

# MC's for top physics:

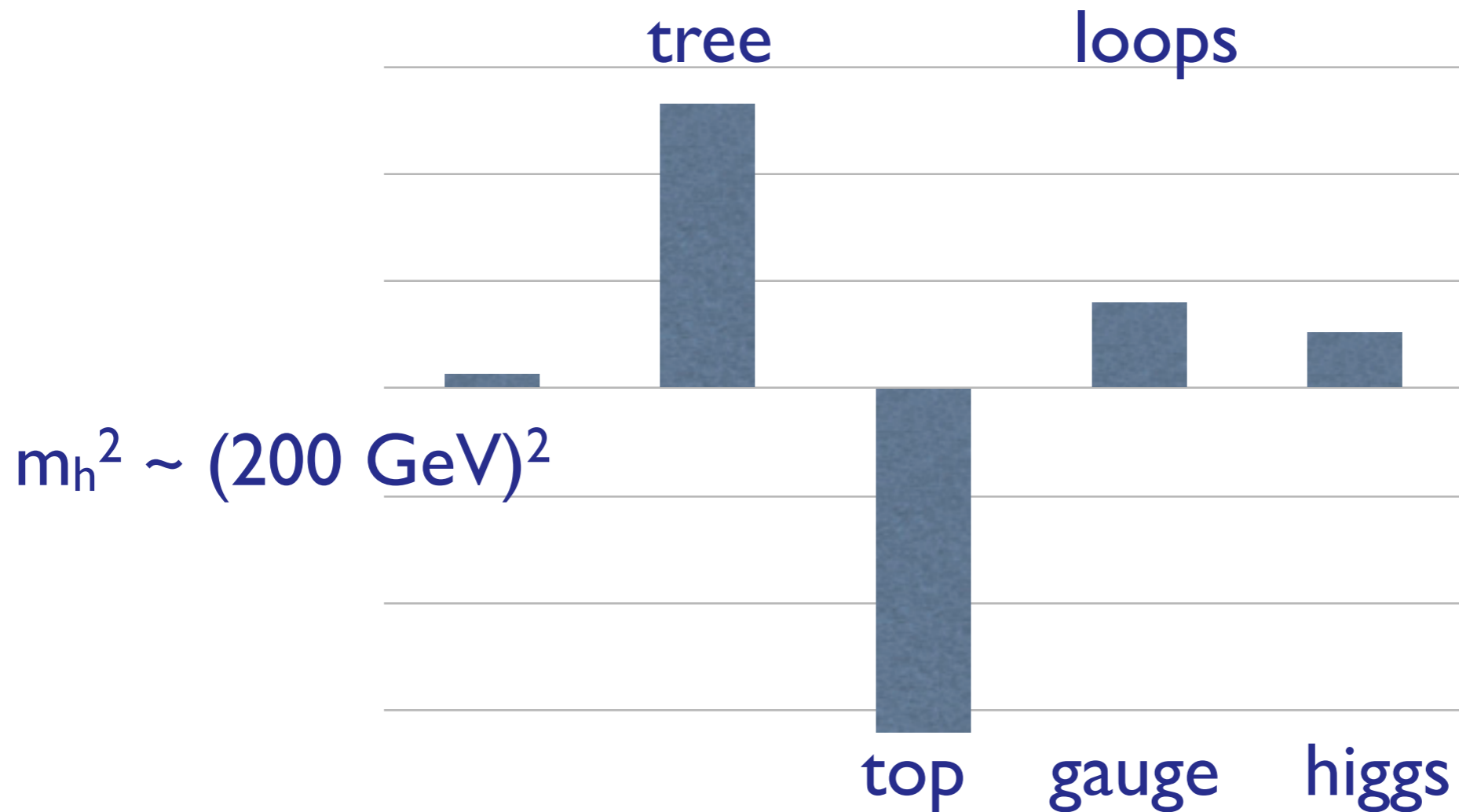
## Matrix Elements and NLO with Parton Showers

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# Why top?

# Top as a link to BSM



$$(200 \text{ GeV})^2 = m_{H_0}^2 + \left[ -(2 \text{ TeV})^2 + (700 \text{ GeV})^2 + (500 \text{ GeV})^2 \right] \left( \frac{\Lambda_{t,W,H}}{10 \text{ TeV}} \right)^2$$

Definition of naturalness: less than 90% cancellation:

$$\Lambda_t \lesssim 3 \text{ TeV} \quad \Lambda_W \lesssim 9 \text{ TeV} \quad \Lambda_H \lesssim 12 \text{ TeV}$$

\* One can actually prove that this case in model independent way, i.e. that the scale associated with top mass generation is very close to that of EWSB.

# Top as a link to BSM

Our **AIM** is twofold:

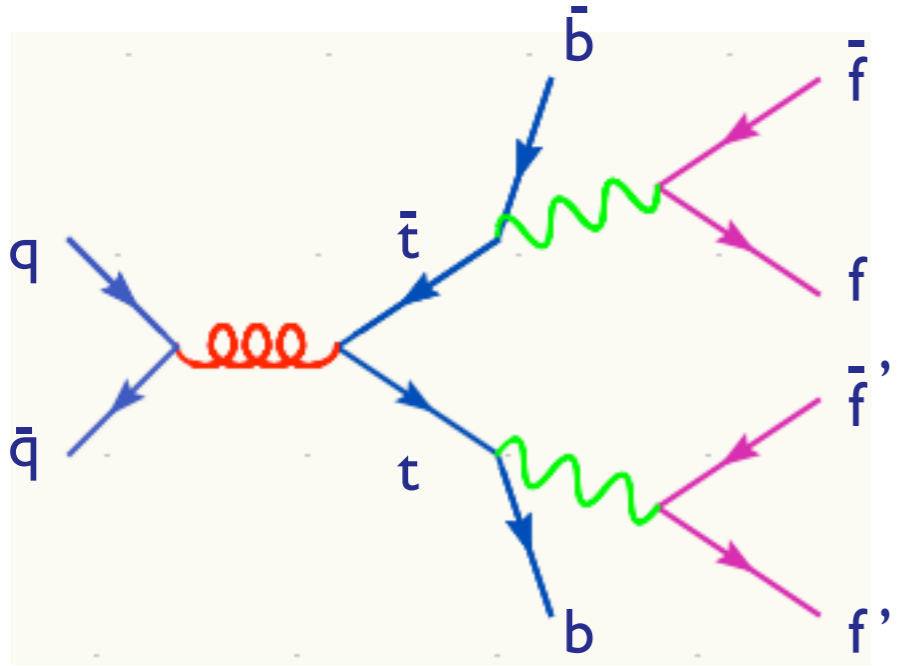
I. Measure all properties (mass, couplings, spin) to establish **indirect** evidence for SM and BSM physics.

Examples: precision EW and QCD ( $m_{\text{top}}, \sigma(t\bar{t}), \sigma(t)$ );  
Rare decays and anomalous couplings. CP violation.

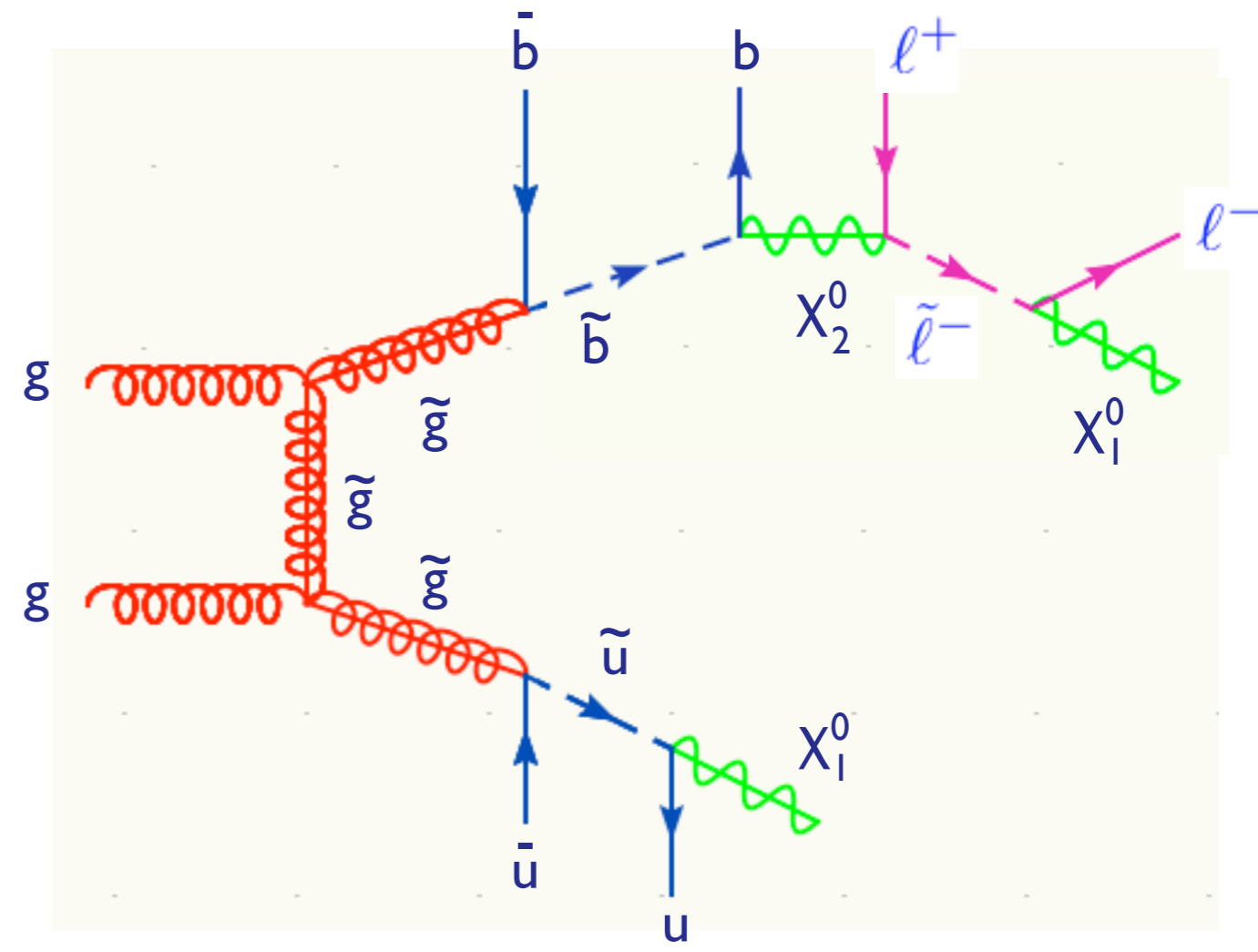
II. Use top as **direct** probe of the EWSB sector and BSM physics

Examples: SM  $t\bar{t}H$ ; BSM:  $Z'$  and  $W'$  resonances; SUSY:  $tH^+$  and  $t \rightarrow bH^+$  or  $\text{stop} \rightarrow t X$ .

# Top as a template

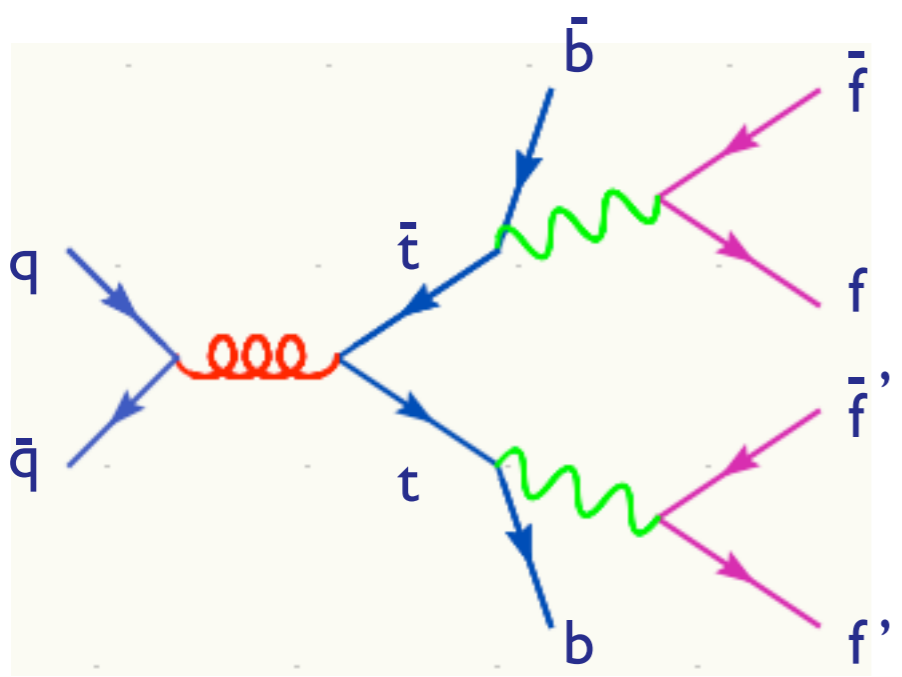


VS

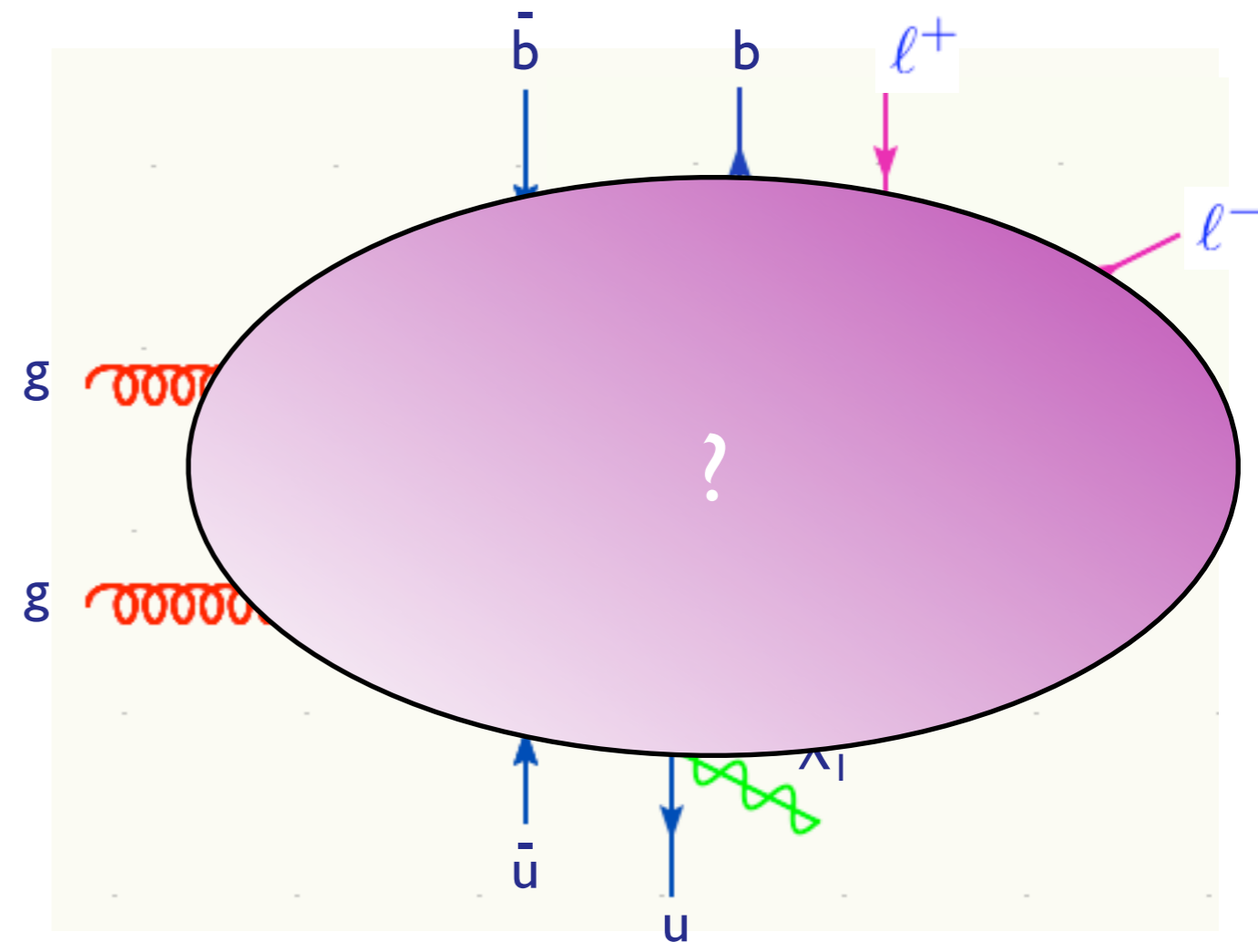


Both involve production of heavy colored states decaying through a chain into jets, leptons and  $\cancel{E}_T$ .

# Top as a template



VS



Both involve production of heavy colored states decaying through a chain into jets, leptons and  $\cancel{E}_T$ .

# Top as background

At the LHC, many measurements will need a good understanding and control of  $t\bar{t}$  and single top events.

A few examples:

- $gg \rightarrow H$  and  $qq \rightarrow Hqq$  with  $H \rightarrow WW$
- $t\bar{t}$  in single top measurements
- $t\bar{t} + \text{jets}$  and  $t\bar{t}bb$  in  $t\bar{t}H$
- $t\bar{t} + \text{jets}$  in SUSY/UED searches (gluino pairs, stop pairs,  $tH^+$ ....)
- .....

# Why MC's?



# The role of MC's in the path to discovery

1. Find excess(es) over SM backgrounds

2. Identify a finite set of coarse models compatible with the excess(es).

3. Look for “predicted excesses” in other channels.

4. Refine

5. Perform more detailed studies to measure mass spectrum, quantum numbers, couplings.

6. Refine

1. Fully exclusive description for rich and energetic final states (multi-jets + EW and QCD particles (W,Z, photon,b,t) with flexible MC to be validated and tuned to control samples. Accurate predictions (NLO,NNLO) for standard candles SM cross sections.

2. Inverse problem tools (Ex: OSET)

3. Simulation of any BSM signature: from models to events in an easy and fast way.

4. Accurate predictions for cross sections of selected models (Ex: SUSY) to identify couplings. Accurate predictions for primary couplings (Ex: spectra calculators).

5. Accurate ME based description for final state distributions which keeps all the relevant information (Ex. decay chain with spin).

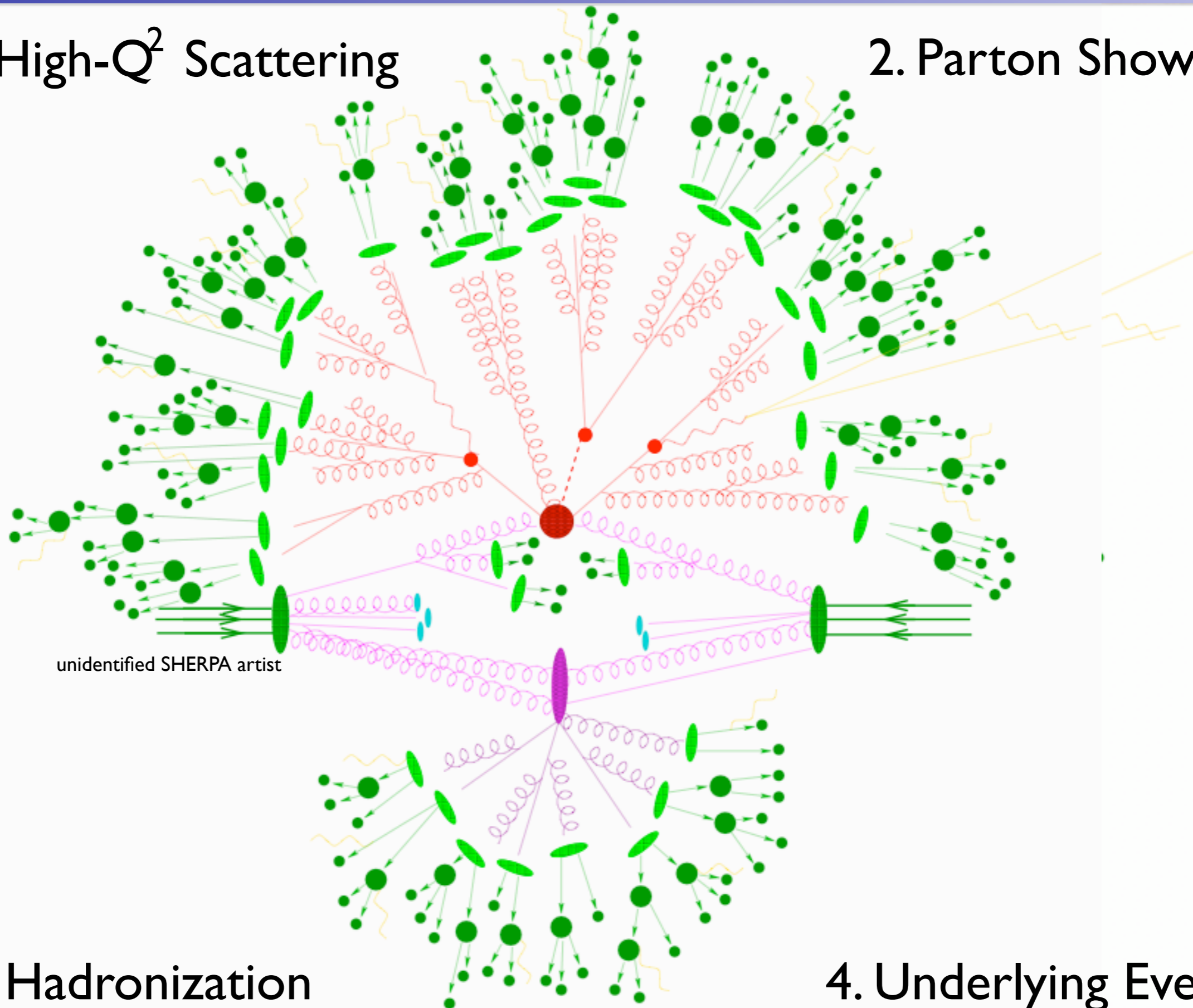
6. Off-shell effects, Matrix Element methods, Global fits (Ex: Sfitter)

# Plan

- Overview
- Matrix Element + Parton Shower
- NLO + Parton Shower
- MC for New Physics
- Conclusions

# 1. High- $Q^2$ Scattering

# 2. Parton Shower

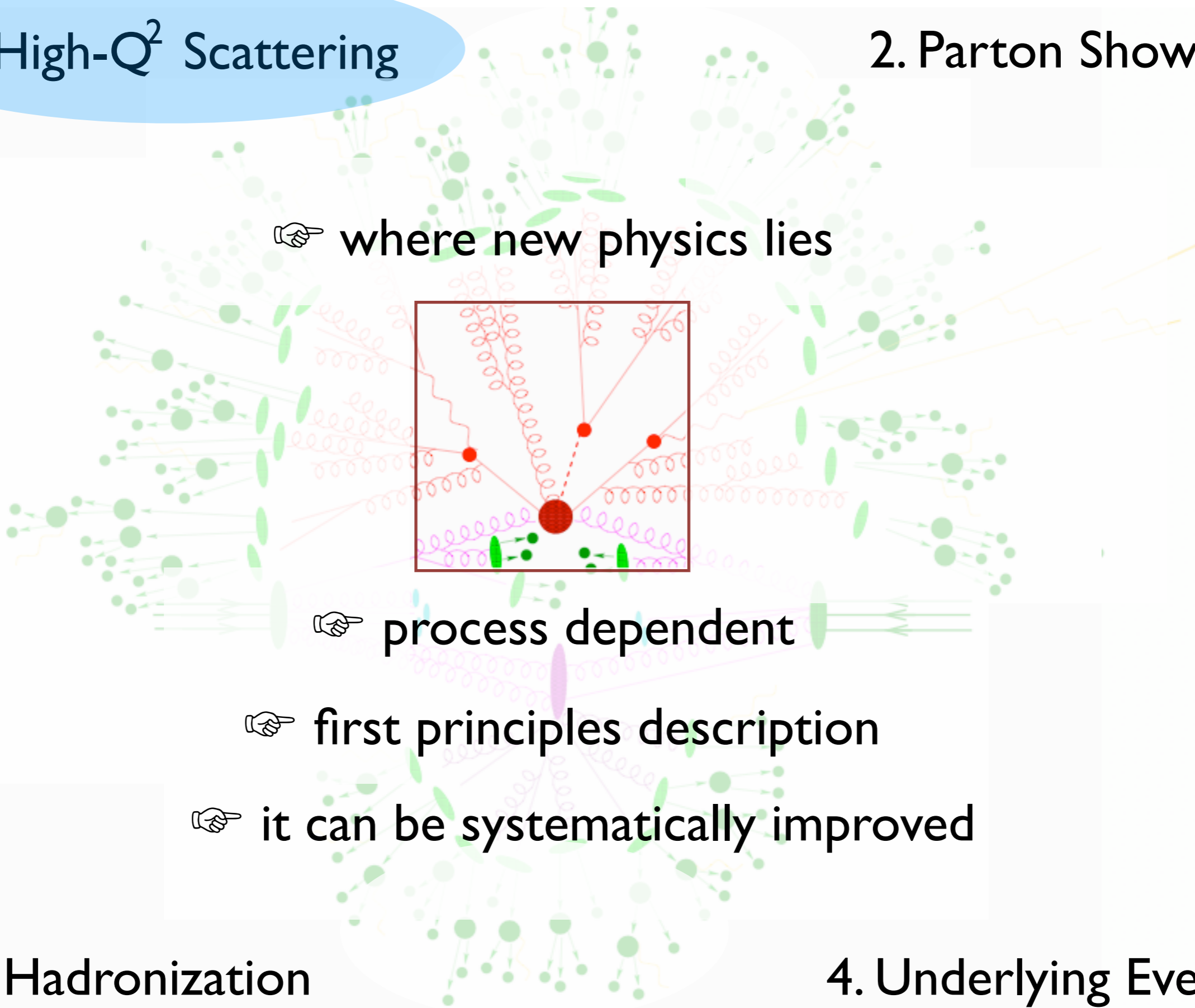


# 3. Hadronization

# 4. Underlying Event

# I. High- $Q^2$ Scattering

# 2. Parton Shower



👉 where new physics lies

👉 process dependent

👉 first principles description

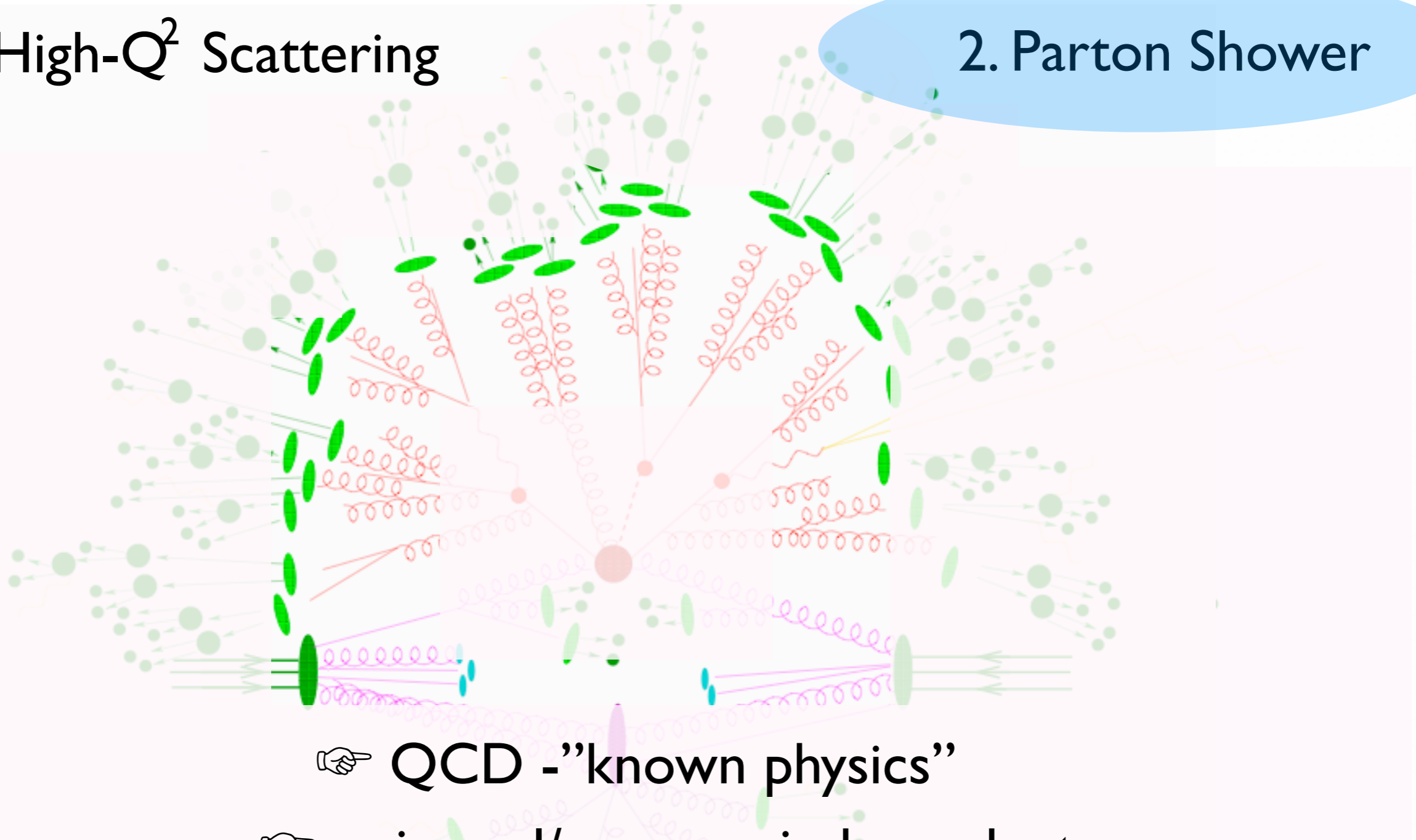
👉 it can be systematically improved

# 3. Hadronization

# 4. Underlying Event

# I. High- $Q^2$ Scattering

# 2. Parton Shower



☞ QCD - "known physics"

☞ universal/ process independent

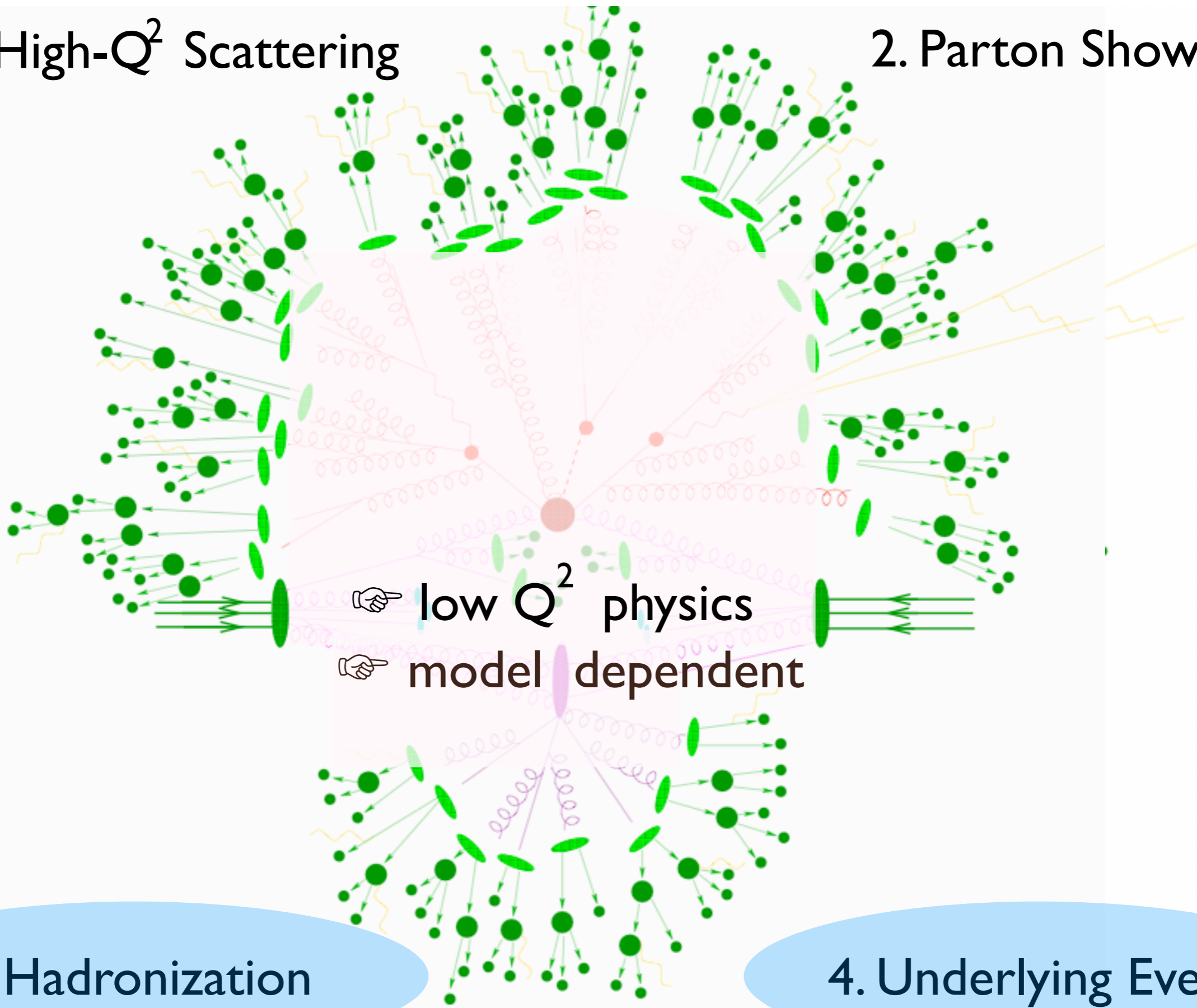
☞ first principles description

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# 3. Hadronization

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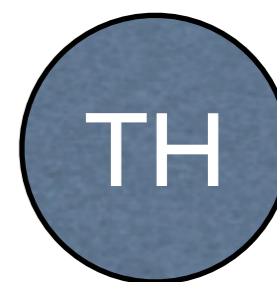
# How we (used to) make predictions?

## First way:

- Include higher order terms in our fixed-order calculations (LO → NLO → NNLO...)

$$\Rightarrow \hat{\sigma}_{ab \rightarrow X} = \sigma_0 + \alpha_S \sigma_1 + \alpha_S^2 \sigma_2 + \dots$$

- Just use the tree-level results for many partons final states



## Comments:

1. The theoretical errors systematically decrease.
2. Pure theoretical point of view.
3. A lot of new techniques and universal algorithms are developed.
4. Final description only in terms of partons and calculation of IR safe observables ⇒ not directly useful for simulations

# How we (used to) make predictions?

## Second way:

- Describe final states with high multiplicities starting from  $2 \rightarrow 1$  or  $2$  procs, using parton showers, and then an hadronization model.



## Comments:

1. Fully exclusive final state description for detector simulations.
2. Normalization is very uncertain
3. Very crude kinematic distributions for multi-parton final states.
4. Improvements are only at the model level.



# How to improve our predictions?

New trend:

TH & EXP

Match fixed-order calculations and parton showers to obtain the most accurate predictions in a detector simulation friendly way!

Two directions:

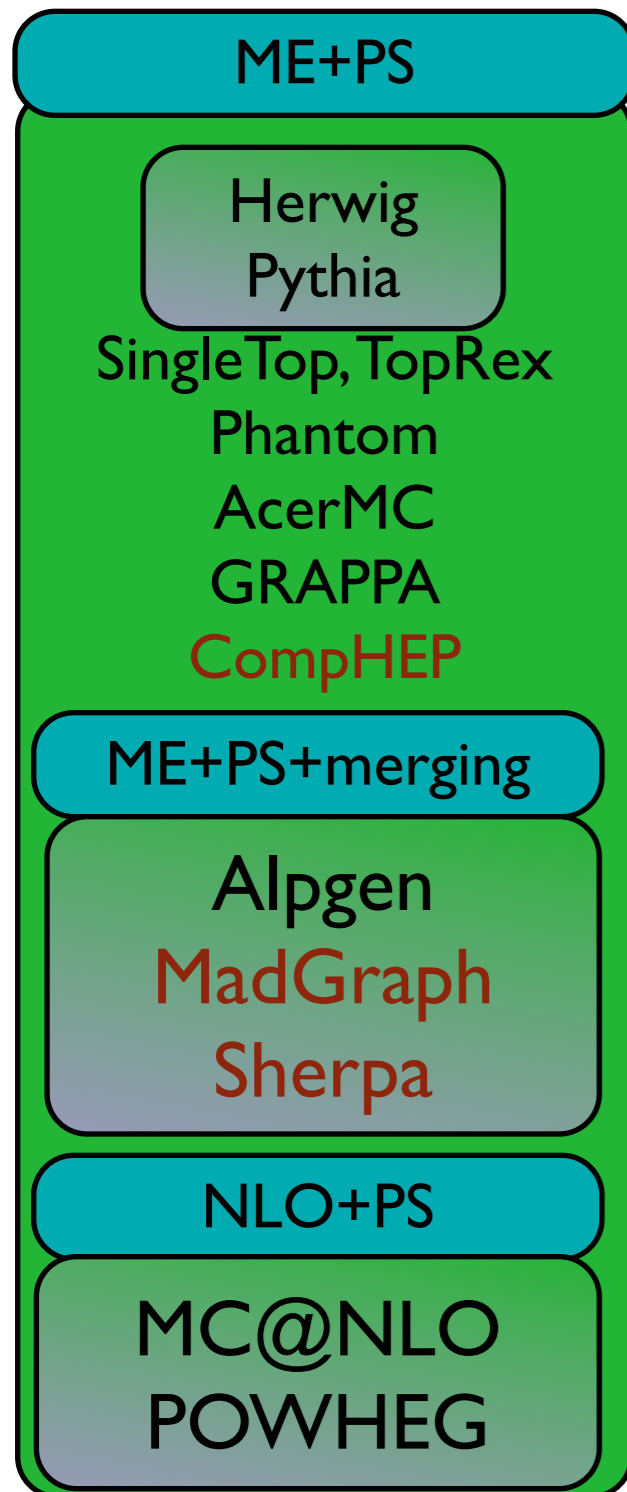
1. Get fully exclusive description of many parton events correct at LO (LL) in all the phase space.

ME+PS

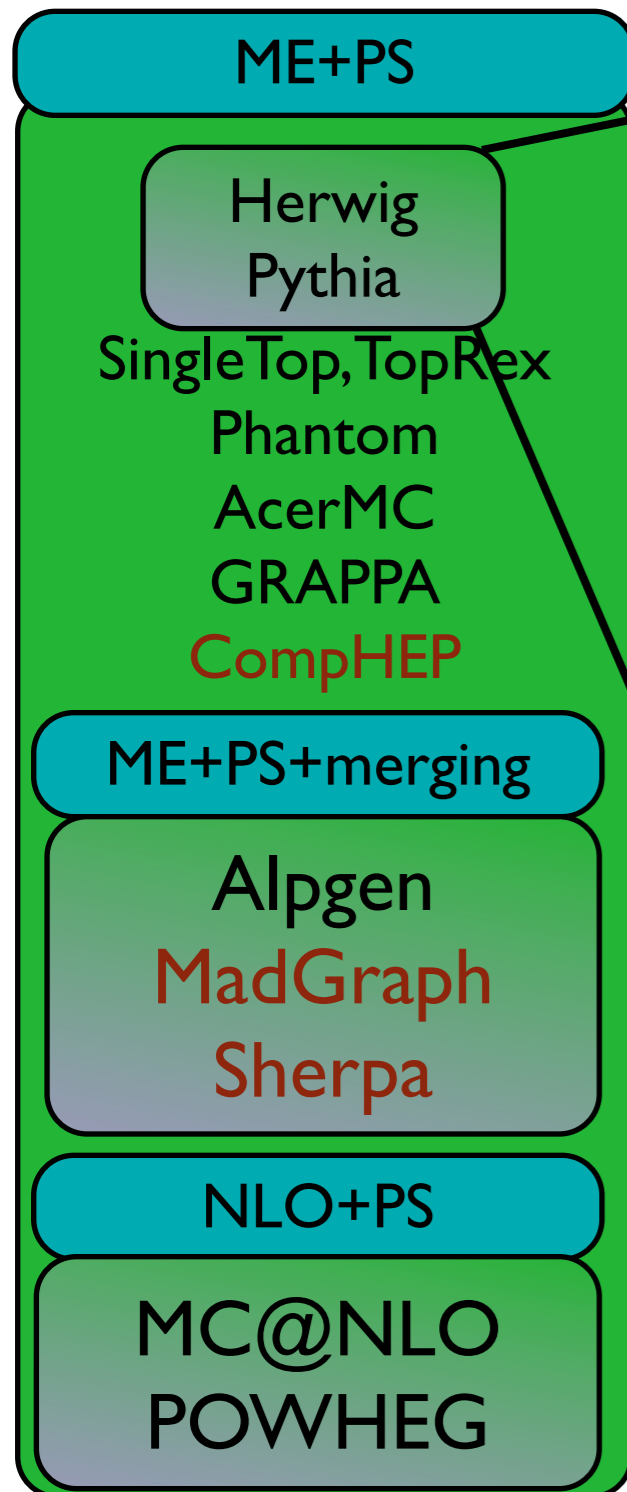
2. Get fully exclusive description of events correct at NLO in the normalization and distributions.

NLO<sub>w</sub>PS

# MC tools for top physics in a nutshell



# MC tools for top physics in a nutshell



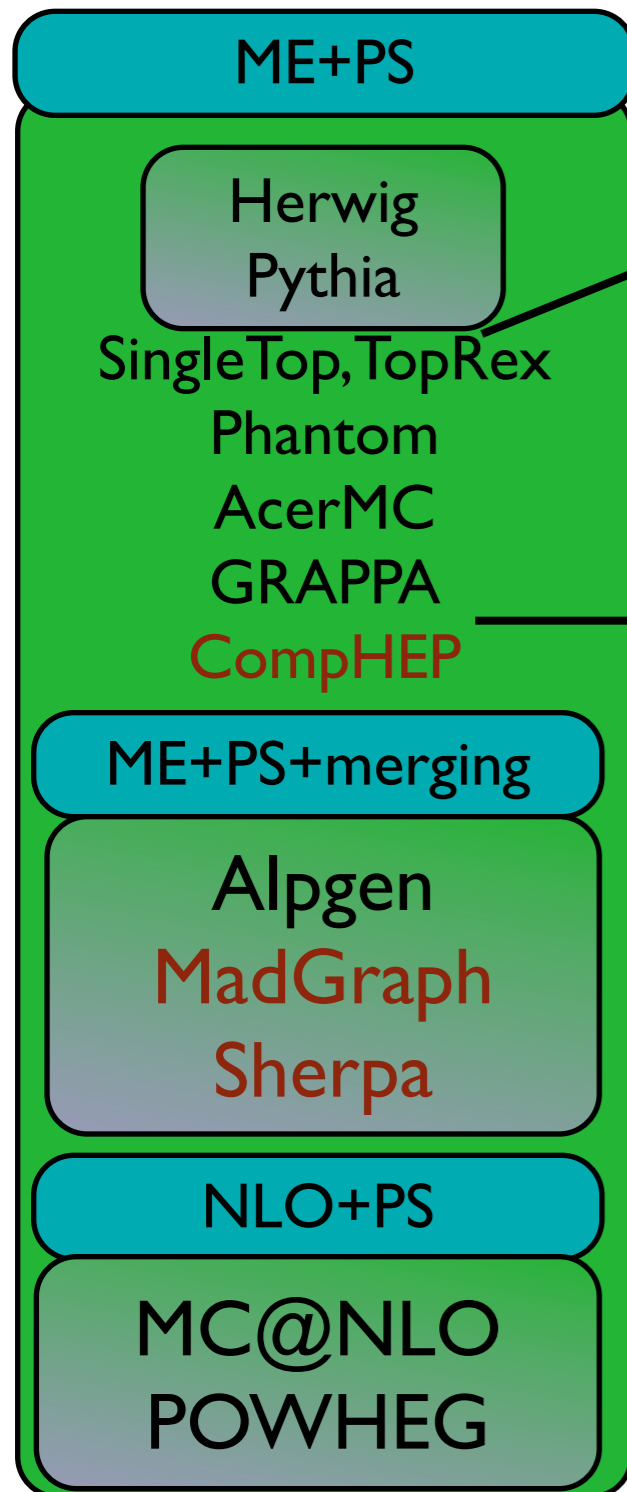
Always the **FIRST** tools. Main purpose is to provide an easily tunable description of the data.

Complete exclusive description of the events: hard scattering, showering & hadronization, underlying event

Significant and intense progress in the development of new showering algorithms with the final aim to go at (N?)NLO [Bauer, Schwartz, 2006; Giele, Kosower, Skands, 2007; Krauss, Schumman, 2007; Bauer, Tackman, Thaler, 2008] . Example: Working implementation for ARIADNE based on dipoles.

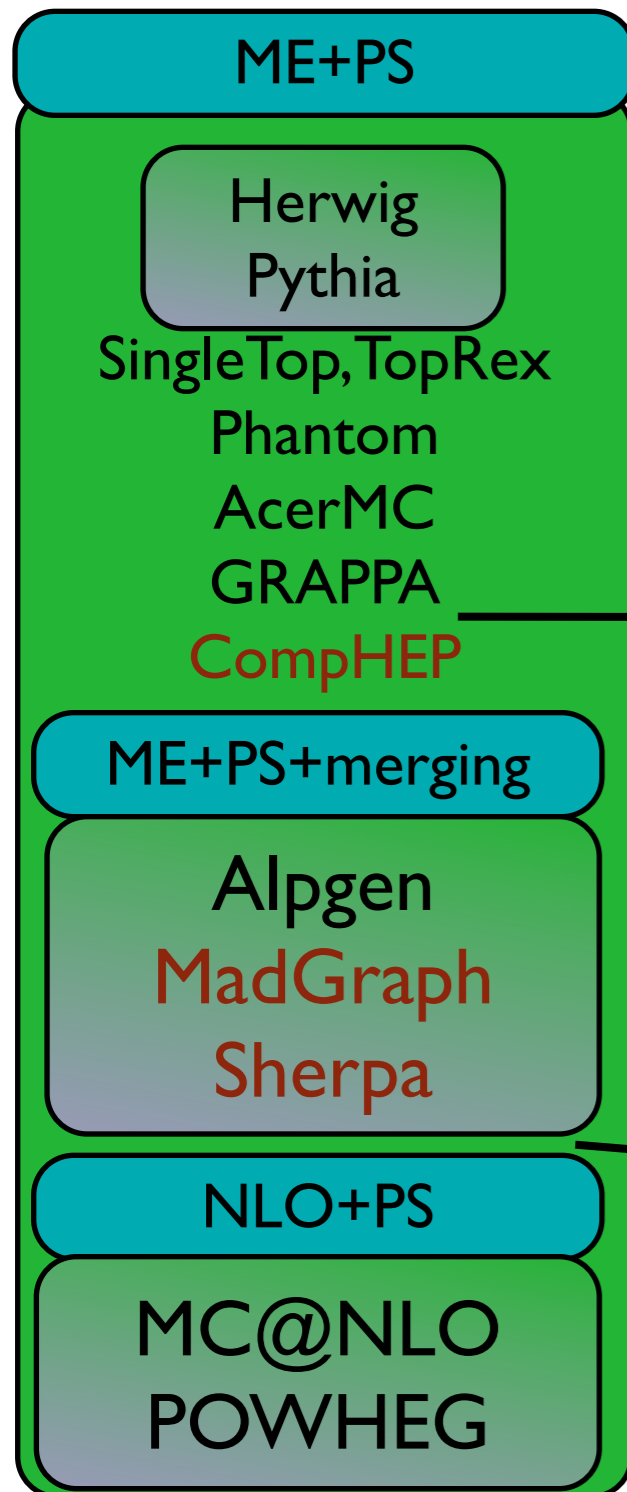
N.B. Apart from SHERPA, which provides its own PS, all other ME codes rely on Pythia or HERWIG for PS, hadro and UE. This is true also for NLO codes (MC@NLO and POWHEG).

# MC tools for top physics in a nutshell



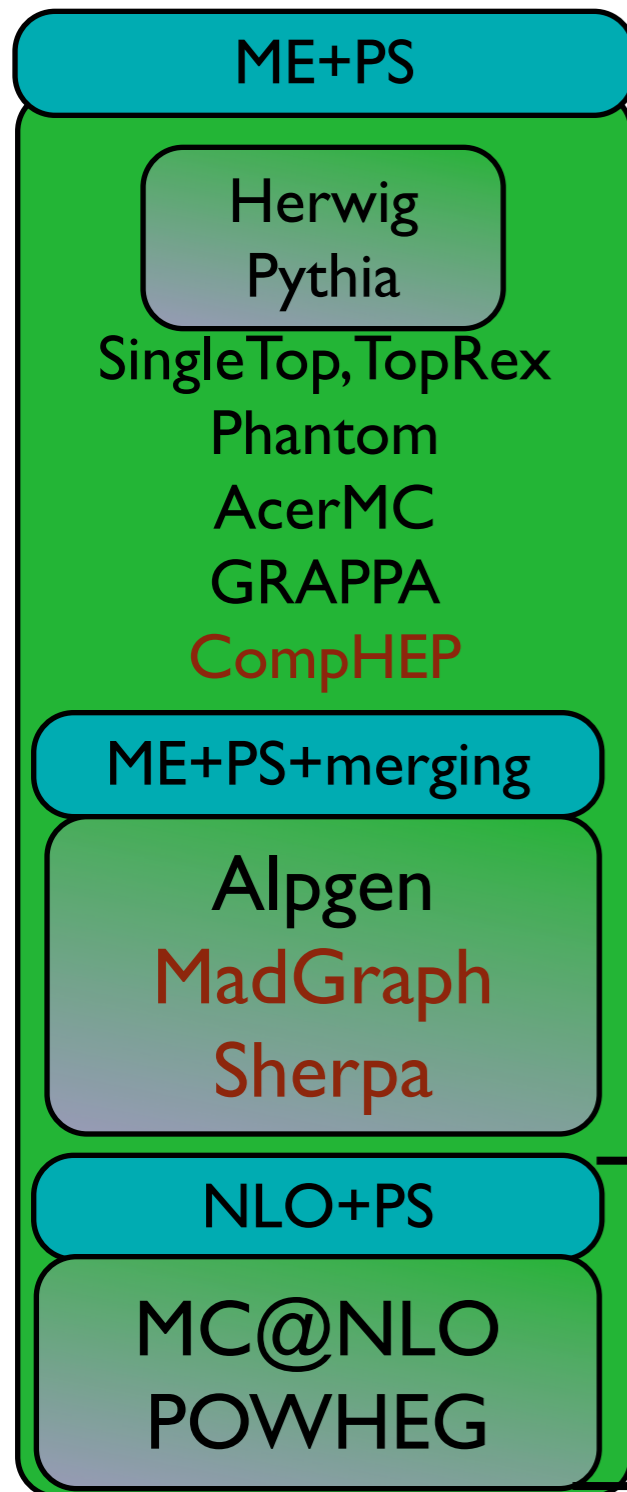
Optimized for a few or several processes. Based on a library of matrix elements (analytic or numeric). Limited in breath but easy to use, and optimal for specific studies. For example,  $pp \rightarrow tt + 6f$  studies. Some of them are localized in a specific exp collaboration.

# MC tools for top physics in a nutshell



Multipurpose Matrix Element creators and generators. Calculations are automatic at tree level. Matching is performed with the parton shower to produce inclusive multi-jet samples. Some codes (Alpgen) are optimized for multi-parton ME. Some codes are suitable for BSM physics.

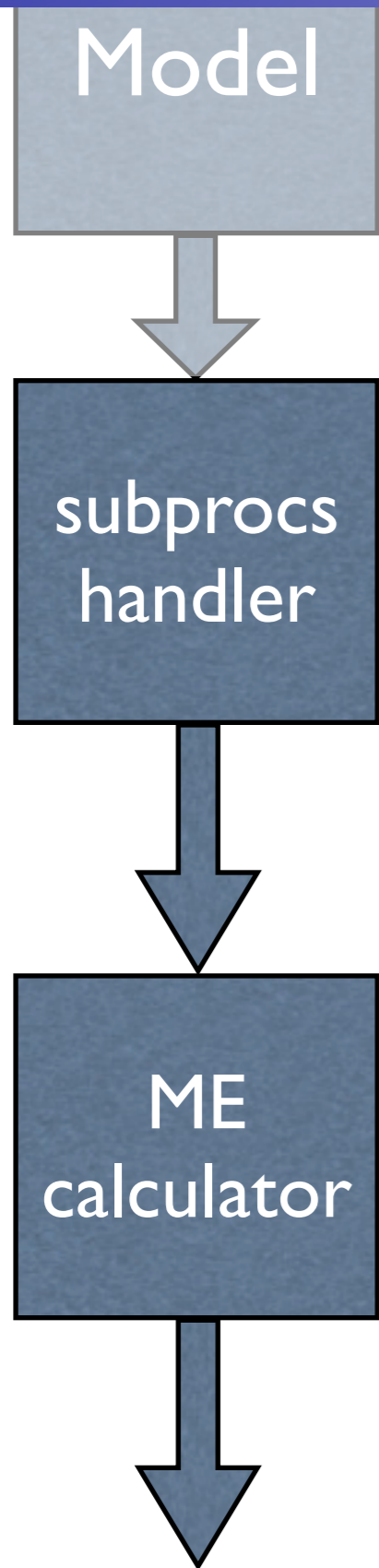
# MC tools for top physics in a nutshell



Combine NLO accuracy in normalization and shapes of hard radiation with parton shower. “Best” tools when NLO calculation is available (i.e. low jet multiplicity). Current limitation is manual work  $\Rightarrow$  small libraries. Only SM.

ME+PS

# Matrix Element based MC's

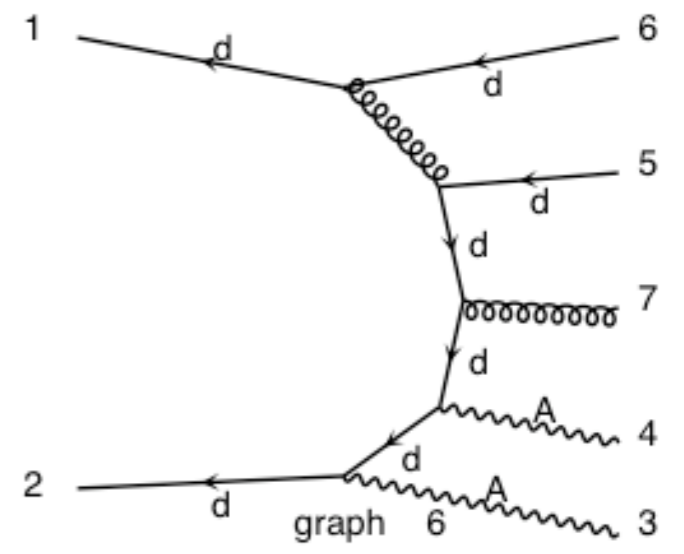


Includes all possible subprocess leading to a given multi-jet final state automatically or manually (done once for all)

“Automatically” generates a code to calculate  $|M|^2$  for arbitrary processes with many partons in the final state.

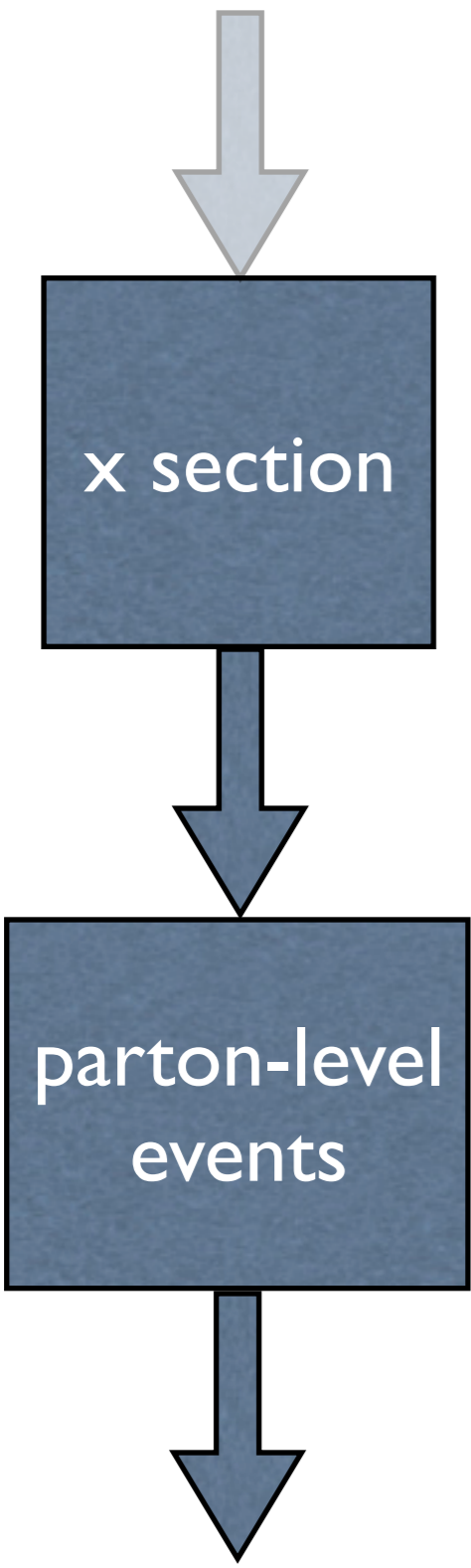
Most use Feynman diagrams w/ tricks to reduce the factorial growth [MadGraph, SHERPA], others have recursive relations to reduce the complexity to exponential [Alpgen, HELAC, Comix].

- $d \sim d \rightarrow a a u u \sim g$
- $d \sim d \rightarrow a a c c \sim g$
- $s \sim s \rightarrow a a u u \sim g$
- $s \sim s \rightarrow a a c c \sim g$



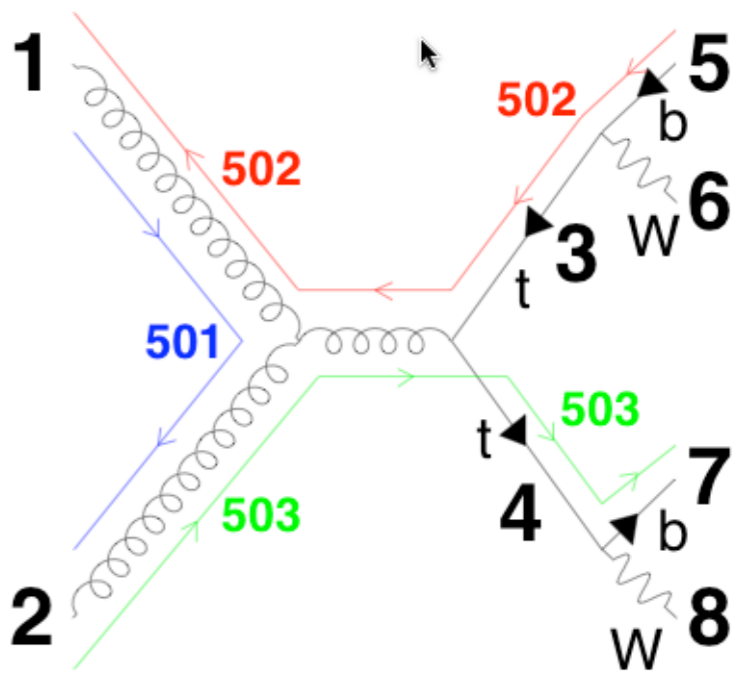
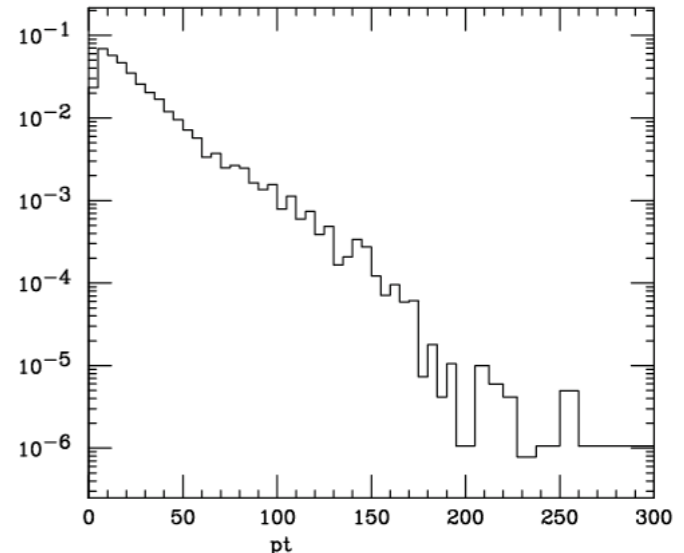


# General structure

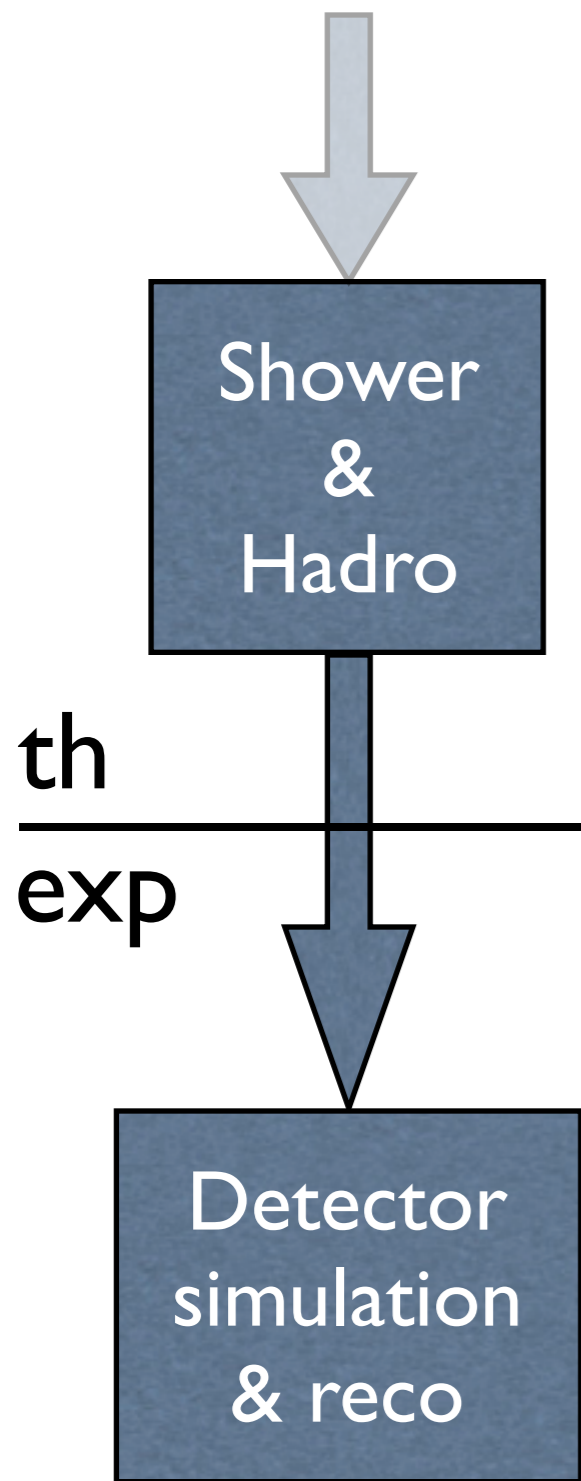


Integrate the matrix element over the phase space using a multi-channel technique and using parton-level cuts.

Events are obtained by unweighting. These are at the parton-level. Information on particle id, momenta, spin, color and mother-daughter is given in the Les Houches format.

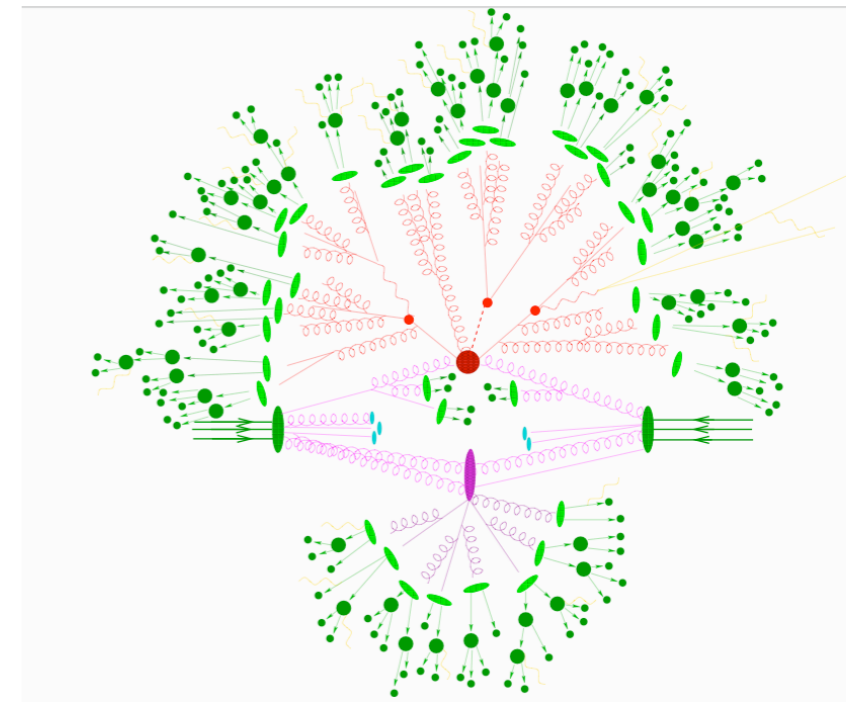


# General structure

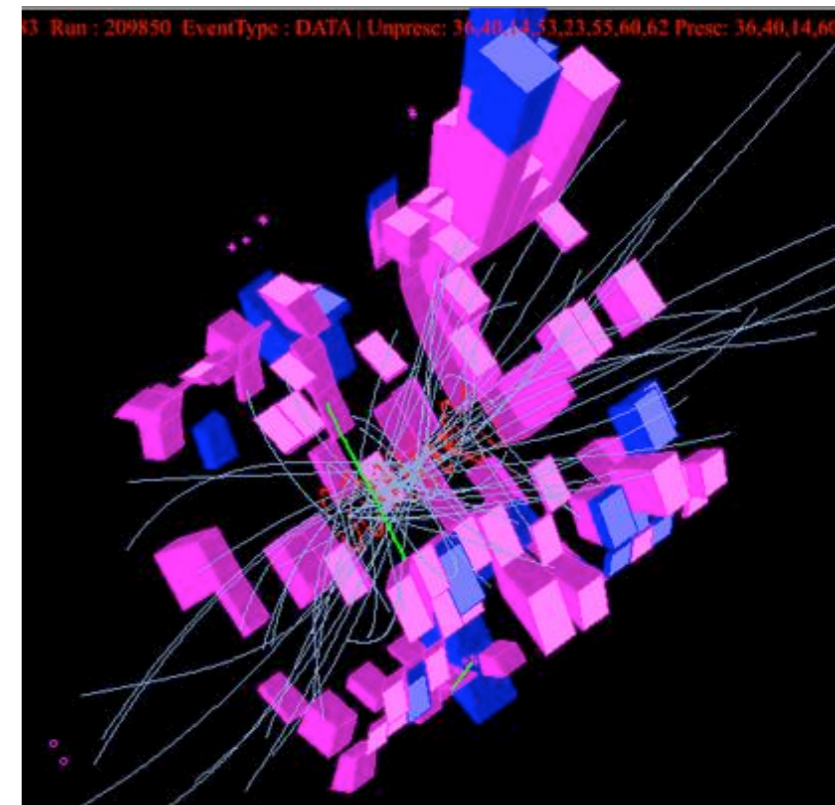


Events in the LH format are passed to the showering and hadronization  $\Rightarrow$  high multiplicity hadron-level events

Parton-Jet merging (MLM or CKKW) happens here!

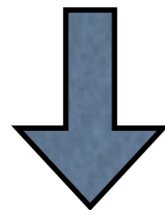


Events in stdhep format are passed through fast or full simulation, and physical objects (leptons, photons, jet, b-jets, taus) are reconstructed.



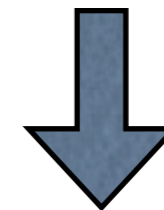
# ME/PS matching

ME



1. parton-level description
2. fixed order calculation
3. quantum interference exact
4. valid when partons are hard and well separated
5. needed for multi-jet description

Shower MC

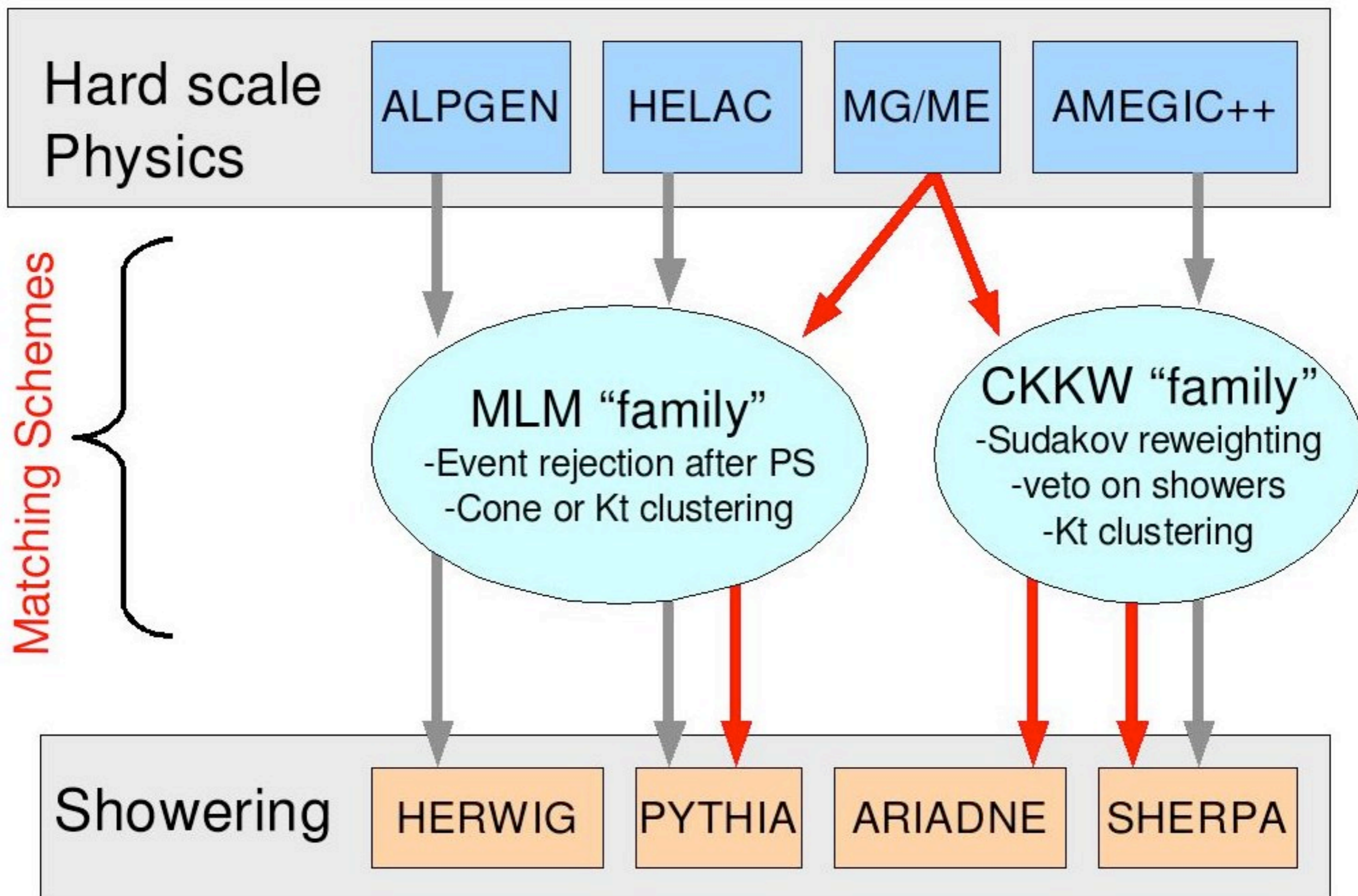


1. hadron-level description
2. resums large logs
3. quantum interference through AA
4. valid when partons are collinear and/or soft
5. needed for realistic studies

**Approaches are complementary!**

**Two recipes available: CKKW and MLM**

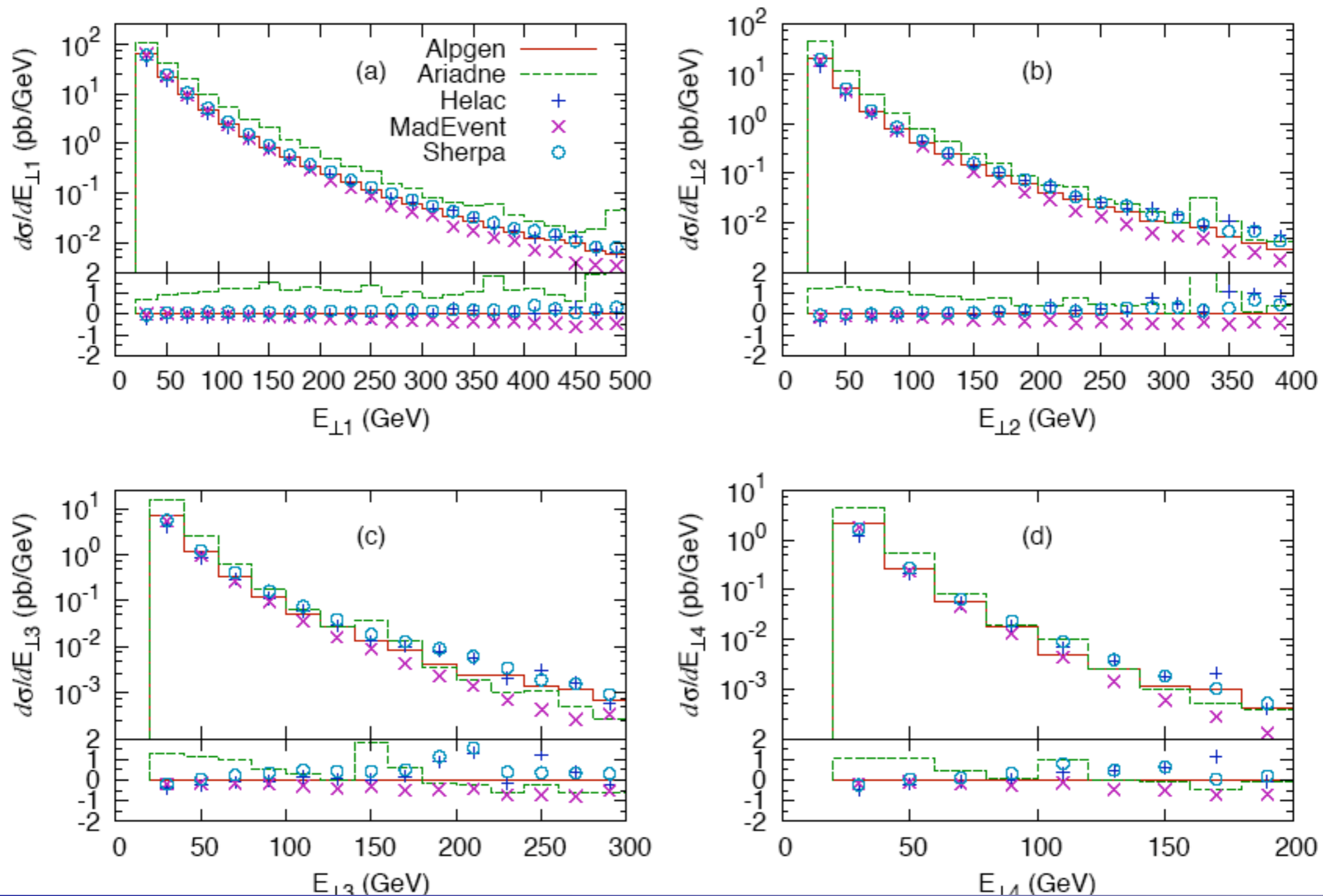
# Available Codes with ME+PS matching



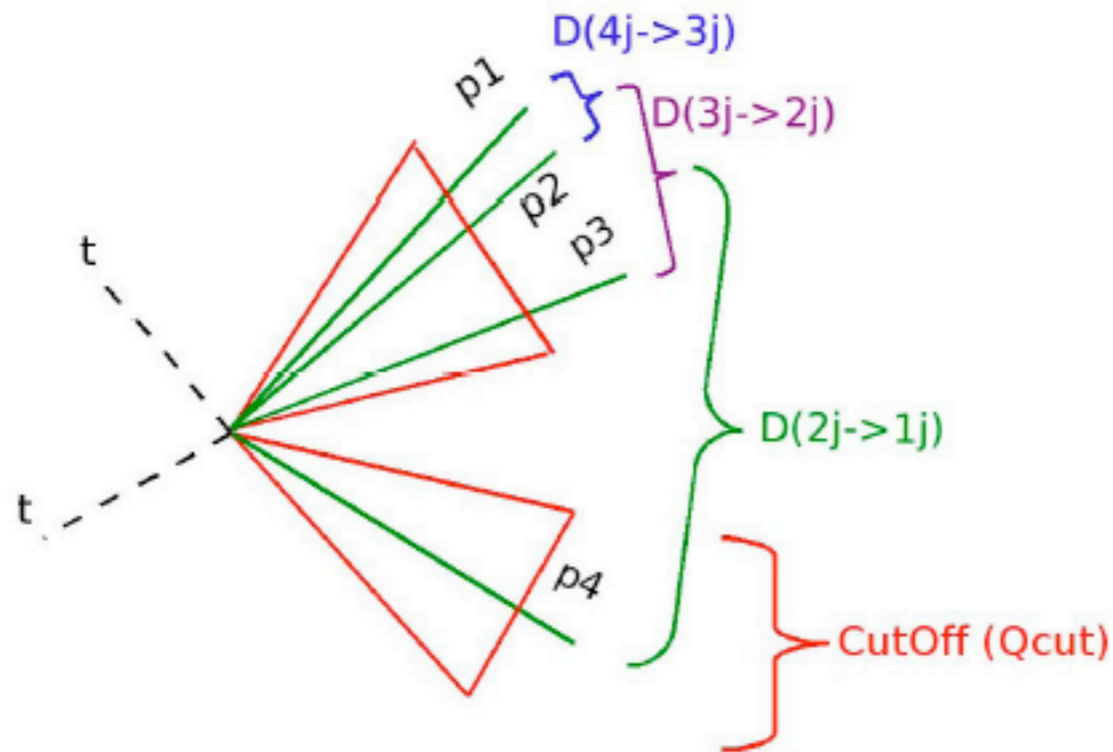
# W+ jets: first comparison

$W^\pm$  + jets comparison plots: Jet  $E_T$  for LHC

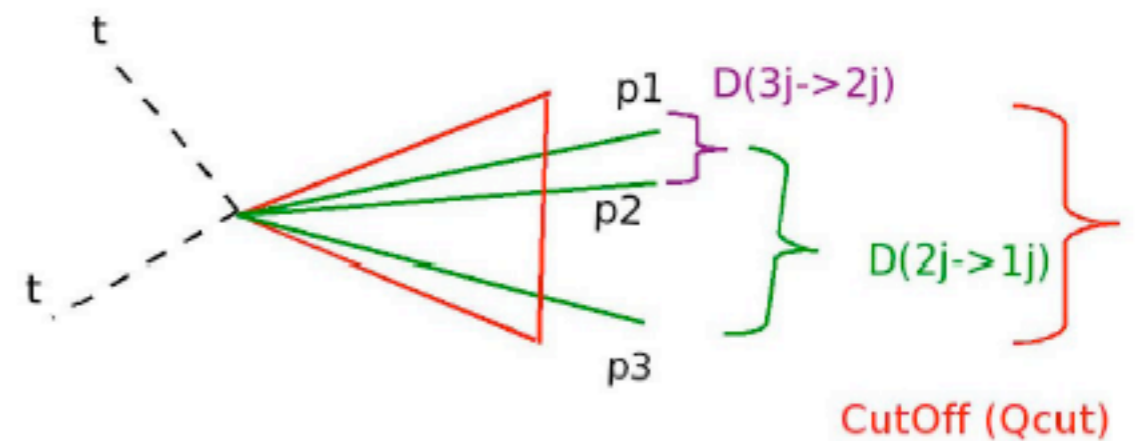
[J. Alwall et al., arXiv:0706.2569]



# Sanity checks: differential jet rates



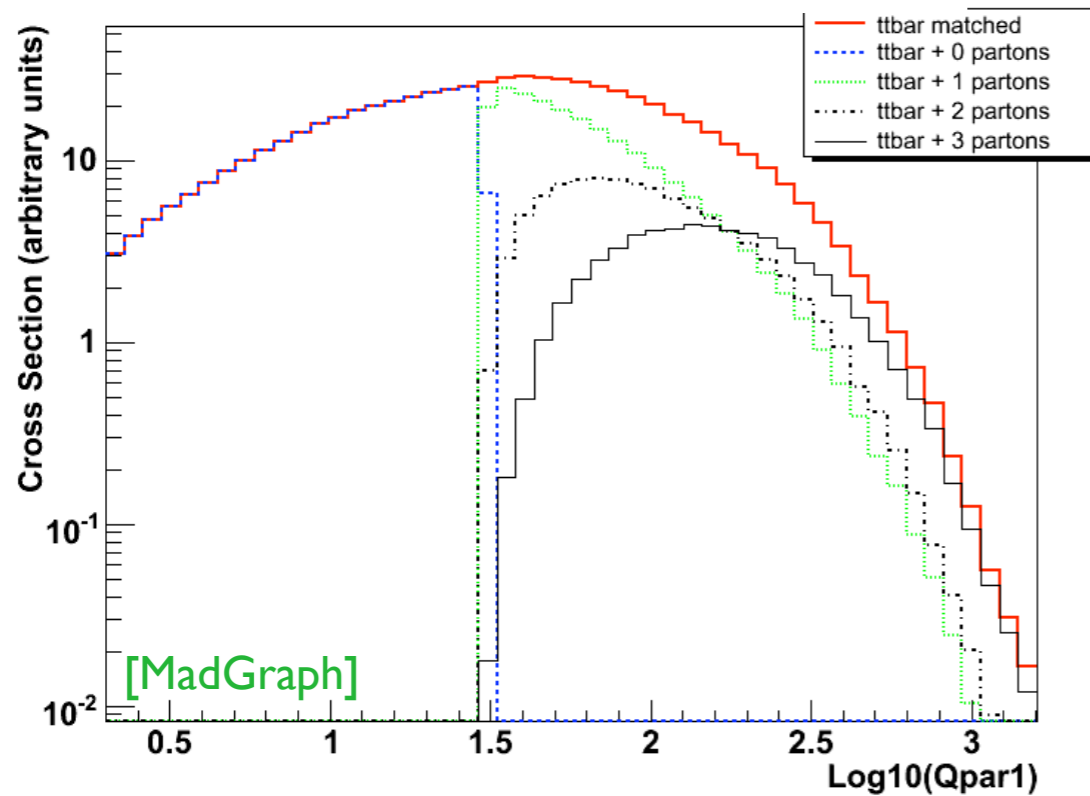
$tt + 2 \text{ partons} + \text{PS}$   
 $D(2 \rightarrow 1) > Q_{\text{cut}}$



$tt + 1 \text{ parton} + \text{PS}$   
 $D(2 \rightarrow 1) < Q_{\text{cut}}$

# Sanity checks: differential jet rates

Diff  $l \rightarrow 0$  jet rates for  $pp \rightarrow t\bar{t} + \text{jets}$  at the LHC

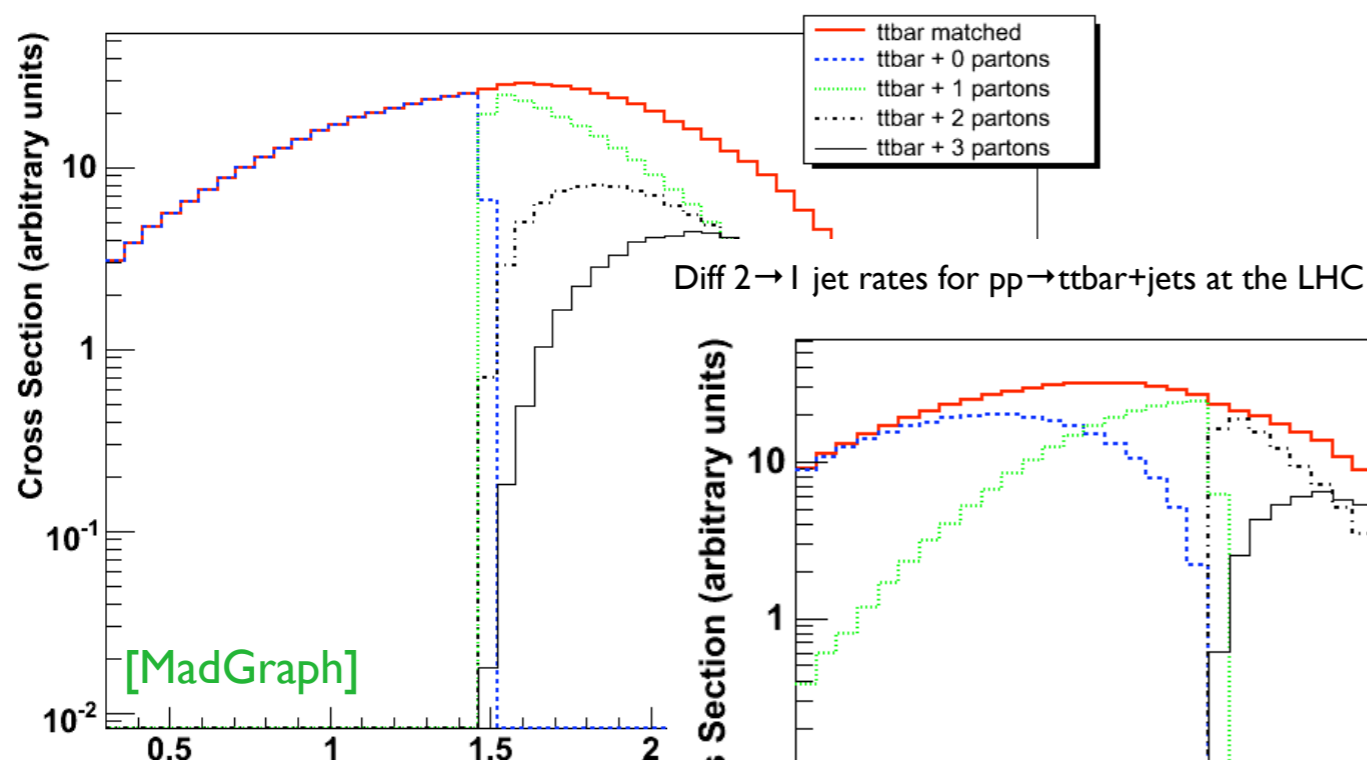


Jet rates are:

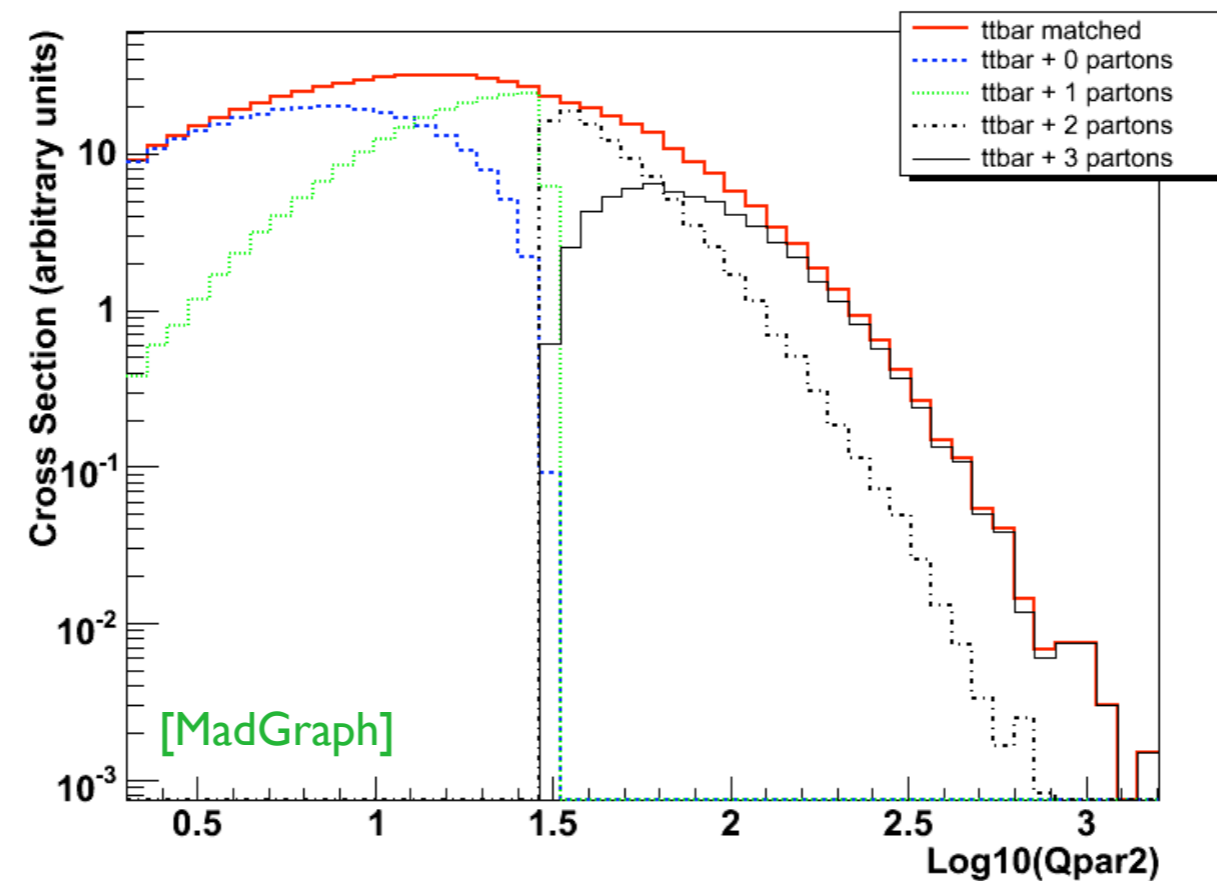
- \* smooth at the cutoff scale
- \* independent of the cutoff scale

# Sanity checks: differential jet rates

Diff 1 $\rightarrow$ 0 jet rates for  $pp\rightarrow t\bar{t}+jets$  at the LHC



Diff 2 $\rightarrow$ 1 jet rates for  $pp\rightarrow t\bar{t}+jets$  at the LHC



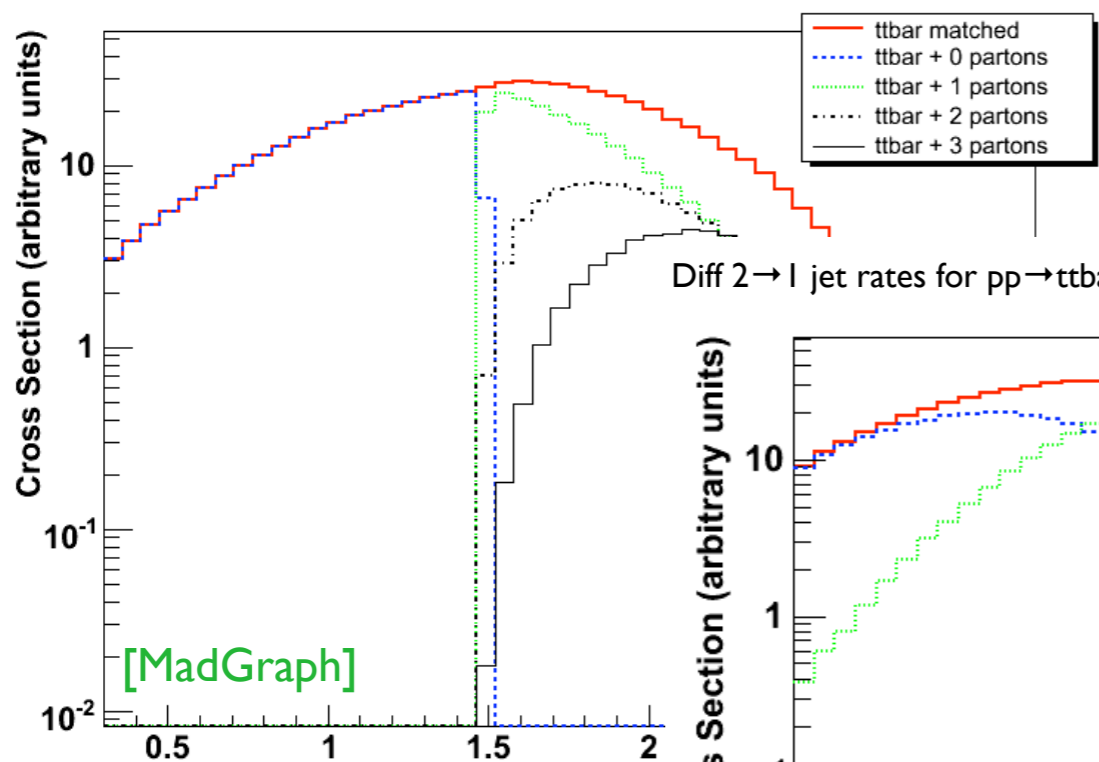
Jet rates are:

- \* smooth at the cutoff scale
- \* independent of the cutoff scale

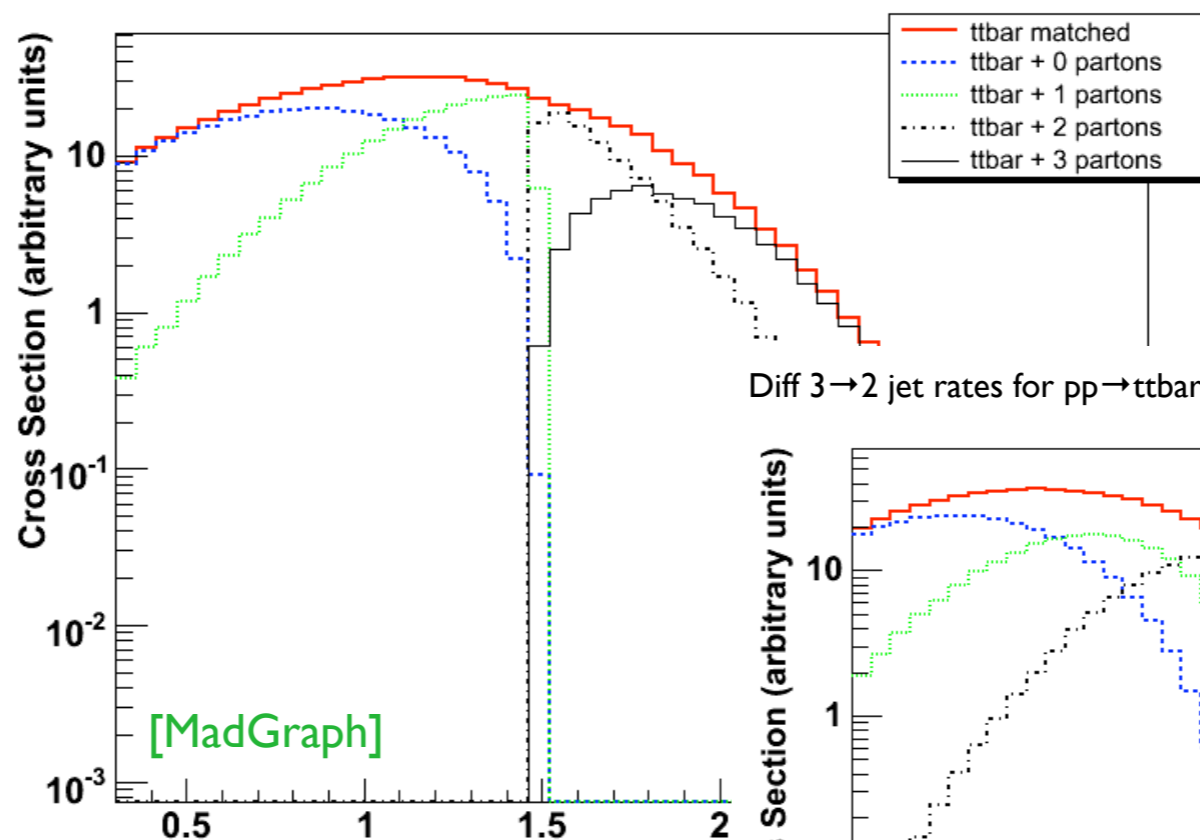


# Sanity checks: differential jet rates

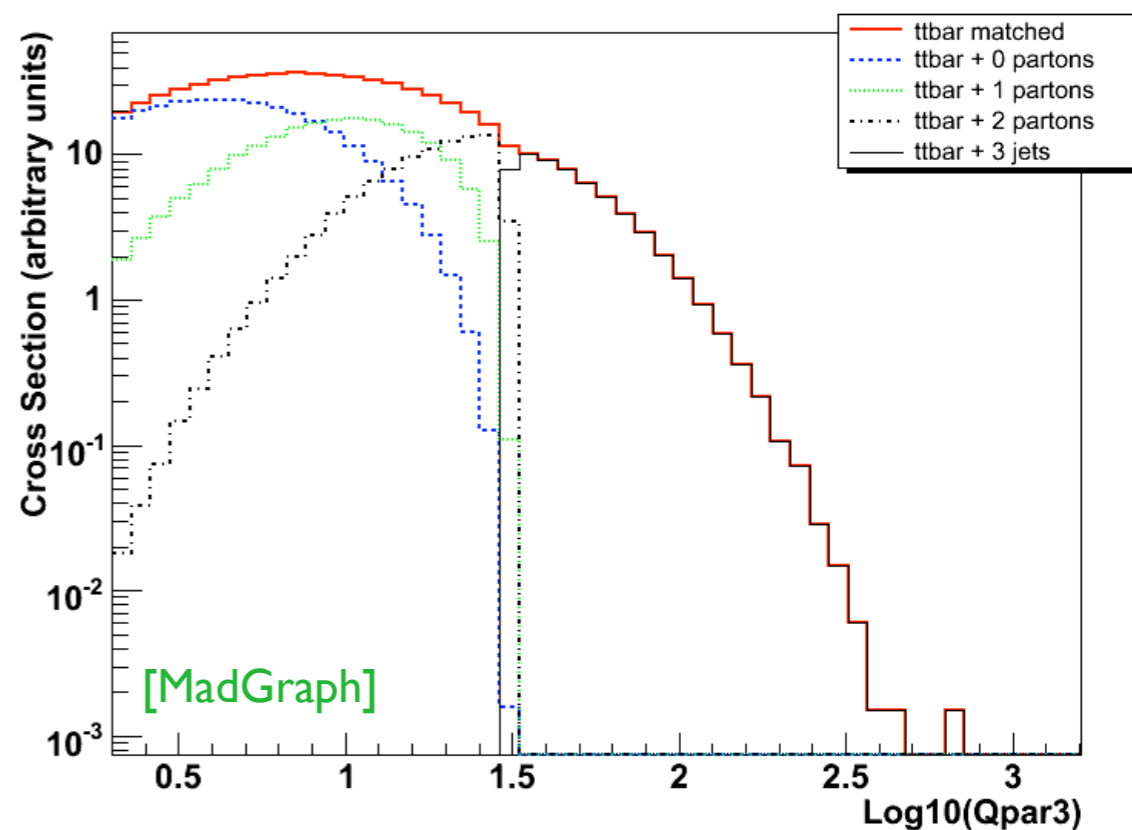
Diff 1→0 jet rates for  $pp \rightarrow t\bar{t} + \text{jets}$  at the LHC



Diff 2→1 jet rates for  $pp \rightarrow t\bar{t} + \text{jets}$  at the LHC



Diff 3→2 jet rates for  $pp \rightarrow t\bar{t} + \text{jets}$  at the LHC

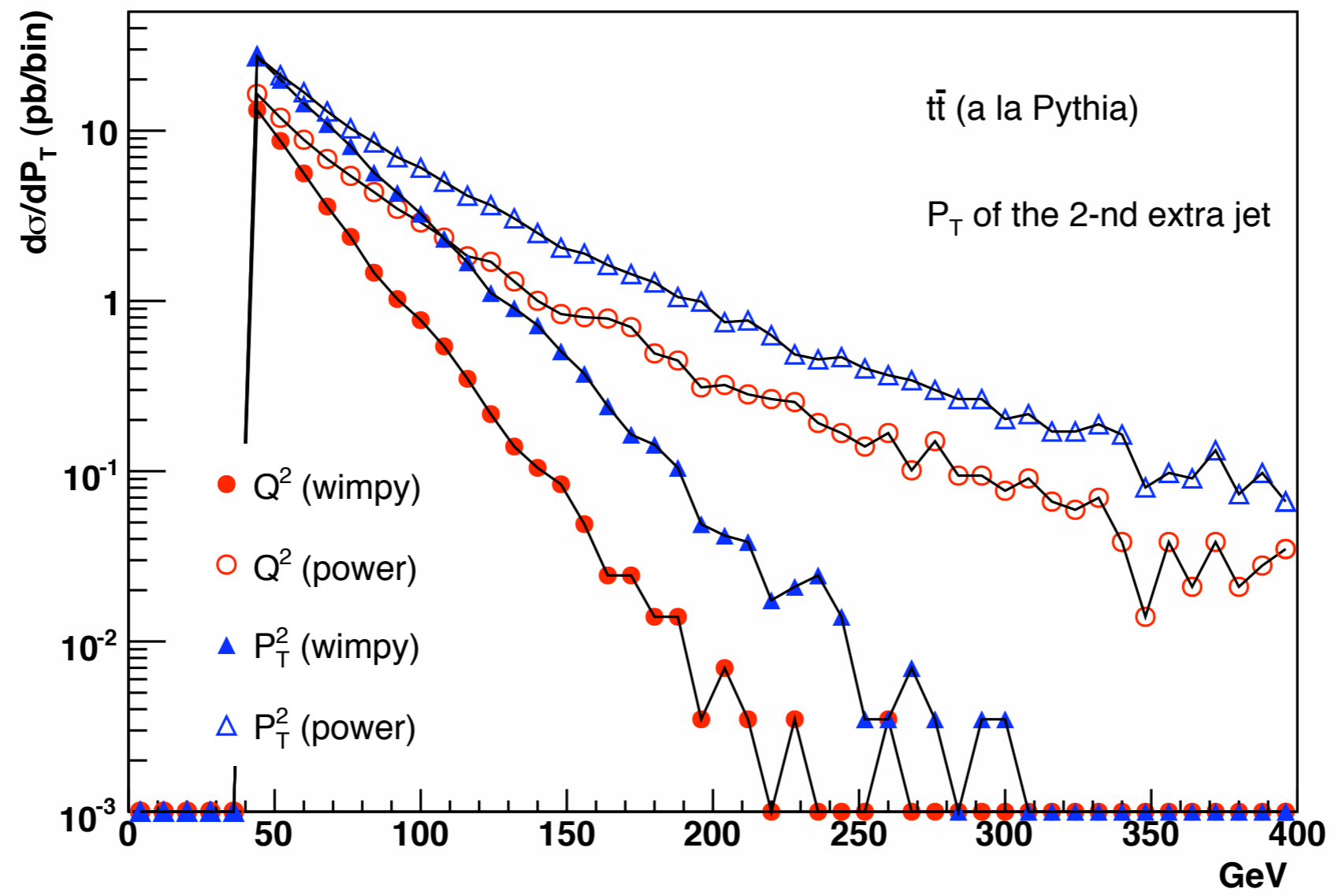


Jet rates are:

- \* smooth at the cutoff scale
- \* independent of the cutoff scale

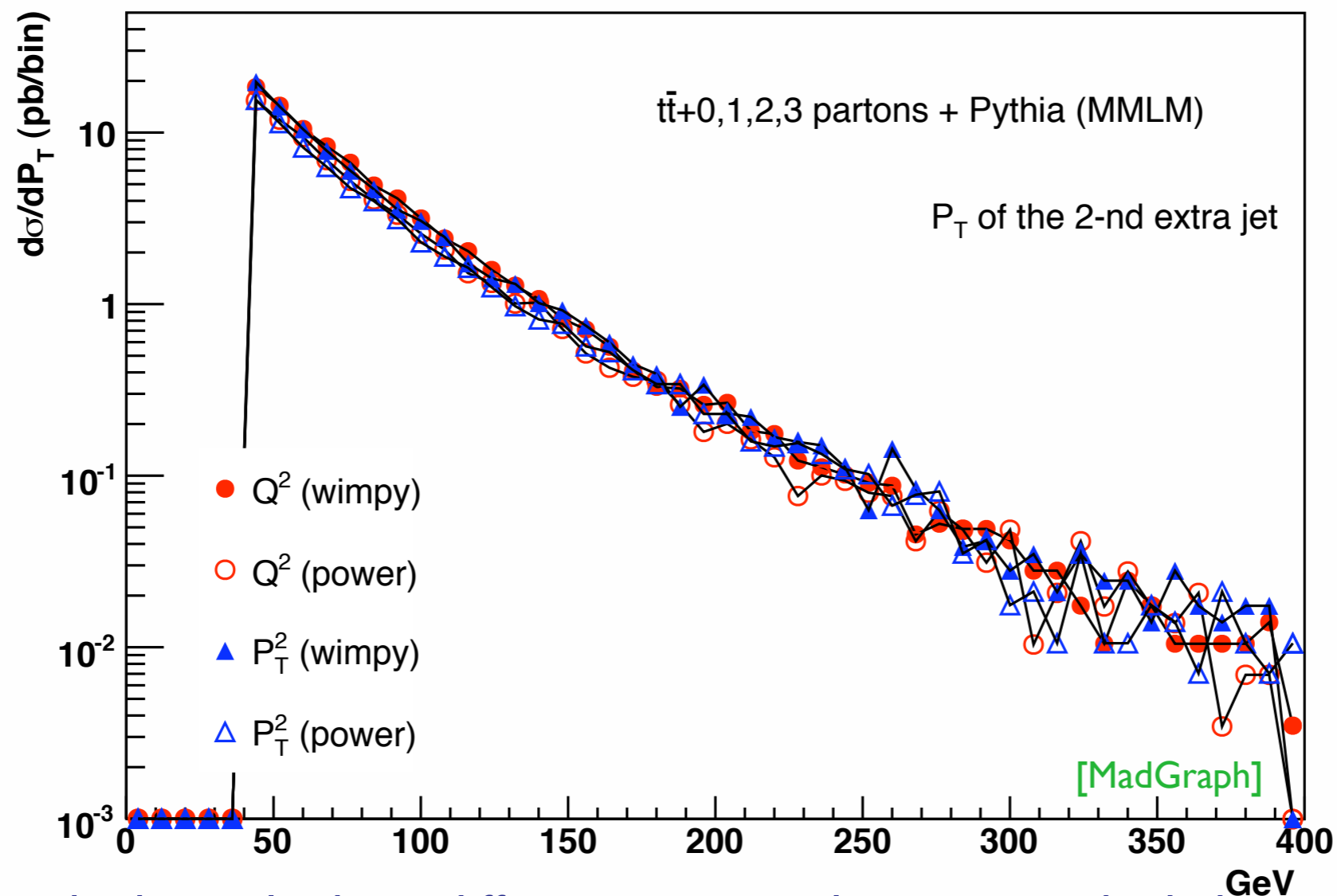
# PS alone vs matched samples

A MC Shower like Pythia is by construction a highly tunable tool. Consider for instance the  $p_T$  distribution of the second jet with different settings:



# PS alone vs matched samples

A MC Shower like Pythia is by construction a highly tunable tool. Consider for instance the  $p_T$  distribution of the second jet with different settings:

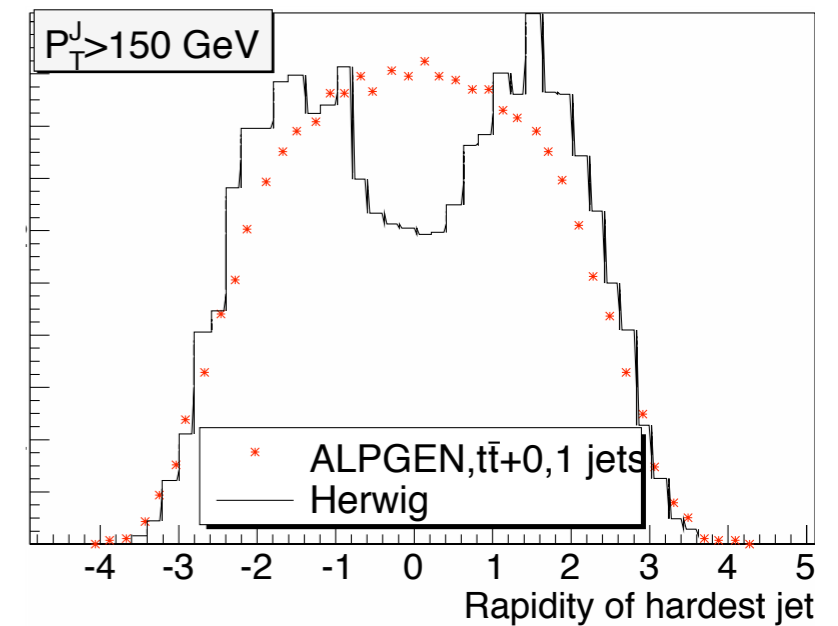
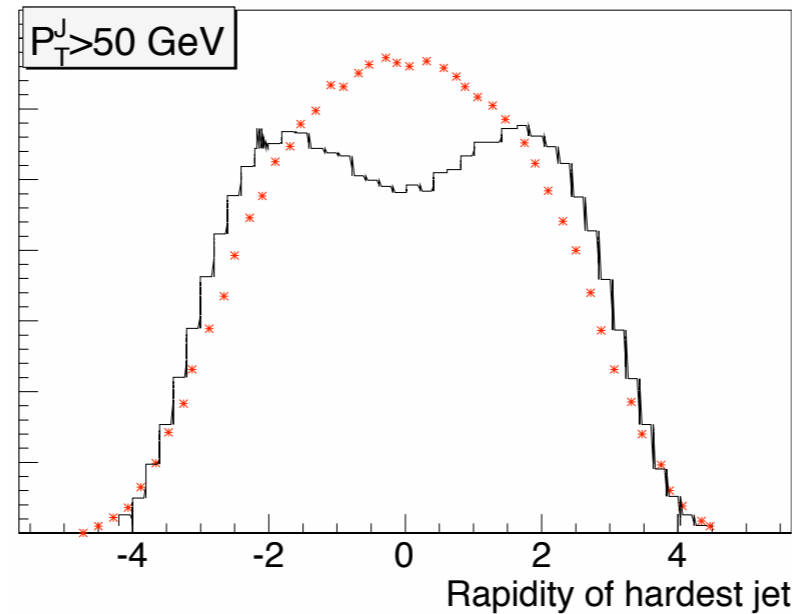
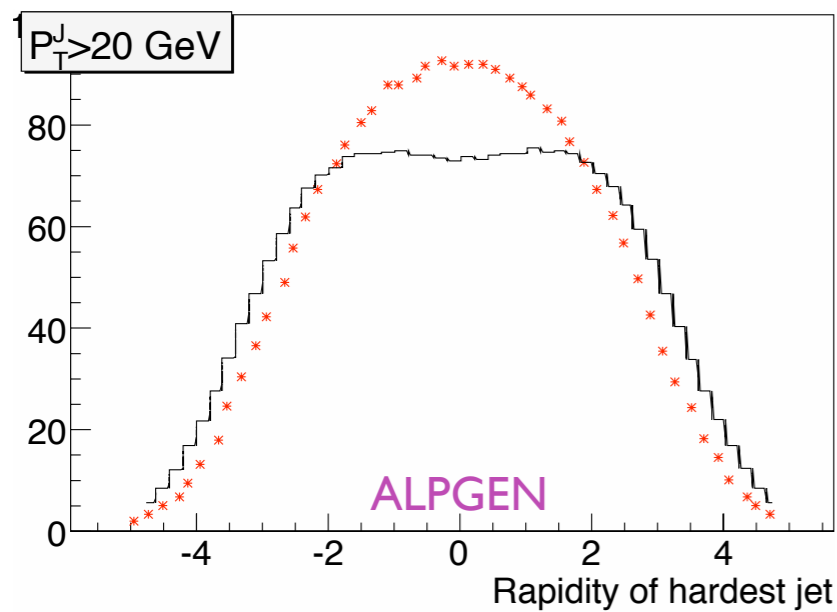


In a matched sample these differences are irrelevant since the behaviour at high  $p_T$  is described by the matrix element  $\Rightarrow$  more predictive power (= less flexibility...)

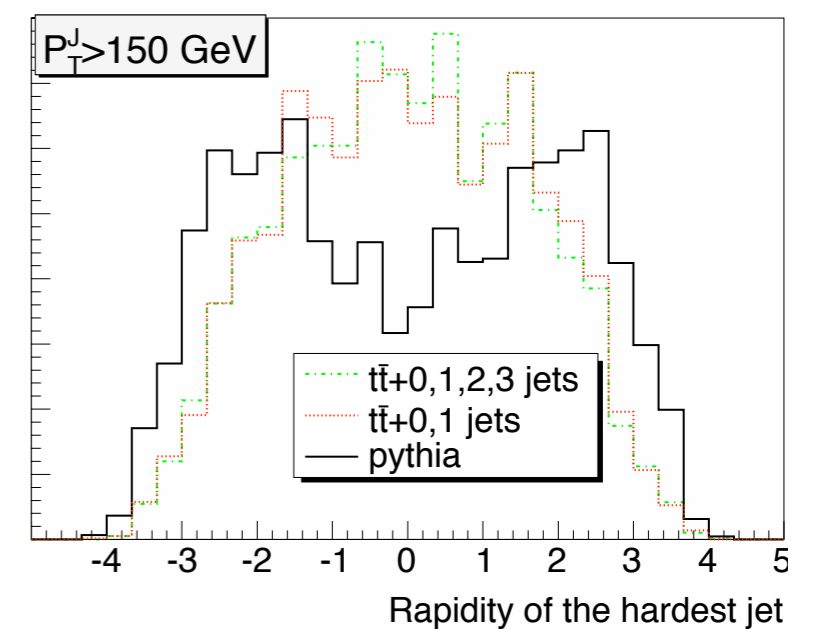
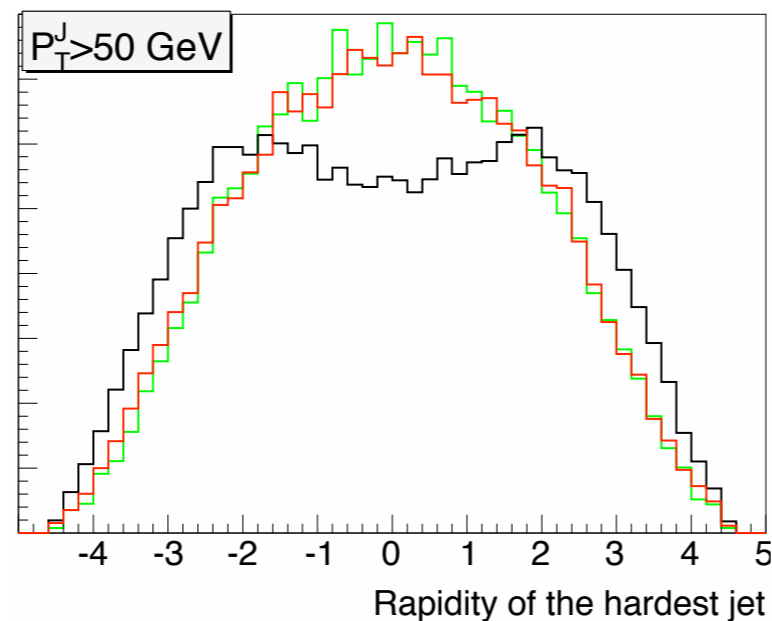
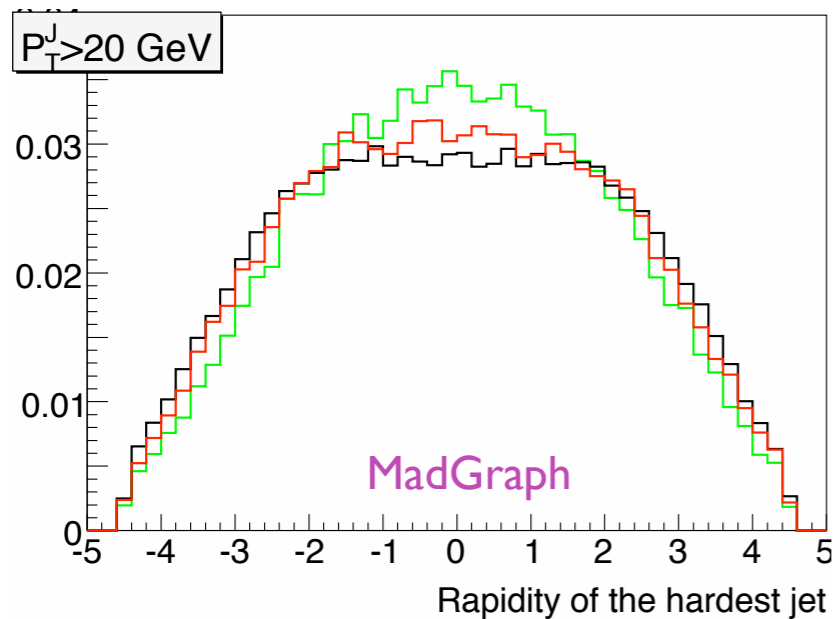
Uncertainties in the matching itself not included.

# Comparisons: 1st jet rapidity

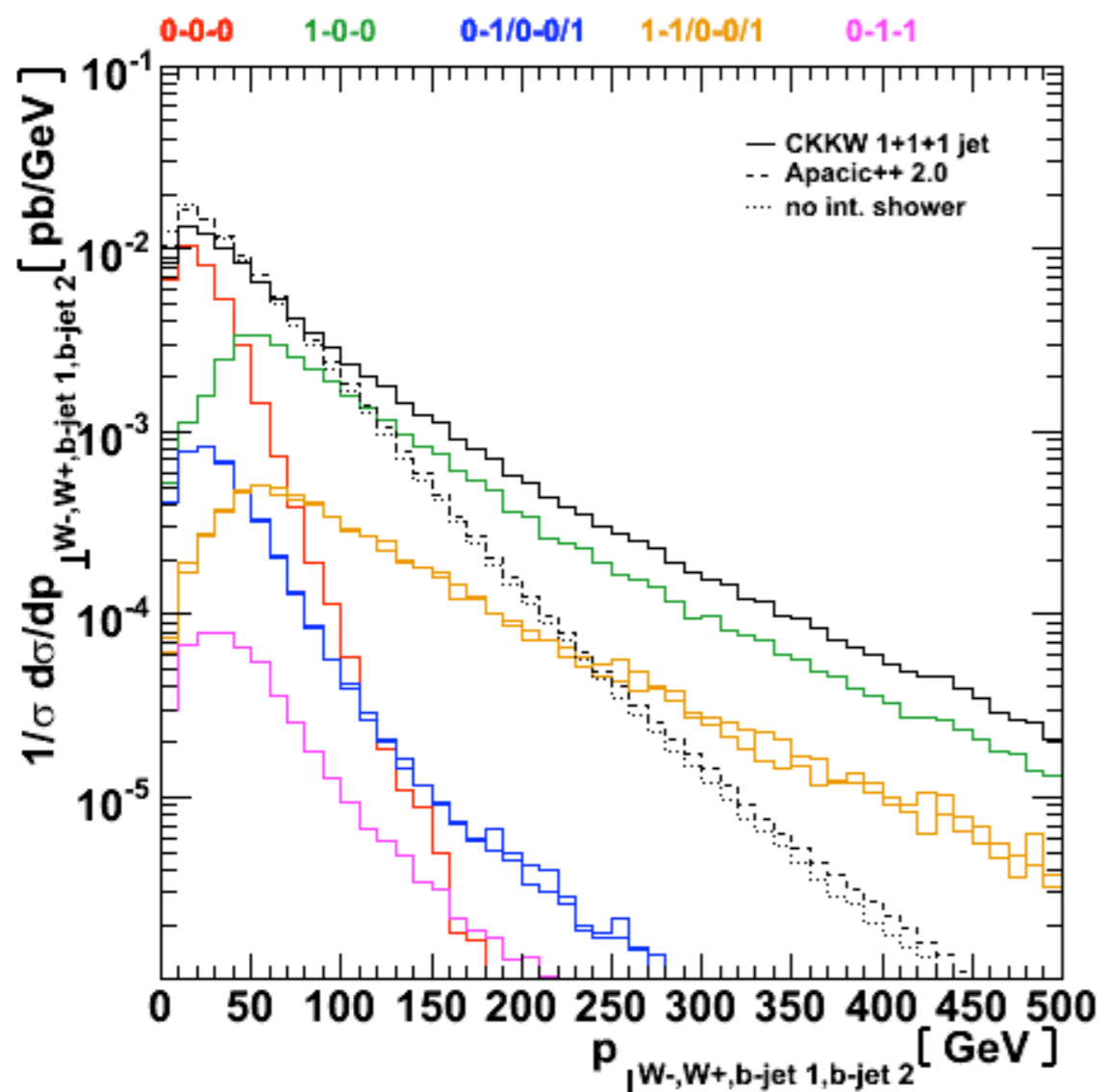
[Mangano, Moretti, Piccinini, Treccani, 2007]



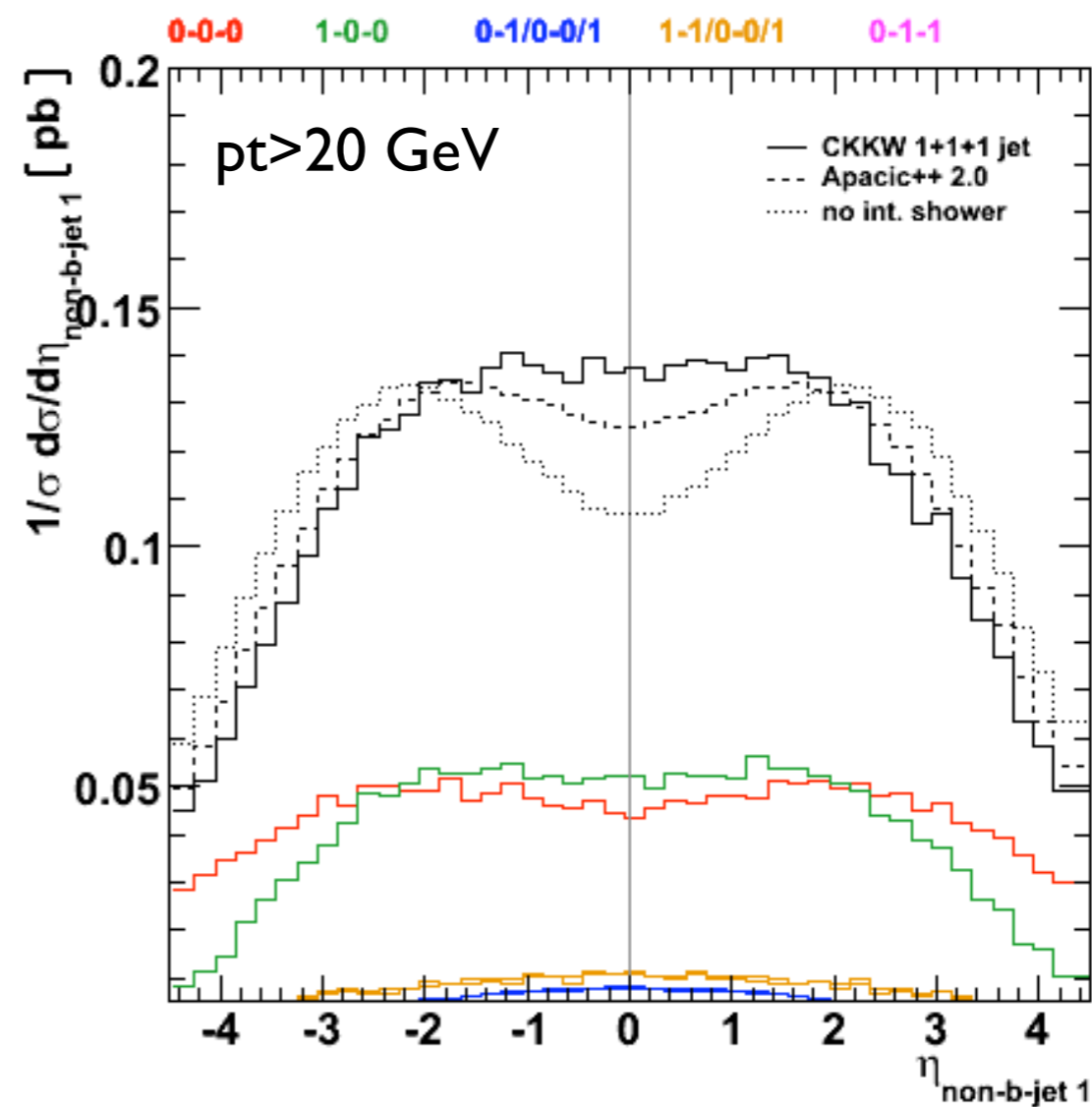
[Alwall, de Visscher, 2007]



# Comparisons: 1st jet rapidity

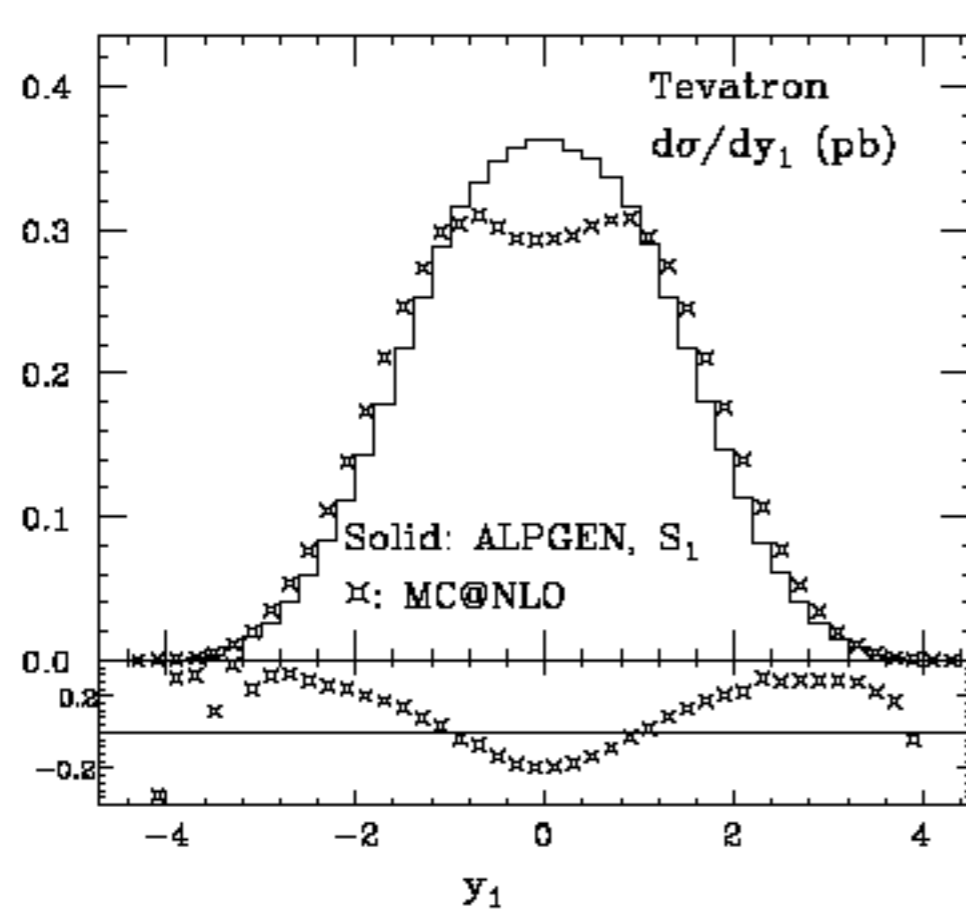


[Sherpa]

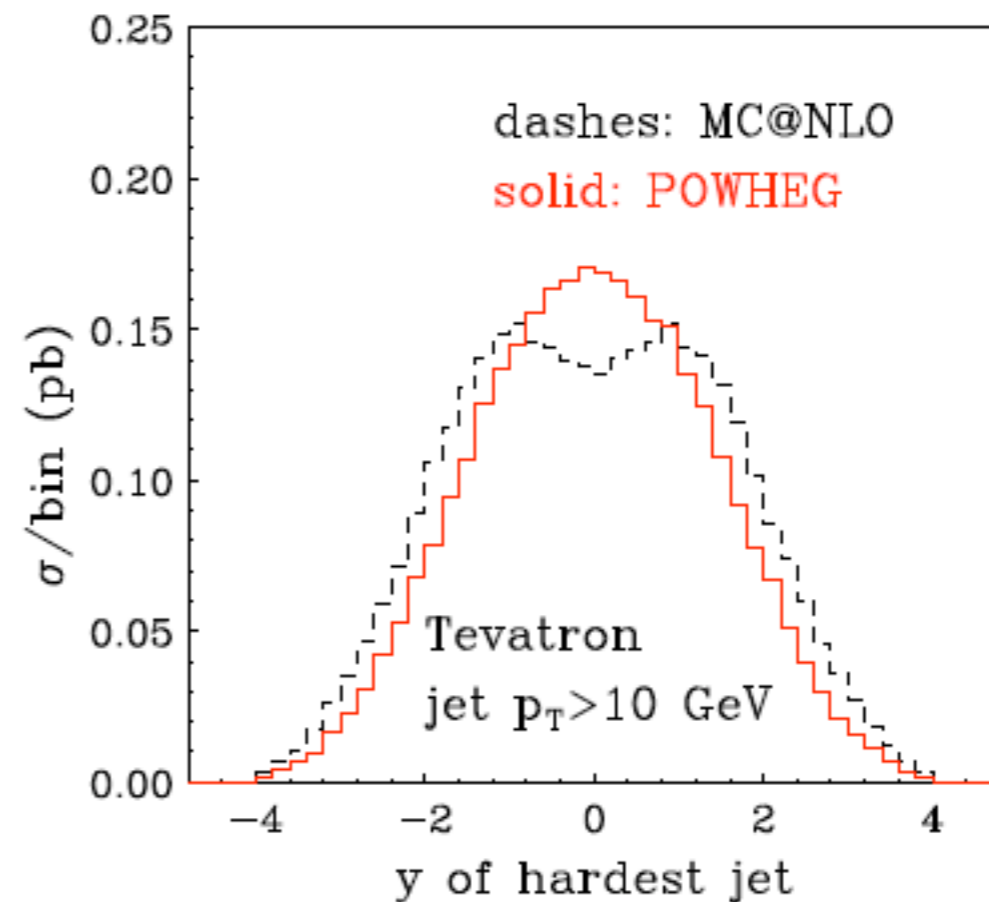


[Sherpa]

# Comparisons: 1st jet rapidity



[Mangano, Moretti, Piccinini, Treccani 2007]



[Frixione, Nason, Ridolfi 2007]

It seems that indeed both Pythia and Herwig develop a deep in the central rapidity region for high-pt jets, which is filled by ME+PS. Hard radiation in MC@NLO is not able to fill it, while POWHEG as a similar behaviour as ME+PS. It will be interesting to see what  $t\bar{t} + 1\text{jet}$  at NLO predicts (ongoing..)

Still a lot to learn by comparing different approaches!

# To remember about the ME-PS matching

- It provides an algorithm to generate multi-jet inclusive samples, that are accurate in all the areas of the phase space avoiding double-counting.
- The matching (à la CKKW) has been rigorously proved in  $e^+e^-$  collisions and it is believed to be true also in  $pp$  collisions. The MLM matching is problematic in  $e^+e^-$  and just a prescription in  $pp$ , where, however, seems to work really well.
- At this stage there is a fair amount of tuning/checking/validation to make it work and evaluation of the systematics is still subject study.
- Since no exact virtual contributions are included the normalization of the cross section for each jet multiplicity is formally LO and therefore uncertain. Normalization has to be obtained from a NLO calculation.
- On the other hand, shapes and often jet rates are (so far) in very good agreement with NLO.
- Fast progress : new studies/proposals/developments every day..

# NLOwPS



## NLO<sub>w</sub>PS

Problem of double counting becomes even more severe at NLO

- \* Real emission from NLO and PS has to be counted once
- \* Virtual contributions in the NLO and Sudakov should not overlap

Current available (and working) solutions:

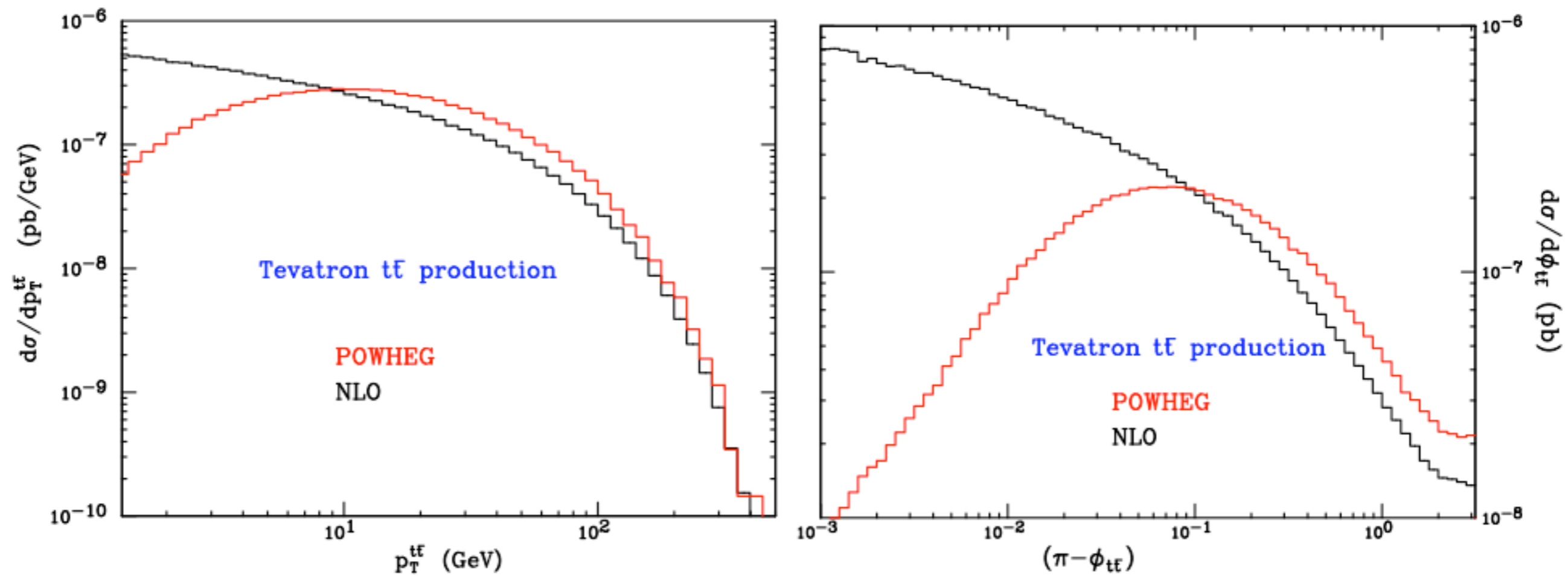
**MC@NLO** [Frixione, Webber, 2003; Frixione, Nason, Webber, 2003]

- Matches NLO to HERWIG angular-ordered PS.
- “Some” work to interface an NLO calculation to HERWIG.  
Uses only FKS subtraction scheme.
- Some events have negative weights.
- Sizable library of procs now.

**POWHEG** [Nason 2004; Frixione, Nason, Oleari, 2007]

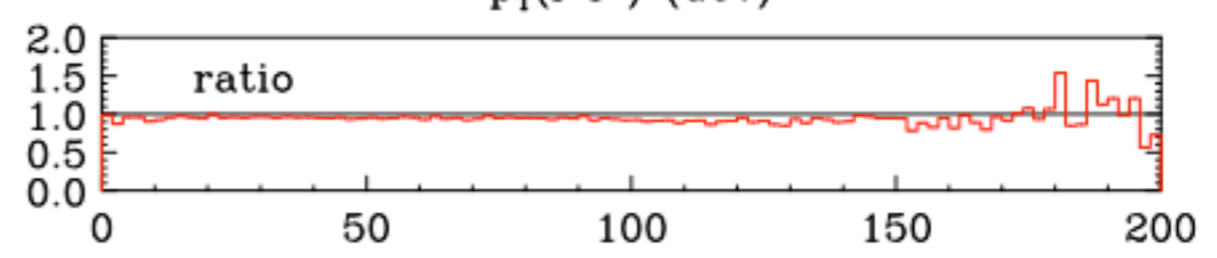
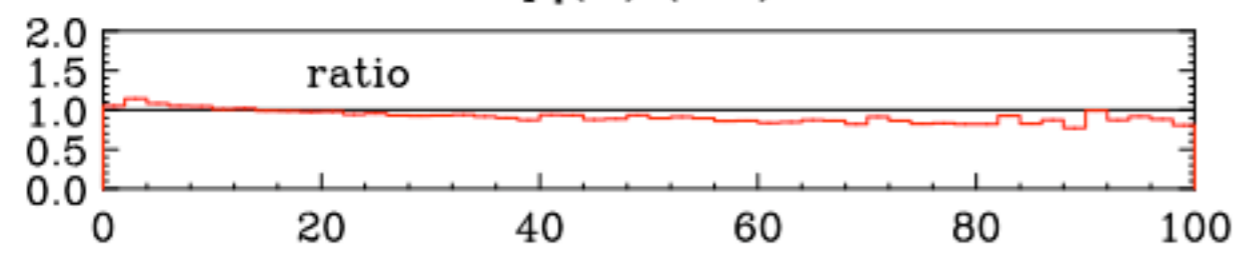
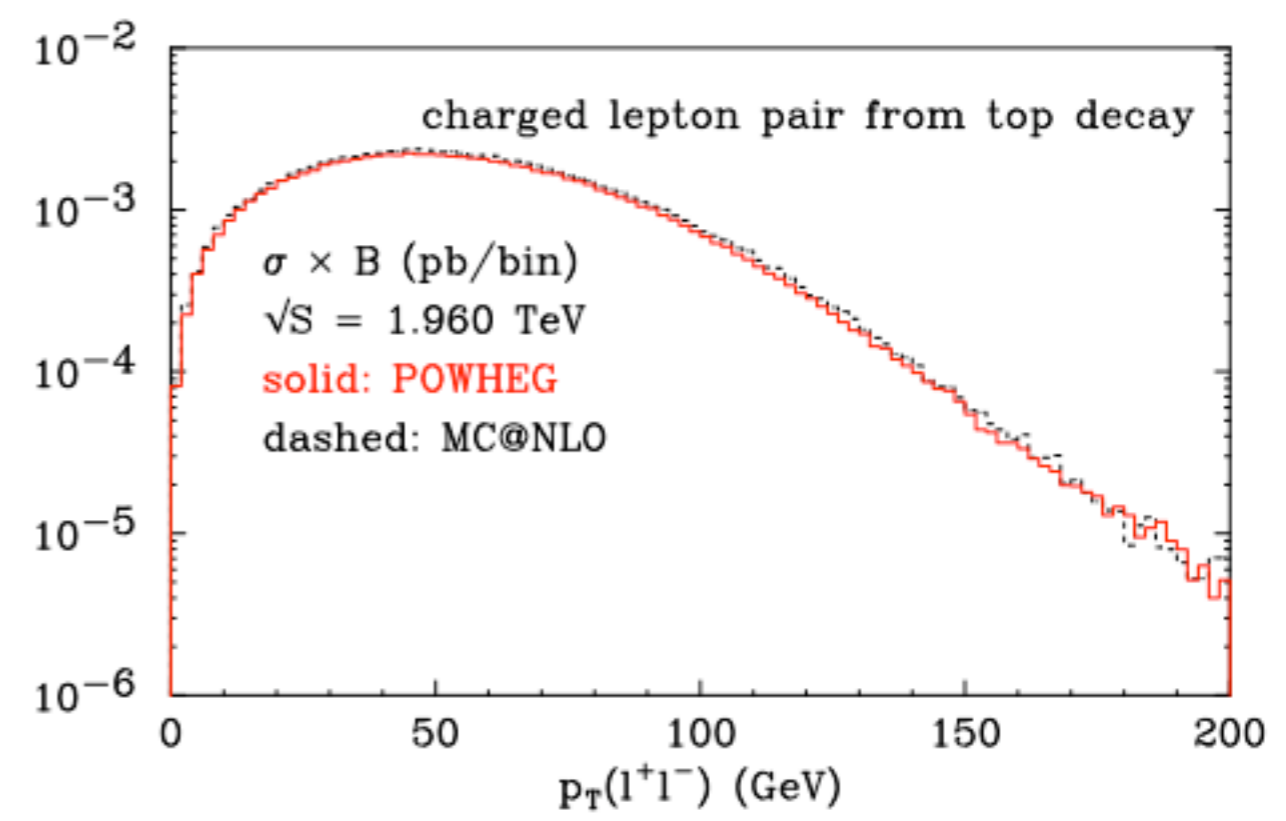
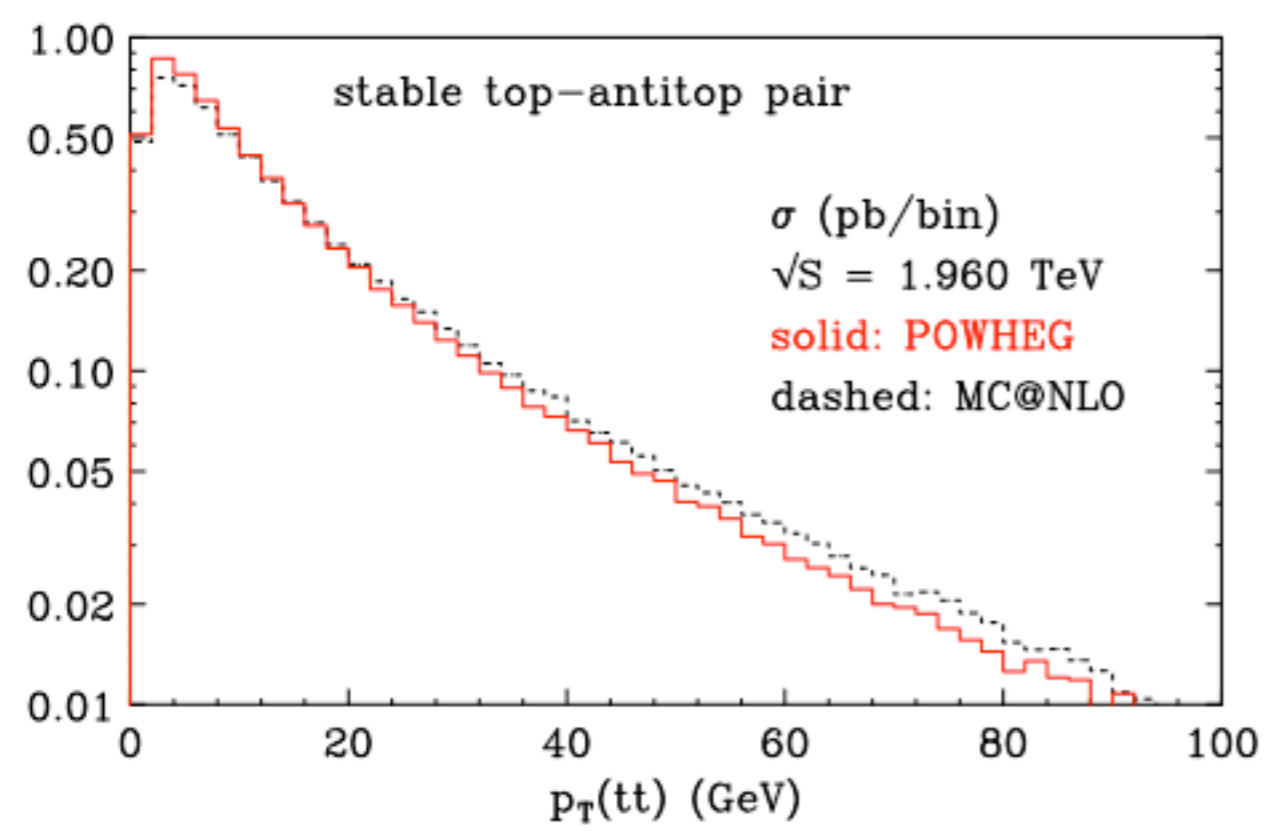
- Is independent from the PS. It can be interfaced to PYTHIA or HERWIG.
- Can use existing NLO results.
- Generates only positive unit weights.
- For top only ttbar (with spin correlations) is available so far.

# $t\bar{t}$ : NLO<sub>w</sub>PS vs NLO



- \* Soft/Collinear resummation of the  $p_T(tt) \rightarrow 0$  region.
- \* At high  $p_T(tt)$  it approaches the  $t\bar{t}$ +parton (tree-level) result.
- \* When  $\Phi(tt) \rightarrow 0$  the emitted radiation is hard and  $t\bar{t}$ -parton result.

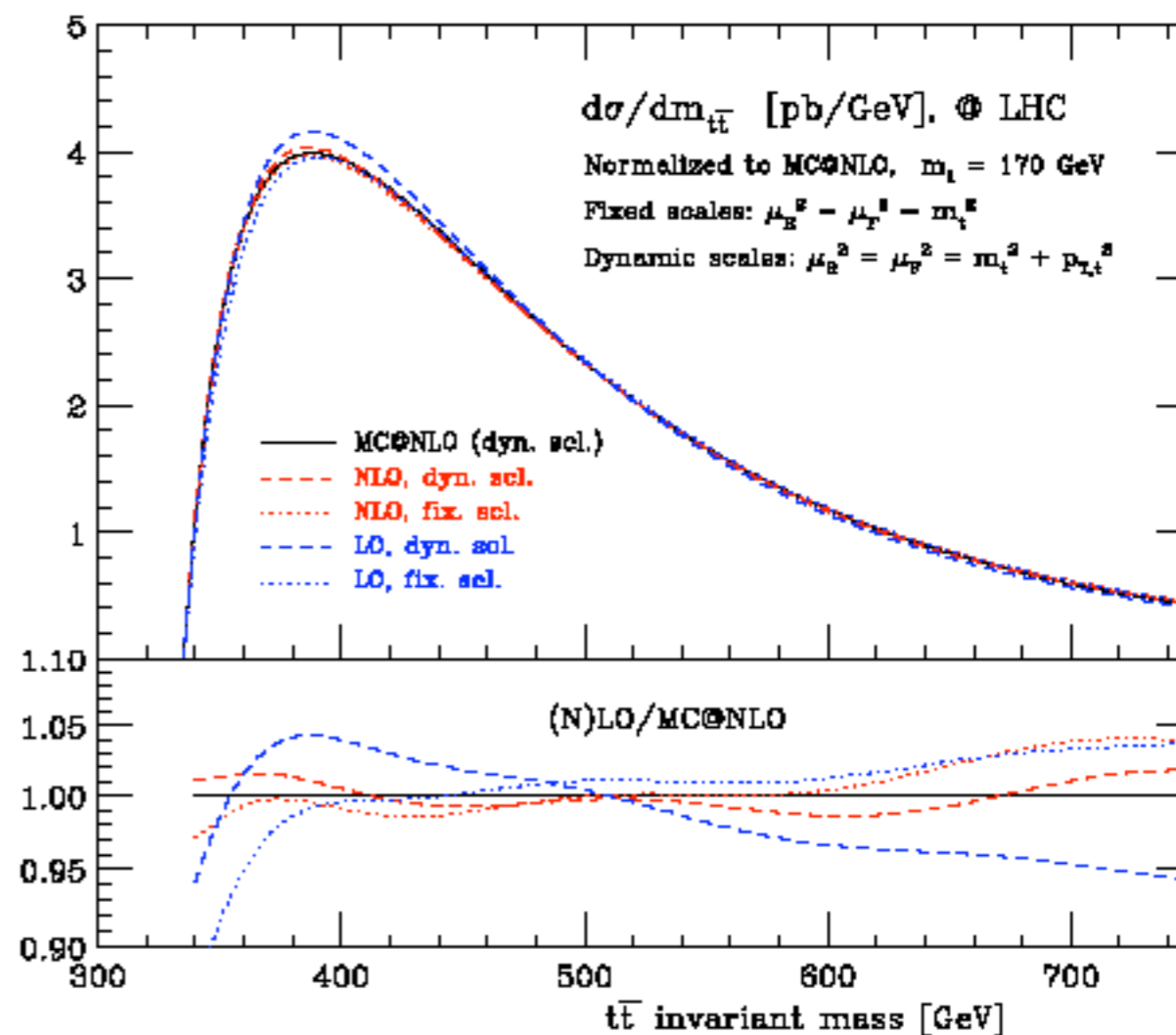
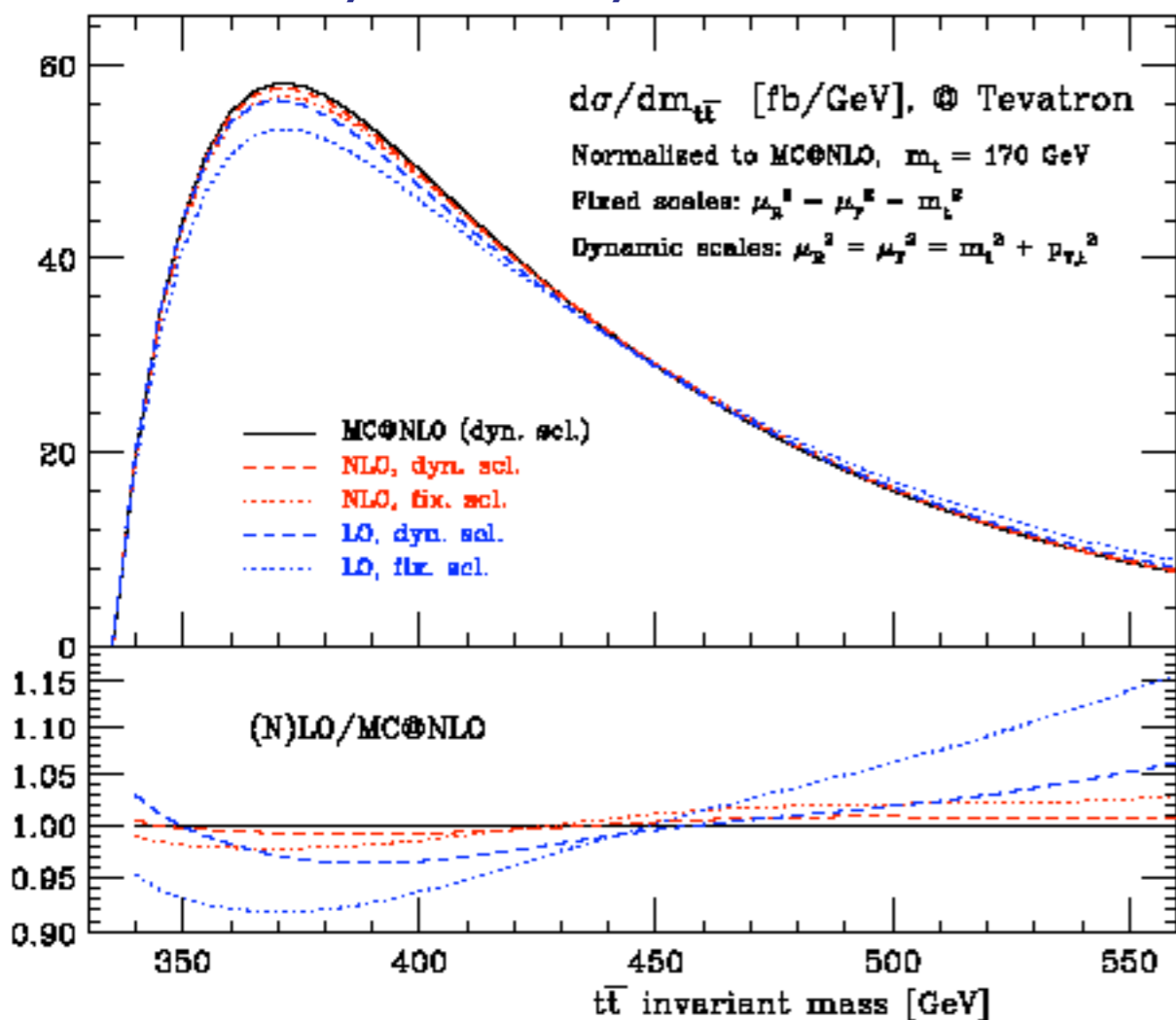
# ttbar: MC@NLO vs POWHEG



Good agreement for all the observables. Sometimes noticeable differences appear that can be ascribed to the different treatment of the higher terms. More exhaustive studies would be welcome.

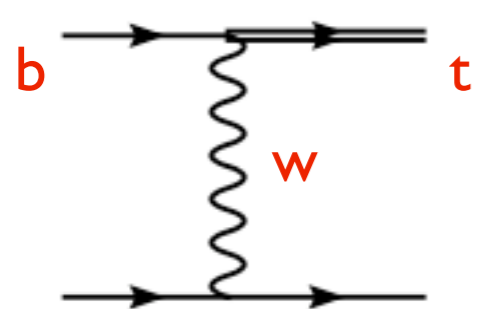
# $m_{t\bar{t}}$ spectrum: low mass

Tree-level production with a dynamical scale reproduces the shape MC@NLO result extremely well. Very stable observable.

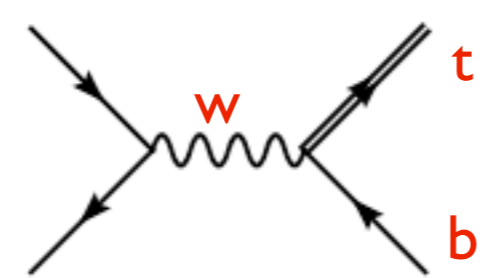


It is always important to validate the MC used in an analysis with the best tool available so far!

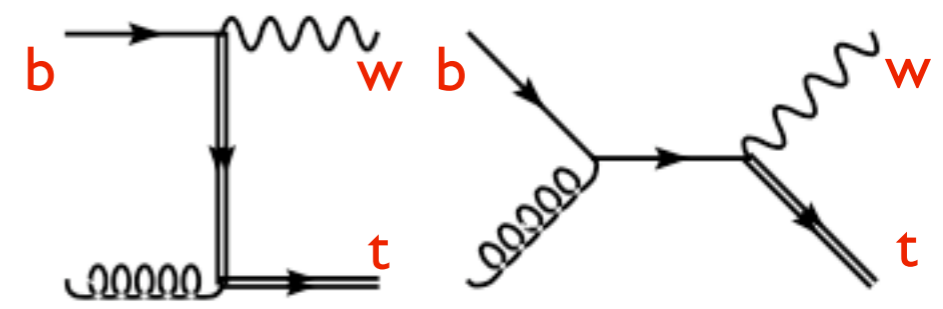
# MC challenges in single top



t-channel



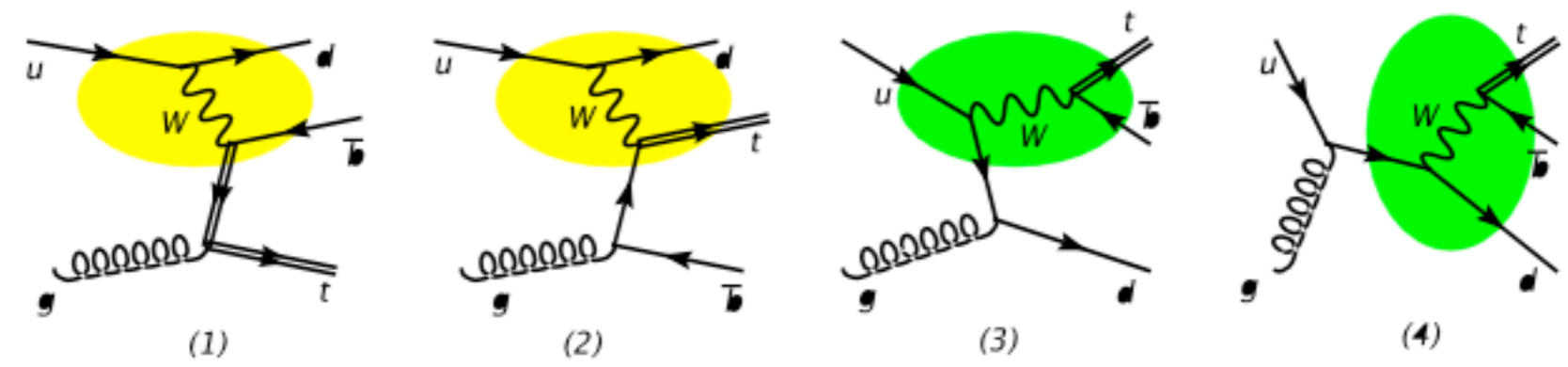
s-channel



Wt-associated

\* All three processes available in MC@NLO, for arbitrary CKM matrix.

\* At NLO t- and s- channel become intertwined :

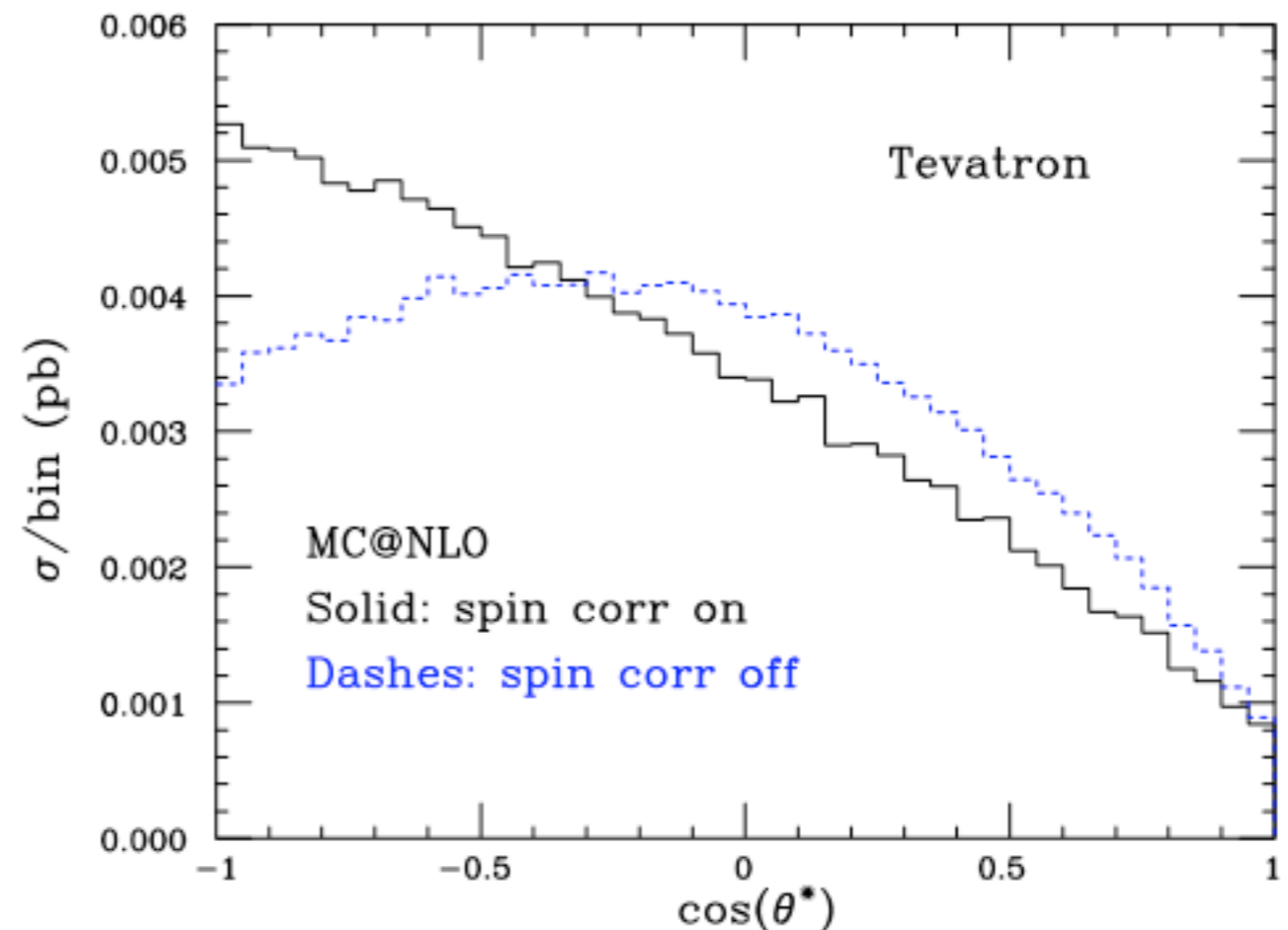
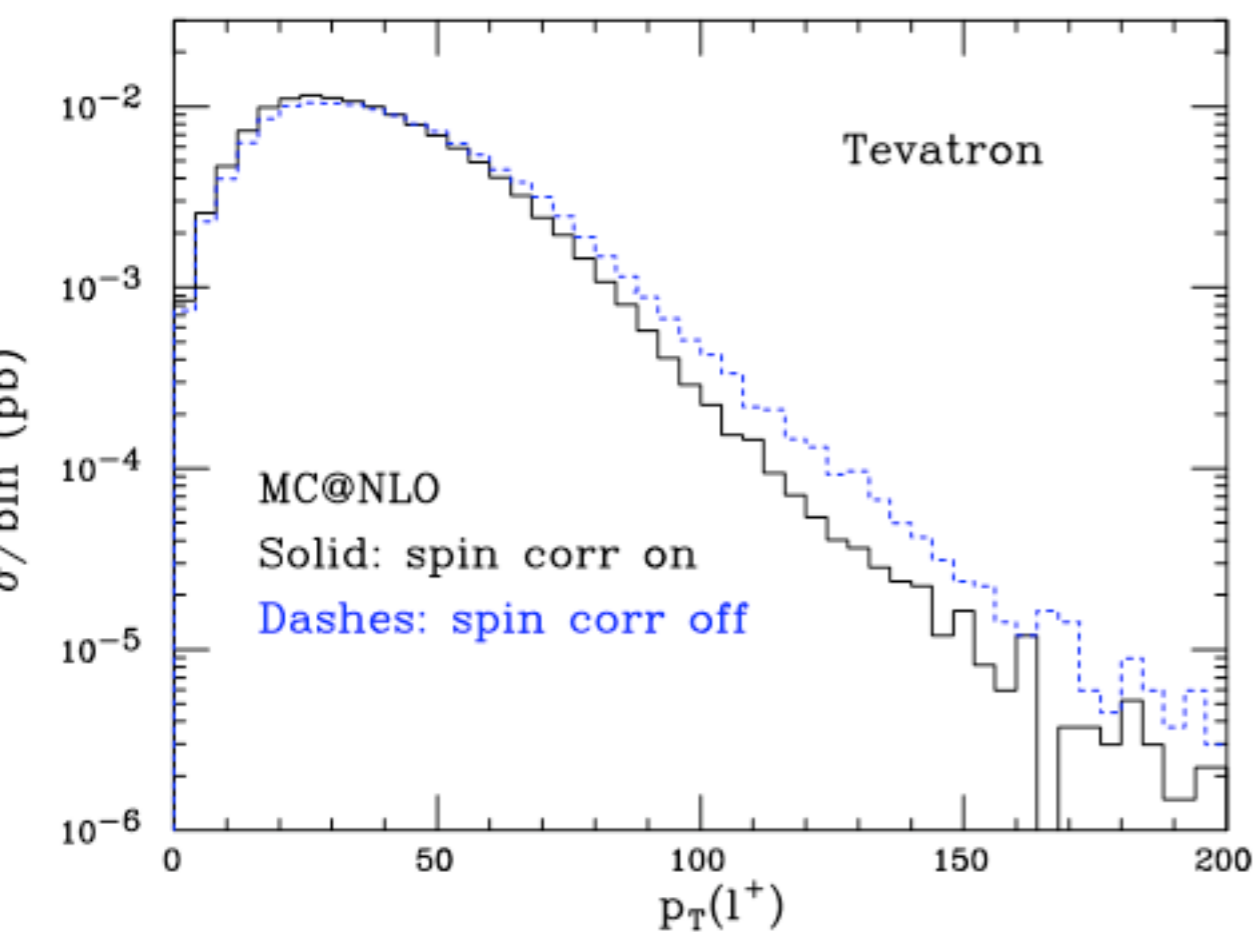


However, since the interference vanishes, their definition as independent procs poses no problems at NLO. Notice that (1) and (4) are not generated by a standard (=Pythia or Herwig) PS approach.

\* Wtb coupling might hide interesting BSM physics [ see Aguilar-Saavedra's talk yesterday ] => very important to keep spin correlations between production and decay

\* Accurate description of the softer b in the t-channel and tW associated production is extremely important.

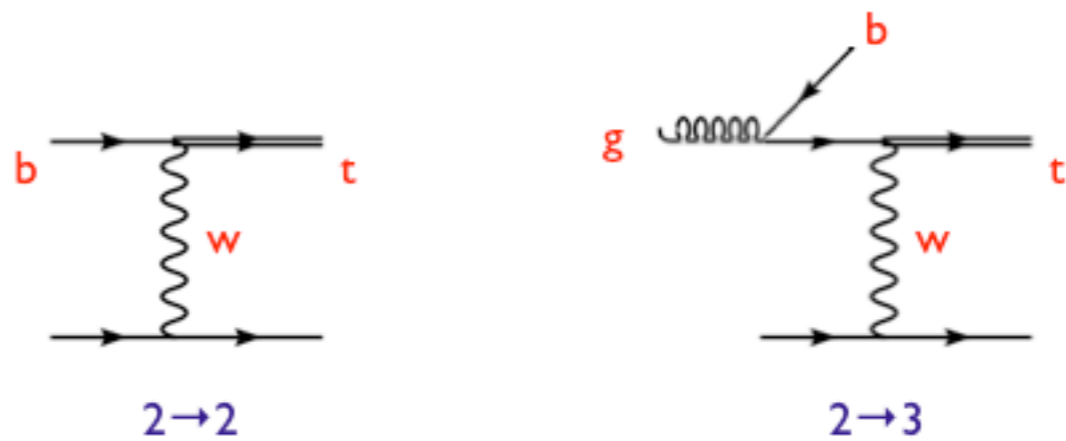
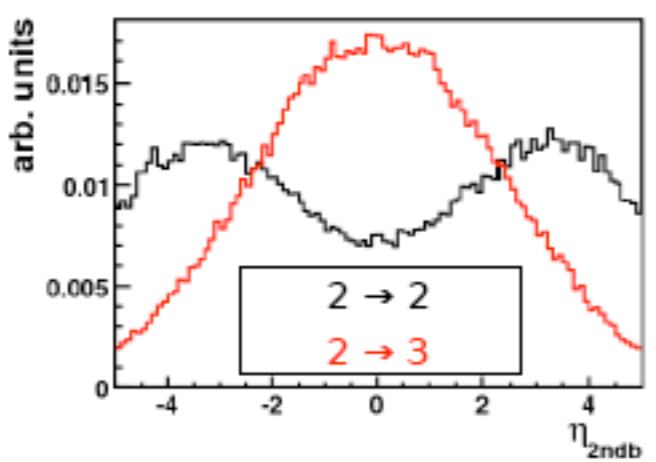
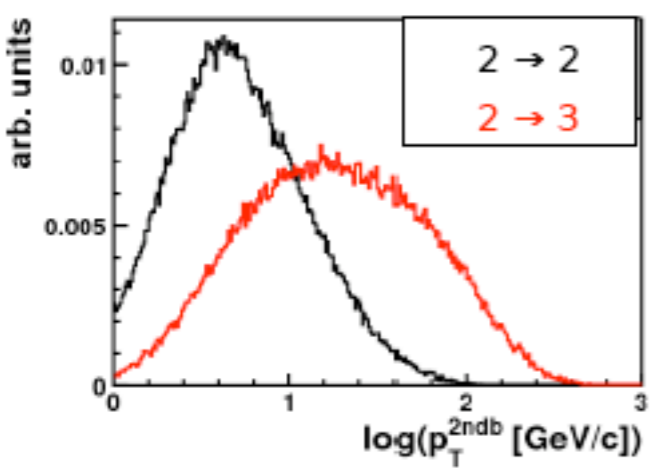
# Single top in MC@NLO



Spin correlation effects are non-negligible even in observables that are not specifically designed to measure them!

# t- channel : pt of the 2<sup>nd</sup> b

A correct prediction for the shape of the 2<sup>nd</sup> b is extremely important to reject ttbar events, by imposing a jet veto



Observation:

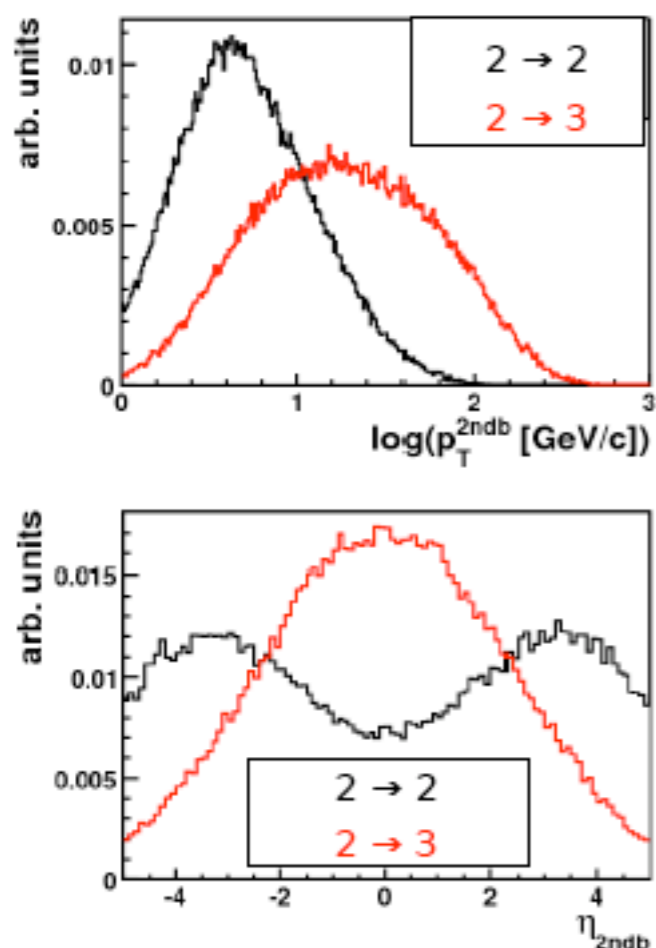
- \* Collinear gluon splitting better handled by the shower (explicit resummation)
- \* high pt tail better described by the ME 2→3

Pragmatic solution : [Boos et al., 2000]

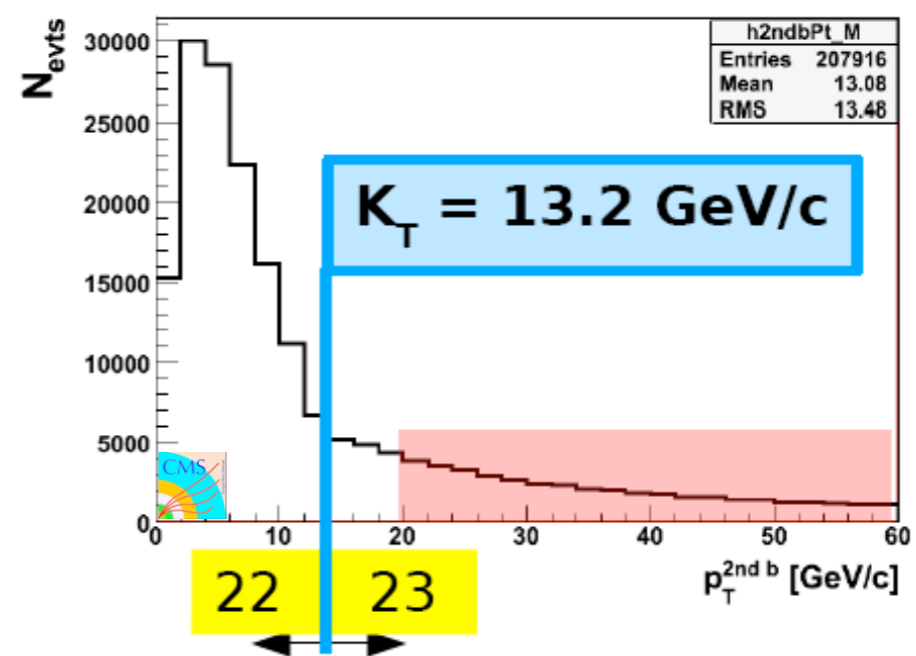
[Muller, Sturm, Wagner, Weinelt, CMS 2008]

# t- channel : $p_T$ of the 2<sup>nd</sup> b

A correct prediction for the shape of the 2<sup>nd</sup> b is extremely important to reject ttbar events, by imposing a jet veto



[Muller, Sturm, Wagner, Weinelt, CMS 2008]

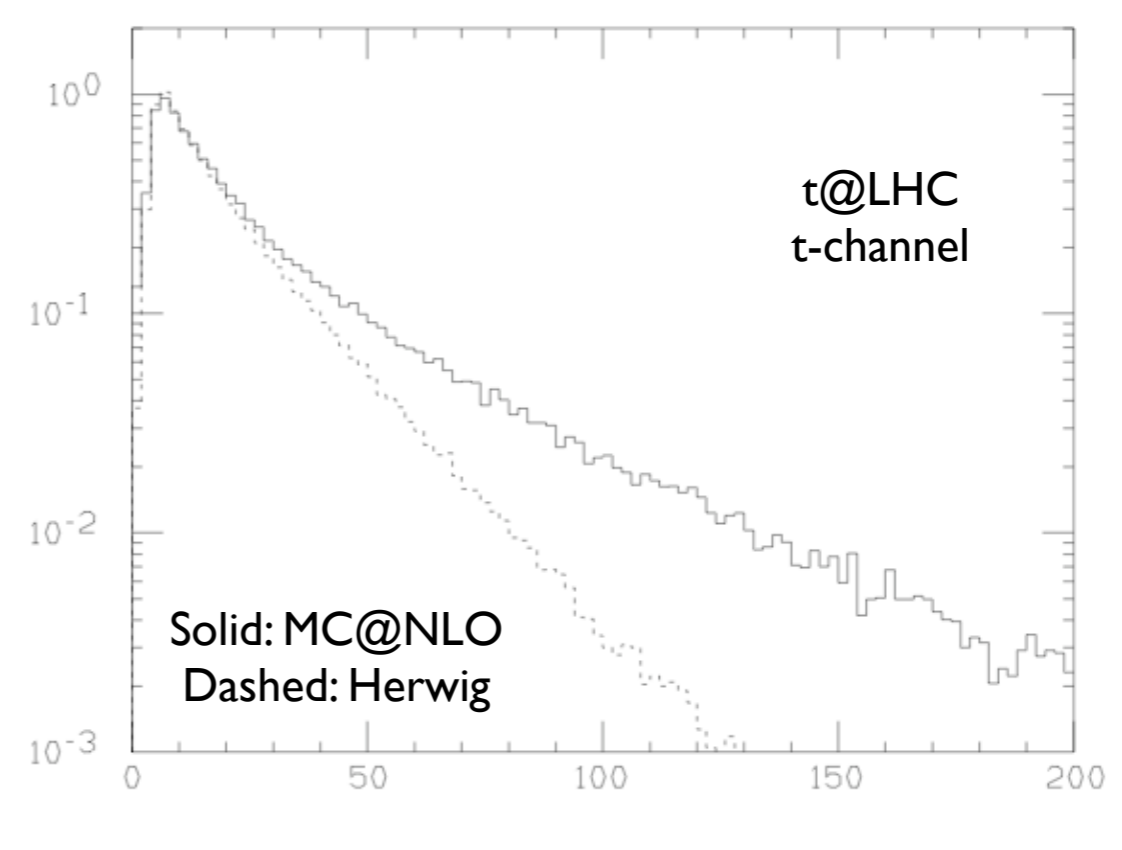
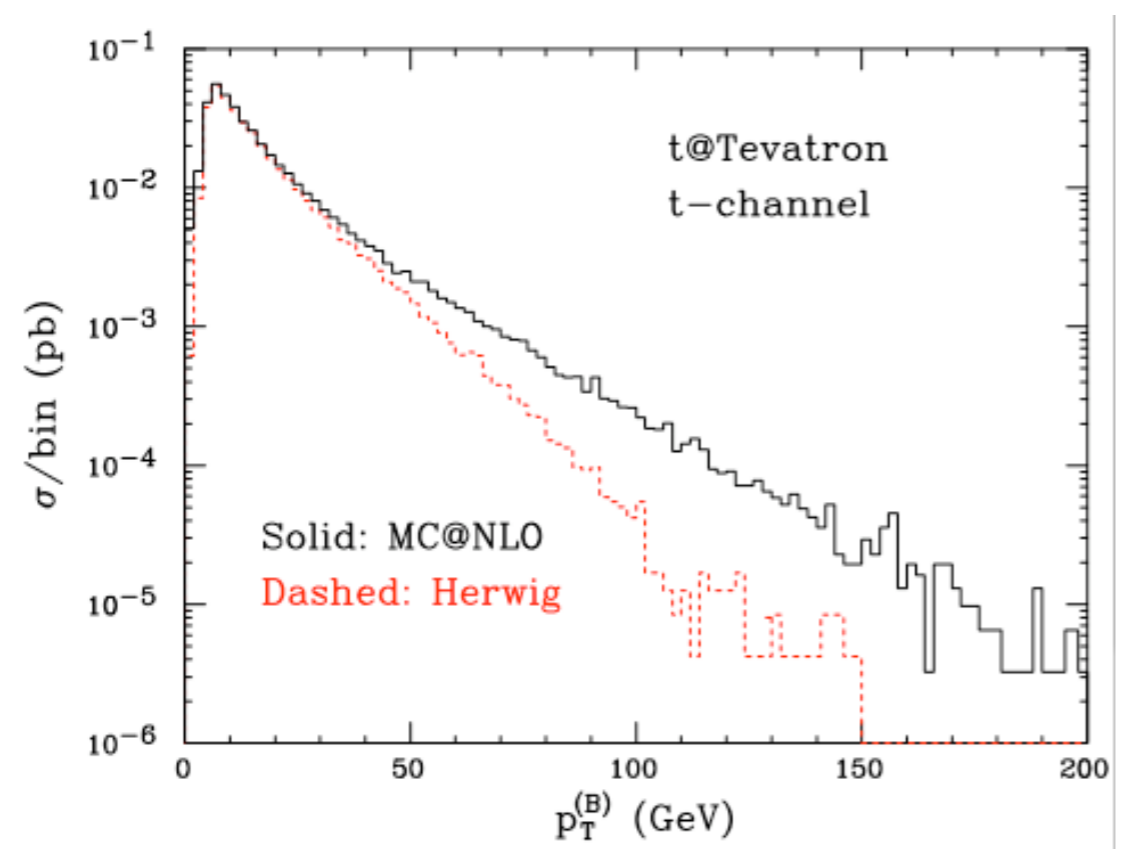


However, this kind of “hard cutoff” matching is clearly very crude and theoretically inconsistent (strong cutoff dependence, other distributions affected,..)



# t- channel : pt of the 2<sup>nd</sup> b

MC@NLO provides a consistent matching between  $2 \rightarrow 2$  and  $2 \rightarrow 3$ .  
This is done assuming the b initially massless.



Expected behaviour reproduced : small pt PS, high pt ME.

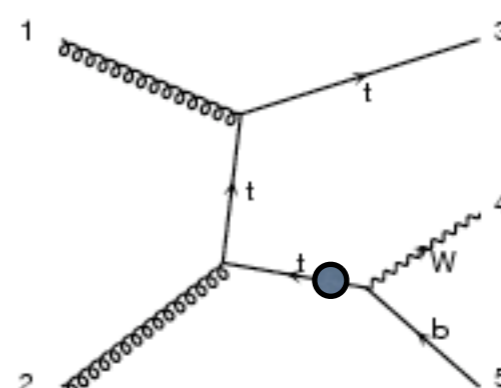
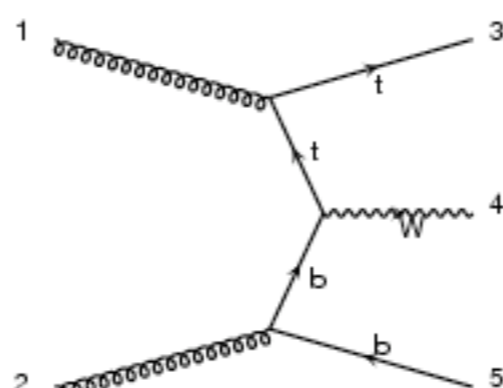
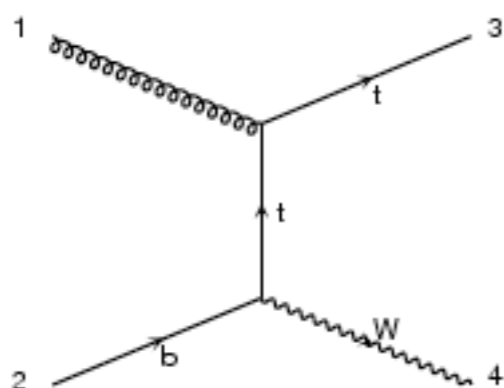
Alternative consistent approach (at LO) in AcerMC. [Hinchliffe and Keservan, 2006]

However, this procedure builds on the assumption that the massless PS describes well the IS collinear splitting into a massive quark pair at low pt....

On-going study on the  $2 \rightarrow 3$  NLO massive calculation [see F. Tramontano's talk]

# tW in MC@NLO

[Frixione, Laenen, Motylinski, Webber, White, YESTERDAY]



b-initiated process like t-channel  $\Rightarrow$  same issue with the  $p_T$  of the 2<sup>nd</sup> b.

Interference with tt at NLO  $\Rightarrow$  non trivial problem : definition of the process is at stake [Tim Tait: (2000), A. Belyaev & E. Boos (2001)].

First MC viable solution proposed [Campbell, FM, Willenbrock, LH 2005] and implemented in MCFM [Campbell, Tramontano, 2006].

However, interference is tamed with a (b-)jet veto  $\Rightarrow$  sensitivity to low  $p_T$  partons  $\Rightarrow$  soft resummation

$\Rightarrow$  MC with PS and with NLO needed.

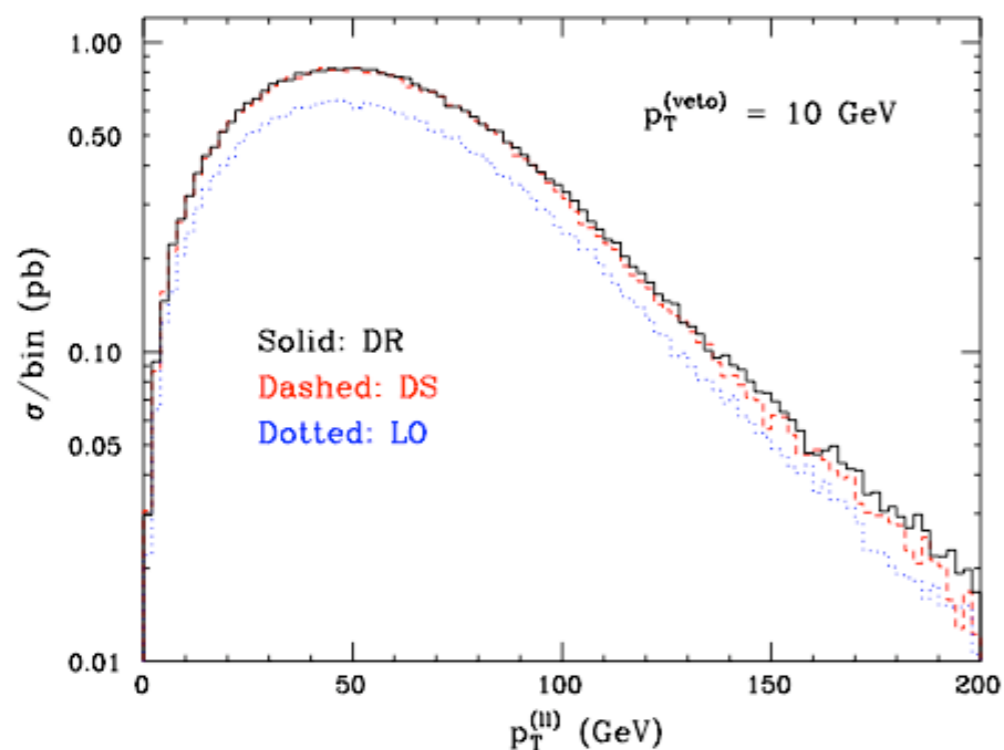


Diagram Removal :  $\hat{S}_{\alpha\beta}$

Diagram Subtraction :  $\left( \hat{S}_{\alpha\beta} + \mathcal{I}_{\alpha\beta} + \mathcal{D}_{\alpha\beta} - \tilde{\mathcal{D}}_{\alpha\beta} \right)$

Result: tW can be defined in

- \* a MC-friendly way
- \* (de facto) non-ambiguous way.

First pheno consistent LHC study in progress.

# To remember about the NLOwPS

MC@NLO [Frixione, Nason, Webber, 2003] is the standard code with a library of SM processes.

POWHEG [Frixione, Nason, Oleari, 2007] is a recent addition with a much smaller library at the moment.

“Best” tools when NLO calculation is available (i.e. low jet multiplicity).

## Main points:

1. NLOwPS provide a consistent way to include K-factors into MC's
2. Scale dependence is meaningful
3. Allows a correct estimate of the PDF errors.
4. Non-trivial dynamics beyond LO included for the first time.

N.B.: The above is true for observables which are at NLO to start with!!!

**Current limitations:** Considerable manual work for the implementation of a new process and only SM.

**Outlook:** Automatization for the real contributions proven feasible. Automatization for the virtuals in sight. General matching procedure available and shower independent. Full automatization in sight.

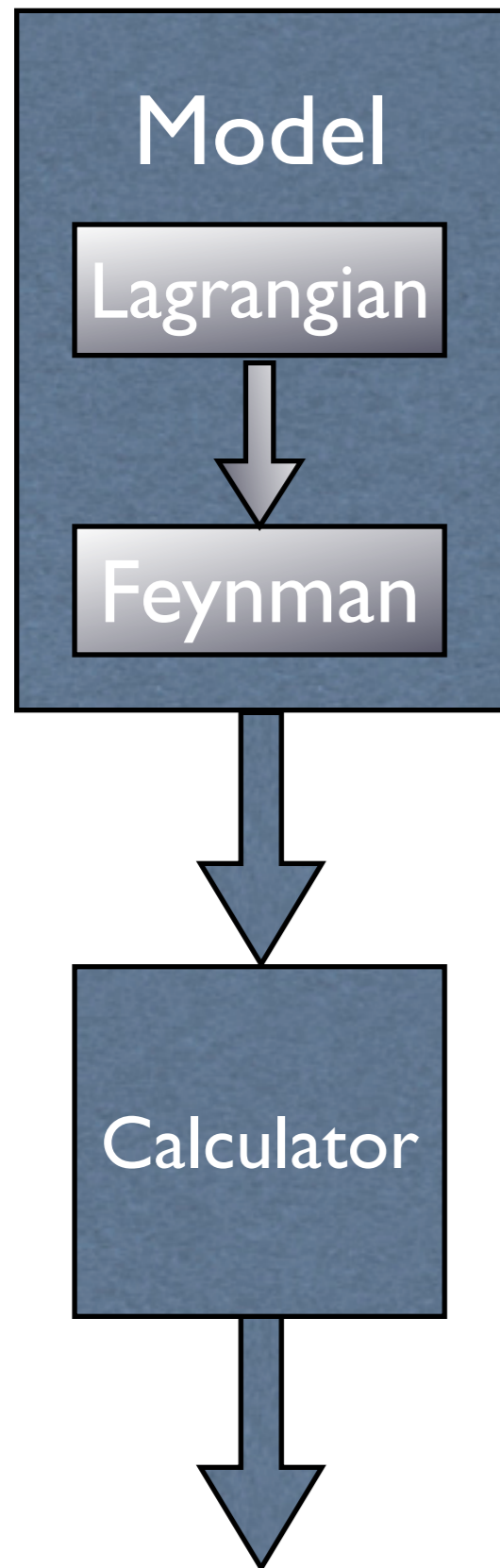
$\mathcal{L} + ME + PS$

# Preliminary observation

- For new physics associated to top, two approaches are possible:
  - ▶ top-down (e.g., model parameter scanning)
  - ▶ bottom-up (e.g., inverse problem, OSET)
- Different strategies lead to different MC tools.
- Some of the MC tools of the new generation (e.g. MadGraph and SHERPA) allow to tackle both by featuring:
  1. Main templates for BSM models available (MSSM, 2HDM, UED,...)
  2. Easy implementation of new models (from Feynman rules)
  3. Any tree-level process available\* (automatic code creation).
  4. Multi-jets samples with matching for QCD background (and possibly BSM signals).

\*within reason.

# Add-on for BSM



Invent a model, renormalizable or not, with new physics. Write the Lagrangian and get the Feynman Rules.

The particles content, the type of interactions and the analytic form of the couplings in the Feynman rules define the model at tree level.

SUSY, Little Higgs, Higgsless, GUT, Extra dimensions (flat, warped, universal,...)

Parameters Calculator.

Given the “primary” couplings, all relevant quantities are calculated: masses, widths and the values of the couplings in the Feynman rules.

Caution: tree-level relations have to be satisfied to avoid gauge violations and/or wrong branching ratios.

FeynHiggs, ISAJET, NMHDecay, SOFTSUSY, SPHENO, SUSPECT, SDECAY...

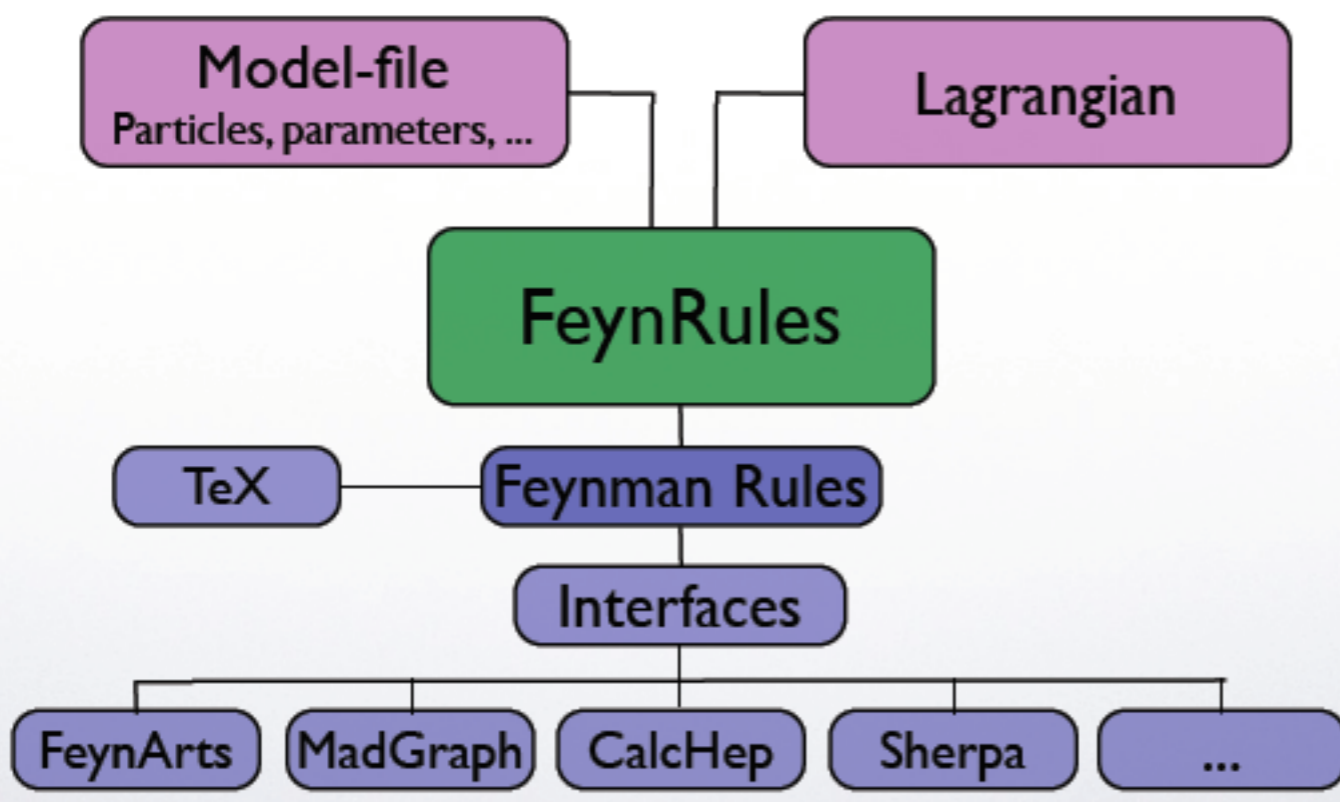
Les Houches interface

# From Lagrangians to Events

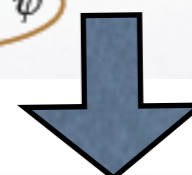
Tools exist to derive Feynman rules and couplings directly from the Lagrangian and effortlessly implement in any MC.

LanHEP [Semenov] : Interfaces to FeynArts (for theorists) and CompHEP/CalcHEP.

FeynRules [Duhr and Christensen] : New tool providing a complete model building development framework and interfaces to several MC's.

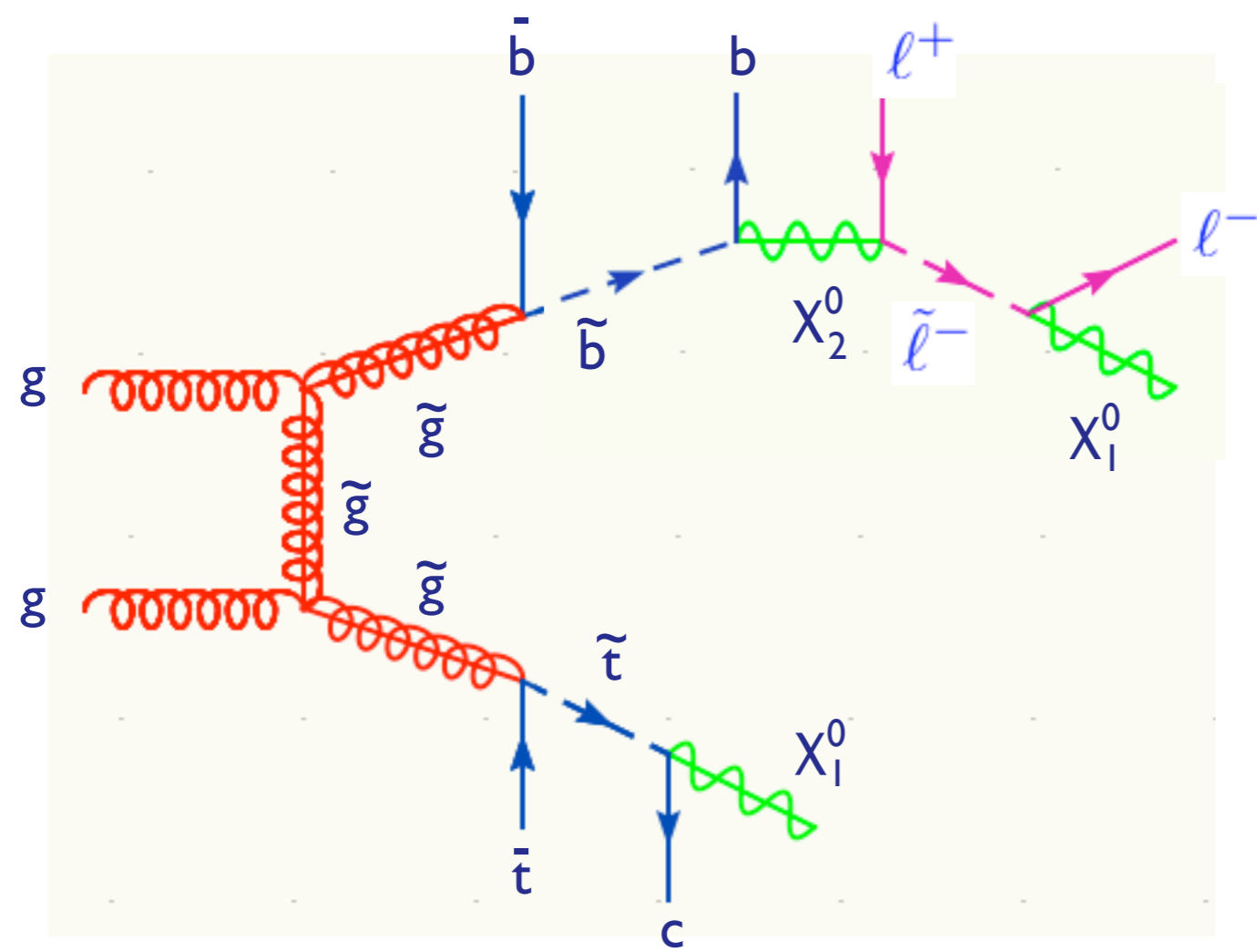


$$\kappa^{-1} \mathcal{L}_F^{\tilde{n}}(\kappa) = \frac{1}{2} \left[ \tilde{h}^{\tilde{n}} \eta^{\mu\nu} - \tilde{h}^{\mu\nu, \tilde{n}} \right] \bar{\psi} i \gamma_{\mu} D_{\nu} \psi - m_{\psi} \tilde{h}^{\tilde{n}} \bar{\psi} \psi + \frac{1}{2} \bar{\psi} i \gamma^{\mu} (\partial_{\mu} \tilde{h}^{\tilde{n}} - \partial^{\nu} \tilde{h}_{\mu\nu}^{\tilde{n}}) \psi + \frac{3\omega}{2} \tilde{\phi}^{\tilde{n}} \bar{\psi} i \gamma^{\mu} D_{\mu} \psi - 2\omega m_{\psi} \tilde{\phi}^{\tilde{n}} \bar{\psi} \psi + \frac{3\omega}{4} \partial_{\mu} \tilde{\phi}^{\tilde{n}} \bar{\psi} i \gamma^{\mu} \psi$$



$$k \left( -2 m_{\psi} \text{om } \phi \psi^{\dagger} \cdot \psi + \frac{3}{4} i \text{om } \partial_{\mu}(\phi) \psi^{\dagger} \cdot \gamma^{\mu} \cdot \psi + \frac{3}{2} \text{om } \phi (i \psi^{\dagger} \cdot \gamma^{\mu} \cdot \partial_{\mu}(\psi) - g \psi^{\dagger} \cdot \gamma^{\mu} \cdot T^a \cdot \psi G_{\mu,a}) + \frac{1}{2} \left( \frac{1}{2} i (\partial_{\mu}(h_{\nu,\nu}) - \partial_{\nu}(h_{\mu,\nu})) \psi^{\dagger} \cdot \gamma^{\mu} \cdot \psi - m_{\psi} \psi^{\dagger} \cdot \psi h_{\mu,\mu} + (i \psi^{\dagger} \cdot \gamma^{\mu} \cdot \partial_{\nu}(\psi) - g \psi^{\dagger} \cdot \gamma^{\mu} \cdot T^a \cdot \psi G_{\nu,a}) (h_{\rho,\rho} \eta_{\mu,\nu} - h_{\mu,\nu}) \right) \right)$$

# Decay chains



Beware that most of the MC's\* make some of or all the following simplifications:

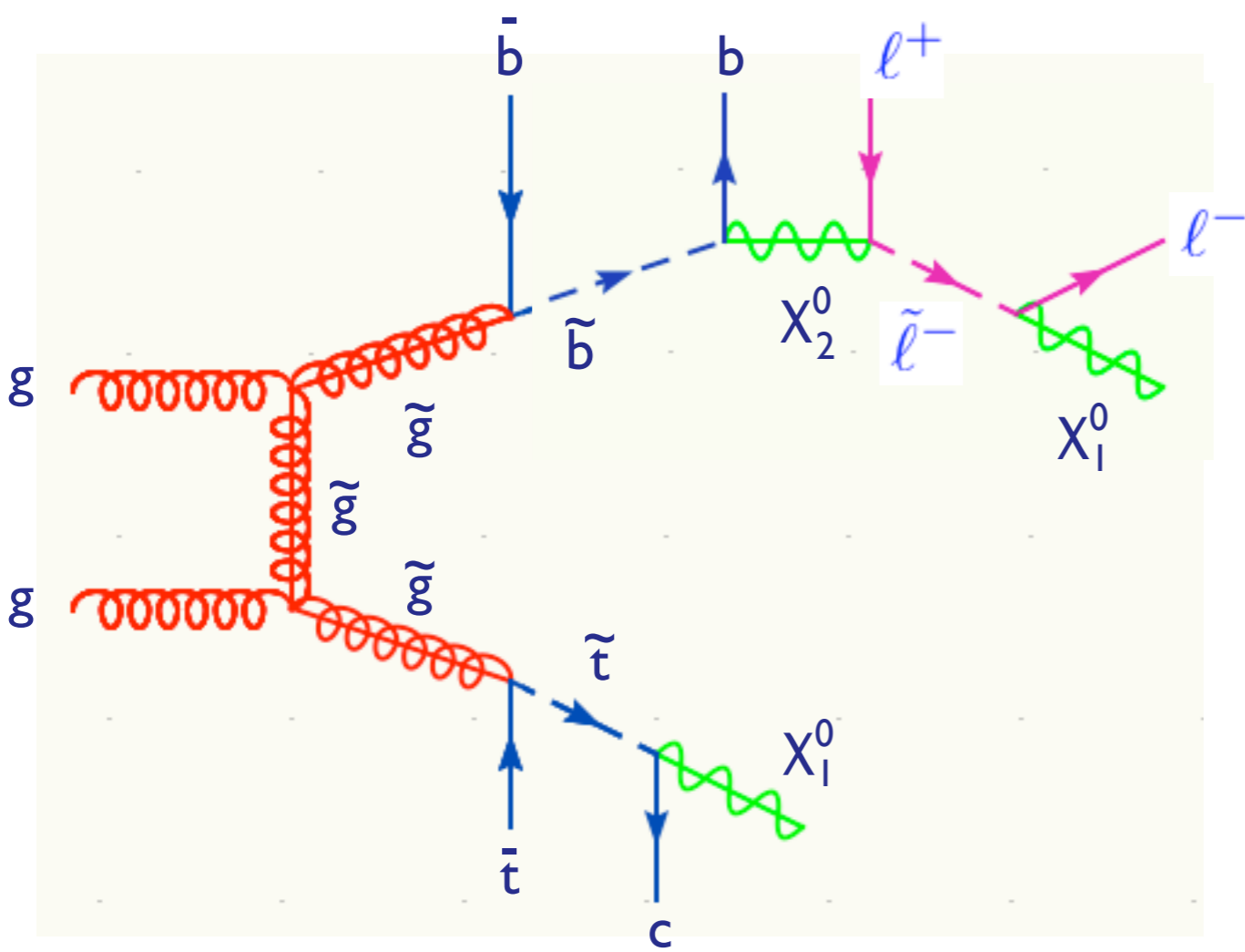
1. production and decay are factorized.
2. Spin is ignored.
3. Chains proceed only through  $1 \rightarrow 2$  decays.
4. The narrow width approximation is employed.
5. Non-resonant diagrams are ignored.

\* Spin correlations between production and decay are correctly accounted for by HERWIG in the narrow width approximation [Richardson 2005; Gigg & Richardson, 2008]



# Decay chains

$$gg \rightarrow (g_0 \rightarrow t \sim (t_1 \rightarrow c \ n_1)) (g_0 \rightarrow b \sim (b_1 \rightarrow (b \ (n_2 \rightarrow \mu^+ \ (\mu_{l-} \rightarrow \mu^- \ n_1))))))$$



Most matrix element based MC's\*:

1. Full matrix element is obtained which includes correlations between production and decays.
2. Spin of the intermediate states is kept.
3. One can go beyond  $1 \rightarrow 2$  decays.
4. Resonances have BW.
5. Non-resonant contributions can be systematically included only where relevant.
6. BW information passed to the event record.

\* As for instance in SHERPA [Krauss et al., ] and in MadGraph [Alwall and Stelzer, 2007]

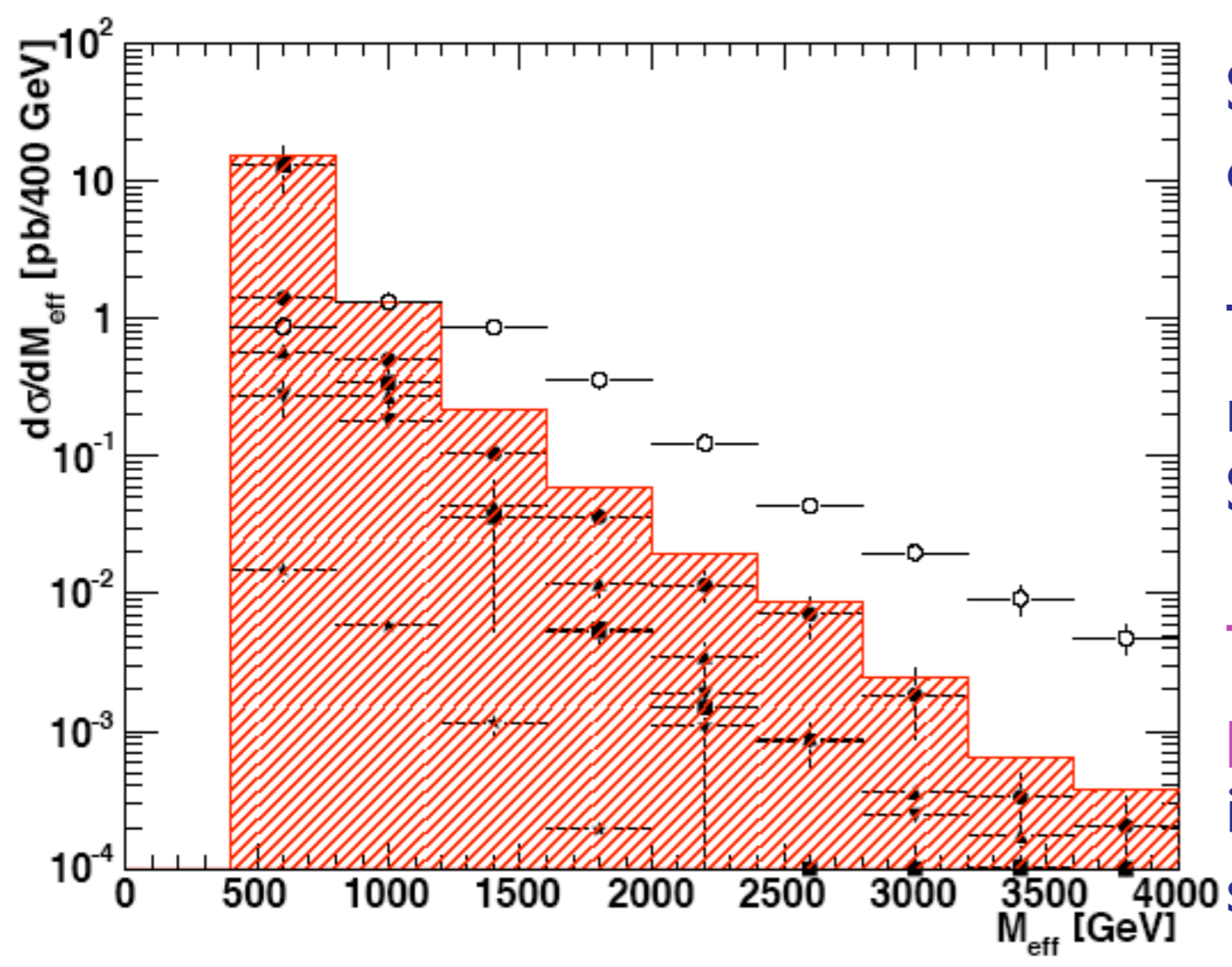
# Example I: SUSY

[Kraml, Raklev, 2006]

$$\tilde{g}\tilde{g} \rightarrow t\bar{t}\tilde{t}_1\tilde{t}_1^*, t\bar{t}\tilde{t}_1^*\tilde{t}_1^*, \bar{t}\tilde{t}_1\tilde{t}_1$$

$$m_{\tilde{t}_1} < m_t$$

$$pp \rightarrow \tilde{g}\tilde{g} \rightarrow bbl^\pm l^\pm + \text{jets} + E_{\text{miss}}^T \quad \tilde{t} \rightarrow c\tilde{\chi}_1^0$$



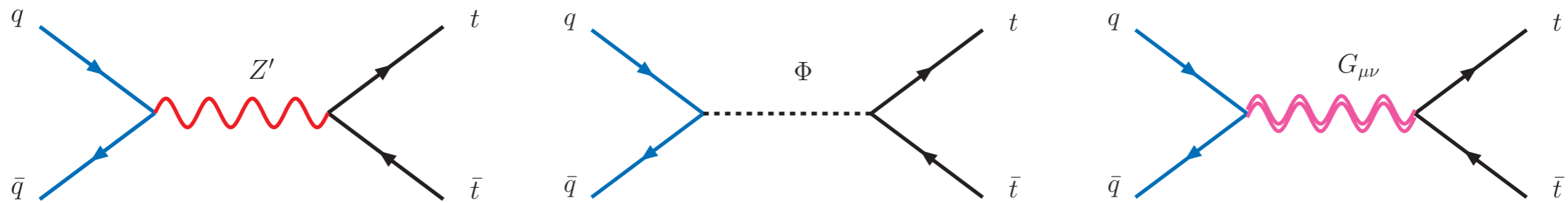
Same-sign top quarks as a signature of light stops.

Typical SUSY inclusive signature: need for a very good control of the SM backgrounds (here Pythia).

The whole analysis can be now performed within one MC including matched samples for the signal and backgrounds.

# Example II: New resonances

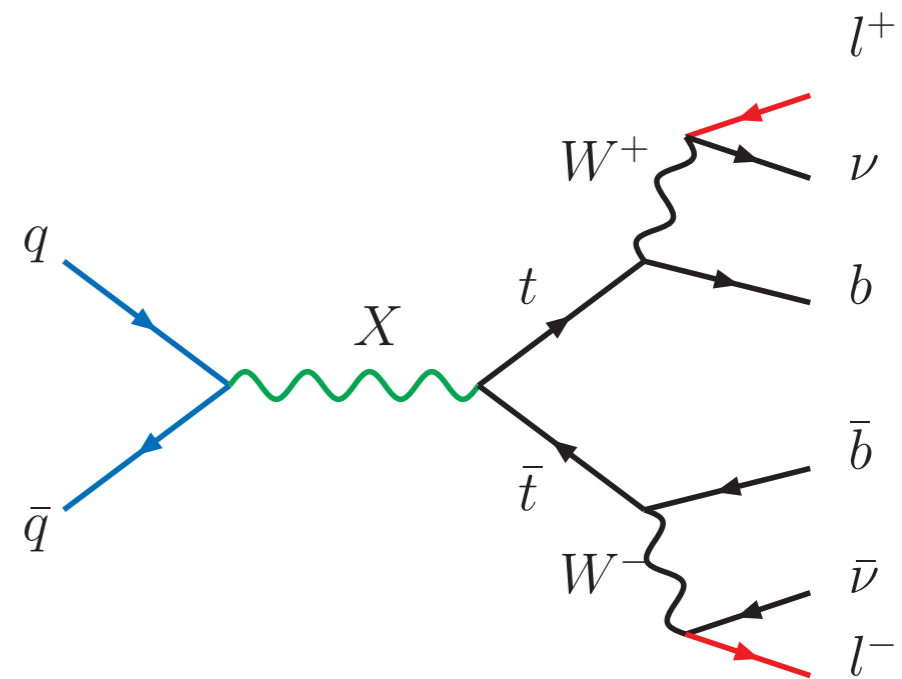
In many scenarios for EWSB new resonances show up, some of which preferably couple to 3rd generation quarks.



Given the large number of models, in this case is more efficient to adopt a “model independent” search and try to get as much information as possible on the quantum numbers and coupling of the resonance.

To access the spin of the intermediate resonance spin correlations should be measured.

It therefore mandatory for such cases to have MC samples where spin correlations are kept and the full matrix element  $\langle pp | X | tt \rangle$  is used.

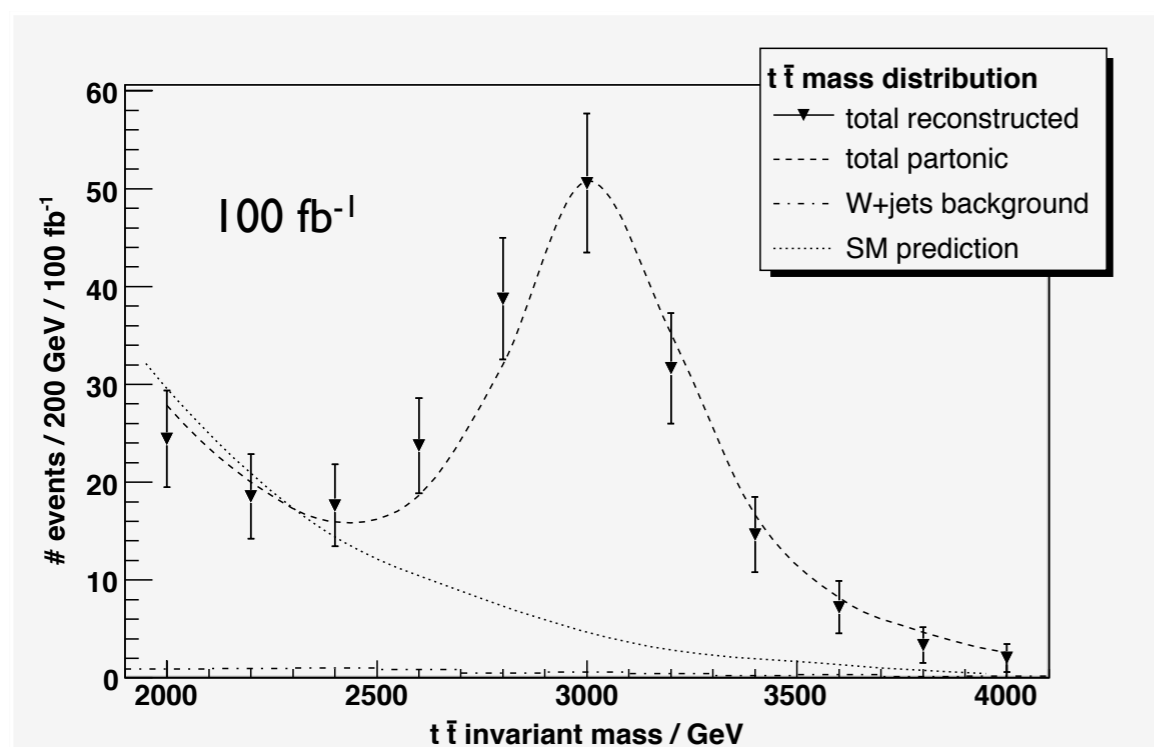


# Example II : KK gluons with SHERPA

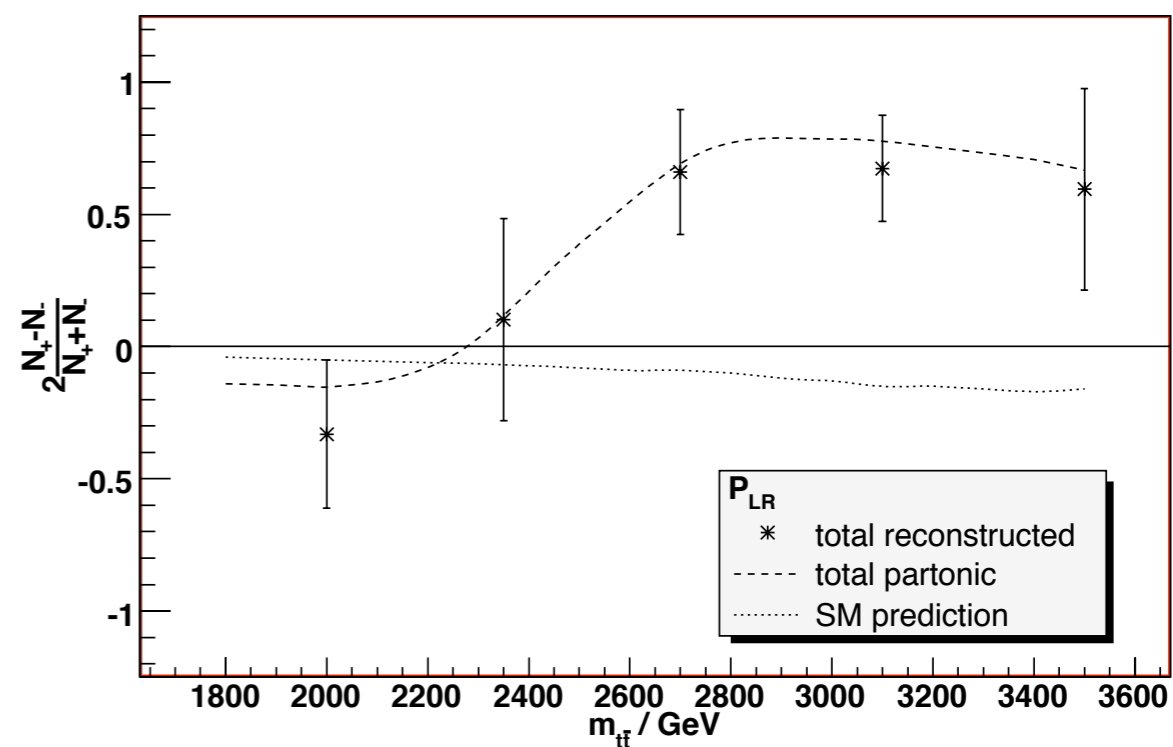
Model from [Agashe et al., hep-ph/0612015]:

RandakII-Sundrum model, with different profiles for SM fermions.

Non-universal couplings to SM particles : suppressed to light quarks,  $\sim 1$  to  $t_L$ , enhanced to  $t_R$ .



MC trivial



MC non-trivial

# Conclusions

- New MC tools are available that can provide an accurate description of both SM and BSM signals and backgrounds involving top:
  - ★ Impressive progress in fixed-order and parton-shower matching both at LO (inclusive  $t\bar{t}$ +jets samples) and NLO ( $t\bar{t}$  in MC@NLO and POWHEG). First systematics comparison available.
  - ★ Progress in the simulation of basically any new physics scenario's involving top (MSSM, new resonances, vector-like partners, anomalous couplings,...)
- New and exciting possibilities of interaction between TH's and EXP's...

Take a look around  
At all people everywhere  
So much energy and excitement in the air  
And the time is right  
To get together with the people you know  
So sing out loud and clear don't be afraid  
**To let the LHC start**

Chorus:

Are you ready for it?  
Rockin' steady for it  
Are you ready for it?  
Rockin' steady for it  
**Are you ready for it?**

**Madonna**

# Credits

Many thanks for providing material, physics insights or  
both to this talk:

Simon de Visscher, Stefano Frixione, Stefan Hoeche,  
Frank Krauss, Carlo Oleari, Steffen Schuman