

Advanced FeynRules topics

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Recap of previous lecture

- Lecture 1 & 2: FeynRules

- ➔ implement a model into FeynRules.

- ➔ export the model to MadGraph 5.

- ➔ check a model implementation.

- Lecture 3: UFO & ALOHA

- ➔ what happens when going from FeynRules to MadGraph.

- This lecture:

- ➔ Superfields in FeynRules.

- ➔ Going beyond tree level: InSurGe & FR@NLO.

- ➔ Galileo.

Superfields in FeynRules

The lifecycle of SUSY pheno

- Example: SUSY model

$$\mathcal{L} = \Phi^\dagger e^{-2gV} \Phi|_{\theta^2\bar{\theta}^2} + \frac{1}{16g^2\tau_{\mathcal{R}}} \text{Tr}(W^\alpha W_\alpha)|_{\theta^2} + \frac{1}{16g^2\tau_{\mathcal{R}}} \text{Tr}(\bar{W}_{\dot{\alpha}} \bar{W}^{\dot{\alpha}})|_{\bar{\theta}^2} + W(\Phi)|_{\theta^2} + W^*(\Phi^\dagger)|_{\bar{\theta}^2} + \mathcal{L}_{\text{soft}}$$

- Very easy ‘theory description’

- ➔ Choose a gauge group (+ additional internal symmetries).
- ➔ Choose the matter content (= chiral superfields in some representation).
- ➔ Write down the most general superpotential.
- ➔ Write down the soft-SUSY breaking terms.
- ➔ (+ check validity of the model)

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- ‘Monte Carlo description’

- ➔ Express superfields in terms of component fields.
- ➔ Express everything in terms of 4-component fermions (beware of the Majoranas!).
- ➔ Express everything in terms of mass eigenstates.
- ➔ Integrate out D and F terms.
- ➔ Implement vertices one-by-one (beware of factors of i , *etc!*)

Supersymmetric models

- FeynRules allows to use the superfield formalism for supersymmetric theories.
- The code then
 - ➔ expands the superfields in the Grassmann variables and integrates them out.
 - ➔ Weyl fermions are transformed into 4-component spinors.
 - ➔ auxiliary fields are integrated out.
- As a result, we obtain a Lagrangian that can be exported to matrix element generators!

Supersymmetric models

- Example: SUSY QCD

- ➔ 1 octet vector superfield

$$V^a = (\tilde{g}^a, G_\mu^a, D^a)$$

- ➔ 1 triplet left-handed chiral superfield

$$Q_L^i = (\tilde{q}_L^i, \chi^i, F_L^i)$$

- ➔ 1 triplet right-handed chiral superfield

$$Q_R^i = (\tilde{q}_R^i, \bar{\xi}^i, F_R^i)$$

- The physical spectrum contains

- ➔ a gauge boson, the gluon

- ➔ two complex triplet scalars

- ➔ an octet Majorana fermion

- ➔ a triplet Dirac fermion, the quark

$$q^i = (\chi^i, \bar{\xi}^i)$$

Supersymmetric models

- Interactions (almost) entirely fixed by SUSY

$$Q_L^\dagger e^{-2g_s V} Q_L + Q_R^\dagger e^{-2g_s V} Q_R + \frac{1}{8g_s^2} \text{Tr}(W^\alpha W_\alpha) + \frac{1}{8g_s^2} \text{Tr}(\bar{W}_{\dot{\alpha}} \bar{W}^{\dot{\alpha}}) \\ + W(Q_L, Q_R^\dagger) + W^*(Q_L^\dagger, Q_R)$$

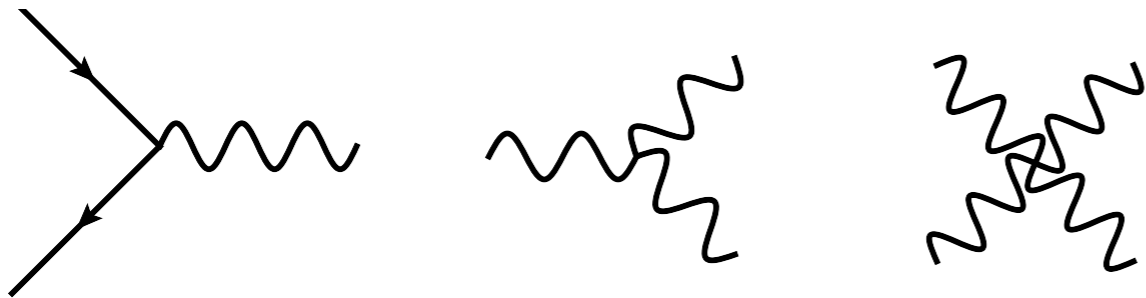
- The gauge sector is already rather complicated in terms of component fields...

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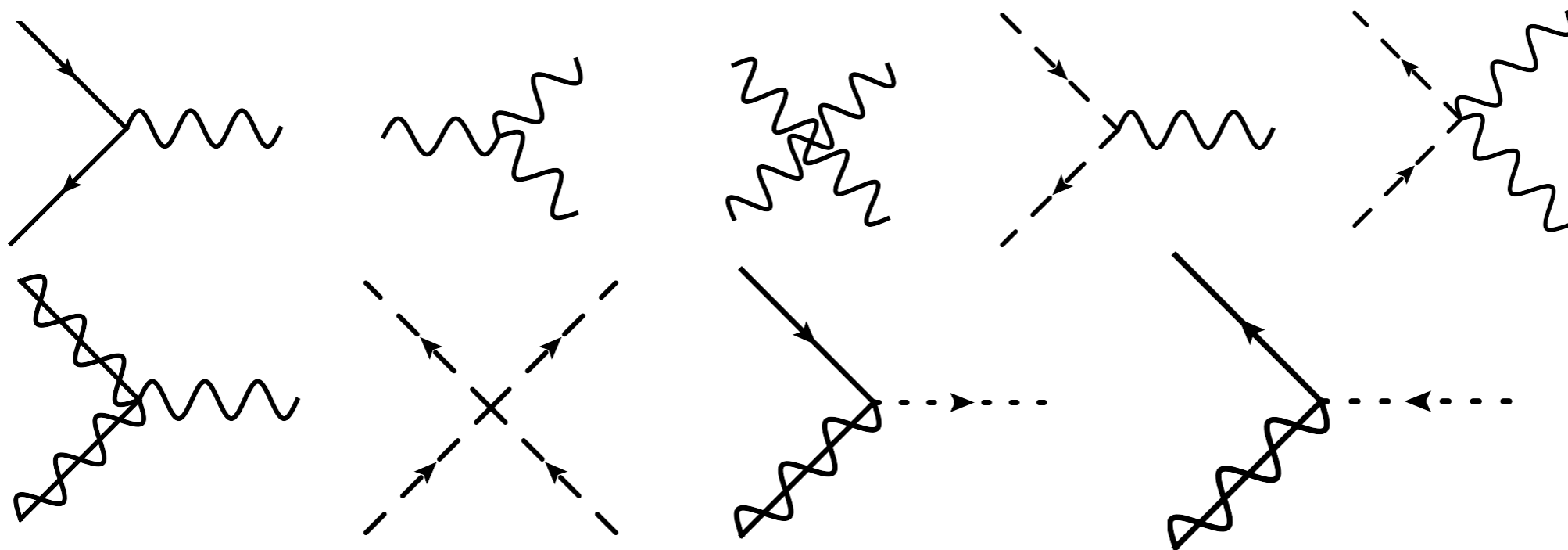


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Defining superfields

```
VSF[1] == { ClassName -> GSF,  
             GaugeBoson -> G,  
             Gaugino -> gow,  
             Indices -> {Index[Gluon]}}
```

$$V^a = (\tilde{g}^a, G_\mu^a, D^a)$$

```
CSF[1] == { ClassName -> QL,  
            Chirality -> Left,  
            Weyl -> qLw,  
            Scalar -> QLs,  
            Indices->{Index[Colour]}}
```

$$Q_L^i = (\tilde{q}_L^i, \chi^i, F_L^i)$$

- The component fields are defined separately.
- Auxiliary F and D fields could be added, but can be left out, and are created on the fly.

Using superfields

$WS = \dots$

$SL = VSFKineticTerms[] + CSFKineticTerms[] + WS + HC[WS];$

- A set of functions allows to transform the superspace action into a component field Lagrangian.
 - ➔ **SF2Components**: expansion in the Grassmann parameters
 - ➔ **ThetaThetabarComponent** etc.: selects the desired coefficient in the Grassmann expansion.
 - ➔ **SolveEqMotionF/SolveEqMotionD**: solves the equations of motion for the F and D terms.
 - ➔ **WeylToDirac**: Transforms Weyl fermions into 4-component fermions.

Using superfields

$$\begin{aligned}
 & -4 \int_{\text{Gluon}S1, \text{Gluon}S2, \text{Gluon}S3} \int_{\text{Gluon}S3, \text{Gluon}S4, \text{Gluon}S5} G_{\text{mu}S1, \text{Gluon}S1} G_{\text{mu}S1, \text{Gluon}S4} G_{\text{mu}S2, \text{Gluon}S2} G_{\text{mu}S2, \text{Gluon}S5} g^4 + \\
 & 16 \partial_{\text{mu}S2} (G_{\text{mu}S1, \text{Gluon}S1}) \int_{\text{Gluon}S1, \text{Gluon}S2, \text{Gluon}S3} G_{\text{mu}S1, \text{Gluon}S2} G_{\text{mu}S2, \text{Gluon}S3} g^3 + 8 i \bar{g}_{\text{O}rS28685, \text{Gluon}S2} \bar{g}_{\text{O}rS28698, \text{Gluon}S1} \int_{\text{Gluon}S1, \text{Gluon}S2, \text{Gluon}S3} G_{\text{mu}S1, \text{Gluon}S3} \gamma^{\text{mu}S1} P_{-rS28685, rS28698} g^3 - \\
 & 8 i \bar{g}_{\text{O}rS28698, \text{Gluon}S1} \bar{g}_{\text{O}rS28685, \text{Gluon}S2} \int_{\text{Gluon}S1, \text{Gluon}S2, \text{Gluon}S3} G_{\text{mu}S1, \text{Gluon}S3} \gamma^{\text{mu}S1} P_{+rS28698, rS28685} g^3 - 8 \partial_{\text{mu}S2} (G_{\text{mu}S1, \text{Gluon}S1})^2 g^2 + \\
 & 8 \partial_{\text{mu}S2} (G_{\text{mu}S1, \text{Gluon}S1}) \partial_{\text{mu}S1} (G_{\text{mu}S2, \text{Gluon}S1}) g^2 + G_{\text{mu}S1, \text{Gluon}S1} G_{\text{mu}S1, \text{Gluon}S2} \text{QLsColour}S1 \text{QLsColour}S2 T_{\text{Colour}S2, \text{Colour}S3}^{\text{Gluon}S1} T_{\text{Colour}S3, \text{Colour}S1}^{\text{Gluon}S2} g^2 + \\
 & G_{\text{mu}S1, \text{Gluon}S1} G_{\text{mu}S1, \text{Gluon}S2} \text{QRsColour}S1 \text{QRsColour}S2 T_{\text{Colour}S2, \text{Colour}S3}^{\text{Gluon}S1} T_{\text{Colour}S3, \text{Colour}S1}^{\text{Gluon}S2} g^2 + 4 i \bar{g}_{\text{O}rS28679, \text{Gluon}S1} \partial_{\text{mu}S1} (g_{\text{O}rS28692, \text{Gluon}S1}) \gamma^{\text{mu}S1} P_{-rS28679, rS28692} g^2 - \\
 & 4 i \partial_{\text{mu}S1} (g_{\text{O}rS28682, \text{Gluon}S1}) \bar{g}_{\text{O}rS28695, \text{Gluon}S1} \gamma^{\text{mu}S1} P_{-rS28682, rS28695} g^2 - 4 i \partial_{\text{mu}S1} (g_{\text{O}rS28692, \text{Gluon}S1}) \bar{g}_{\text{O}rS28679, \text{Gluon}S1} \gamma^{\text{mu}S1} P_{+rS28692, rS28679} g^2 + \\
 & 4 i \bar{g}_{\text{O}rS28695, \text{Gluon}S1} \partial_{\text{mu}S1} (g_{\text{O}rS28682, \text{Gluon}S1}) \gamma^{\text{mu}S1} P_{+rS28695, rS28682} g^2 - i \partial_{\text{mu}S1} (\text{QLsColour}S1) G_{\text{mu}S1, \text{Gluon}S1} \text{QLsColour}S2 T_{\text{Colour}S1, \text{Colour}S2}^{\text{Gluon}S1} g^2 + \\
 & i \sqrt{2} \bar{q}_{rS28675, \text{Colour}S1} \bar{g}_{\text{O}rS28688, \text{Gluon}S1} P_{+rS28675, rS28688} \text{QLsColour}S2 T_{\text{Colour}S1, \text{Colour}S2}^{\text{Gluon}S1} g^2 + i \partial_{\text{mu}S1} (\text{QRsColour}S1) G_{\text{mu}S1, \text{Gluon}S1} \text{QRsColour}S2 T_{\text{Colour}S1, \text{Colour}S2}^{\text{Gluon}S1} g^2 - \\
 & i \sqrt{2} \bar{q}_{rS28689, \text{Colour}S1} \bar{g}_{\text{O}rS28676, \text{Gluon}S1} P_{-rS28689, rS28676} \text{QRsColour}S2 T_{\text{Colour}S1, \text{Colour}S2}^{\text{Gluon}S1} g^2 + i \partial_{\text{mu}S1} (\text{QLsColour}S1) G_{\text{mu}S1, \text{Gluon}S1} \text{QLsColour}S2 T_{\text{Colour}S2, \text{Colour}S1}^{\text{Gluon}S1} g^2 - \\
 & i \sqrt{2} \bar{g}_{\text{O}rS28677, \text{Gluon}S1} \bar{q}_{rS28690, \text{Colour}S1} P_{-rS28677, rS28690} \text{QLsColour}S2 T_{\text{Colour}S2, \text{Colour}S1}^{\text{Gluon}S1} g^2 - i \partial_{\text{mu}S1} (\text{QRsColour}S1) G_{\text{mu}S1, \text{Gluon}S1} \text{QRsColour}S2 T_{\text{Colour}S2, \text{Colour}S1}^{\text{Gluon}S1} g^2 + \\
 & i \sqrt{2} \bar{g}_{\text{O}rS28691, \text{Gluon}S1} \bar{q}_{rS28678, \text{Colour}S1} P_{+rS28691, rS28678} \text{QRsColour}S2 T_{\text{Colour}S2, \text{Colour}S1}^{\text{Gluon}S1} g^2 + \bar{q}_{rS28687, \text{Colour}S2} \bar{q}_{rS28700, \text{Colour}S1} G_{\text{mu}S1, \text{Gluon}S1} T_{\text{Colour}S2, \text{Colour}S1}^{\text{Gluon}S1} \gamma^{\text{mu}S1} P_{-rS28687, rS28700} g^2 - \\
 & \bar{q}_{rS28699, \text{Colour}S1} \bar{q}_{rS28686, \text{Colour}S2} G_{\text{mu}S1, \text{Gluon}S1} T_{\text{Colour}S1, \text{Colour}S2}^{\text{Gluon}S1} \gamma^{\text{mu}S1} P_{+rS28699, rS28686} g^2 + \frac{1}{2} \partial_{\text{mu}S1} (\text{QLsColour}S1) \partial_{\text{mu}S1} (\text{QLsColour}S1) + \\
 & \frac{1}{2} \partial_{\text{mu}S1} (\text{QRsColour}S1) \partial_{\text{mu}S1} (\text{QRsColour}S1) - \frac{1}{4} \partial_{\text{mu}S1} (\partial_{\text{mu}S1} (\text{QLsColour}S1)) \text{QLsColour}S1 - \frac{1}{4} \partial_{\text{mu}S1} (\partial_{\text{mu}S1} (\text{QLsColour}S1)) \text{QLsColour}S1 - \frac{1}{4} \partial_{\text{mu}S1} (\partial_{\text{mu}S1} (\text{QRsColour}S1)) \text{QRsColour}S1 - \\
 & \frac{1}{4} \partial_{\text{mu}S1} (\partial_{\text{mu}S1} (\text{QRsColour}S1)) \text{QRsColour}S1 + \frac{1}{32} \text{QLsColour}S1S28637 \text{QLsColour}S1S28639 \text{QLsColour}S2S28637 \text{QLsColour}S2S28639 T_{\text{Colour}S2S28637, \text{Colour}S1S28637}^{\text{Gluon}S1} T_{\text{Colour}S2S28639, \text{Colour}S1S28639}^{\text{Gluon}S1} + \\
 & \frac{1}{32} \text{QLsColour}S1S28639 \text{QLsColour}S2S28639 \text{QRsColour}S1S28638 \text{QRsColour}S2S28638 T_{\text{Colour}S2S28638, \text{Colour}S1S28638}^{\text{Gluon}S1} T_{\text{Colour}S2S28639, \text{Colour}S1S28639}^{\text{Gluon}S1} + \\
 & \frac{1}{32} \text{QLsColour}S1S28637 \text{QLsColour}S2S28637 \text{QRsColour}S1S28640 \text{QRsColour}S2S28640 T_{\text{Colour}S2S28637, \text{Colour}S1S28637}^{\text{Gluon}S1} T_{\text{Colour}S2S28640, \text{Colour}S1S28640}^{\text{Gluon}S1} + \\
 & \frac{1}{32} \text{QRsColour}S1S28638 \text{QRsColour}S1S28640 \text{QRsColour}S2S28638 \text{QRsColour}S2S28640 T_{\text{Colour}S2S28638, \text{Colour}S1S28638}^{\text{Gluon}S1} T_{\text{Colour}S2S28640, \text{Colour}S1S28640}^{\text{Gluon}S1} - \\
 & \frac{1}{16} \text{QLsColour}S1 \text{QLsColour}S1S28641 \text{QLsColour}S2 \text{QLsColour}S2S28641 T_{\text{Colour}S2, \text{Colour}S1}^{\text{Gluon}S1} T_{\text{Colour}S2S28641, \text{Colour}S1S28641}^{\text{Gluon}S1} - \\
 & \frac{1}{16} \text{QLsColour}S1 \text{QLsColour}S2 \text{QRsColour}S1S28642 \text{QRsColour}S2S28642 T_{\text{Colour}S2, \text{Colour}S1}^{\text{Gluon}S1} T_{\text{Colour}S2S28642, \text{Colour}S1S28642}^{\text{Gluon}S1} - \\
 & \frac{1}{16} \text{QLsColour}S1S28643 \text{QLsColour}S2S28643 \text{QRsColour}S1 \text{QRsColour}S2 T_{\text{Colour}S2, \text{Colour}S1}^{\text{Gluon}S1} T_{\text{Colour}S2S28643, \text{Colour}S1S28643}^{\text{Gluon}S1} - \\
 & \frac{1}{16} \text{QRsColour}S1 \text{QRsColour}S1S28644 \text{QRsColour}S2 \text{QRsColour}S2S28644 T_{\text{Colour}S2, \text{Colour}S1}^{\text{Gluon}S1} T_{\text{Colour}S2S28644, \text{Colour}S1S28644}^{\text{Gluon}S1} + \frac{1}{2} i \bar{q}_{rS28680, \text{Colour}S1} \partial_{\text{mu}S1} (q_{rS28693, \text{Colour}S1}) \gamma^{\text{mu}S1} P_{-rS28680, rS28693} - \\
 & \frac{1}{2} i \partial_{\text{mu}S1} (\bar{q}_{rS28683, \text{Colour}S1}) \bar{q}_{rS28696, \text{Colour}S1} \gamma^{\text{mu}S1} P_{-rS28683, rS28696} - \frac{1}{2} i \partial_{\text{mu}S1} (\bar{q}_{rS28694, \text{Colour}S1}) \bar{q}_{rS28681, \text{Colour}S1} \gamma^{\text{mu}S1} P_{+rS28694, rS28681} + \frac{1}{2} i \bar{q}_{rS28697, \text{Colour}S1} \partial_{\text{mu}S1} (q_{rS28684, \text{Colour}S1}) \gamma^{\text{mu}S1} P_{+rS28697, rS28684}
 \end{aligned}$$

Model database

We encourage model builders writing order to make them useful to a comm FeynRules model database, please see

- [✉ claude.duhr@durham.ac.uk](mailto:claude.duhr@durham.ac.uk)
- [✉ neil@hep.wisc.edu](mailto:neil@hep.wisc.edu)
- [✉ fuks@cern.ch](mailto:fuks@cern.ch)

Available models

[Standard Model](#)

[Simple extensions of the SM \(9\)](#)

[Supersymmetric Models \(4\)](#)

[Extra-dimensional Models \(4\)](#)

[Strongly coupled and effective field theories \(4\)](#)

[Miscellaneous \(0\)](#)

Model	Contact
MSSM	✉ B. Fuks
NMSSM	✉ B. Fuks
RPV-MSSM	✉ B. Fuks
R-MSSM	✉ B. Fuks

Summary

- FeynRules allows to use the superfield formalism for supersymmetric theories.
 - ➔ expands the superfields in the Grassmann variables and integrates them out.
 - ➔ Weyl fermions are transformed into 4-component spinors.
 - ➔ auxiliary fields are integrated out.
- SUSY phenomenology however has other aspects as well, like solving RG equations to obtain the low energy mass spectrum from some high scale input.
- This needs input beyond tree level.

To infinity and beyond...



- We have reached the point that we have a complete and fully automatized chain from the Lagrangian to the events.

- This chain is however restricted to **tree level** so far.
- **Next goal:** extend this chain to beyond-tree-level information.
 - ➔ **Development branch 1: InSurGe**
'spectrum generator generator'.
 - ➔ **Development branch 2: FR@NLO**
getting ready for automated BSM NLO computation.

Beyond tree-level

InSurGe

SUSY RG equations

- The fully generic MSSM depends on 105 free parameters.
 - ➔ impossible to do phenomenology with such a huge parameter space.
- Very often one assumes an organizing principle are at some high UV scale, where, e.g., couplings and/or masses unify.
- The parameters at the weak scale are then fixed by the renormalization group flow.

$$\beta_g^{(1)} = g^3 [S(\mathcal{R}) - 3C(G)] \quad \beta_M^{(1)} = 2g^2 M [S(\mathcal{R}) - 3C(G)]$$

- **Classical example:** mSUGRA

$$m_0, m_{1/2}, A_0, \tan \beta, \text{sign}(\mu)$$

InSurGe

- The development version of FeynRules allows to extract the one-loop renormalization group equations for generic SUSY models. [A. Alloul, K. de Causmaecker, B. Fuks]

InSurGe

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```
RGE[LSoft, SuperW]
```

InSurGe

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- Starting from the superspace action, the FeynRules computes the RGE's

RGE[LSoft, SuperW]

$$\begin{aligned} \frac{d\mu}{dt} &= \mu \left[-\frac{3g'^2}{80\pi^2} - \frac{3g_w^2}{16\pi^2} + \frac{3}{16\pi^2} \text{Tr}[\mathbf{y}^{d\dagger} \mathbf{y}^d] + \frac{3}{16\pi^2} \text{Tr}[\mathbf{y}^{u\dagger} \mathbf{y}^u] + \frac{1}{16\pi^2} \text{Tr}[\mathbf{y}^{e\dagger} \mathbf{y}^e] \right] \\ \frac{db}{dt} &= b \left[-\frac{3g'^2}{80\pi^2} - \frac{3g_w^2}{16\pi^2} + \frac{3}{16\pi^2} \text{Tr}[\mathbf{y}^{d\dagger} \mathbf{y}^d] + \frac{3}{16\pi^2} \text{Tr}[\mathbf{y}^{u\dagger} \mathbf{y}^u] + \frac{1}{16\pi^2} \text{Tr}[\mathbf{y}^{e\dagger} \mathbf{y}^e] \right] \\ &+ \mu \left[\frac{3g'^2 M_1}{40\pi^2} + \frac{3g_w^2 M_2}{8\pi^2} + \frac{3}{8\pi^2} \text{Tr}[\mathbf{y}^{d\dagger} \mathbf{T}^d] + \frac{3}{8\pi^2} \text{Tr}[\mathbf{y}^{u\dagger} \mathbf{T}^u] + \frac{1}{8\pi^2} \text{Tr}[\mathbf{y}^{e\dagger} \mathbf{T}^e] \right] \end{aligned}$$

InSurGe

- Right now, FeynRules/InSurGe can extract the one and two-loop RGE's for arbitrary (renormalizable) SUSY models.
- These RGE's can then be fed into spectrum generators to obtain the mass spectrum for arbitrary SUSY models.
- In this way, FeynRules/InSurGe turns into a 'spectrum generator generator'.

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- These RGE's can then be fed into spectrum generators to obtain the mass spectrum for arbitrary SUSY models.
- In this way, FeynRules/InSurGe turns into a 'spectrum generator generator'.
- As an example, an interface to SuSpect 3 is being developed that will allow to inject the FeynRules/InSurGe RGE's directly into the SuSpect framework.

```
WriteSuSpectOutput[LSoft, SuperW]
```

InSurGe

- In parallel, InSurGe will come with its own RGE solver.
- Under development right now:
 - ➔ Finding the minimum of the Higgs potential.
 - ➔ Diagonalizing the mass matrices.
- At the same time this solves another problem...

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- Under development right now:
 - ➔ Finding the minimum of the Higgs potential.
 - ➔ Diagonalizing the mass matrices.
- At the same time this solves another problem...
- Remember that FeynRules does not allow to diagonalize the mass matrices automatically.
- InSurGe will solve this problem at the same time (not only for SUSY models).

Mass diagonalization with InSurGe

- In the future, FeynRules will extract the mass matrices automatically from the Lagrangian.

Mass diagonalization with InSurGe

Preliminary results

```
In[57]:= GetMassMatrices[lagr]
```

The Lagrangian you entered contains physical fields, make sure these don't come from a previous field rotation. Calculations have been however performed.

Calculations took 35.8097

seconds to be performed. Results are stored in the variable `MassMatrices`

```
In[58]:= MassMatrices[[7, 2]].MatrixForm[MassMatrices[[7, 3]]].MassMatrices[[7, 4]]
```

```
Out[58]= {bow, wow3, hdw1, huw2} .
```

$$\begin{pmatrix} M_1 & 0 & -\frac{1}{2} i g' v_d & \frac{1}{2} i g' v_u \\ 0 & M_2 & \frac{1}{2} i g_w v_d & -\frac{1}{2} i g_w v_u \\ -\frac{1}{2} i g' v_d & \frac{1}{2} i g_w v_d & 0 & -\mu \\ \frac{1}{2} i g' v_u & -\frac{1}{2} i g_w v_u & -\mu & 0 \end{pmatrix} \cdot \{\mathbf{bow}, \mathbf{wow}_3, \mathbf{hdw}_1, \mathbf{huw}_2\}$$

Mass diagonalization with InSurGe

- In the future, FeynRules will extract the mass matrices automatically from the Lagrangian.
- The mass matrix is then passed on to InSurGe:
 - ➔ a C++ code is written out that diagonalizes the mass matrices and returns a `param_card.dat` with the numerical values of the masses and mixing matrices.
- The C++ code is standalone and can be run outside and independently of Mathematica.

Workflow

$$\mathcal{L} = (D_\mu \phi)^\dagger D^\mu \phi - V(\phi) - \frac{1}{4} F_{\mu\nu} F^{\mu\nu}$$

FeynRules

UFO

Mass matrices
SUSY RGE's

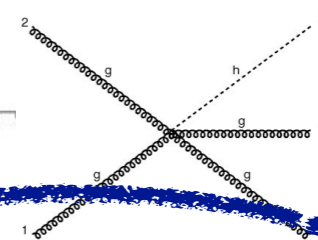
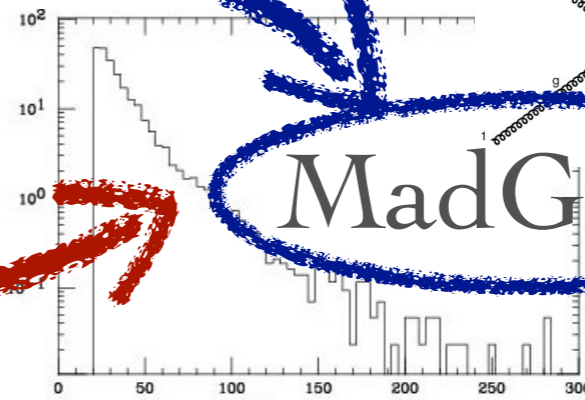
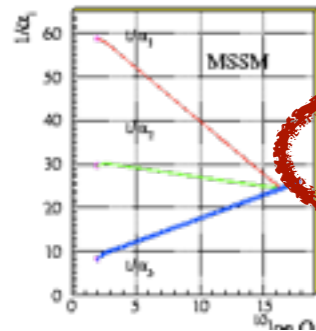
ALOHA

ALOHA
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PYTHON
programming

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MadGraph



Beyond tree-level

FR@NLO

Towards NLO

- We are slowly getting to the point that we have automated tools for NLO computations:
 - ➔ Blackhat
 - ➔ GoSam
 - ➔ FeynArts/FormCalc
 - ➔ Helac-NLO
 - ➔ MadLoops
 - ➔ Rocket
- Most of these codes only do SM processes so far.
- Reason: Beyond LO, we do not only need tree-level Feynman rules, but in addition we need
 - ➔ UV counterterms.
 - ➔ Feynman rules for non cut-constructible part of the amplitude ('R2 terms').

Extraction of UV counterterms

- The (not public) development version of FeynRules already allows to extract counterterm Feynman rules.

```
ExtractCounterterms[1[s,f],{aS,1}]
```

$$\blacktriangleright I_{sf} \rightarrow I_{sf} + \frac{\alpha_s}{4\pi} \left[(\delta Z_{\parallel}^{L(1)})_{ff'} (P_L)_{ss'} + (\delta Z_{\parallel}^{R(1)})_{ff'} (P_R)_{ss'} \right] I_{s'f'}$$

```
ExtractCounterterms[ydo,{{aS,2},{aEW,1}}]
```

$$\blacktriangleright y_d \rightarrow y_d + \frac{\alpha_s}{2\pi} \delta y_d^{(1,0)} + \frac{\alpha}{2\pi} \delta y_d^{(0,1)} + \frac{\alpha_s^2}{4\pi^2} \delta y_d^{(2,0)} + \frac{\alpha_s \alpha}{4\pi^2} \delta y_d^{(1,1)} + \frac{\alpha_s^2 \alpha}{8\pi^3} \delta y_d^{(2,1)}$$

- At the moment, the values of the counterterms for the independent parameters and the fields must still be given by hand.
- In the future: FeynRules will call FeynArts/FormCalc on the fly to obtain the values of the counterterms.

Extraction of counterterms

- **First milestone:** include counterterms into the FeynArts interface.

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➔ FeynArts output **without** counterterms:

$$C[V[1], V[1], V[3], -V[3]] == \{ \{(-I)*gc8, \mathbf{0}\}, \\ \{(-I)*gc8, \mathbf{0}\}, \\ \{(2*I)*gc8, \mathbf{0}\} \}$$

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$$C[V[1], V[1], V[3], -V[3]] == \left\{ \begin{array}{l} \{(-I) * gc8, \mathbf{0}\}, \\ \{(-I) * gc8, \mathbf{0}\}, \\ \{(2 * I) * gc8, \mathbf{0}\} \end{array} \right.$$

➔ FeynArts output **with** counterterms:

$$C[V[1], V[1], V[3], -V[3]] == \left\{ \begin{array}{l} \{(-I) * gc8, \mathbf{(I * (aEW * deltagw01 + aS * deltagw10) * gw * sw^2) / Pi}\}, \\ \{(-I) * gc8, \mathbf{(I * (aEW * deltagw01 + aS * deltagw10) * gw * sw^2) / Pi}\}, \\ \{(2 * I) * gc8, \mathbf{((-2 * I) * (aEW * deltagw01 + aS * deltagw10) * gw * sw^2) / Pi}\} \end{array} \right.$$

Extraction of counterterms

- **First milestone:** include counterterms into the FeynArts interface.
- **Next:**
 - ➔ on-the-fly use of FeynArts to compute the analytical values of the independent one-loop counterterms.
 - ➔ write counterterms out to the UFO.

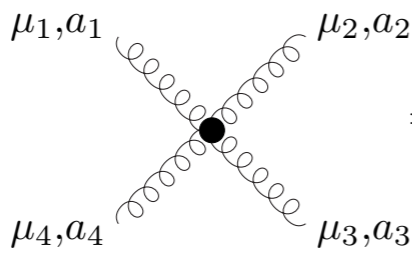
Extraction of counterterms

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```
V_R24G = CTVertex(name = 'V_R24G',
  particles = [ P.G, P.G, P.G, P.G ],
  color = [ 'Tr(1,2)*Tr(3,4)', 'Tr(1,3)*Tr(2,4)', 'Tr(1,4)*Tr(2,3)', \
    'd(-1,1,2)*d(-1,3,4)', 'd(-1,1,3)*d(-1,2,4)', 'd(-1,1,4)*d(-1,2,3)'],
  lorentz = [ L.R2_4G_1234, L.R2_4G_1324, L.R2_4G_1423 ],
  loop_particles = [ [[P.G]], [[P.u],[P.d],[P.c],[P.s]] ],
  couplings = {(0,0,0):C.GC_4GR2_Gluon_delta5,(0,1,0):C.GC_4GR2_Gluon_delta7,(0,2,0):C.GC_4GR2_Gluon_delta7, \
    (1,0,0):C.GC_4GR2_Gluon_delta7,(1,1,0):C.GC_4GR2_Gluon_delta5,(1,2,0):C.GC_4GR2_Gluon_delta7, \
    (2,0,0):C.GC_4GR2_Gluon_delta7,(2,1,0):C.GC_4GR2_Gluon_delta7,(2,2,0):C.GC_4GR2_Gluon_delta5, \
    (3,0,0):C.GC_4GR2_4Struct,(3,1,0):C.GC_4GR2_2Struct,(3,2,0):C.GC_4GR2_2Struct, \
    (4,0,0):C.GC_4GR2_2Struct,(4,1,0):C.GC_4GR2_4Struct,(4,2,0):C.GC_4GR2_2Struct, \
    (5,0,0):C.GC_4GR2_2Struct,(5,1,0):C.GC_4GR2_2Struct,(5,2,0):C.GC_4GR2_4Struct, \
    (0,0,1):C.GC_4GR2_Fermion_delta11,(0,1,1):C.GC_4GR2_Fermion_delta5,(0,2,1):C.GC_4GR2_Fermion_delta5, \
    (1,0,1):C.GC_4GR2_Fermion_delta5,(1,1,1):C.GC_4GR2_Fermion_delta11,(1,2,1):C.GC_4GR2_Fermion_delta5, \
    (2,0,1):C.GC_4GR2_Fermion_delta5,(2,1,1):C.GC_4GR2_Fermion_delta5,(2,2,1):C.GC_4GR2_Fermion_delta11, \
    (3,0,1):C.GC_4GR2_11Struct,(3,1,1):C.GC_4GR2_5Struct,(3,2,1):C.GC_4GR2_5Struct, \
    (4,0,1):C.GC_4GR2_5Struct,(4,1,1):C.GC_4GR2_11Struct,(4,2,1):C.GC_4GR2_5Struct, \
    (5,0,1):C.GC_4GR2_5Struct,(5,1,1):C.GC_4GR2_5Struct,(5,2,1):C.GC_4GR2_11Struct },
  type = 'R2')
```

R2 terms

- All the automatized NLO codes are based, in one way or another, on some unitary-based approach.
- Unitarity, however, does not provide everything, but misses the rational pieces (without cuts).
- Some can be obtained, others (R2) need a different approach.
- R2 terms can be obtained via effective tree-level Feynman rules.



The diagram shows a four-point interaction between gluons. The external lines are labeled with momenta μ_1, a_1 , μ_2, a_2 , μ_3, a_3 , and μ_4, a_4 . The interaction is represented by a central black dot where the four lines meet.

$$\begin{aligned}
 &= -\frac{ig^4 N_{col}}{96\pi^2} \sum_{P(234)} \left\{ \left[\frac{\delta_{a_1 a_2} \delta_{a_3 a_4} + \delta_{a_1 a_3} \delta_{a_4 a_2} + \delta_{a_1 a_4} \delta_{a_2 a_3}}{N_{col}} \right. \right. \\
 &\quad \left. \left. + 4 \text{Tr}(t^{a_1} t^{a_3} t^{a_2} t^{a_4} + t^{a_1} t^{a_4} t^{a_2} t^{a_3}) (3 + \lambda_{HV}) \right. \right. \\
 &\quad \left. \left. - \text{Tr}(\{t^{a_1} t^{a_2}\} \{t^{a_3} t^{a_4}\}) (5 + 2\lambda_{HV}) \right] g_{\mu_1 \mu_2} g_{\mu_3 \mu_4} \right. \\
 &\quad \left. + 12 \frac{N_f}{N_{col}} \text{Tr}(t^{a_1} t^{a_2} t^{a_3} t^{a_4}) \left(\frac{5}{3} g_{\mu_1 \mu_3} g_{\mu_2 \mu_4} - g_{\mu_1 \mu_2} g_{\mu_3 \mu_4} - g_{\mu_2 \mu_3} g_{\mu_1 \mu_4} \right) \right\}
 \end{aligned}$$

[Draggiotis, Garzelli, Papadopoulos, Pittau;
Garzelli, Malamos, Pittau]

What are the R_2 rational terms?

$$\bar{A}(\bar{q}) = \frac{1}{(2\pi)^4} \int d^d \bar{q} \frac{\bar{N}(\bar{q})}{\bar{D}_0 \bar{D}_1 \dots \bar{D}_{m-1}}, \quad \bar{D}_i = (\bar{q} + p_i)^2 - m_i^2$$

$$\bar{N}(\bar{q}) = N(q) + \tilde{N}(\tilde{q}, q, \epsilon)$$

where \bar{X} lives in d dimension, X in 4, \tilde{X} in ϵ .

R_2 definition

$$R_2 \equiv \frac{1}{(2\pi)^4} \int d^d \bar{q} \frac{\tilde{N}(\tilde{q}, q, \epsilon)}{\bar{D}_0 \bar{D}_1 \dots \bar{D}_{m-1}}$$

Finite (< 4 legs) set of vertices computed once for all!

1 SM

- takes approx. 10min

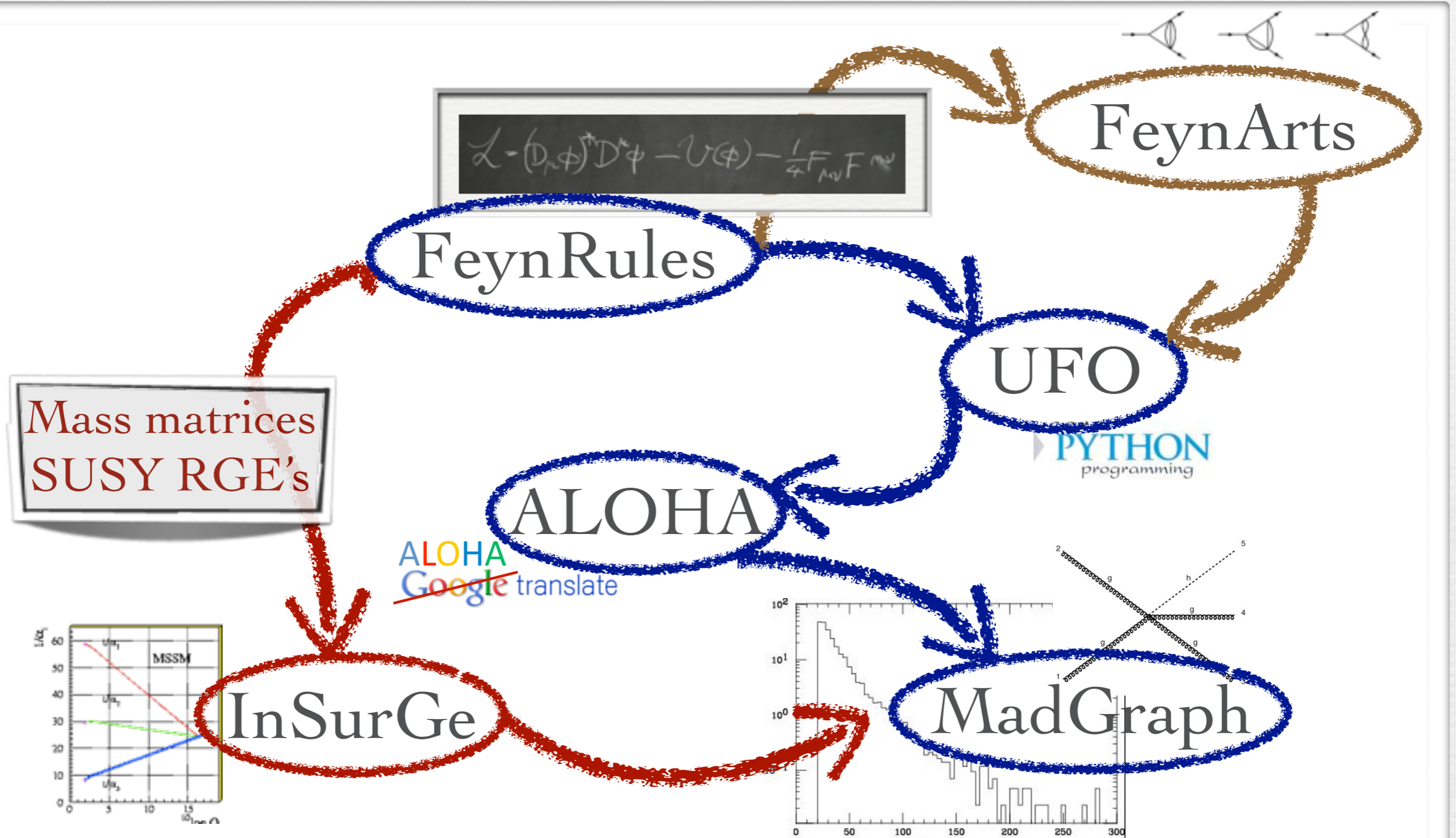
- agree up to $-1/i$ factors (definition of A, ϕ^\pm) with

QCD : Draggiotis, P. et al. JHEP 0904 (2009) 072.

EW : Garzelli, M.V. et al. JHEP 1001 (2010) 040, Erratum-ibid.
1010 (2010) 097.

2 MSSM : On going

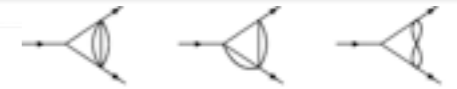
Workflow



Workflow

The only thing left to do is to write a FeynRules model file

$$\mathcal{L} = (D_\mu \phi)^\dagger D^\mu \phi - V(\phi) - \frac{1}{4} F_{\mu\nu} F^{\mu\nu}$$



FeynArts

FeynRules

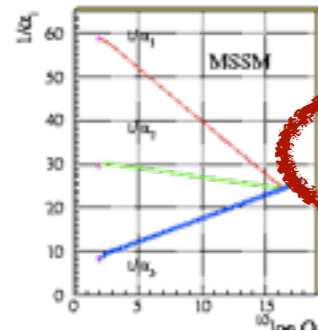
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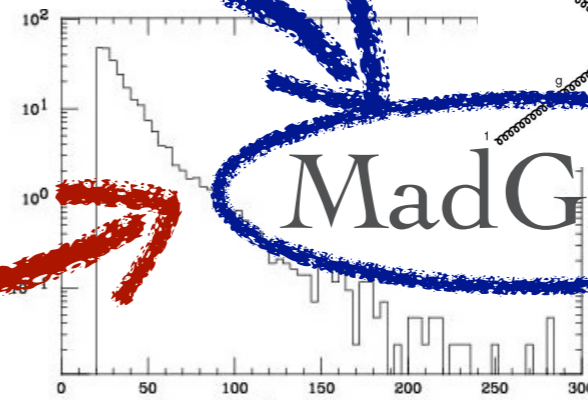
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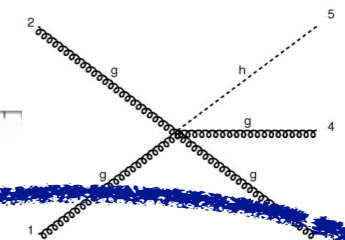
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Galileo

Model building with FeynRules

- All the steps from the Lagrangian to phenomenology are/will be automatized.
- The only ‘bottleneck’ right now is the writing of the FeynRules model file.
 - ➔ Avoid typos in the FeynRules syntax.
 - ➔ A lot of information to provide for complicated models.
- In the future, the Galileo package will alleviate this task.

[N. D. Christensen, D. Salmon, N. Setzer
C. Speckner, Stefanus]
- Galileo will enable the user to specify the symmetry groups and the particle content of the model, and the FeynRules model file is generated automatically.

Galileo

Core Components: Compact Lie Algebras

Compact Lie Algebras: Finished and Tested

- Understands any Compact Lie Algebra, $U(1)$, and \mathbb{Z}_n
 $A_n, B_n, C_n, D_n, G_2, F_4, E_6, E_7, E_8, U(1), \mathbb{Z}_n$
 - Arbitrary Representations
 - All Generators of any Representations
- Product States
 - Gives Singlets

Galileo

Core Components: Fields

Fields: Mostly Finished and Tested

- Understands Spin 0, $\frac{1}{2}$, and 1 fields (fundamental objects)
- Understands Superfields
 - Implemented as a container
 - Chiral Superfields stored as collection of spin-0 and spin- $\frac{1}{2}$
- Output
 - to MathML
 - NEED to FeynRules

Galileo

Core Components: Model

Model: Mostly Finished and Partly Tested

- Understands any direct product group (of Known Groups)
 - Local or Global (\mathbb{Z}_n only global)
- Lagrangian
 - Completely and automatically generated given fields
 - Handles all operators up to dimension 6
 - For SUSY Generates Kahler potential and Superpotential
 - Output to MathML
- Needs
 - Save/Read Model - XML file
 - Write complet FeynRules `.fr` file

Galileo

SM 1 Higgs (Dim 4)

SM with 1 Higgs up to Dimension 4

$$\begin{aligned} & \bar{Q}_l \gamma^\mu D_\mu Q_l + \bar{U}_r \gamma^\mu D_\mu U_r + \bar{D}_r \gamma^\mu D_\mu D_r + \bar{L}_l \gamma^\mu D_\mu L_l + \bar{E}_r \gamma^\mu D_\mu E_r + D_\mu \Phi^* D^\mu \Phi + G_{C\mu\nu} G_C^{\mu\nu} + \theta_0 \epsilon_{\alpha\beta\gamma\delta} G_C^{\alpha\beta} G_C^{\gamma\delta} \\ & + W_{L\mu\nu} W_L^{\mu\nu} + \theta_1 \epsilon_{\alpha\beta\gamma\delta} W_L^{\alpha\beta} W_L^{\gamma\delta} + B_{Y\mu\nu} B_Y^{\mu\nu} \end{aligned}$$

$$\mu_{r0} \Phi \Phi^*$$

$$\lambda_{r1} \Phi \Phi \Phi^* \Phi^*$$

$$y_2 \Phi \bar{U}_r Q_l + y_2^* \Phi^* \bar{Q}_l U_r + y_3 \Phi^* \bar{D}_r Q_l + y_3^* \Phi \bar{Q}_l D_r + y_4 \Phi^* \bar{E}_r L_l + y_4^* \Phi \bar{L}_l E_r$$

Galileo

SM 1 Higgs (Dim 6)

$$\bar{Q}_l \gamma^\mu D_\mu Q_l + \bar{U}_r \gamma^\mu D_\mu U_r + \bar{D}_r \gamma^\mu D_\mu D_r + \bar{L}_l \gamma^\mu D_\mu L_l + \bar{E}_r \gamma^\mu D_\mu E_r + D_\mu \Phi^* D^\mu \Phi + G_{C\mu\nu} G_C^{\mu\nu} + \theta_0 \epsilon_{\alpha\beta\gamma\delta} G_C^{\alpha\beta} G_C^{\gamma\delta} + W_{L\mu\nu} W_L^{\mu\nu} + \theta_1 \epsilon_{\alpha\beta\gamma\delta} W_L^{\alpha\beta} W_L^{\gamma\delta} + B_{Y\mu\nu} B_Y^{\mu\nu}$$

$$\mu_{r0} \Phi \Phi^*$$

$$\lambda_{r1} \Phi \Phi \Phi^* \Phi^*$$

$$y_2 \Phi \bar{U}_r Q_l + y_2^* \Phi^* \bar{Q}_l U_r + y_3 \Phi^* \bar{D}_r Q_l + y_3^* \Phi \bar{Q}_l D_r + y_4 \Phi^* \bar{E}_r L_l + y_4^* \Phi \bar{L}_l E_r +$$

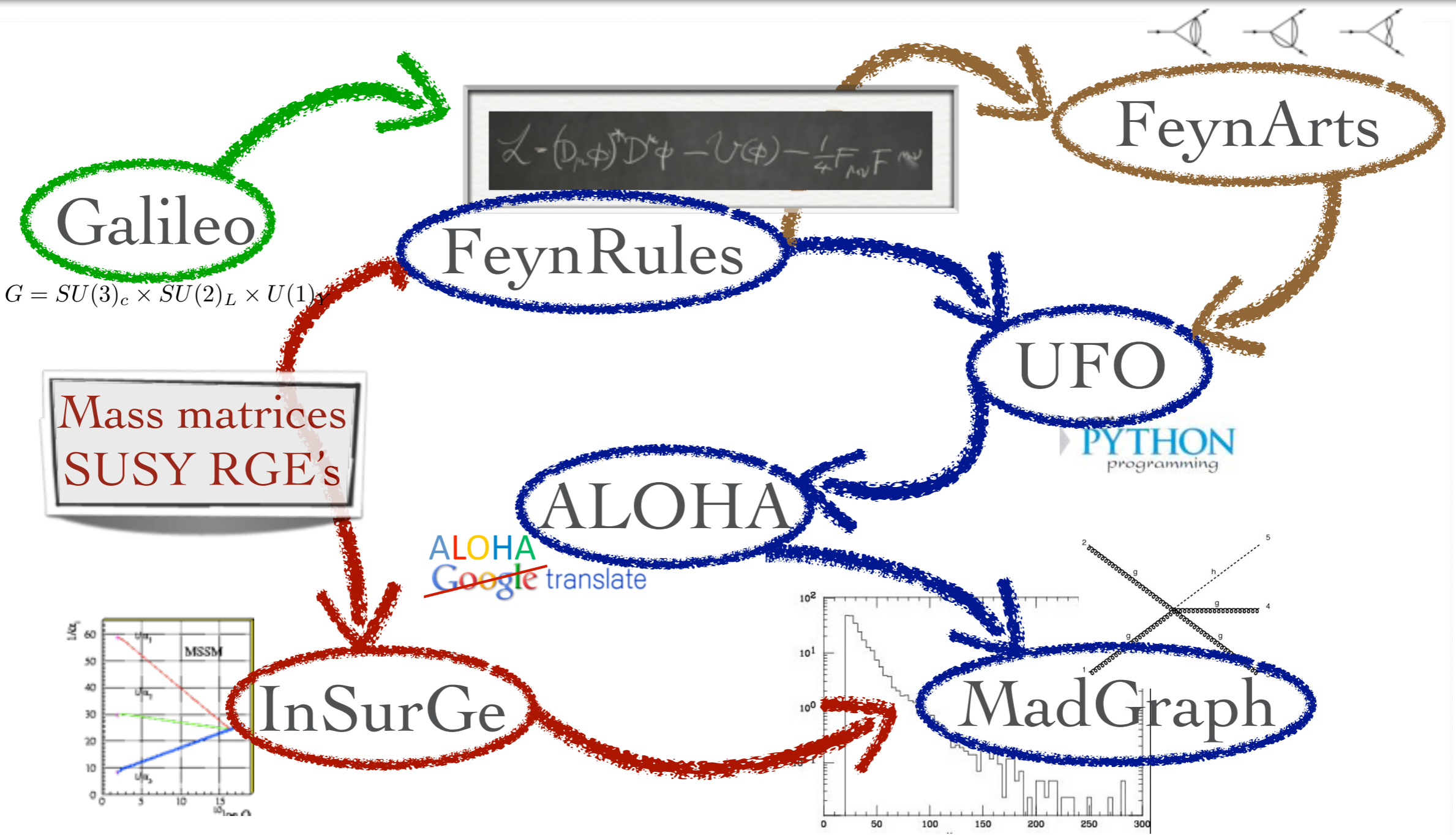
$$y_5 \Phi \Phi \bar{L}_l L_l^c + y_5^* \Phi^* \Phi^* \bar{L}_l^c L_l +$$

$$y_6 \Phi^* \Phi \Phi \bar{U}_r Q_l + y_6^* \Phi \Phi^* \Phi^* \bar{Q}_l U_r + y_7 \Phi^* \Phi^* \Phi \bar{D}_r Q_l + y_7^* \Phi \Phi \Phi^* \bar{Q}_l D_r + y_8 \Phi^* \Phi^* \Phi \bar{E}_r L_l + y_8^* \Phi \Phi \Phi^* \bar{L}_l E_r$$

Galileo

$$\begin{aligned}
& g_{r9} \bar{Q}_l \gamma_\mu Q_l \bar{Q}_l \gamma^\mu Q_l + g_{r10} \bar{Q}_l^c \sigma_{\mu\nu} Q_l \bar{Q}_l \sigma^{\mu\nu} Q_l^c + g_{r11} \bar{Q}_l^c Q_l \bar{Q}_l Q_l^c + g_{r12} \bar{Q}_l \gamma_\mu Q_l \bar{U}_r \gamma^\mu U_r + g_{r13} \bar{Q}_l^c \gamma_\mu U_r \bar{U}_r \gamma^\mu Q_l^c + g_{r14} \bar{Q}_l \sigma_{\mu\nu} U_r \bar{U}_r \\
& \sigma^{\mu\nu} Q_l + g_{r15} \bar{Q}_l U_r \bar{U}_r Q_l + g_{r16} \bar{Q}_l \gamma_\mu Q_l \bar{D}_r \gamma^\mu D_r + g_{r17} \bar{Q}_l^c \gamma_\mu D_r \bar{D}_r \gamma^\mu Q_l^c + g_{r18} \bar{Q}_l \sigma_{\mu\nu} D_r \bar{D}_r \sigma^{\mu\nu} Q_l + g_{r19} \bar{Q}_l D_r \bar{D}_r Q_l + g_{r20} \bar{Q}_l \gamma_\mu L_l \\
& \bar{L}_l \gamma^\mu Q_l + g_{r21} \bar{Q}_l \gamma_\mu Q_l \bar{L}_l \gamma^\mu L_l + g_{r22} \bar{Q}_l^c \sigma_{\mu\nu} L_l \bar{L}_l \sigma^{\mu\nu} Q_l^c + g_{r23} \bar{Q}_l^c L_l \bar{L}_l Q_l^c + g_{r24} \bar{Q}_l \gamma_\mu Q_l \bar{E}_r \gamma^\mu E_r + g_{r25} \bar{Q}_l^c \gamma_\mu E_r \bar{E}_r \gamma^\mu Q_l^c + g_{r26} \bar{Q}_l \\
& \sigma_{\mu\nu} E_r \bar{E}_r \sigma^{\mu\nu} Q_l + g_{r27} \bar{Q}_l E_r \bar{E}_r Q_l + g_{r28} \bar{Q}_l \sigma_{\mu\nu} Q_l^c \bar{L}_l \sigma^{\mu\nu} Q_l^c + g_{r28}^* \bar{Q}_l^c \sigma_{\mu\nu} Q_l \bar{Q}_l^c \sigma^{\mu\nu} L_l + g_{r29} \bar{Q}_l Q_l^c \bar{L}_l Q_l^c + g_{r29}^* \bar{Q}_l^c Q_l \bar{Q}_l^c L_l + g_{r30} \bar{D}_r \\
& \sigma_{\mu\nu} Q_l \bar{U}_r \sigma^{\mu\nu} Q_l + g_{r30}^* \bar{Q}_l \sigma_{\mu\nu} D_r \bar{Q}_l \sigma^{\mu\nu} U_r + g_{r31} \bar{D}_r Q_l \bar{U}_r Q_l + g_{r31}^* \bar{Q}_l D_r \bar{Q}_l U_r + g_{r32} \bar{U}_r \sigma_{\mu\nu} Q_l \bar{D}_r \sigma^{\mu\nu} Q_l + g_{r32}^* \bar{Q}_l \sigma_{\mu\nu} U_r \bar{Q}_l \sigma^{\mu\nu} D_r + \\
& g_{r33} \bar{U}_r Q_l \bar{D}_r Q_l + g_{r33}^* \bar{Q}_l U_r \bar{Q}_l D_r + g_{r34} \bar{Q}_l \sigma_{\mu\nu} Q_l^c \bar{U}_r^c \sigma^{\mu\nu} D_r + g_{r34}^* \bar{Q}_l^c \sigma_{\mu\nu} Q_l \bar{D}_r \sigma^{\mu\nu} U_r^c + g_{r35} \bar{Q}_l Q_l^c \bar{U}_r^c D_r + g_{r35}^* \bar{Q}_l^c Q_l \bar{D}_r U_r^c + g_{r36} \bar{E}_r \\
& \gamma_\mu Q_l^c \bar{U}_r \gamma^\mu Q_l^c + g_{r36}^* \bar{Q}_l^c \gamma_\mu E_r \bar{Q}_l^c \gamma^\mu U_r + g_{r37} \bar{U}_r \gamma_\mu Q_l^c \bar{E}_r \gamma^\mu Q_l^c + g_{r37}^* \bar{Q}_l^c \gamma_\mu U_r \bar{Q}_l^c \gamma^\mu E_r + g_{r38} \bar{Q}_l \sigma_{\mu\nu} Q_l^c \bar{E}_r \sigma^{\mu\nu} U_r^c + g_{r38}^* \bar{Q}_l^c \sigma_{\mu\nu} Q_l \\
& \bar{U}_r^c \sigma^{\mu\nu} E_r + g_{r39} \bar{Q}_l Q_l^c \bar{E}_r U_r^c + g_{r39}^* \bar{Q}_l^c Q_l \bar{U}_r^c E_r + g_{r40} \bar{E}_r \sigma_{\mu\nu} Q_l \bar{U}_r \sigma^{\mu\nu} L_l + g_{r40}^* \bar{Q}_l \sigma_{\mu\nu} E_r \bar{L}_l \sigma^{\mu\nu} U_r + g_{r41} \bar{E}_r Q_l \bar{U}_r L_l + g_{r41}^* \bar{Q}_l E_r \bar{L}_l U_r \\
& + g_{r42} \bar{U}_r \sigma_{\mu\nu} Q_l \bar{E}_r \sigma^{\mu\nu} L_l + g_{r42}^* \bar{Q}_l \sigma_{\mu\nu} U_r \bar{L}_l \sigma^{\mu\nu} E_r + g_{r43} \bar{U}_r Q_l \bar{E}_r L_l + g_{r43}^* \bar{Q}_l U_r \bar{L}_l E_r + g_{r44} \bar{L}_l \sigma_{\mu\nu} Q_l^c \bar{U}_r^c \sigma^{\mu\nu} E_r + g_{r44}^* \bar{Q}_l^c \sigma_{\mu\nu} L_l \bar{E}_r \sigma^{\mu\nu} \\
& U_r^c + g_{r45} \bar{L}_l Q_l^c \bar{U}_r^c E_r + g_{r45}^* \bar{Q}_l^c L_l \bar{E}_r U_r^c + g_{r46} \bar{D}_r \gamma_\mu Q_l^c \bar{L}_l \gamma^\mu U_r^c + g_{r46}^* \bar{Q}_l^c \gamma_\mu D_r \bar{U}_r^c \gamma^\mu L_l + g_{r47} \bar{U}_r \gamma_\mu Q_l^c \bar{L}_l \gamma^\mu D_r^c + g_{r47}^* \bar{Q}_l^c \gamma_\mu U_r \bar{D}_r^c \gamma^\mu L_l \\
& + g_{r48} \bar{L}_l \sigma_{\mu\nu} Q_l^c \bar{D}_r \sigma^{\mu\nu} U_r^c + g_{r48}^* \bar{Q}_l^c \sigma_{\mu\nu} L_l \bar{U}_r^c \sigma^{\mu\nu} D_r + g_{r49} \bar{L}_l Q_l^c \bar{D}_r U_r^c + g_{r49}^* \bar{Q}_l^c L_l \bar{U}_r^c D_r + g_{r50} \bar{L}_l \gamma_\mu Q_l \bar{D}_r \gamma^\mu E_r + g_{r50}^* \bar{Q}_l \gamma_\mu L_l \bar{E}_r \gamma^\mu D_r \\
& + g_{r51} \bar{E}_r \gamma_\mu Q_l^c \bar{D}_r \gamma^\mu L_l + g_{r51}^* \bar{Q}_l^c \gamma_\mu E_r \bar{L}_l \gamma^\mu D_r^c + g_{r52} \bar{D}_r \sigma_{\mu\nu} Q_l \bar{L}_l \sigma^{\mu\nu} E_r + g_{r52}^* \bar{Q}_l \sigma_{\mu\nu} D_r \bar{E}_r \sigma^{\mu\nu} L_l + g_{r53} \bar{D}_r Q_l \bar{L}_l E_r + g_{r53}^* \bar{Q}_l D_r \bar{E}_r L_l + \\
& g_{r54} \bar{U}_r \gamma_\mu U_r \bar{U}_r \gamma^\mu U_r + g_{r55} \bar{U}_r^c \sigma_{\mu\nu} U_r \bar{U}_r \sigma^{\mu\nu} U_r^c + g_{r56} \bar{U}_r U_r \bar{U}_r U_r + g_{r57} \bar{U}_r \gamma_\mu D_r \bar{D}_r \gamma^\mu U_r + g_{r58} \bar{U}_r \gamma_\mu U_r \bar{D}_r \gamma^\mu D_r + g_{r59} \\
& \bar{U}_r^c \sigma_{\mu\nu} D_r \bar{D}_r \sigma^{\mu\nu} U_r^c + g_{r60} \bar{U}_r^c D_r \bar{D}_r U_r^c + g_{r61} \bar{U}_r \gamma_\mu U_r \bar{L}_l \gamma^\mu L_l + g_{r62} \bar{U}_r^c \gamma_\mu L_l \bar{L}_l \gamma^\mu U_r^c + g_{r63} \bar{U}_r \sigma_{\mu\nu} L_l \bar{L}_l \sigma^{\mu\nu} U_r + g_{r64} \bar{U}_r L_l \bar{L}_l U_r + \\
& g_{r65} \bar{U}_r \gamma_\mu E_r \bar{E}_r \gamma^\mu U_r + g_{r66} \bar{U}_r \gamma_\mu U_r \bar{E}_r \gamma^\mu E_r + g_{r67} \bar{U}_r^c \sigma_{\mu\nu} E_r \bar{E}_r \sigma^{\mu\nu} U_r^c + g_{r68} \bar{U}_r^c E_r \bar{E}_r U_r^c + g_{r69} \bar{E}_r \sigma_{\mu\nu} U_r^c \bar{D}_r \sigma^{\mu\nu} U_r^c + g_{r69}^* \bar{U}_r^c \sigma_{\mu\nu} E_r \\
& \bar{U}_r^c \sigma^{\mu\nu} D_r + g_{r70} \bar{E}_r U_r^c \bar{D}_r U_r^c + g_{r70}^* \bar{U}_r^c E_r \bar{U}_r^c D_r + g_{r71} \bar{D}_r \sigma_{\mu\nu} U_r^c \bar{E}_r \sigma^{\mu\nu} U_r^c + g_{r71}^* \bar{U}_r^c \sigma_{\mu\nu} D_r \bar{U}_r^c \sigma^{\mu\nu} E_r + g_{r72} \bar{D}_r U_r^c \bar{E}_r U_r^c + g_{r72}^* \bar{U}_r^c D_r \\
& \bar{U}_r^c E_r + g_{r73} \bar{U}_r \sigma_{\mu\nu} U_r^c \bar{E}_r \sigma^{\mu\nu} D_r^c + g_{r73}^* \bar{U}_r^c \sigma_{\mu\nu} U_r \bar{D}_r^c \sigma^{\mu\nu} E_r + g_{r74} \bar{U}_r U_r^c \bar{E}_r D_r^c + g_{r74}^* \bar{U}_r^c U_r \bar{D}_r^c E_r + g_{r75} \bar{D}_r \gamma_\mu D_r \bar{D}_r \gamma^\mu D_r + g_{r76} \\
& \bar{D}_r^c \sigma_{\mu\nu} D_r \bar{D}_r \sigma^{\mu\nu} D_r^c + g_{r77} \bar{D}_r^c D_r \bar{D}_r D_r^c + g_{r78} \bar{D}_r \gamma_\mu D_r \bar{L}_l \gamma^\mu L_l + g_{r79} \bar{D}_r^c \gamma_\mu L_l \bar{L}_l \gamma^\mu D_r^c + g_{r80} \bar{D}_r \sigma_{\mu\nu} L_l \bar{L}_l \sigma^{\mu\nu} D_r + g_{r81} \bar{D}_r L_l \bar{L}_l D_r + \\
& g_{r82} \bar{D}_r \gamma_\mu E_r \bar{E}_r \gamma^\mu D_r + g_{r83} \bar{D}_r \gamma_\mu D_r \bar{E}_r \gamma^\mu E_r + g_{r84} \bar{D}_r^c \sigma_{\mu\nu} E_r \bar{E}_r \sigma^{\mu\nu} D_r^c + g_{r85} \bar{D}_r^c E_r \bar{E}_r D_r^c + g_{r86} \bar{L}_l \gamma_\mu L_l \bar{L}_l \gamma^\mu L_l + g_{r87} \bar{L}_l^c \sigma_{\mu\nu} L_l \bar{L}_l \sigma^{\mu\nu} \\
& L_l^c + g_{r88} \bar{L}_l^c L_l \bar{L}_l L_l^c + g_{r89} \bar{L}_l \gamma_\mu L_l \bar{E}_r \gamma^\mu E_r + g_{r90} \bar{L}_l^c \gamma_\mu E_r \bar{E}_r \gamma^\mu L_l^c + g_{r91} \bar{L}_l \sigma_{\mu\nu} E_r \bar{E}_r \sigma^{\mu\nu} L_l + g_{r92} \bar{L}_l E_r \bar{E}_r L_l + g_{r93} \bar{E}_r \gamma_\mu E_r \bar{E}_r \gamma^\mu E_r + \\
& g_{r94} \bar{E}_r^c \sigma_{\mu\nu} E_r \bar{E}_r \sigma^{\mu\nu} E_r^c + g_{r95} \bar{E}_r^c E_r \bar{E}_r E_r^c + g_{r96} \Phi \Phi \Phi \Phi^* \Phi^* \Phi^*
\end{aligned}$$

Workflow



Galileo

$$G = SU(3)_c \times SU(2)_L \times U(1)$$

FeynRules

$$\mathcal{L} = (D_\mu \phi)^\dagger D^\mu \phi - V(\phi) - \frac{1}{4} F_{\mu\nu} F^{\mu\nu}$$

FeynArts

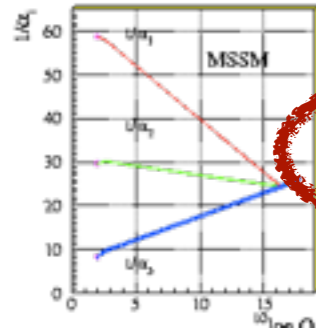
UFO

PYTHON
programming

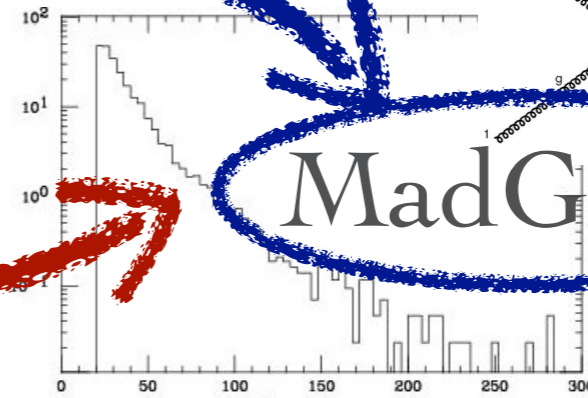
ALOHA

ALOHA
Google translate

Mass matrices
SUSY RGE's



InSurGe



MadGraph