



# Beyond the Standard Model phenomenology with FEYNRULES

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# Outline

1. FEYNRULES in a nutshell
2. Implementing supersymmetric QCD in FEYNRULES
3. Using FEYNRULES with the supersymmetric QCD model
4. Advanced model implementation techniques
5. Summary

# Monte Carlo tools and discoveries at the LHC (I)

## Assumption

There is some new physics to be discovered

## A BSM LHC story

### ◆ *A priori* preparation

- ♣ Viable model building (top-down & bottom-up)
- ♣ Phenomenological studies
- ♣ Prospective collider analyses

### ◆ *A posteriori* reactions to announcements

- ♣ Model building (top-down & bottom-up)
- ♣ Recasting experimental analyses
- ♣ Measurements (precision predictions)

## Predictions for the LHC

### ◆ Option 1: handmade calculations

- ♣ Factorial growth of the number of diagrams
- ♣ Tedious and error prone

### ◆ Option 2: Monte Carlo simulations

- ♣ Easy to use
- ♣ Can include the full collision environment

# Monte Carlo tools and discoveries at the LHC (2)

## Predictions for the LHC

- ◆ ~~Option 1: handmade calculations~~
  - ♣ Factorial growth of the number of diagrams
  - ♣ Tedious and error prone

- ◆ Option 2: Monte Carlo simulations
  - ♣ Easy to use
  - ♣ Can include the full collision environment

### ◆ How to implement a new physics model in a Monte Carlo program?

- ★ Model definition: particles, parameters & vertices ( $\equiv$  Lagrangian)
- ★ To be translated in a programming language, following some conventions, etc.
- ★ Tedious, time-consuming, error prone
- ★ Iterations for all considered tools and models
- ★ Beware of the restrictions of each tool (Lorentz structures, color structures)

- ★ **Highly redundant** (each tool, each model)
- ★ **No-brainer task** (from Feynman rules to code)

FEYNRULES

Systematization  
& automation

# FEYNRULES in a nutshell

## ◆ What is FEYNRULES?

- ❖ A framework to **develop new physics models**
- ❖ **Automatic export** to several Monte Carlo event generators

- ⇒ Facilitate phenomenological investigations of BSM models
- ⇒ Facilitate the confrontation of BSM models to data

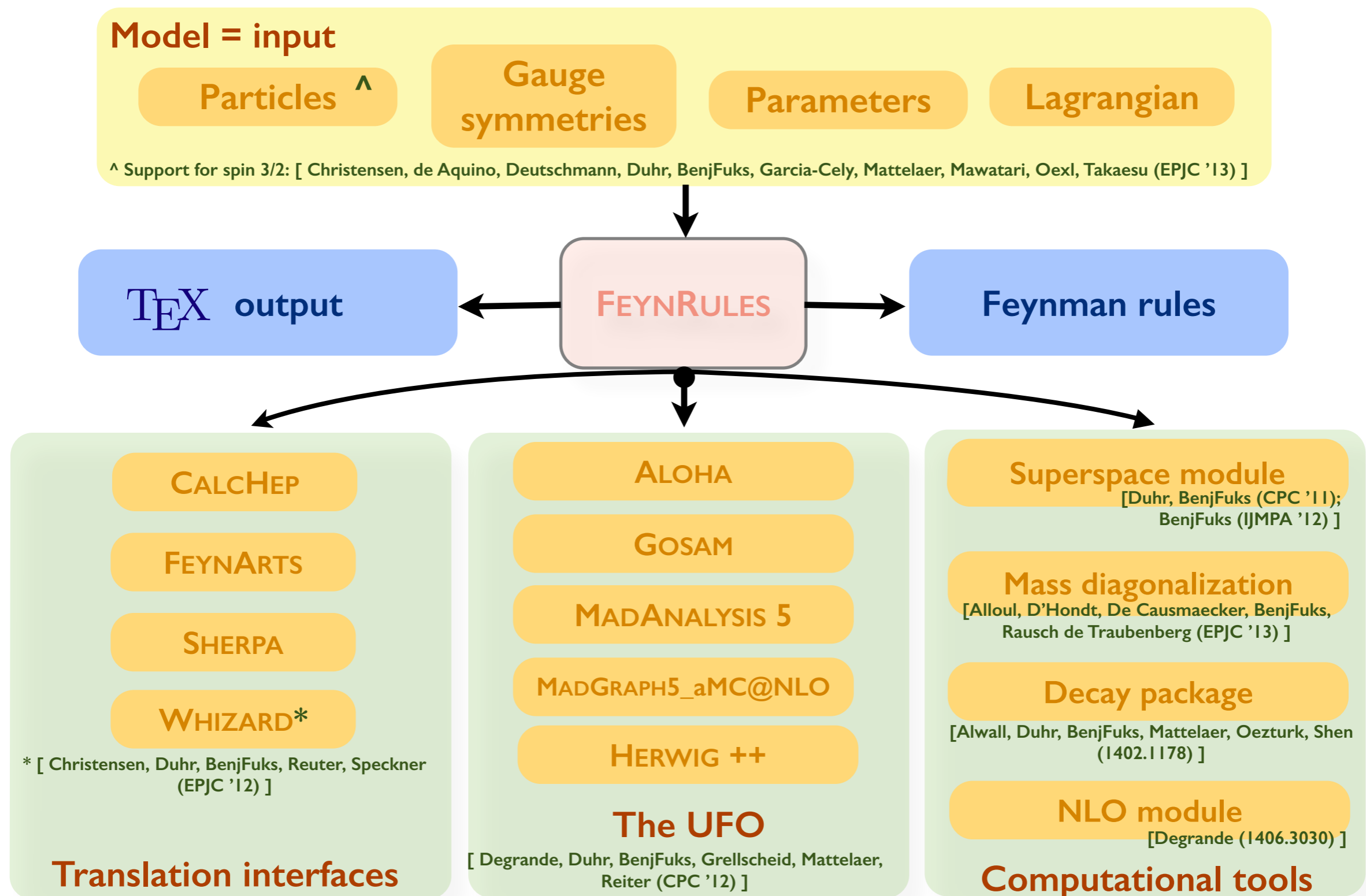
- ❖ **Validation** of an implementation using several of the linked Monte Carlo programs

## ◆ Main features

- ❖ **MATHEMATICA** package
- ❖ Core function: **derives Feynman rules from a Lagrangian**
- ❖ **Requirements**: locality, Lorentz and gauge invariance
- ❖ **Supported fields**: scalar, (two- and four-component) fermion, vector (and ghost), spin-3/2, tensor, superfield

# From FEYNRULES to Monte Carlo tools

[ Christensen, Duhr (CPC '09); Alloul, Christensen, Degrande, Duhr, BenjFuks (CPC '14) ]



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# (Broken) supersymmetric QCD: the model

## ◆ Particle content (simplified for the scope of the lecture)

- ❖ **Two matter supermultiplets** in the fundamental representation of  $SU(3)_c$ 
  - ★ One massive Dirac fermion: a **quark**
  - ★ Two massive scalar fields: a left-handed and a right-handed **squark**
- ❖ **One  $(SU(3)_c)$  gauge supermultiplet**
  - ★ One massive Majorana fermion: a **gluino**
  - ★ One massless gauge boson: the **gluon**

## ◆ The dynamics of the model is embedded in the Lagrangian

$$\begin{aligned}
 \mathcal{L} = & -\frac{1}{4}g_{\mu\nu}g^{\mu\nu} + \frac{i}{2}\bar{g}\not{D}\tilde{g} + D_\mu\tilde{q}_L^\dagger D^\mu\tilde{q}_L + D_\mu\tilde{q}_R^\dagger D^\mu\tilde{q}_R + i\bar{q}\not{D}q \\
 & - m_{\tilde{q}_L}^2\tilde{q}_L^\dagger\tilde{q}_L - m_{\tilde{q}_R}^2\tilde{q}_R^\dagger\tilde{q}_R - m_q\bar{q}q - \frac{1}{2}m_{\tilde{g}}\bar{\tilde{g}}\tilde{g} \\
 & - \frac{g_s^2}{2}\left[-\tilde{q}_L^\dagger T^a\tilde{q}_L + \tilde{q}_R^\dagger T^a\tilde{q}_R\right]\left[-\tilde{q}_L^\dagger T^a\tilde{q}_L + \tilde{q}_R^\dagger T^a\tilde{q}_R\right] \\
 & + \sqrt{2}g_s\left[-\tilde{q}_L^\dagger T^a(\bar{g}^a P_L q) + (\bar{q}P_L\tilde{g}^a)T^a\tilde{q}_R + \text{h.c.}\right]
 \end{aligned}$$

- ❖ Kinetic terms for all fields (first line)
- ❖ Mass terms for the quark, squarks and gluino (second line)
- ❖ (supersymmetric) gauge interactions for all fields (the last two lines of the Lagrangian)



# How to write a FEYNRULES model file?

◆ A FEYNRULES model file is compliant with the MATHEMATICA syntax

◆ It is a `.fr` file containing:

❖ A **preamble**

- ★ Author information
- ★ Model information
- ★ Index definitions

❖ The declaration of the **gauge group**

- ★ Abelian or not
- ★ Representation matrices
- ★ Structure constants
- ★ Coupling constant
- ★ Gauge boson or vector superfield

❖ The declaration of the **fields**

- ★ Names, spins, PDG codes
- ★ Indices, quantum numbers
- ★ Masses, widths
- ★ **Classes and class members**

❖ The declaration of the **parameters**

- ★ External and internal
- ★ Scalar and tensor

❖ A Lagrangian

# The preamble of the model file: general information

## ◆ An electronic signature for the model implementation

- ♣ Important for traceability, documentation, contact with the model authors, etc.
- ♣ Reference publications used can be added
- ♣ Webpage information can be added

```
(* ***** *)
(* ***** *)
(* ***** FeynRules model file: SUSY-QCD ***** *)
(* ***** Author: B. Fuks ***** *)
(* ***** *)
(* ***** *)

(* ***** *)
(* ***** Information ***** *)
(* ***** *)
M$ModelName = "SUSYQCD";

M$Information = {
  Authors    -> {"Benjamin Fuks"},
  Date       -> "16.06.14",
  Version    -> "1.0.0",
  Institutions -> {"CERN / IPHC Strasbourg / U. of Strasbourg"},
  Emails     -> {"benjamin.fuks@iphc.cnrs.fr"}
};
```

# The preamble of the model file: indices

## ◆ The dimension of the indices must be declared

### ✿ In our SUSY-QCD model:

- ★ Fundamental  $SU(3)_C$  indices for the squarks and quark: *Colour*, dimension 3
- ★ Adjoint  $SU(3)_C$  indices for the gluon and gluino: *Gluon*, dimension 8
- ★ Lorentz and spin indices are automatically handled

```
(* ***** *)
(* ***** Indices ***** *)
(* ***** *)
IndexRange[Index[Gluon]] = NoUnfold[Range[8]];
IndexRange[Index[Colour]] = NoUnfold[Range[3]];
```

### ✿ QCD has a special role for Monte Carlo event generators ➤ many special names

- ★ *Colour* and *Gluon* for the indices
- ★ *G* for the gluon field
- ★ *T* for the fundamental representation matrices, *f* and *d* for the structure constants
- ★ *G* (gs will be used) and *aS* for the coupling constants

## ◆ The style of the indices can be specified

### ✿ In our SUSY-QCD model:

- ★ Fundamental  $SU(3)_C$  indices: starts with the letter *m*
- ★ Adjoint  $SU(3)_C$  indices: starts with the letter *a*

```
IndexStyle[Colour, m];
IndexStyle[Gluon, a];
```

# The declaration of the gauge group

## ◆ Each direct factor of the group is declared as an element of the $M\$GaugeGroups$ list

❖ A declaration  $\equiv$  a set of MATHEMATICA replacement rules

❖ In our SUSY-QCD model:

★ We must only declare  $SU(3)_C$ : we choose the name  $SU3C$

```
(* ***** *)
(* ***** Gauge groups ***** *)
(* ***** *)
M$GaugeGroups = {
  SU3C == {
    Abelian          -> False,
    GaugeBoson       -> G,
    CouplingConstant -> gs,
    StructureConstant -> f,
    Representations  -> { {T, Colour} }
  }
};
```

❖ Each rule represents one group property (reminder for QCD: special names exist)

★ *Abelian*: abelian or non-abelian gauge group

★ *GaugeBoson*: the associated gauge boson

★ *CouplingConstant*, *StructureConstant*: coupling and the structure constants

★ *Representation*: list of 2-tuples linking an index (*Colour*) to the symbol of a representation matrix (*T*)

## ◆ Advantages of a proper gauge group declaration

❖ Render the writing of the Lagrangian easier:

★ *Covariant derivatives* ( $DC[field, Lorentz\ index]$ )

★ *Field strength tensors* ( $FS[field, Lorentz\ index\ 1, Lorentz\ index\ 2]$ )

★ Useful for Lagrangian building in superspace (very briefly covered in the last part of this lecture)

➤ Duhr, BenjFuks [CPC 182 (2011) 2404]; BenjFuks [IJMPA 27 (2012) 1230007]

See the manual for more details on gauge groups

# Declaring the gluon field

## ◆ Each field is declared as an element of the $M\$ClassesDescription$ list

- ♣ A declaration  $\equiv$  a set of MATHEMATICA replacement rules
- ♣ In our SUSY-QCD model, we first declare the  $SU(3)_c$  gauge boson:  $G$

```
(* *****)
(* ***** Fields ***** *)
(* *****)
M$ClassesDescription = {
  V[1] == {
    ClassName      -> G,
    SelfConjugate  -> True,
    Indices        -> {Index[Gluon]},
    Mass           -> 0,
    Width          -> 0,
    PDG            -> 21
  },
  ...
}
```

- ♣ Each rule represents a property of the declared field
  - ★ Vector field  $\triangleright$  the label is  $V[I]$  (with  $V$ , and not  $F, S, R, T$ , etc.)
  - ★ **Classname**: defines the symbol to use in the Lagrangian  $\triangleright G$
  - ★ **SelfConjugate**: the gluon is its own antiparticle  $\triangleright True$
  - ★ **Indices**: the gluon lies in the adjoint representation of  $SU(3)_c$ 
    - $\triangleright$  The gluon has been previously set as the gauge boson of  $SU(3)_c$
    - $\triangleright$  Its index (*Gluon*) is internally linked to the adjoint representation of the group
  - ★ Other properties: vanishing mass and widths, PDG code set to 21
  - ★ Not used options: *Unphysical, Definitions, PropagatorLabel, PropagatorType, PropagatorArrow, ParticleName, AntiParticleName, QuantumNumbers*

See the manual for more details on field declarations

# Declaring the gluino field

## ◆ A second element in the $M\$ClassesDescription$ list

```
F[1] == {
  ClassName      -> go,
  SelfConjugate  -> True,
  Indices        -> {Index[Gluon]},
  Mass           -> {Mgo,500},
  Width          -> {Wgo,10},
  PDG            -> 1000021
},
```

### ❖ Differences with the gluon field declaration

- ★ Four-component fermionic field ➤ the label is **F[I]** (with an F)
- ★ **Classname** in the fermion case: defines two symbols to use in the Lagrangian ➤ *go* and *gobar*
- ★ **Mass and width**: we define two symbols and their associated numerical values

## ◆ Second option: two-component spinors (Lagrangians are sometimes easier to write)

```
W[1] == {
  ClassName      -> go|w,
  Unphysical     -> True,
  Chirality      -> Left,
  SelfConjugate  -> False,
  Indices        -> {Index[Gluon]},
},
```

```
F[1] == {
  ClassName      -> go,
  WeylComponents -> gow,
  SelfConjugate  -> True,
  Indices        -> {Index[Gluon]},
  Mass           -> {Mgo,500},
  Width          -> {Wgo,10},
  PDG            -> 1000021
},
```

- ★ Two-component and four-component fermionic fields ➤ the labels are **WV[I]** (with a W) and **F[I]**
- ★ Weyl fermions are **unphysical** and linked to a four-component fermion (**WeylComponents**)
- ★ Several symbols are defined ➤ *go* and *gobar*; *gow* and *gowbar*
- ★ The **chirality** of the Weyl fermion can be specified

See the manual for more details on field declarations

# Declaring the (top) quark field

## ◆ A third element in the $M\$ClassesDescription$ list

```
F[2] == {
  ClassName      -> q,
  SelfConjugate  -> False,
  Indices        -> {Index[Colour]},
  Mass           -> {Mq, 173},
  Width          -> {Wq, 1.50833649},
  PDG            -> 6
},
```

- ✿ Nothing special compared to the other fields
  - ★ Fundamental QCD indices are specified

## ◆ Parenthesis: three generations of up-type quarks ( $Gen$ being the generation index)

```
F[3] == {
  ClassName      -> uq,
  ClassMembers   -> {u, c, t},
  SelfConjugate  -> False,
  Indices        -> {Index[Gen], Index[Colour]},
  FlavorIndex    -> Gen,
  Mass           -> {Mu, {MU, 2.55*^-3}, {MC, 1.42}, {MT, 173}},
  Width          -> {0, 0, {WT, 1.50833649}},
  PDG            -> {2, 4, 6}
},
```

- ★ **ClassMembers**: specify all the members of the class
- ★ **FlavorIndex**: defines which index is the flavor index (efficiency of the code)
- ★ **Mass, Width, PDG**: one for each class member (plus a generic mass symbol)

See the manual for more details on field declarations

# Declaring the (top) squark fields

## ◆ Extra elements in the $M\$ClassesDescription$ list

```
S[1] == {
  ClassName      -> sqL,
  SelfConjugate  -> False,
  Indices        -> {Index[Colour]},
  Unphysical     -> True,
  Definitions    -> {sqL[c_] -> Cos[theta] sq1[c] - Sin[theta] sq2[c]}
},
S[2] == {
  ClassName      -> sqR,
  SelfConjugate  -> False,
  Indices        -> {Index[Colour]},
  Unphysical     -> True,
  Definitions    -> {sqR[c_] -> Sin[theta] sq1[c] + Cos[theta] sq2[c]}
},
```

```
S[3] == {
  ClassName      -> sq1,
  SelfConjugate  -> False,
  Indices        -> {Index[Colour]},
  Mass          -> {Msq1,300},
  Width         -> {Wsq1,10},
  PDG           -> 1000006
},
S[4] == {
  ClassName      -> sq2,
  SelfConjugate  -> False,
  Indices        -> {Index[Colour]},
  Mass          -> {Msq2,800},
  Width         -> {Wsq2,2},
  PDG           -> 2000006
}
```

## ❖ Squark fields mix:

$$\begin{pmatrix} \tilde{q}_1 \\ \tilde{q}_2 \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \tilde{q}_L \\ \tilde{q}_R \end{pmatrix}$$

- ★ Left and right-handed squarks are declared as **unphysical**
- ★ **Definitions** linking gauge- and mass-eigenstates are provided (inversion of the relation above)
  - The rotations will be performed **automatically** by FEYNRULES
  - The Lagrangian can be written in the gauge basis (easier)

See the manual for more details on field declarations



# Parameter declaration

## ◆ The model contains three parameters:

- ❖ The strong coupling constants:  $g_s, \alpha_s$  are both needed (required by the Monte Carlo tools)
- ❖ The squark mixing angle  $\theta$
- ❖ Masses and widths are handled automatically

## ◆ Parameters are declared as elements of the list $M\$Parameters$

- ❖ A declaration  $\equiv$  a set of MATHEMATICA replacement rules
- ❖ In our SUSY-QCD model, we have:

```
(* ***** *)
(* ***** Parameters ***** *)
(* ***** *)
M$Parameters = {
  aS == {
    ParameterType -> External,
    Value         -> 0.1184,
    InteractionOrder -> {QCD, 2}
  },
  gs == {
    ParameterType -> Internal,
    Value         -> Sqrt[4 Pi aS],
    InteractionOrder -> {QCD, 1},
    ParameterName  -> G
  },
  theta == {
    ParameterType -> External,
    Value         -> Pi/4.
  }
};
```

- ★ We have **Internal** and **External** parameters
  - **External**  $\equiv$  free model parameter  $\Rightarrow$  numerical value
  - **Internal**  $\equiv$  dependent model parameter  $\Rightarrow$  formula
- ★ **InteractionOrder**: specific to MADGRAPH  
(more efficient diagram generation)
- ★ **ParameterName**: specific to Monte Carlo tools
- ★ Not used options: *TeX, Definitions, ComplexParameter, Description, BlockName, OrderBlock*
- ★ Specific to tensor parameters: *Indices, Unitary*

See the manual for more details on parameter declarations

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# Implementing the vector Lagrangian

## ◆ Temporary restriction to the gauge content of the theory

- ❖ One  $(SU(3)_c)$  gauge supermultiplet
  - ★ One massive Majorana fermion: a **gluino**
  - ★ One massless gauge boson: the **gluon**

## ◆ The dynamics of the model is embedded in the vector Lagrangian

$$\mathcal{L}_{\text{vector}} = -\frac{1}{4}g_{\mu\nu}g^{\mu\nu} + \frac{i}{2}\bar{\tilde{g}}\not{D}\tilde{g} - \frac{1}{2}m_{\tilde{g}}\bar{\tilde{g}}\tilde{g}$$

- ❖ Kinetic terms for all the gluino and gluon fields
- ❖ Mass terms for the gluino
- ❖ Gauge interactions for both fields (embedded into gauge covariant objects)

## ◆ The implementation in FEYNRULES is easy (cf. predefined functions linked to the gauge group)

- ❖ Option 1 (left): all non-Lorentz indices understood (FEYNRULES takes care of reintroducing them)
- ❖ Option 2 (right): all indices explicit

```
LVector1 := -1/4 FS[G,mu,nu,a] FS[G,mu,nu,a] +
            I/2 gobar.Ga[mu].DC[go,mu] -
            1/2 Mgo gobar.go;
```

```
LVector2 := -1/4 FS[G,mu,nu,a] FS[G,mu,nu,a] +
            I/2 Ga[mu,s1,s2] gobar[s1,a].DC[go[s2,a],mu] -
            1/2 Mgo gobar[s1,a].go[s1,a];
```

# Starting a MATHEMATICA session (I)

## ◆ Step I: loading FEYNRULES into MATHEMATICA

- ♣ Setting up the FEYNRULES path
- ♣ Loading FEYNRULES itself

```
In[2]:= $FeynRulesPath = SetDirectory["~/Work/tools/FeynRules/branch/feynrules-current"];
<< FeynRules`

- FeynRules -
Version: 2.0.25 (22 April 2014).
Authors: A. Alloul, N. Christensen, C. Degrande, C. Duhr, B. Fuks

Please cite:
- arXiv:1310.1921;
- Comput.Phys.Commun.180:1614-1641,2009 (arXiv:0806.4194).

http://feynrules.phys.ucl.ac.be

The FeynRules palette can be opened using the command FRPalette[].
```

## ◆ The output

- ♣ Information on FEYNRULES, the authors, the version number, references, etc.

# Starting a MATHEMATICA session (2)

## ◆ Step2: loading the SUSY-QCD implementation into FEYNRULES

- ♣ Moving to the right directory
- ♣ Loading the model itself

```
In[4]:= SetDirectory[NotebookDirectory[]];
LoadModel["susyqcd_nthu.fr"];

This model implementation was created by
Benjamin Fuks
Model Version: 1.0.0
For more information, type ModelInformation[].

- Loading particle classes.
- Loading gauge group classes.
- Loading parameter classes.

Model SUSYQCD loaded.
```

## ◆ The output

- ♣ Information of the model file preamble are printed to the screen

# Checking the implementation: comparing with books

## ◆ Printing the vector Lagrangian and comparing with the textbook expression

$$\mathcal{L}_{\text{vector}} = -\frac{1}{4}g_{\mu\nu}g^{\mu\nu} + \frac{i}{2}\bar{g}\not{D}g - \frac{1}{2}m_{\bar{g}}\bar{g}g$$

```
LVector1 := -1/4 FS[G,mu,nu,a] FS[G,mu,nu,a] +
            I/2 gobar.Ga[mu].DC[go,mu] -
            1/2 Mgo gobar.go;
```

```
In[8]:= LVector1
```

```
Out[8]= -1/2 Mgo go . go + 1/2 i go . Y^mu . (partial_mu [go] - i gs FSU3C^a$471 . go G_mu,a$471) -
        1/4 (-partial_nu [G_mu,a] + partial_mu [G_nu,a] + gs f_a,bb$469,cc$469 G_mu,bb$469 G_nu,cc$469)
        (-partial_nu [G_mu,a] + partial_mu [G_nu,a] + gs f_a,bb$470,cc$470 G_mu,bb$470 G_nu,cc$470)
```

## ◆ The output

♣ Covariant derivatives and field strength tensors have been automatically evaluated

## ◆ Restoring all indices automatically

```
In[9]:= ExpandIndices[LVector1]
```

```
Out[9]= -1/4 partial_nu [G_mu,a]^2 + 1/2 partial_nu [G_mu,a] partial_mu [G_nu,a] - 1/4 partial_mu [G_nu,a]^2 - 1/2 Mgo go_{i1$478,i2$478} . go_{i1$478,i2$478} + 1/4 gs partial_nu [G_mu,a] f_a,bb$473,cc$473 G_mu,bb$473 G_nu,cc$473 -
        1/4 gs partial_mu [G_nu,a] f_a,bb$473,cc$473 G_mu,bb$473 G_nu,cc$473 + 1/4 gs partial_nu [G_mu,a] f_a,bb$474,cc$474 G_mu,bb$474 G_nu,cc$474 -
        1/4 gs partial_mu [G_nu,a] f_a,bb$474,cc$474 G_mu,bb$474 G_nu,cc$474 - 1/4 gs^2 f_a,bb$473,cc$473 f_a,bb$474,cc$474 G_mu,bb$473 G_mu,bb$474 G_nu,cc$473 G_nu,cc$474 +
        1/4 i go_{i$478,i1$480} . partial_mu [go_{j$478,i1$480}] Y_{i$478,j$478}^mu - 1/4 i gs go_{i$481,i$480} . go_{j$481,j$480} f_a$475,i$480,j$480 G_mu,a$475 Y_{i$481,j$481}^mu -
        1/4 i partial_mu [go_{j$478,i1$480}] . go_{i$478,i1$480} Y_{j$478,i$478}^mu + 1/4 i gs go_{j$481,j$480} . go_{i$481,i$480} f_a$475,i$480,j$480 G_mu,a$475 Y_{j$481,i$481}^mu
```

# Checking the implementation: hermiticity, normalization

## ◆ The Lagrangian must be Hermitian

$$\mathcal{L}_{\text{vector}} = -\frac{1}{4}g_{\mu\nu}g^{\mu\nu} + \frac{i}{2}\bar{\tilde{g}}\not{D}\tilde{g} - \frac{1}{2}m_{\tilde{g}}\bar{\tilde{g}}\tilde{g}$$

```
In[10]:= CheckHermiticity[LVector1];
```

Checking for hermiticity by calculating the Feynman rules contained in L-HC[L].

If the lagrangian is hermitian, then the number of vertices should be zero.

**Starting Feynman rule calculation.**

Expanding the Lagrangian...

No vertices found.

0 vertices obtained.

The lagrangian is hermitian.

## ◆ The Lagrangian must be canonically normalized

```
In[11]:= CheckKineticTermNormalisation[LVector1];
```

Neglecting all terms with more than 2 particles.

All kinetic terms are diagonal.

All kinetic terms are correctly normalized.

- ❖ Other methods: *CheckDiagonalQuadraticTerms*, *CheckDiagonalKineticTerms*, *CheckDiagonalMassTerms*

# Checking the implementation: the mass spectrum

## ◆ Checks at the level of the masses can be performed

- ♣ The masses can be extracted from the Lagrangian
- ♣ The masses are given when particles are declared
- ♣ They should match

```
In[12]:= CheckMassSpectrum[LVector1]
```

```
Neglecting all terms with more than 2 particles.
```

```
All mass terms are diagonal.
```

```
Getting mass spectrum.
```

```
Checking for less than 0.1% agreement with model file values.
```

```
Out[12]//TableForm=
```

Particle	Analytic value	Numerical value	Model-file value
go	Mgo	500.	500.



# The Feynman rules (I)

## ◆ Extract all N-point interactions from the Lagrangian (with N>2)

```
In[14]:= FeynmanRules[LVector1, ScreenOutput -> True];
Starting Feynman rule calculation.
Expanding the Lagrangian...
Collecting the different structures that enter the vertex.
3 possible non-zero vertices have been found -> starting the computation: 3 / 3.
3 vertices obtained.
( * * * * * )
Vertex 1
Particle 1 : Vector , G
Particle 2 : Vector , G
Particle 3 : Vector , G
Vertex:
gs f_{a_1, a_2, a_3} p_1^{\mu_3} \eta_{\mu_1, \mu_2} - gs f_{a_1, a_2, a_3} p_2^{\mu_3} \eta_{\mu_1, \mu_2} - gs f_{a_1, a_2, a_3} p_1^{\mu_2} \eta_{\mu_1, \mu_3} +
gs f_{a_1, a_2, a_3} p_3^{\mu_2} \eta_{\mu_1, \mu_3} + gs f_{a_1, a_2, a_3} p_2^{\mu_1} \eta_{\mu_2, \mu_3} - gs f_{a_1, a_2, a_3} p_3^{\mu_1} \eta_{\mu_2, \mu_3}
```

- ❖ Three vertices are found
- ❖ The expression above consists of the **triple gluon vertex**
- ❖ The index  $a_i$  is related to the  $i^{\text{th}}$  particle  
(reminder:  $a$  is the index style for the adjoint  $SU(3)_C$  indices)
- ❖ The index  $\mu_i$  is the Lorentz index of the  $i^{\text{th}}$  (vector) particle

# The Feynman rules (2)

## ◆ The second vertex: the quartic gluon interaction

Vertex 2

Particle 1 : Vector , G

Particle 2 : Vector , G

Particle 3 : Vector , G

Particle 4 : Vector , G

Vertex:

$$\begin{aligned}
 & i g_s^2 f_{a_1, a_3, \text{Gluon}\$l} f_{a_2, a_4, \text{Gluon}\$l} \eta_{\mu_1, \mu_4} \eta_{\mu_2, \mu_3} + i g_s^2 f_{a_1, a_2, \text{Gluon}\$l} f_{a_3, a_4, \text{Gluon}\$l} \eta_{\mu_1, \mu_4} \eta_{\mu_2, \mu_3} + \\
 & i g_s^2 f_{a_1, a_4, \text{Gluon}\$l} f_{a_2, a_3, \text{Gluon}\$l} \eta_{\mu_1, \mu_3} \eta_{\mu_2, \mu_4} - i g_s^2 f_{a_1, a_2, \text{Gluon}\$l} f_{a_3, a_4, \text{Gluon}\$l} \eta_{\mu_1, \mu_3} \eta_{\mu_2, \mu_4} - \\
 & i g_s^2 f_{a_1, a_4, \text{Gluon}\$l} f_{a_2, a_3, \text{Gluon}\$l} \eta_{\mu_1, \mu_2} \eta_{\mu_3, \mu_4} - i g_s^2 f_{a_1, a_3, \text{Gluon}\$l} f_{a_2, a_4, \text{Gluon}\$l} \eta_{\mu_1, \mu_2} \eta_{\mu_3, \mu_4}
 \end{aligned}$$

- ❖ The index  $a_i$  is related to the  $i^{\text{th}}$  particle  
(reminder:  $a$  is the index style for the adjoint  $SU(3)_C$  indices)
- ❖ The index  $\text{Gluon}\$l$  is repeated  $\Rightarrow$  **summed index**
- ❖ The index  $\mu_i$  is the Lorentz index of the  $i^{\text{th}}$  (vector) particle

# The Feynman rules (3)

## ◆ The third vertex: the gluon-guino interaction

Vertex 3

Particle 1 : Majorana , go

Particle 2 : Majorana , go

Particle 3 : Vector , G

Vertex:

$$g_s f_{a_1, a_2, a_3} \gamma_{s_1, s_2}^{\mu_3}$$

- ❖ The index  $a_i$  is related to the  $i^{\text{th}}$  particle  
(reminder:  $a$  is the index style for the adjoint  $SU(3)_C$  indices)
- ❖ The index  $\mu_i$  is the Lorentz index of the  $i^{\text{th}}$  (vector) particle
- ❖ The index  $s_i$  is the spin index of the  $i^{\text{th}}$  (fermionic) particle

## ◆ The function `FeynmanRules` has many options

- ❖ Restriction on the interactions (`MaxParticles`, `MaxCanonicalDimension`, etc.)
- ❖ Selection of specific particles (`Free`, `Contains`, etc.)
- ❖ `ScreenOutput`: displaying the vertices to the screen or not
- ❖ `FlavorExpand`: perform a flavor expansion (otherwise, classes are used)

See the manual for more details on the function `FeynmanRules`

# Implementing the matter Lagrangian (I)

## ◆ Matter fields

❖ **Two matter supermultiplets** in the fundamental representation of  $SU(3)_c$

★ One massive Dirac fermion: a **quark**

★ Two mixing massive scalar fields: two **squark**

★ Gauge interactions with the  $SU(3)_c$  **gauge supermultiplet**

$$\begin{pmatrix} \tilde{q}_1 \\ \tilde{q}_2 \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \tilde{q}_L \\ \tilde{q}_R \end{pmatrix}$$

## ◆ The dynamics of the model is embedded in the matter Lagrangian

$$\begin{aligned} \mathcal{L}_{\text{matter}} = & D_\mu \tilde{q}_L^\dagger D^\mu \tilde{q}_L + D_\mu \tilde{q}_R^\dagger D^\mu \tilde{q}_R + i \bar{q} \not{D} q - m_{\tilde{q}_i}^2 \tilde{q}_i^\dagger \tilde{q}_i - m_q \bar{q} q \\ & - \frac{g_s^2}{2} \left[ -\tilde{q}_L^\dagger T^a \tilde{q}_L + \tilde{q}_R^\dagger T^a \tilde{q}_R \right] \left[ -\tilde{q}_L^\dagger T^a \tilde{q}_L + \tilde{q}_R^\dagger T^a \tilde{q}_R \right] \\ & + \sqrt{2} g_s \left[ -\tilde{q}_L^\dagger T^a (\tilde{g}^a P_L q) + (\bar{q} P_L \tilde{g}^a) T^a \tilde{q}_R \right] + \text{h.c.} \end{aligned}$$

❖ Kinetic and gauge interaction terms (half of the first line) in the gauge basis

❖ Mass terms (second half of the first line) in the mass basis (because easier to implement)

❖ D-terms (second line)

❖ Supersymmetric gauge quark-squark-gluino interactions (fourth line)

# Implementing the matter Lagrangian (2)

## ◆ Kinetic and gauge interactions

$$\mathcal{L}_{\text{matter}} = D_\mu \tilde{q}_L^\dagger D^\mu \tilde{q}_L + D_\mu \tilde{q}_R^\dagger D^\mu \tilde{q}_R + i \bar{q} \not{D} q - m_{\tilde{q}_i}^2 \tilde{q}_i^\dagger \tilde{q}_i - m_q \bar{q} q$$

## ◆ The implementation in FEYNRULES is easy (cf. predefined functions linked to the gauge group)

```
Lkin := DC[sqLbar[cc],mu] DC[sqL[cc],mu] +
        DC[sqRbar[cc],mu] DC[sqR[cc],mu] +
        I qbar.Ga[mu].DC[q,mu] -
        Mq qbar.q -
        Msq1^2 sq1bar[cc] sq1[cc] -
        Msq2^2 sq2bar[cc] sq2[cc];
```

❖ Repeated indices are summed

## ◆ We can compute the Feynman rules

```
In[9]:= Simplify[FeynmanRules[Lkin]] // MatrixForm
```

Starting Feynman rule calculation.

Expanding the Lagrangian...

Collecting the different structures that enter the vertex.

9 possible non-zero vertices have been found -> starting the computation: 9 / 9.

5 vertices obtained.

Out[9]/MatrixForm=

$$\left( \begin{array}{l} \{G, 1\}, \{sq1, 2\}, \{sq1^\dagger, 3\} \\ \{G, 1\}, \{sq2, 2\}, \{sq2^\dagger, 3\} \\ \{G, 1\}, \{G, 2\}, \{sq1, 3\}, \{sq1^\dagger, 4\} \\ \{G, 1\}, \{G, 2\}, \{sq2, 3\}, \{sq2^\dagger, 4\} \\ \{\bar{q}, 1\}, \{q, 2\}, \{G, 3\} \end{array} \right) \begin{array}{l} i g s (p_2^{\mu_1} - p_3^{\mu_1}) T_{m_3, m_2}^{a_1} \\ i g s (p_2^{\mu_1} - p_3^{\mu_1}) T_{m_3, m_2}^{a_1} \\ i g s^2 \eta_{\mu_1, \mu_2} (T_{m_4, \text{Colour}\$1}^{a_1} T_{\text{Colour}\$1, m_3}^{a_2} + T_{\text{Colour}\$1, m_3}^{a_1} T_{m_4, \text{Colour}\$1}^{a_2}) \\ i g s^2 \eta_{\mu_1, \mu_2} (T_{m_4, \text{Colour}\$1}^{a_1} T_{\text{Colour}\$1, m_3}^{a_2} + T_{\text{Colour}\$1, m_3}^{a_1} T_{m_4, \text{Colour}\$1}^{a_2}) \\ i g s \gamma_{s_1, s_2}^{\mu_3} T_{m_1, m_2}^{a_3} \end{array}$$

❖ Field rotations have been performed automatically

# Implementing the matter Lagrangian (3)

## ◆ The D-terms

$$\mathcal{L}_{\text{matter}} = -\frac{g_s^2}{2} \left[ -\tilde{q}_L^\dagger T^a \tilde{q}_L + \tilde{q}_R^\dagger T^a \tilde{q}_R \right] \left[ -\tilde{q}_L^\dagger T^a \tilde{q}_L + \tilde{q}_R^\dagger T^a \tilde{q}_R \right]$$

## ◆ The implementation in FEYNRULES is easy

```
LD := -1/2 gs^2 *
      (sqRbar[cc1] T[a,cc1,cc2] sqR[cc2] - sqLbar[cc1] T[a,cc1,cc2] sqL[cc2]) *
      (sqRbar[cc3] T[a,cc3,cc4] sqR[cc4] - sqLbar[cc3] T[a,cc3,cc4] sqL[cc4]);
```

- ❖ Repeated indices (cc1, cc2, cc3, cc4) are summed
- ❖ A single index can only be used twice

## ◆ We can compute the Feynman rules

$\{\{sq1, 1\}, \{sq1, 2\}, \{sq1^\dagger, 3\}, \{sq1^\dagger, 4\}\}$	$-i g_s^2 \cos[2 \theta]^2 (T_{m_3, m_2}^{\text{Gluon} \$1} T_{m_4, m_1}^{\text{Gluon} \$1} + T_{m_3, m_1}^{\text{Gluon} \$1} T_{m_4, m_2}^{\text{Gluon} \$1})$
$\{\{sq1, 1\}, \{sq1^\dagger, 2\}, \{sq1^\dagger, 3\}, \{sq2, 4\}\}$	$\frac{1}{2} i g_s^2 \sin[4 \theta] (T_{m_2, m_4}^{\text{Gluon} \$1} T_{m_3, m_1}^{\text{Gluon} \$1} + T_{m_2, m_1}^{\text{Gluon} \$1} T_{m_3, m_4}^{\text{Gluon} \$1})$
$\{\{sq1^\dagger, 1\}, \{sq1^\dagger, 2\}, \{sq2, 3\}, \{sq2, 4\}\}$	$-i g_s^2 \sin[2 \theta]^2 (T_{m_1, m_4}^{\text{Gluon} \$1} T_{m_2, m_3}^{\text{Gluon} \$1} + T_{m_1, m_3}^{\text{Gluon} \$1} T_{m_2, m_4}^{\text{Gluon} \$1})$
$\{\{sq1, 1\}, \{sq1, 2\}, \{sq1^\dagger, 3\}, \{sq2^\dagger, 4\}\}$	$\frac{1}{2} i g_s^2 \sin[4 \theta] (T_{m_3, m_2}^{\text{Gluon} \$1} T_{m_4, m_1}^{\text{Gluon} \$1} + T_{m_3, m_1}^{\text{Gluon} \$1} T_{m_4, m_2}^{\text{Gluon} \$1})$
$\{\{sq1, 1\}, \{sq1^\dagger, 2\}, \{sq2, 3\}, \{sq2^\dagger, 4\}\}$	$-i g_s^2 (\sin[2 \theta]^2 T_{m_2, m_3}^{\text{Gluon} \$1} T_{m_4, m_1}^{\text{Gluon} \$1} - \cos[2 \theta]^2 T_{m_2, m_1}^{\text{Gluon} \$1} T_{m_4, m_3}^{\text{Gluon} \$1})$
$\{\{sq1^\dagger, 1\}, \{sq2, 2\}, \{sq2, 3\}, \{sq2^\dagger, 4\}\}$	$-\frac{1}{2} i g_s^2 \sin[4 \theta] (T_{m_1, m_3}^{\text{Gluon} \$1} T_{m_4, m_2}^{\text{Gluon} \$1} + T_{m_1, m_2}^{\text{Gluon} \$1} T_{m_4, m_3}^{\text{Gluon} \$1})$
$\{\{sq1, 1\}, \{sq1, 2\}, \{sq2^\dagger, 3\}, \{sq2^\dagger, 4\}\}$	$-i g_s^2 \sin[2 \theta]^2 (T_{m_3, m_2}^{\text{Gluon} \$1} T_{m_4, m_1}^{\text{Gluon} \$1} + T_{m_3, m_1}^{\text{Gluon} \$1} T_{m_4, m_2}^{\text{Gluon} \$1})$
$\{\{sq1, 1\}, \{sq2, 2\}, \{sq2^\dagger, 3\}, \{sq2^\dagger, 4\}\}$	$-\frac{1}{2} i g_s^2 \sin[4 \theta] (T_{m_3, m_2}^{\text{Gluon} \$1} T_{m_4, m_1}^{\text{Gluon} \$1} + T_{m_3, m_1}^{\text{Gluon} \$1} T_{m_4, m_2}^{\text{Gluon} \$1})$
$\{\{sq2, 1\}, \{sq2, 2\}, \{sq2^\dagger, 3\}, \{sq2^\dagger, 4\}\}$	$-i g_s^2 \cos[2 \theta]^2 (T_{m_3, m_2}^{\text{Gluon} \$1} T_{m_4, m_1}^{\text{Gluon} \$1} + T_{m_3, m_1}^{\text{Gluon} \$1} T_{m_4, m_2}^{\text{Gluon} \$1})$

- ❖ All nine vertices automatically derived from the (very compact) Lagrangian

# Implementing the matter Lagrangian (4)

## ◆ The squark-quark-gluino interactions

$$\mathcal{L}_{\text{matter}} = \sqrt{2}g_s \left[ -\tilde{q}_L^\dagger T^a (\tilde{g}^a P_L q) + (\bar{q} P_L \tilde{g}^a) T^a \tilde{q}_R \right] + \text{h.c.}$$

## ◆ The implementation in FEYNRULES is again easy

```
Lgosqq := Sqrt[2] gs ProjM[s1,s2] * (
  - sqLbar[cc1] T[a,cc1,cc2] gobar[s1,a].q[s2,cc2] +
  qbar[s1,cc1].go[s2,a] T[a,cc1,cc2] sqR[cc2]);
```

- ❖ All indices must this time be explicit (the scalars cannot be included in a fermion chain)
- ❖ *ProjM* denotes the left-handed chirality projector (*ProjP* is the right-handed one)

## ◆ We can compute the Feynman rules (including the Hermitian conjugate pieces, with HC)

```
In[37]:= Simplify[FeynmanRules[Lgosqq + HC[Lgosqq]]] // MatrixForm
```

Starting Feynman rule calculation.

Expanding the Lagrangian...

Collecting the different structures that enter the vertex.

4 possible non-zero vertices have been found -> starting the computation: 4 / 4.

4 vertices obtained.

Out[37]//MatrixForm=

$$\left( \begin{array}{l} \left\{ \{\bar{q}, 1\}, \{go, 2\}, \{sq1, 3\} \right\} \quad -i \sqrt{2} g_s \left( \cos[\theta] P_{+s_1, s_2} - P_{-s_1, s_2} \sin[\theta] \right) T_{m_1, m_3}^{a_2} \\ \left\{ \{go, 1\}, \{q, 2\}, \{sq1^\dagger, 3\} \right\} \quad -i \sqrt{2} g_s \left( \cos[\theta] P_{-s_1, s_2} - P_{+s_1, s_2} \sin[\theta] \right) T_{m_3, m_2}^{a_1} \\ \left\{ \{\bar{q}, 1\}, \{go, 2\}, \{sq2, 3\} \right\} \quad i \sqrt{2} g_s \left( \cos[\theta] P_{-s_1, s_2} + P_{+s_1, s_2} \sin[\theta] \right) T_{m_1, m_3}^{a_2} \\ \left\{ \{go, 1\}, \{q, 2\}, \{sq2^\dagger, 3\} \right\} \quad i \sqrt{2} g_s \left( \cos[\theta] P_{+s_1, s_2} + P_{-s_1, s_2} \sin[\theta] \right) T_{m_3, m_2}^{a_1} \end{array} \right)$$

# From FEYNRULES to phenomenology

- ◆ The SUSY-QCD model has been implemented into FEYNRULES

```
LMatter := Lkin + LD + Lgosqq + HC[Lgosqq];  
LSUSYQCD := LVector1 + LMatter;
```

- ◆ The Feynman rules have been extracted (and checked)
- ◆ We are now ready to export the model to one or several Monte Carlo tools
  - ♣ CALCHEP / COMPHEP
  - ♣ FEYNARTS / FORMCALC
  - ♣ SHERPA
  - ♣ UFO ➤ MADGRAPH5\_AMC@NLO
  - ♣ WHIZARD / OMEGA

```
WriteCHOutput [LSUSYQCD]  
WriteFeynArtsOutput [LSUSYQCD]  
WriteSHOutput [LSUSYQCD]  
WriteUFO [LSUSYQCD]  
WriteWOOutput [LSUSYQCD]
```



# Limitations and fine prints

## ◆ Particle / parameter names

- ❖ The strong interaction has a special role
  - ★ Name for the gluon field, the coupling constant, etc.
  - ★ Which parameter is internal/external
  - ★ The numerical value of  $\alpha_s$  at the Z-pole
- ❖ For some generators, the electroweak interaction has also a special role
  - ★ Name for the Fermi coupling, the Z-boson mass
  - ★ Which parameter is external/internal
  - ★ At which scale must the numerical values be given

## ◆ Color structures: not supported in full generality

- ❖ The interfaces discard the non-supported vertices
- ❖ Representations handled by the FEYNRULES interfaces:
  - ★ CALCHEP: 1, 3, 8 (but limited)
  - ★ FEYNARTS: all
  - ★ MADGRAPH5\_AMC@NLO: 1, 3, 6, 8
  - ★ SHERPA: 1, 3, 8
  - ★ WHIZARD: 1, 3, 8

## ◆ Lorentz structures and spins: not supported in full generality

- ❖ The interfaces discard the non-supported vertices
- ❖ Representations handled by the FEYNRULES interfaces:
  - ★ CALCHEP: MSSM-like Lorentz structures; spin 0, 1/2, 1, 3/2, 2
  - ★ FEYNARTS: all Lorentz structures; spin 0, 1/2, 1
  - ★ MADGRAPH5\_AMC@NLO: all Lorentz structures; spin 0, 1/2, 1, 3/2, 2
  - ★ SHERPA: SM-like Lorentz structures; spin 0, 1/2, 1
  - ★ WHIZARD: MSSM-like Lorentz structures; spin 0, 1/2, 1, 2

# From FEYNRULES to MADGRAPH 5: the UFO (I)

## ◆ The UFO in a nutshell

- ❖ UFO  $\equiv$  Universal FEYNRULES output (**not tied** to any Monte Carlo tool)
- ❖ Allows the model to contain **generic** color and Lorentz structures
- ❖ The FEYNRULES interface creates a **PYTHON module** to be linked to any code
- ❖ This module contains **all** the model information
- ❖ More information  $\triangleright$  Degrande, Duhr, BenjFuks, Grellscheid, Mattelaer, Reiter [CPC 183 (2012) 1201]

## ◆ The UFO version of the model can be used by several programs

- ❖ ALOHA
- ❖ GOSAM
- ❖ HERWIG++
- ❖ MADANALYSIS 5
- ❖ MADGRAPH5\_AMC@NLO

## ◆ Used with a copy-paste

```
In[6]:= WriteUFO[LSUSYQCD]
      --- Universal FeynRules Output (UFO) v 1.1 ---
      Warning: no electric charge defined. Putting all electric charges to zero.
      Starting Feynman rule calculation.
      Expanding the Lagrangian...
      Collecting the different structures that enter the vertex.
      25 possible non-zero vertices have been found -> starting the computation: 25 / 25.
      21 vertices obtained.
      Flavor expansion of the vertices: 21 / 21
      - Saved vertices in InterfaceRun[ 1 ].
      Computing the squared matrix elements relevant for the 1->2 decays:
      4 / 4
      Squared matrix elent compute in 0.112652 seconds.
      Decay widths computed in 0.002791 seconds.
      Preparing Python output.
      - Splitting vertices into building blocks.
      - Optimizing: 21/21 .
      - Writing files.
      Done!
```

# From FEYNRULES to MADGRAPH 5: the UFO (2)

## ◆ Particles

- ❖ Similar to the FEYNRULES model format, but with different names for the attributes (masses, widths, PDG codes, etc.)
- ❖ Antiparticles are automatically derived from the knowledge of the corresponding particle
- ❖ Special keyword for the (numerical value) zero
- ❖ Note: the embedding of the color and spin representation

```
G = Particle(pdg_code = 21,
             name = 'G',
             antiname = 'G',
             spin = 3,
             color = 8,
             mass = Param.ZERO,
             width = Param.ZERO,
             texname = 'G',
             antitexname = 'G',
             charge = 0)
```

```
go = Particle(pdg_code = 1000021,
              name = 'go',
              antiname = 'go',
              spin = 2,
              color = 8,
              mass = Param.Mgo,
              width = Param.Wgo,
              texname = 'go',
              antitexname = 'go',
              charge = 0)
```

```
sq1 = Particle(pdg_code = 1000006,
               name = 'sq1',
               antiname = 'sq1~',
               spin = 1,
               color = 3,
               mass = Param.Msq1,
               width = Param.Wsq1,
               texname = 'sq1',
               antitexname = 'sq1~',
               charge = 0)
```

```
sq1__tilde__ = sq1.anti()
```

```
sq2 = Particle(pdg_code = 2000006,
               name = 'sq2',
               antiname = 'sq2~',
               spin = 1,
               color = 3,
               mass = Param.Msq2,
               width = Param.Wsq2,
               texname = 'sq2',
               antitexname = 'sq2~',
               charge = 0)
```

```
sq2__tilde__ = sq2.anti()
```

```
q = Particle(pdg_code = 6,
             name = 'q',
             antiname = 'q~',
             spin = 2,
             color = 3,
             mass = Param.Mq,
             width = Param.Wq,
             texname = 'q',
             antitexname = 'q~',
             charge = 0)
```

```
q__tilde__ = q.anti()
```

# From FEYNRULES to MADGRAPH 5: the UFO (3)

## ◆ Parameters

- ❖ Similar to the FEYNRULES model format, but with different names
- ❖ Note: the automated generation of a LesHouches structure for the external parameters (in particular for the masses and widths)
- ❖ PYTHON-compliant formula for the internal parameters

```
aS = Parameter(name = 'aS',
               nature = 'external',
               type = 'real',
               value = 0.1184,
               texname = '\\text{aS}',
               lhablock = 'FRBlock',
               lhacode = [ 1 ])
```

```
theta = Parameter(name = 'theta',
                  nature = 'external',
                  type = 'real',
                  value = 0.7853981633974483,
                  texname = '\\theta',
                  lhablock = 'FRBlock',
                  lhacode = [ 2 ])
```

```
Mgo = Parameter(name = 'Mgo',
                 nature = 'external',
                 type = 'real',
                 value = 500,
                 texname = '\\text{Mgo}',
                 lhablock = 'MASS',
                 lhacode = [ 1000021 ])
```

```
Wq = Parameter(name = 'Wq',
                nature = 'external',
                type = 'real',
                value = 1.50833649,
                texname = '\\text{Wq}',
                lhablock = 'DECAY',
                lhacode = [ 6 ])
```

```
G = Parameter(name = 'G',
               nature = 'internal',
               type = 'real',
               value = '2*cmath.sqrt(aS)*cmath.sqrt(cmath.pi)',
               texname = 'G')
```

# From FEYNRULES to MADGRAPH 5: the UFO (4)

## ◆ Vertices

### ♣ Decomposed in a **spin x color** basis

★ For instance, the quartic gluon vertex can be written as

$$\begin{aligned}
 & ig_s^2 f^{a_1 a_2 b} f^{b a_3 a_4} (\eta^{\mu_1 \mu_4} \eta^{\mu_2 \mu_3} - \eta^{\mu_1 \mu_3} \eta^{\mu_2 \mu_4}) \\
 & + ig_s^2 f^{a_1 a_3 b} f^{b a_2 a_4} (\eta^{\mu_1 \mu_4} \eta^{\mu_2 \mu_3} - \eta^{\mu_1 \mu_2} \eta^{\mu_3 \mu_4}) \\
 & + ig_s^2 f^{a_1 a_4 b} f^{b a_2 a_3} (\eta^{\mu_1 \mu_3} \eta^{\mu_2 \mu_4} - \eta^{\mu_1 \mu_2} \eta^{\mu_3 \mu_4})
 \end{aligned}
 \Rightarrow
 \begin{aligned}
 & (f^{a_1 a_2 b} f^{b a_3 a_4}, f^{a_1 a_3 b} f^{b a_2 a_4}, f^{a_1 a_4 b} f^{b a_2 a_3}) \\
 & \times \begin{pmatrix} ig_s^2 & 0 & 0 \\ 0 & ig_s^2 & 0 \\ 0 & 0 & ig_s^2 \end{pmatrix} \begin{pmatrix} \eta^{\mu_1 \mu_4} \eta^{\mu_2 \mu_3} - \eta^{\mu_1 \mu_3} \eta^{\mu_2 \mu_4} \\ \eta^{\mu_1 \mu_4} \eta^{\mu_2 \mu_3} - \eta^{\mu_1 \mu_2} \eta^{\mu_3 \mu_4} \\ \eta^{\mu_1 \mu_3} \eta^{\mu_2 \mu_4} - \eta^{\mu_1 \mu_2} \eta^{\mu_3 \mu_4} \end{pmatrix}
 \end{aligned}$$

### ♣ Each element of the decomposition is stored separately

★ **vertices.py**: contains the decomposition, the color basis and generic elements for the couplings and the spin basis

```
V_2 = Vertex(name = 'V_2',
             particles = [ P.G, P.G, P.G, P.G ],
             color = [ 'f(-1,1,2)*f(3,4,-1)', 'f(-1,1,3)*f(2,4,-1)', 'f(-1,1,4)*f(2,3,-1)' ],
             lorentz = [ L.VVVV1, L.VVVV2, L.VVVV3 ],
             couplings = {(1,1):C.GC_4, (0,0):C.GC_4, (2,2):C.GC_4})
```

★ **lorentz.py**: contains each necessary element of the spin basis (reused across vertices as much as possible)

```
VVVV1 = Lorentz(name = 'VVVV1',
                spins = [ 3, 3, 3, 3 ],
                structure = 'Metric(1,4)*Metric(2,3) - Metric(1,3)*Metric(2,4)')

VVVV2 = Lorentz(name = 'VVVV2',
                spins = [ 3, 3, 3, 3 ],
                structure = 'Metric(1,4)*Metric(2,3) - Metric(1,2)*Metric(3,4)')

VVVV3 = Lorentz(name = 'VVVV3',
                spins = [ 3, 3, 3, 3 ],
                structure = 'Metric(1,3)*Metric(2,4) - Metric(1,2)*Metric(3,4)')
```

★ **couplings.py**: the coupling strengths (reused across vertices as much as possible)

```
GC_4 = Coupling(name = 'GC_4',
                 value = 'complex(0,1)*G**2',
                 order = {'QCD':2})
```

# Outline

1. FEYNRULES in a nutshell
2. Implementing supersymmetric QCD in FEYNRULES
3. Using FEYNRULES with the supersymmetric QCD model
4. Advanced model implementation techniques
5. Summary

# Advanced techniques for implementing models in FEYNRULES

- ◆ Extension / restriction of existing models
- ◆ The superspace module of FEYNRULES
- ◆ Mass diagonalization
- ◆ Two-body decays
- ◆ Next-to-leading order module

# Merging and extending model implementations

## ◆ Many BSM models of interest are simple extensions of another model

### ❖ FEYNRULES allows one to start from a given model

- ★ Add new particles, parameters, Lagrangian terms
- ★ Modify existing particles, parameters, Lagrangian terms
- ★ Remove some particles, parameters, Lagrangian terms

### ❖ Special cases very relevant for LHC physics: Simplified Model Spectra

- ★ The Standard Model + one or two new particles
- ★ Often inspired by the MSSM
- ★ Example: the SM + lightest stop and neutralino + the relevant subset of MSSM interactions

## ◆ The merged FEYNRULES model contains two .fr files

### ❖ The parent model implementation

### ❖ One extra file with the modifications

### ❖ They must be loaded together (the parent model first)

```
LoadModel["SM.fr", "stops.fr"];|
```

- ★ No need to re-implement what is common (gauge groups, etc.)

## ◆ One can start from any of the models available on the FEYNRULES model database

<http://feynrules.irmp.ucl.ac.be>



# The FEYNRULES model database

## ◆ About 40 models are available online

### ❖ Simple extensions of the Standard Model

- ★ Simplified model spectra
- ★ Four generation models
- ★ Vector-like quarks
- ★ Two-Higgs-Doublet Models, Hidden Abelian Higgs
- ★ etc.

### ❖ Supersymmetric models

- ★ MSSM with and without R-parity
- ★ The NMSSM
- ★ R-symmetric supersymmetric models
- ★ Left-right supersymmetric models

### ❖ Extra-dimensional models

- ★ Universal extra-dimensions
- ★ Large extra-dimensions
- ★ Heidi, Minimal Higgsless models
- ★ Randall-Sundrum

### ❖ Strongly coupled and effective field theories

- ★ Technicolor
- ★ Models with dimension-six and dimension-eight operators
- ★ etc.

#### Available models

[Standard Model](#)

[Simple extensions of the SM \(18\)](#)

[Supersymmetric Models \(5\)](#)

[Extra-dimensional Models \(4\)](#)

[Strongly coupled and effective field theories \(8\)](#)

[Miscellaneous \(0\)](#)

# Restricting model implementations

## ◆ Many BSM models of interest are subset of other

- ❖ Equivalent to the parent model, but with some parameters set to specific values (0 or 1)
  - ★ Example 1: the massless version of a model (massless light quarks in the Standard Model)
  - ★ Example 2: a mixing matrix set to the identity (no-CKM matrix in the Standard Model)
- ❖ FEYNRULES allows one to start from a given model
  - ★ Write the restrictions under the form of a list of MATHEMATICA replacement rules (**M\$Restrictions**)
  - ★ Read them into FEYNRULES
  - ★ Apply them before the computation of any Feynman rule

## ◆ The output Feynman rules (and this Monte Carlo model files)

- ❖ Are free from the restricted parameters
- ❖ Smaller files, more efficiency at the Monte Carlo level
  - ★ Example: the general MSSM has more than 10000 vertices; its flavor-conserving version only ~1000

## ◆ One practical example: the Standard Model without CKM-mixing

```
M$Restrictions = {
  CKM[i_,i_] -> 1,
  CKM[i_?NumericQ, j_?NumericQ] :> 0 /; (i != j),
  cabi -> 0
}
```

```
LoadRestriction["DiagonalCKM.rst"]
```

# The supersymmetry module

## ◆ A module dedicated to calculations in superspace

- ❖ Superfield declaration and **links to the component fields**
- ❖ Series expansion in terms of component fields
- ❖ Automatic **derivation of supersymmetric Lagrangians**
- ❖ Solving the equations of motion of the unphysical fields
- ❖ Many extra built-in functions
- ❖ See: Duhr, BenjFuks [CPC 182 (2011) 2404]; BenjFuks [IJMPA 27 (2012) 1230007]

## A left-handed squark superfield

```
CSF[1] == {
  ClassName      -> QL,
  Chirality      -> Left,
  Weyl          -> qLw,
  Scalar        -> sqL,
  Indices       -> {Index[Colour]}
}
```

## ◆ Supersymmetric model implementation

- ❖ Declaration of the model **gauge group**
- ❖ Declaration of all **fields and superfields**
- ❖ Declaration of all model **parameters**
- ❖ **Can the writing of the Lagrangian be simplified?**

## ◆ Supersymmetric Lagrangian in superspace are very compact

$$\mathcal{L} = \Phi^\dagger e^{-2gV} \Phi|_{\theta^2\bar{\theta}^2} + \frac{1}{16g^2\tau_{\mathcal{R}}} \text{Tr}(W^\alpha W_\alpha)|_{\theta^2} + \frac{1}{16g^2\tau_{\mathcal{R}}} \text{Tr}(\bar{W}_{\dot{\alpha}} \bar{W}^{\dot{\alpha}})|_{\bar{\theta}^2} \\ + W(\Phi)|_{\theta^2} + W^*(\Phi^\dagger)|_{\bar{\theta}^2} + \mathcal{L}_{\text{soft}}$$

➤ Model independent ⇒ can be automated

➤ Model dependent ⇒ to be provided

- ❖ First line: kinetic and gauge interaction terms for all fields
- ❖ Second line: model dependent superpotential and supersymmetry breaking Lagrangian

See the manual for more details on the superspace module

# Implementing supersymmetric Lagrangians

## ◆ Supersymmetric Lagrangian in superspace are very compact

$$\mathcal{L} = \Phi^\dagger e^{-2gV} \Phi|_{\theta^2\bar{\theta}^2} + \frac{1}{16g^2\tau_{\mathcal{R}}} \text{Tr}(W^\alpha W_\alpha)|_{\theta^2} + \frac{1}{16g^2\tau_{\mathcal{R}}} \text{Tr}(\bar{W}_{\dot{\alpha}} \bar{W}^{\dot{\alpha}})|_{\bar{\theta}^2} \\ + W(\Phi)|_{\theta^2} + W^*(\Phi^\dagger)|_{\bar{\theta}^2} + \mathcal{L}_{\text{soft}}$$

### ♣ First line: kinetic and gauge interaction terms for all fields

- ★ Model independent  $\Rightarrow$  can be automated (`CSFKineticTerms` and `VSFKineticTerms` functions)
- ★ Dedicated methods to access the components of a superfield (`Theta2Component`, etc.)

```
Lag = Theta2Thetabar2Component[ CSFKineticTerms[] ] +
      Theta2Component[ VSFKineticTerms[] ] +
      Thetabar2Component[ VSFKineticTerms[] ];
```

### ♣ Second line: superpotential and supersymmetry-breaking terms

- ★ Model dependent  $\Rightarrow$  to be provided by the user (`SuperW` and `LSoft`)

```
Lag2 = LSoft + Theta2Component[ SuperW ] + Thetabar2Component[ SuperW ];
```

## ◆ The Lagrangian above must be post-processed

- ♣ Solving the **equation of motion** for the auxiliary fields and inserting the solution back into the Lagrangian
  - ★ Automated (`SolveEqMotionF` and `SolveEqMotionD`)
- ♣ Replacement of **Weyl spinors** in terms of Majorana and Dirac spinors
  - ★ Automated (`WeylToDirac`)
- ♣ Rotation to the mass basis
  - ★ Standard FEYNRULES function (`ExpandIndices`)

See the manual for more details on the superspace module

# Mass matrices and their diagonalization

## ◆ The problematics of the mass matrices

- ❖ Lagrangians are usually easily written in the gauge basis
- ❖ The included mass matrices are thus in general non-diagonal ➤ **diagonalization required**
- ❖ The gauge basis of fields must be rotated to the mass basis where the mass matrices are diagonal
- ❖ **This diagonalization cannot in general be achieved analytically**

## ◆ The ASPERGE package of FEYNRULES

- ❖ A MATHEMATICA module allowing one to **extract the mass matrices from the Lagrangian**
- ❖ A generator of C++ code ➤ **numerical diagonalization** of all the model mass matrices (the generated code can be used in a **standalone way**)
- ❖ See: Alloul, D'Hondt, De Causmaecker, BenjFuks, Rausch de Traubenberg [ EPJC 73 (2013) 2325 ]

## ◆ Example: the Z-boson and photon in the Standard Model

- ❖ Each mixing is declared as a set of replacement rules (in *M\$MixingsDescription*)
- ❖ **Each rule represent a property of the mixing relation**

$$\begin{pmatrix} A_\mu \\ Z_\mu \end{pmatrix} = U_w \begin{pmatrix} B_\mu \\ W_\mu^3 \end{pmatrix} \Rightarrow$$

```
Mix["AZ"] == {
  MassBasis      -> {A, Z},
  GaugeBasis     -> {B, Wi[3]},
  MixingMatrix   -> UW,
  BlockName      -> WEAKMIX
}
```

- ❖ ASPERGE can compute the mass matrices:
- ❖ **ASPERGE can generate its standalone C++ version:**

```
ComputeMassMatrix[Lag];
WriteASperGe[Lag];
```

See the manual for more details on the ASPERGE module

# Two-body decays

## ◆ The problematics of the decay widths and branching ratios

- ❖ Some Monte Carlo tools need the decay table (widths and branching ratios) to decay particles
- ❖ Widths and branching ratios are not independent quantities ➤ **need to be calculated**
- ❖ Some Monte Carlo tools compute these quantities on the fly
  - **the procedure is repeated each time it is needed**
- ❖ FeynRules offers a way to include analytical information on the two-body decay

## ◆ The decay module of FEYNRULES

- ❖ Two body decays can be directly read **from the three-point vertices** ( $\mathcal{V}$ ) of the model

$$\Gamma_{1 \rightarrow 2} = \frac{1}{2|M|S} \int d\Phi_N |\mathcal{M}_{1 \rightarrow 2}|^2 = \frac{\sqrt{\lambda(M^2, m_1^2, m_2^2)}}{16 \pi S |M|^3} \mathcal{V}_{\ell_1 \ell_2 \ell_3}^{a_1 a_2 a_3} \mathcal{P}_1^{\ell_1 \ell'_1} \mathcal{P}_2^{\ell_2 \ell'_2} \mathcal{P}_3^{\ell_3 \ell'_3} (\mathcal{V}^*)_{\ell'_1 \ell'_2 \ell'_3}^{a_1 a_2 a_3}$$

- ★ Partial width for the decay of a particle of mass  $M$  to two particles of masses  $m_1$  and  $m_2$
- ★ Includes a symmetry factor  $S$  and  $\mathcal{P}$  denotes the polarization tensor of each particle
- ❖ FEYNRULES makes use of MATHEMATICA to compute all partial widths of the model
  - ★ Ignores open and closed channels ➤ **benchmark independent**
  - ★ **The information is exported to the UFO** (already used by MADWIDTH, the decay module of MADGRAPH5)

# The decay module of FEYNRULES

## ◆ Automatic decay width computations

- ❖ All two-body decay widths can be easily computed from the Lagrangian

```
verts      = FeynmanRules[Lag];
vertsexp   = FlavorExpansion[verts];
results    = ComputeWidths[vertsexp];
```

- ❖ Many functions available for analytical calculations (**PartialWidth**, **TotWidth**, **BranchingRatio**)
- ❖ The numerical value provided for the particle widths can be updated accordingly (**UpdateWidths**)
- ❖ See: Alwall, Duhr, BenjFuks, Mattelaer, Öztürk, Shen [ arXiv:1402.1178 ]

## ◆ The information is (by default) employed at the UFO interface level

- ❖ Can be turned of: (**AddDecays** → **False**)
- ❖ The UFO contains an extra file *decays.py*
- ❖ Example of the Standard Model UFO: the top quark

```
Decay_t = Decay(name = 'Decay_t',
               particle = P.t,
               partial_widths = {(P.W__plus__, P.d): '((MT**2 - MW**2)*((3*CKM3x1*ee**2*MT**2*complexconjugate(CKM3x1))/(2.*sw**2) +
(3*CKM3x1*ee**2*MT**4*complexconjugate(CKM3x1))/(2.*MW**2*sw**2) - (3*CKM3x1*ee**2*MW**2*complexconjugate(CKM3x1))/sw**2))/
(96.*cmath.pi*abs(MT)**3)',
                               (P.W__plus__, P.s): '((MT**2 - MW**2)*((3*CKM3x2*ee**2*MT**2*complexconjugate(CKM3x2))/(2.*sw**2) +
(3*CKM3x2*ee**2*MT**4*complexconjugate(CKM3x2))/(2.*MW**2*sw**2) - (3*CKM3x2*ee**2*MW**2*complexconjugate(CKM3x2))/sw**2))/
(96.*cmath.pi*abs(MT)**3)',
                               (P.W__plus__, P.b): '(((3*CKM3x3*ee**2*MB**2*complexconjugate(CKM3x3))/(2.*sw**2) +
(3*CKM3x3*ee**2*MT**2*complexconjugate(CKM3x3))/(2.*sw**2) + (3*CKM3x3*ee**2*MB**4*complexconjugate(CKM3x3))/(2.*MW**2*sw**2) -
(3*CKM3x3*ee**2*MB**2*MT**2*complexconjugate(CKM3x3))/(MW**2*sw**2) + (3*CKM3x3*ee**2*MT**4*complexconjugate(CKM3x3))/(2.*MW**2*sw**2) -
(3*CKM3x3*ee**2*MW**2*complexconjugate(CKM3x3))/sw**2)*cmath.sqrt(MB**4 - 2*MB**2*MT**2 + MT**4 - 2*MB**2*MW**2 - 2*MT**2*MW**2 +
MW**4))/(96.*cmath.pi*abs(MT)**3)')})
```

See the manual for more details on the decay module

# FEYNRULES @ NLO

## ◆ Ingredients for a next-to-leading order calculation

- ❖ Tree-level vertices
- ❖ UV counterterms
- ❖  $R_2$  counterterms (depending on the NLO event generator)

## ◆ The $R_2$ counterterms

- ❖ One-loop modules are based on unitarity approaches
- ❖ This requires the computation of the rational parts of the loops ( $R_1$  and  $R_2$ )
- ❖ The  $R_2$  counterterm can be calculated once and for all for each model
  - ★ From the knowledge of the tree-level Lagrangian
  - ★ Using special Feynman rules (numerators in  $\epsilon$  dimensions; denominators in  $D$  dimensions)

## ◆ Technical details for the NLO module of FEYNRULES

- ❖ Automatic renormalization of the Lagrangian
- ❖ Use of the FEYNARTS interface of FEYNRULES
- ❖ Generation of a script for NLO diagram generation
  - $R_2$  and UV counterterms
- ❖ Export to the UFO
- ❖ See: Degrande [ arXiv:1406.3030]



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# Summary

## ◆ The quest for new physics at the LHC has started

- ❖ Rely on **Monte Carlo event generators** for background and signal modeling
- ❖ FEYNRULES facilitates the implementation of new physics models in those tools

## ◆ FEYNRULES: <http://feynrules.irmp.ucl.ac.be>

- ❖ **Straightforward implementation** of new physics model in Monte Carlo tools
  - ★ Interfaces to many programs
- ❖ FEYNRULES is shipped with its **own computational modules**
  - ★ A superspace module
  - ★ A decay package
  - ★ A mass diagonalization module (ASPERGE)
  - ★ A brand new NLO module

Try it on with your  
favorite model!

# From FEYNRULES to event analysis

## 0. Implementation of the model in FEYNRULES and generation of the Monte Carlo model files

### 1. Event generation with any Monte Carlo event generator

- ✦ Both **signal** and **backgrounds**
- ✦ **Precision** in the normalization: (N)NLO inclusive results
- ✦ Generator choice: beware of restrictions (supported Lorentz and color structures)  
→ e.g., **MADGRAPH5\_AMC@NLO**

Next talk

### 2. Parton showering and hadronization

- ✦ **Precision** in the shapes: multiparton matrix-element merging techniques  
→ e.g., **MADGRAPH 5 + PYTHIA (automated merging)**

Next-to-next talk

### 3. Detector simulation

later this week

### 4. Event analysis with MADANALYSIS 5

- ✦ **Parton-level**, **hadron-level** and **reconstructed-level** analyses (and more)

later this week