



QCD BASICS FOR ACCURATE LHC PHYSICS

LECTURE IV

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CLAIMS AND AIMS

Four lectures:

- I. Intro and QCD fundamentals
- 2. QCD in the final state
- 3. From accurate QCD to useful QCD
- 4. Advanced QCD with applications at the LHC





HOW DO IMPROVE?

I. We reach NLO and NNLO accuracy

TH-Accurate

2. We include parton showers

EXP-Useful







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- 2. Computationally expensive
- 3. Limited number of particles
- 4. Valid when partons are hard and well separated
- 5. Quantum interference correct
- 6. Needed for multi-jet description





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Shower MC



- I. Resums logs to all orders
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Approaches are complementary: merge them!





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Approaches are complementary: merge them!

Difficulty: avoid double counting, ensure smooth distributions





TWO METHODS

- ME+PS : CKKW and MLM merging
- NLOwPS : MC@NLO and POWHEG





PS ALONE VS MATCHED SAMPLES

In the soft-collinear approximation of Parton Shower MCs, parameters are used to tune the result \Rightarrow Large variation in results (small prediction power)





PS ALONE VS MATCHED SAMPLES

In a matched sample these differences are irrelevant since the behavior at high pt is dominated by the matrix element.





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GOAL FOR ME-PS MERGING/MATCHING

• Regularization of matrix element divergence



- Regularization of matrix element divergence
- Correction of the parton shower for large momenta



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- Smooth jet distributions



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2nd QCD radiation jet in top pair production at the LHC



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Double counting between ME and PS easily avoided using phase space cut between the two: PS below cutoff, ME above cutoff.





- So double counting no problem, but what about getting smooth distributions that are independent of the precise value of Q^c?
- Below cutoff, distribution is given by PS
 need to make ME look like PS near cutoff
- Let's take another look at the PS!













• How does the PS generate the configuration above?







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- Probability for the splitting at t_I is given by $(\Delta_q(t_1,t_0))^2 rac{lpha_s(t_1)}{2\pi} P_{gq}(z)$







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and for the whole tree

$$\frac{(\Delta_q(t_{\rm cut},t_0))^2}{2\Delta_g(t_2,t_1)} \frac{(\Delta_q(u_1,t_2))^2}{(\Delta_q(u_1,t_2))^2} \frac{\alpha_s(t_1)}{2\pi} \frac{P_{gq}(z)}{2\pi} \frac{\alpha_s(t_2)}{2\pi} \frac{P_{qg}(z')}{2\pi}$$







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Sudakov suppression due to disallowing additional radiation above the scale t_{cut}











 $|\mathcal{M}|^2(\hat{s},p_3,p_4,...)$

- To get an equivalent treatment of the corresponding matrix element, do as follows:
 - I. Cluster the event using some clustering algorithmthis gives us a corresponding "parton shower history"
 - 2. Reweight $\boldsymbol{\alpha}_{s}$ in each clustering vertex with the clustering scale $|\mathcal{M}|^{2} \rightarrow |\mathcal{M}|^{2} \frac{\alpha_{s}(t_{1})}{\alpha_{s}(t_{0})} \frac{\alpha_{s}(t_{2})}{\alpha_{s}(t_{0})}$
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 - 3. Use some algorithm to apply the equivalent Sudakov suppression $(\Delta_q(t_{\text{cut}}, t_0))^2 \Delta_g(t_2, t_1) (\Delta_q(_{\text{cut}}, t_2))^2$



MLM MATCHING

[M.L. Mangano, 2002, 2006] [J.Alwall et al 2007, 2008]

• The simplest way to do the Sudakov suppression is to run the shower on the event, starting from t_0 !



- If hardest shower emission scale $k_{T1} > t_{cut}$, reject the event, if all $k_{T1,2,3} < t_{cut}$, keep the event
- The probability for this is so the internal structure of the shower history is ignored. In practice, this approximation is still pretty good.
- Allows matching with any shower, without modifications!





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- We have a number of choices to make in the above procedure. The most important are:
 - I. The clustering scheme used to determine the parton shower history of the ME event
 - 2. What to use for the scale t_0 (factorization scale)
 - 3. How to divide the phase space between parton showers and matrix elements





- We have a number of choices to make in the above procedure. The most important are:
 - I. The clustering scheme used to determine the parton shower history of the ME event
 - 2. What to use for the scale t_0 (factorization scale)
 - 3. How to divide the phase space between parton showers and matrix elements
- In MadGraph and the MadGraph-Pythia interface, there are three different schemes implemented:
 - a. Cone jet scheme (original MLM scheme from AlpGen)
 - b. k_{T} -jet MLM scheme
 - c. 'Shower- k_T ' scheme





I. The default clustering scheme used inside MadGraph to determine the parton shower history is the Durham k_T scheme. For e^+e^- :

$$k_{Tij}^2 = 2\min(E_i^2, E_j^2)(1 - \cos\theta_{ij})$$

and for hadron collisions, the minimum of:

$$k_{Tibeam} = m_i^2 + p_{Ti}^2 = (E_i + p_{zi})(E_i - p_{zi})$$

and

$$k_{Tij}^2 = \min(p_{Ti}^2, p_{Tj}^2) R_{ij}$$

with

$$R_{ij} = 2[\cosh(y_i - y_j) - \cos(\phi_i - \phi_j)] \simeq (\Delta y)^2 + (\Delta \phi)^2$$

2. Find the smallest k_{Tij} (or k_{Tibeam}), combine partons *i* and *j* (or *i* and the beam), and continue until you reach a $2 \rightarrow 2$ (or $2 \rightarrow 1$)





Additional notes:

- MadGraph only allows clustering according to valid diagrams in the process. This means that, e.g., two quarks or quarkantiquark of different flavor are never clustered, and the clustering always gives a physically allowed parton shower history.
- For on-shell s-channel propagators, the clustering value is the invariant mass.
- If there is an on-shell propagator in the diagram, only clustering according to diagrams with this propagator is allowed.





2. The clustering provides a convenient choice for factorization scale t₀:



Cluster back to the 2 \rightarrow 2 (here qq \rightarrow W⁻g) system, and use the W boson transverse mass in that system.

- Special treatment (still beta) for
 - Processes with final-state b quarks that are considered as heavy particles (the 4-flavor scheme)
 - Processes with t-channel color singlet exchange, e.g. weak boson fusion processes.









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 - a. Cone jet MLM scheme:
 - Use cuts in $p_T (p_T^{ME})$ and ΔR between partons in ME
 - Cluster events after parton shower using a cone jet algorithm with the same ΔR and $p_T^{match} > p_T^{ME}$
 - Keep event if all jets are matched to ME partons (i.e., all ME partons are within ΔR of a jet)



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 - b. k_{T} -jet MLM scheme:
 - Use cut in the Durham $k_{\rm T}$ in ME
 - Cluster events after parton shower using the same k_T clustering algorithm into k_T jets with $k_T^{\text{match}} > k_T^{\text{ME}}$
 - Keep event if all jets are matched to ME partons (i.e., all partons are within k_T^{match} to a jet)



- 3. How to divide the phase space between PS and ME: This is where the different schemes differ!
 - c. Shower- k_T scheme:
 - Use cut in the Durham $k_{\rm T}$ in ME
 - After parton shower, get information from the PS generator about the k_T^{PS} of the hardest shower emission
 - Keep event if $k_T^{PS} < k_T^{match}$





SUMMARY OF MLM ALGORITHM

- I. Generate ME events (with different parton multiplicities) using parton-level cuts ($p_T^{ME}/\Delta R$ or k_T^{ME})
- 2. Cluster each event and reweight α_s and PDFs based on the scales in the clustering vertices
- 3. Run the parton shower with starting scale $t_0 = m_T$.
- 4. Check that the number of jets after parton shower is the same as ME partons, and that all jets after parton shower are matched to the ME partons (using one of the schemes in the last slides) at a scale Q^{match}. If yes, keep the event. If no, reject the event. Q^{match} is called the *matching scale*.

One more subtlety: the highest multiplicity sample





HIGHEST MULTIPLICITY SAMPLE

- For the highest jet multiplicity that we generate with the matrix element, we need to allow additional jets above the matching scale Q^{match}., since we will otherwise not get a jet-inclusive description.
- However, we need to reject events with additional jets above the scale of the softest ME parton to avoid double counting.
- For the k_T MLM and shower-kT schemes, the clustering scales of the ME partons are communicated to Pythia using an additional line in the LHE event file written by MadEvent.



How to do matching in MadGraph +Pythia Example: Simulation of $pp \rightarrow W$ with 0, 1, 2 jets (possible on a laptop!)



Matching automatically done when run through MadEvent and Pythia!



HOW TO DO MATCHING IN MADGRAPH +PYTHIA

- By default, k_T -MLM matching is run if xqcut > 0, with the matching scale QCUT = max(xqcut*1.4, xqcut+10)
- For shower-kT, by default QCUT = xqcut
- If you want to change the Pythia setting for matching scale or switch to shower-k_T matching:

```
In pythia_card.dat:
...
! This sets the matching scale, needs to be > xqcut
QCUT = 30
! This switches from kT-MLM to shower-kT matching
! Note that MSTP(81)>=20 needed (pT-ordered shower)
SHOWERKT = T
```



HOW TO DO VALIDATE THE MATCHING

- The matched cross section is found at the end of the Pythia log file
- The matched cross section (for X+0,1,... jets) should be close to the unmatched cross section for the 0-jet sample
- The matching scale (QCUT) should typically be chosen around 1/6-1/3 x hard scale (so xqcut correspondingly lower)
- The differential jet rate plots should be smooth
- When QCUT is varied (within the region of validity), the matched cross section should not vary significantly



- The "differential jet rates" are simply the clustering scales in k_T jet clustering
- The 0→1 diff. jet rate (DJR1) is the p_T of the last remaining jet after clustering
- The I→2 diff. jet rate (DJR2) is the smallest of the p⊤ of the 2nd last remaining jet and the k⊤ between the 2nd jet and the 1st jet
- Note that only radiated jets (not jets from decays) are included in the jet rate plots







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MATCHING RESULTS

W+jets production at the Tevatron for MadGraph+Pythia (kT-jet MLM scheme)



Jet distributions smooth, and stable when we vary the matching scale!

ThikTank on Physics@LHC, 05-09 Dec 2011



TH/EXP COMPARISON AT THE LHC









MATCHING IN NEW PHYSICS PRODUCTION

Alwall, de Visscher, FM [arXiv:0810.5350]

- We know that matching of ME+PS is vital for jet production in SM backgrounds
- But is it relevant for heavy BSM particle production?
 - ➡ Very hard jets from decays
 - Parton showers expected to be more accurate for larger masses
- Using gluino and squark production as example
- Turns out there are many cases where matching is necessary for precise description!





EXAMPLE

- Example: Gluinos that decay to two quarks+LSP with free ratio of gluino/LSP mass
- Special difficulty when decay products nearly massdegenerate with produced particle
- No (small) missing transverse energy in decay







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- Special difficulty when decay products nearly massdegenerate with produced particle
- No (small) missing transverse energy in decay
 - ➡ Need recoil agains ISR jet!







EXAMPLE



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DOUBLE COUNTING OF DECAYS

• Special difficulty in e.g. SUSY matching: Double counting of decays to jets!







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- This has been solved in recent versions of MadGraph 5 by the new "\$" syntax mg5> import model_v4 mssm mg5> generate p p > dr dr~ j j \$ go
- This removes any on-shell gluinos from the event generation (where on-shell is defined as m ± n·Γ with n set by bwcutoff in the run_card.dat)
- The corresponding region is exactly filled if you run gluino production with gluinos decaying to dr j (using the same bwcutoff).









Invariant mass distributions of d_r squark and d quark



d



EXAMPLE: BSM MULTIJET FINAL STATES

pp→Graviton (ADD&RS) +jets pp→X6 +jets de Aquino, Hagiwara, Li, FM [arXiv:1101.5499] Diquark mass (GeV) pb/bin pb/bin MadGraph MadGraph Summed Contributions pp->sextet, matched -3 jets incl. pp->sextet, unmatched 10-2 ets excl.) jets lexel.: void 10⁻³ 10 10⁻² 10-4 10⁻³ 10⁻⁵ 50 0 100 200 250 300 350 40 150 2500 3000 3500 4000 4500 500 P_{T} of jet H_t(4 jets) with P₂>50

New Physics models can be easily included in Matrix Element generators via FeynRules and results automatically for multi-jet inclusive final state obtained at the same level of accuracy that for the SM.





THE POWER OF MATCHING: LOOP EFFECTS IN H+JETS

Alwall, Li, FM [arXiv:1110.1728]



Matched samples can be obtained easily. They agree pretty well with HqT predictions.





