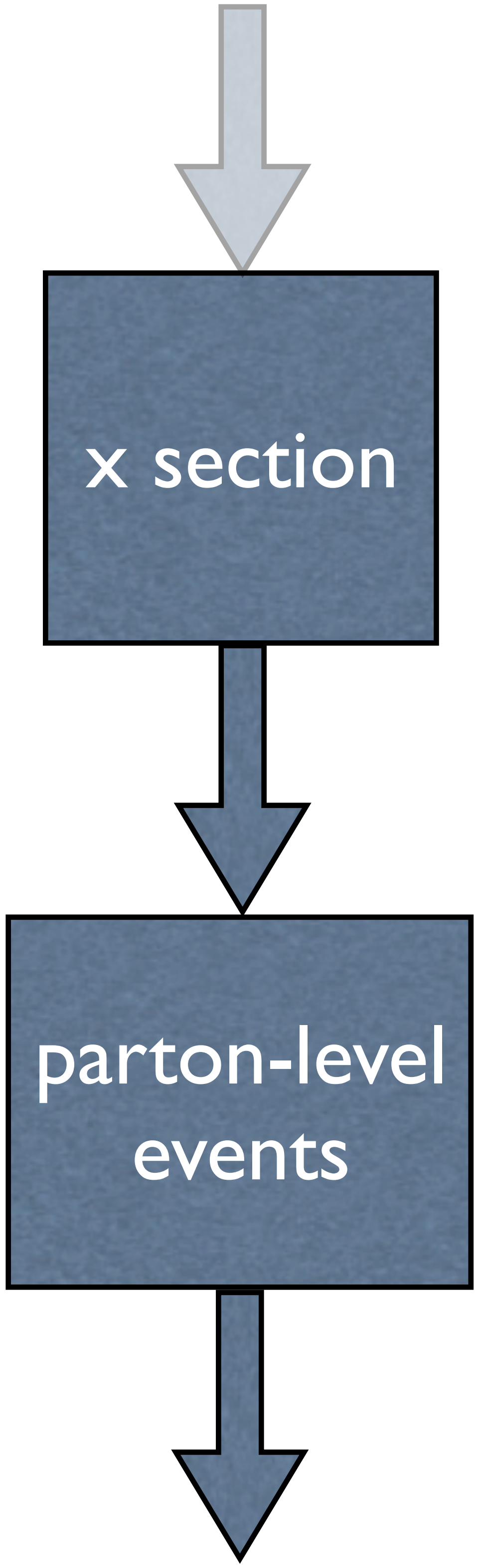


diagram 949 QCD=5, QED=1

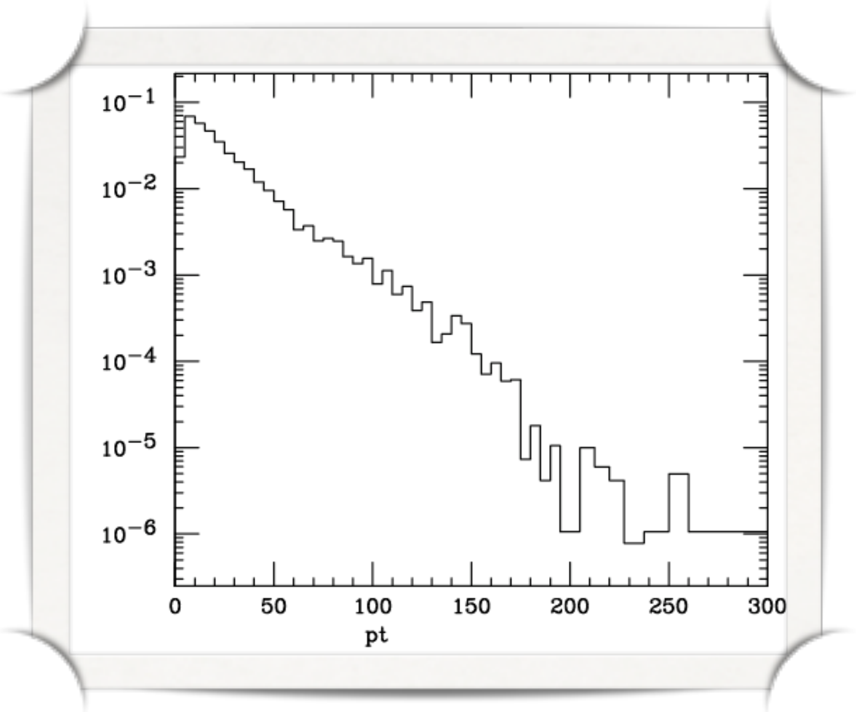
GENERAL STRUCTURE



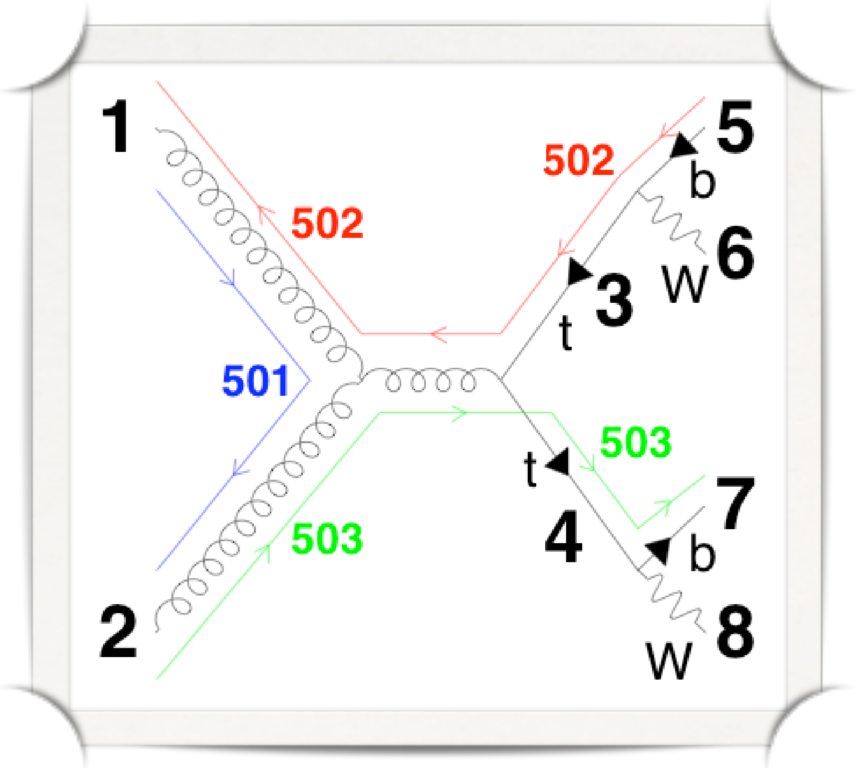
x section

parton-level events

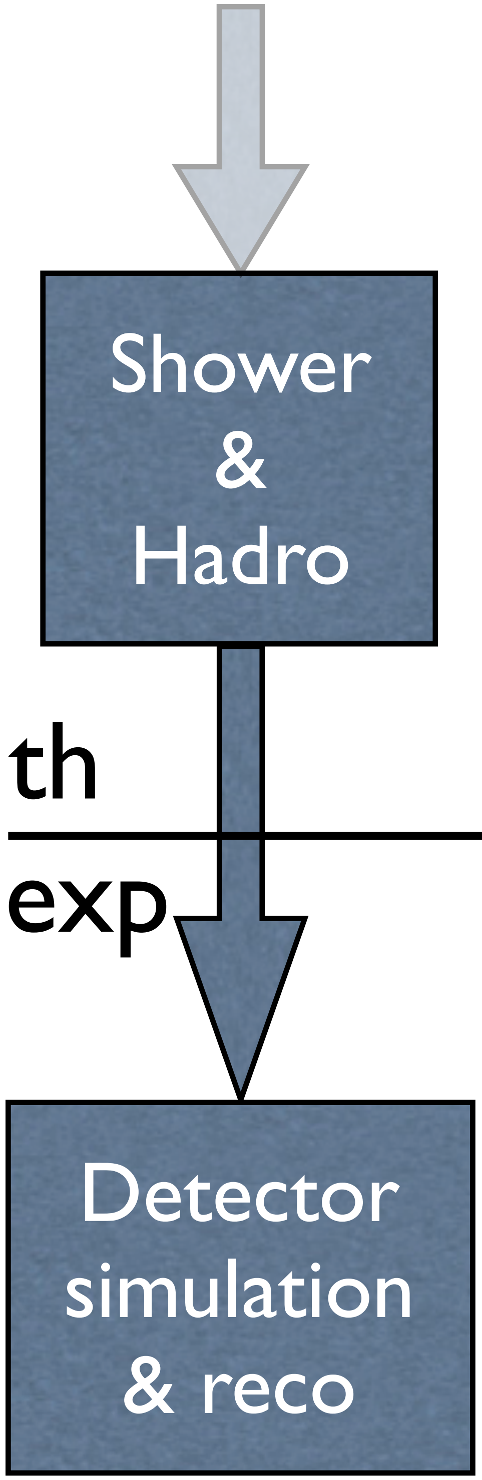
Integrate the matrix element over the phase space using a multi-channel technique and using parton-level cuts.



Events are obtained by unweighting. These are at the parton-level. Information on particle id, momenta, spin, color is given in the Les Houches format.

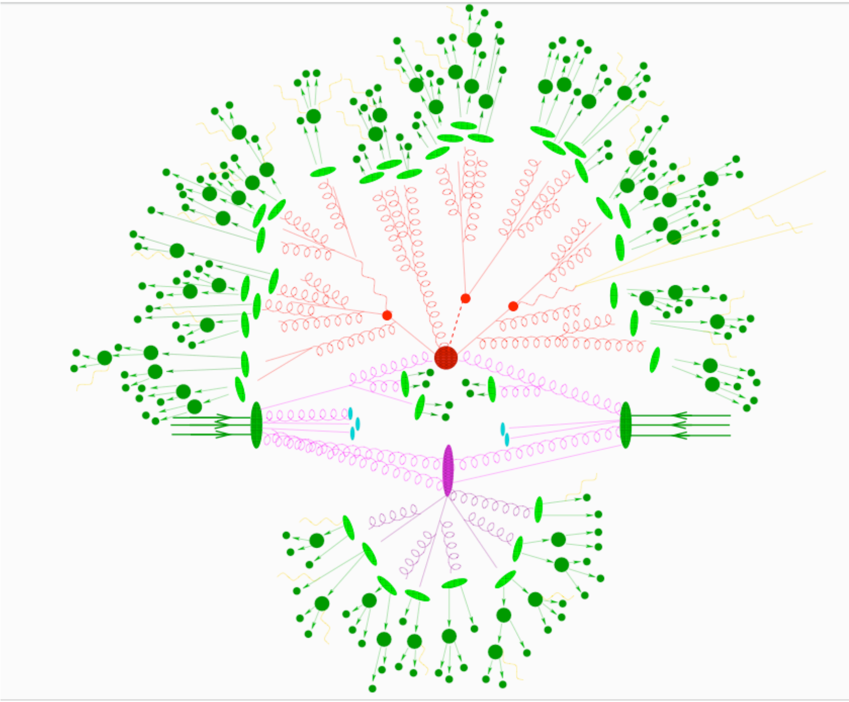


GENERAL STRUCTURE

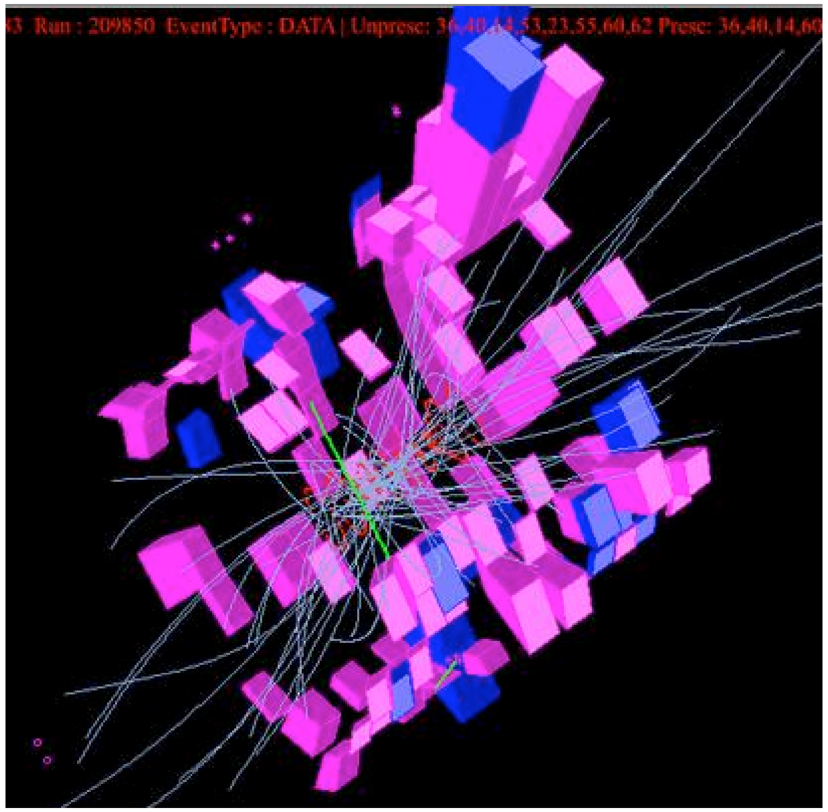


Events in the LH format are passed to the showering and hadronization⇒

high multiplicity hadron-level events



Events in HepMC format are passed through fast or full simulation, and physical objects (leptons, photons, jet, b-jets, taus) are reconstructed.



CODES

- Example of tree-level Monte Carlo codes:
  - Alpgen: fast matrix elements due to use of recursion relations. SM only.
  - Comix (Sherpa): fast matrix elements due to use of recursion relations. Some BSM models implemented (however, e.g. no Majorana particles).
  - MadGraph: Feynman diagrams to generate matrix elements which results in high unweighting efficiency. Virtually all BSM models are (or can be) implemented.
- and more: CalcHEP/CompHEP, Whizard...

Skip FeynRules

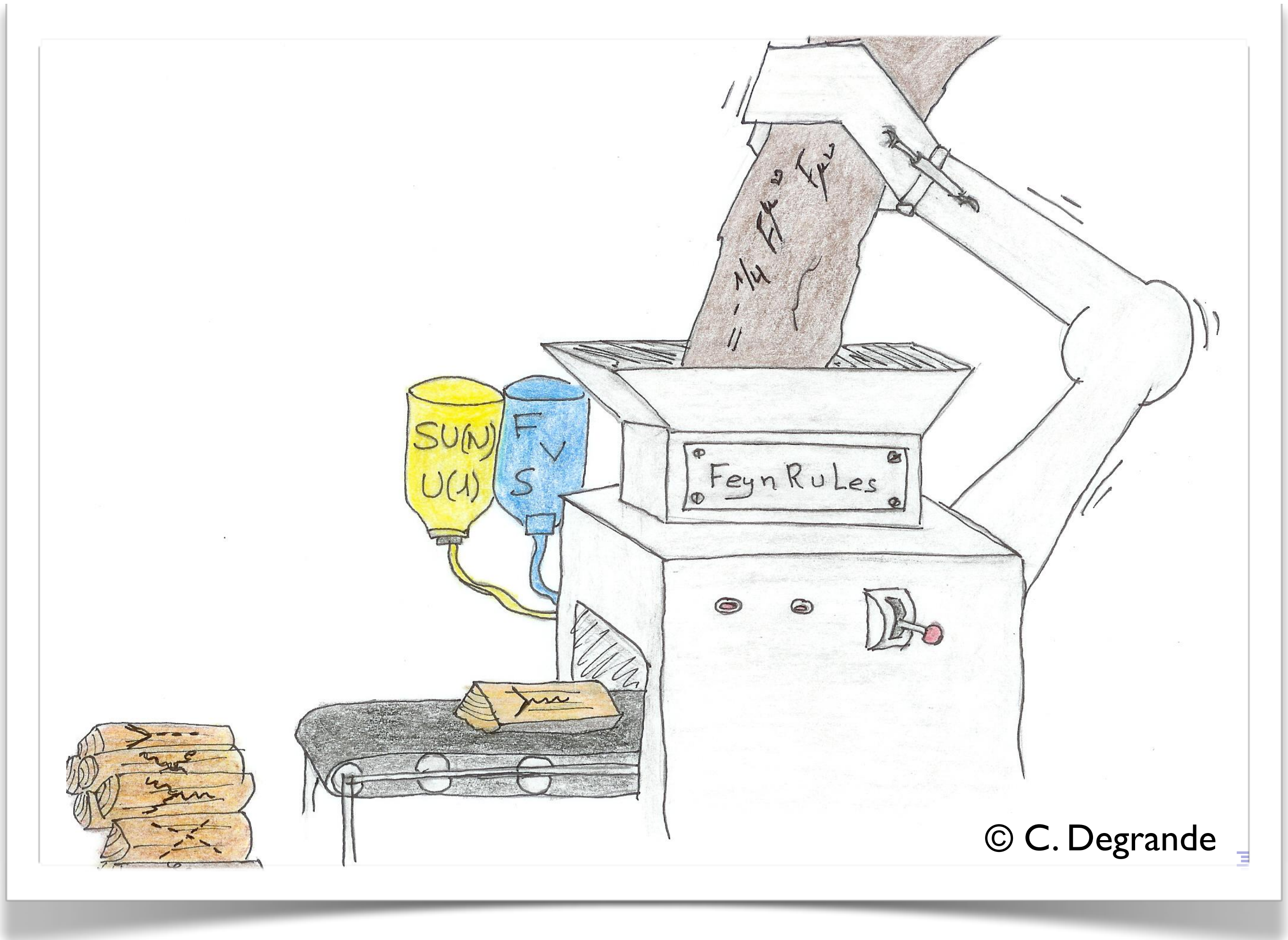
FEYNRULES

- FeynRules is a Mathematica package that allows to derive Feynman rules from a Lagrangian.
- Current public version: 1.6.x.
- The only requirements on the Lagrangian are:
  - ➡ All indices need to be contracted (i.e. Lorentz and gauge invariance)
  - ➡ Locality
  - ➡ Supported field types:

    spin 0, 1/2, 1, 2 & ghosts (3/2 are coming)

FEYNRULES

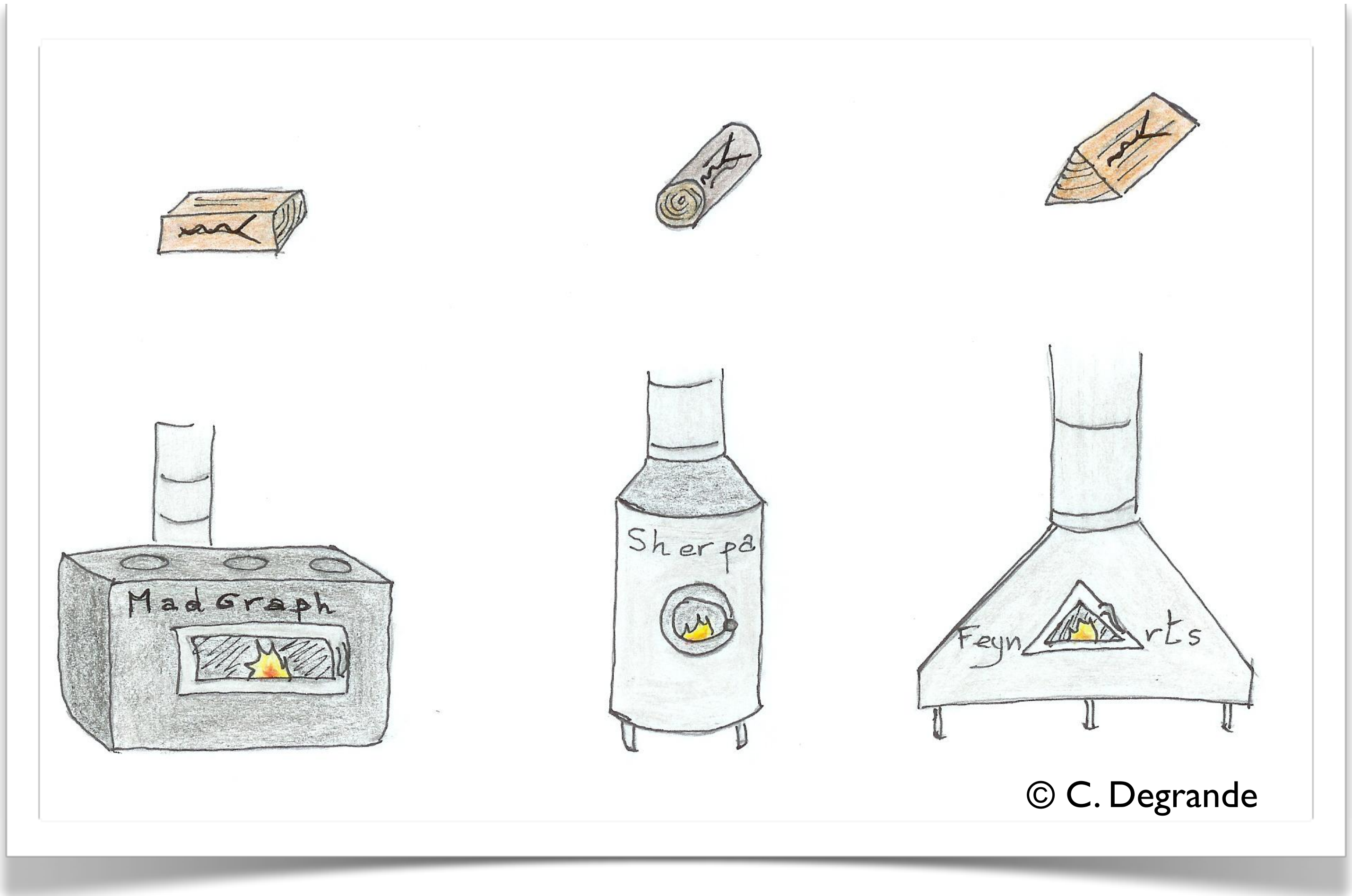
- FeynRules comes with a set of interfaces, that allow to export the Feynman rules to various matrix element generators.
- Interfaces coming with current public version
    - CalcHep / CompHep
    - FeynArts / FormCalc
    - MadGraph
    - Sherpa
    - Whizard / Omega
    - Universal FeynRules Output



© C. Degrande

FEYNRULES

- FeynRules comes with a set of interfaces, that allow to export the Feynman rules to various matrix element generators.
- Interfaces coming with current public version
  - ➡ CalcHep / CompHep
  - ➡ FeynArts / FormCalc
  - ➡ MadGraph
  - ➡ Sherpa
  - ➡ Whizard / Omega
  - ➡ Universal FeynRules Output



© C. Degrande

FEYNRULES

- The input requested form the user is twofold.

- The Model File:
  Definitions of particles and
  parameters (e.g., a quark)

```
F[1]  ==
  {ClassName      ->   q,
    SelfConjugate ->  False,
    Indices           ->  {Index[Colour]},
    Mass              ->  {MQ,  200},
    Width             ->  {WQ, 5}   }
```

- The Lagrangian:

$$\mathcal{L} = -\frac{1}{4} G^a_{\mu\nu}\, G^{\mu\nu}_a + i\bar{q}\,\gamma^\mu\, D_\mu q - M_q\, \bar{q}\, q$$

```
L =
-1/4 FS[G,mu,nu,a] FS[G,mu,nu,a]
+ I qbar.Ga[mu].del[q,mu]
- MQ qbar.q
```

FEYNRULES

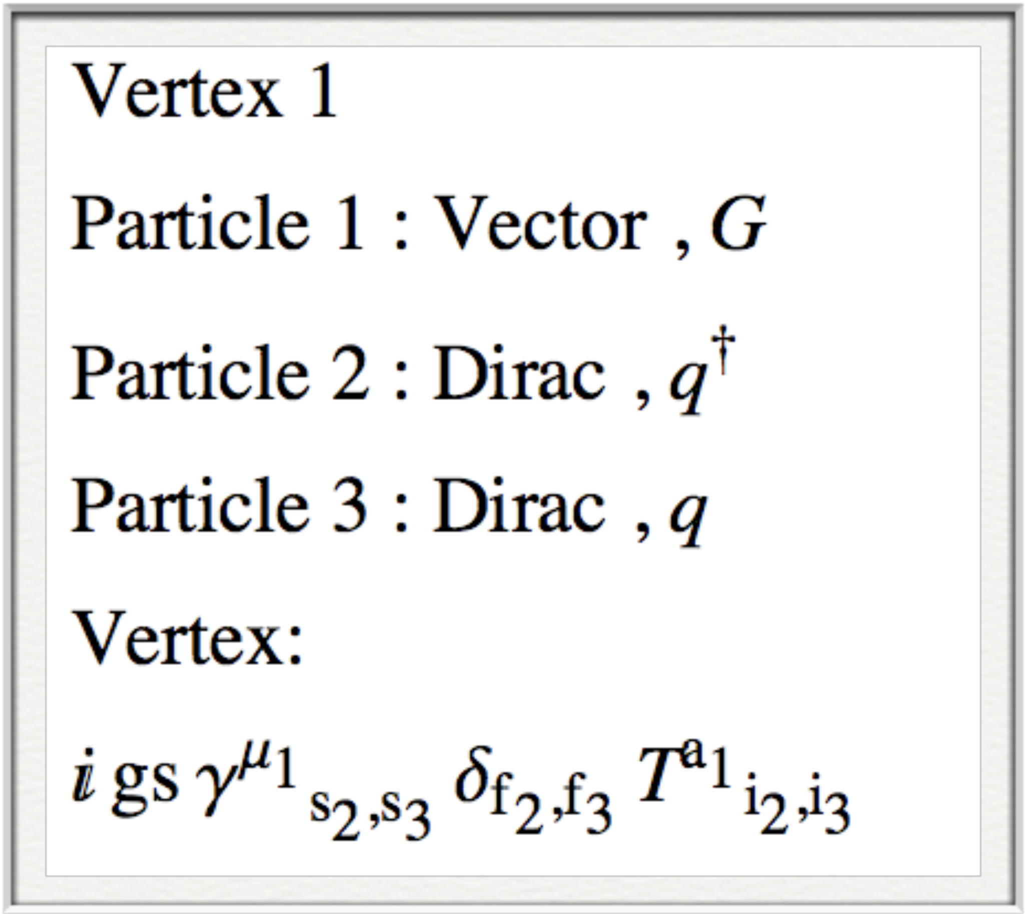
- Once this information has been provided, FeynRules can be used to compute the Feynman rules for the model:

`FeynmanRules[ L ]`

FEYNRULES

- Once this information has been provided, FeynRules can be used to compute the Feynman rules for the model:

FeynmanRules[L]



Vertex 1

Particle 1 : Vector , G

Particle 2 : Dirac , $q^\dagger$

Particle 3 : Dirac , q

Vertex:

$$i\, \mathrm{gs}\, \gamma^{\mu_1}{}_{\mathrm{s}_2,\mathrm{s}_3}\, \delta_{\mathrm{f}_2,\mathrm{f}_3}\, T^{\mathrm{a}_1}{}_{\mathrm{i}_2,\mathrm{i}_3}$$

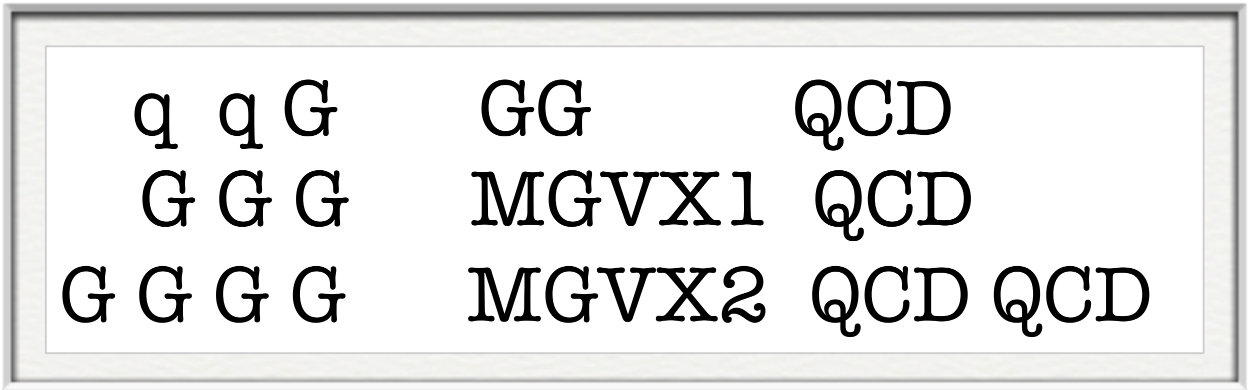
FEYNRULES

- Once we have the Feynman rules, we can export them to a MC event generator via the UFO:

WriteUFOutput[L]

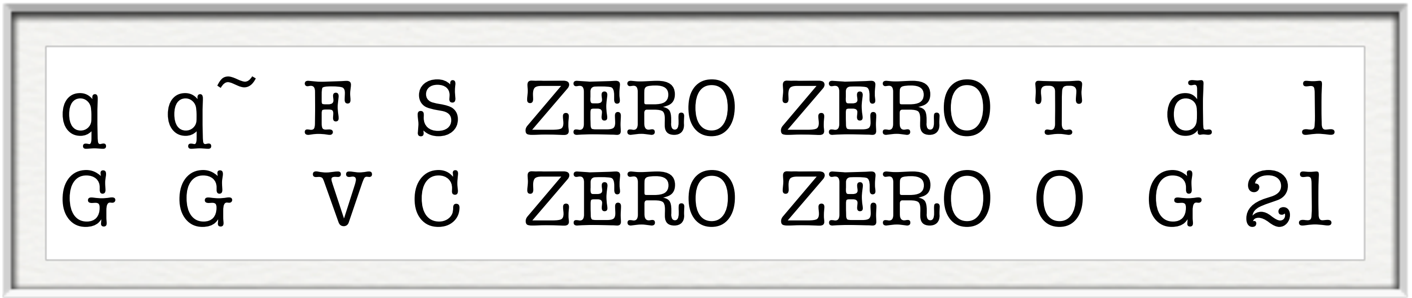
- This produces a set of files that can be directly used in the matrix element generator (“plug ‘n’ play”).

interactions.dat



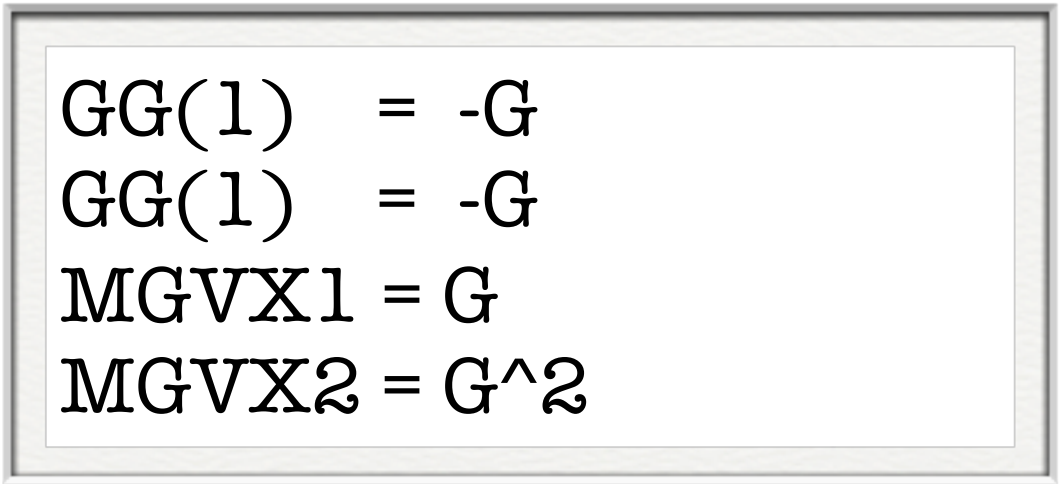
```
  q  q G     GG        QCD
   G G G    MGVX1  QCD
G G G G     MGVX2  QCD QCD
```

particles.dat



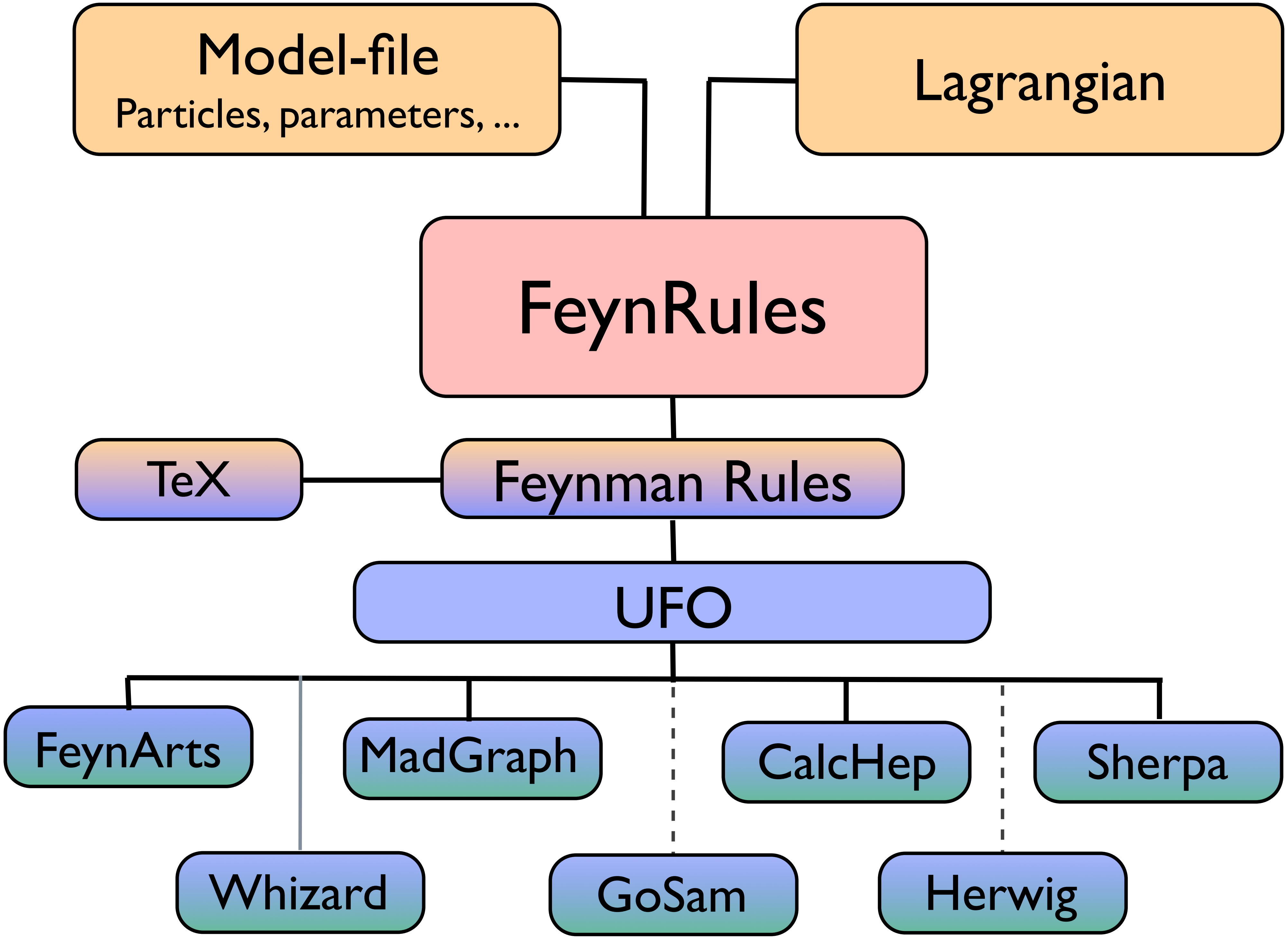
```
q  q~  F  S  ZERO ZERO T  d   1
G  G   V  C  ZERO ZERO O  G  21
```

couplings.dat



```
GG(1)   =  -G
GG(1)   =  -G
MGVX1 = G
MGVX2 = G^2
```

FEYNRULES



Model-file

Particles, parameters, ...

Lagrangian

FeynRules

TeX

Feynman Rules

UFO

FeynArts

MadGraph

CalcHep

Sherpa

Whizard

GoSam

Herwig

FEYNRULES

- Already available models:
  - Standard Model
  - Simple extensions of the SM (4th generation, 2HDM, ...)
  - SUSY models ((N)MSSM, RPV-MSSM, ...)
  - Extra-dimensional models (minimal UED, Large Extra Dimensions, ...)
  - Strongly coupled and effective field theories (Minimal Walking Technicolor, Chiral Perturbation theory, ...)
- Straight-forward to start from a given model and to add extra particles/interactions
- All available models, restrictions, syntax and more information can be found on the FeynRules website:

http://feynrules.phys.ucl.ac.be

LO PREDICTIONS : REMARKS

$$\sigma_X = \sum_{a,b} \int_0^1 dx_1 dx_2\, f_a(x_1,\mu_F^2) f_b(x_2,\mu_F^2) \times \hat{\sigma}_{ab\to X}(x_1,x_2,\alpha_S(\mu_R^2),\frac{Q^2}{\mu_F^2},\frac{Q^2}{\mu_R^2})$$

- By calculating the short distance coefficient at tree-level we obtain the first estimate of rates for inclusive final states.

- **Even at LO extra radiation is included:** it is described by the PDF's in the initial state and by the definition of a final state parton, which at LO represents all possible final state evolutions.

- Due to the above approximations a cross section at LO can strongly depend on the factorization and renormalization scales.

- Predictions can be systematically improved, at NLO and NNLO, by including higher order corrections in the short distance and in the evolution of the PDF's.

SUMMARY

- Having accurate and flexible simulations tools available for the LHC is a necessity (even more now!!)
- At LO event generation is technically challenging, yet conceptually straightforward.

CREDITS

To organize this presentation I have benefited from lectures (and actual slides), talks and discussions with many people.
In particular:

- Mike Seymour (MC basics)
- Claude Duhr (FeynRules)
- Johan Alwall (ME+PS merging)
- Rikkert Frederix, Paolo Torrielli (NLO+PS)
- Stefano Frixione, Michelangelo Mangano, Paolo Nason (for QCD, PS, LO, NLO, and more...)
- ....

Whom I all warmly thank!!