

# Effects of Jet Matching in BSM Signals

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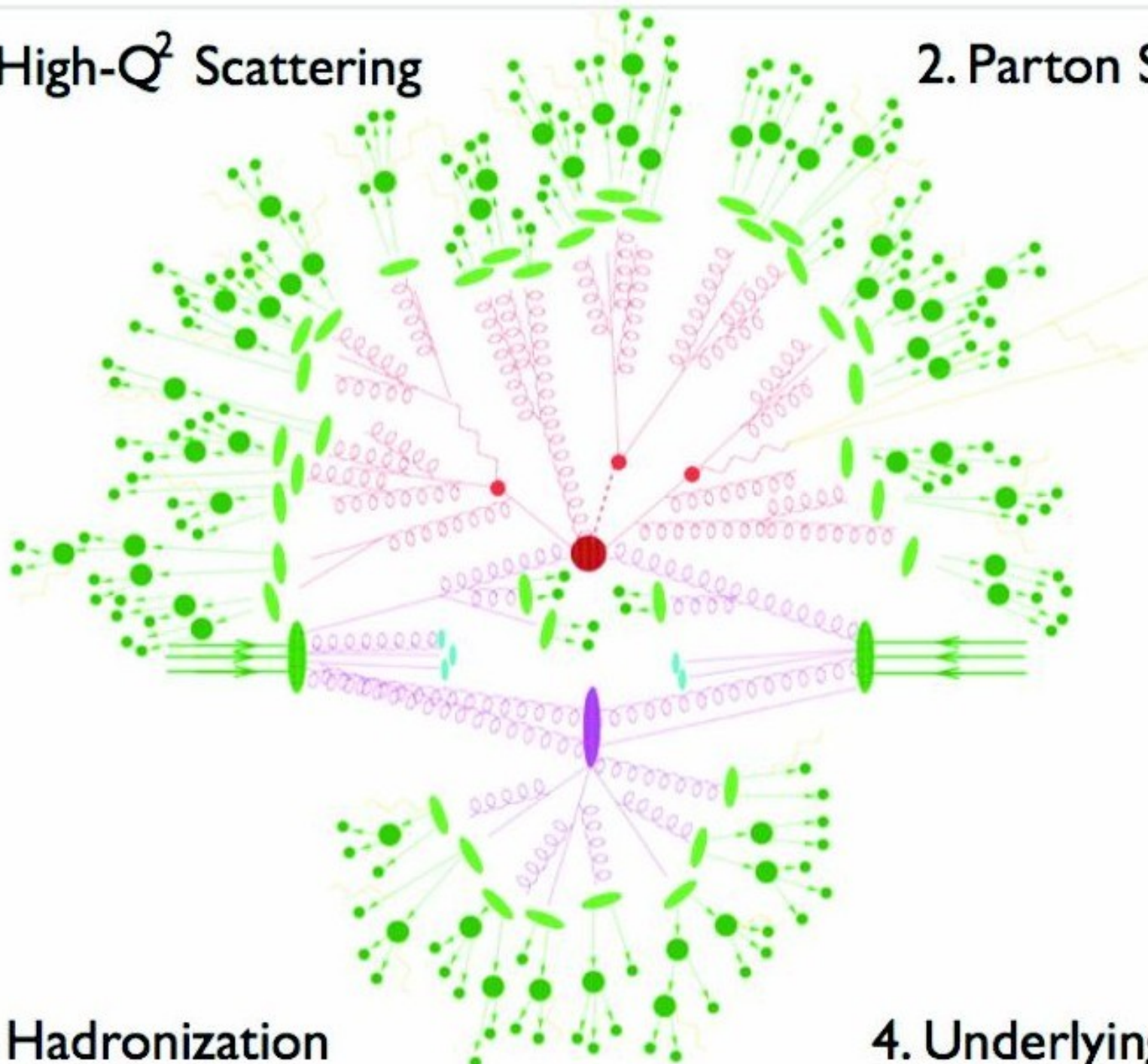
arXiv:0810.xxxx (very soon!)

IPMU Seminar, Tokyo, Japan, 14 Oct 2008

# Introduction – The Big Picture

1. High- $Q^2$  Scattering

2. Parton Shower



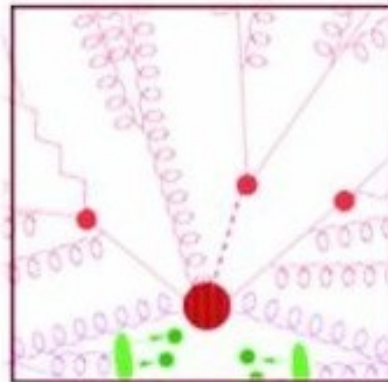
3. Hadronization

4. Underlying Event

# I. High- $Q^2$ Scattering

# 2. Parton Shower

☞ where new physics lies



☞ process dependent

☞ first principles description

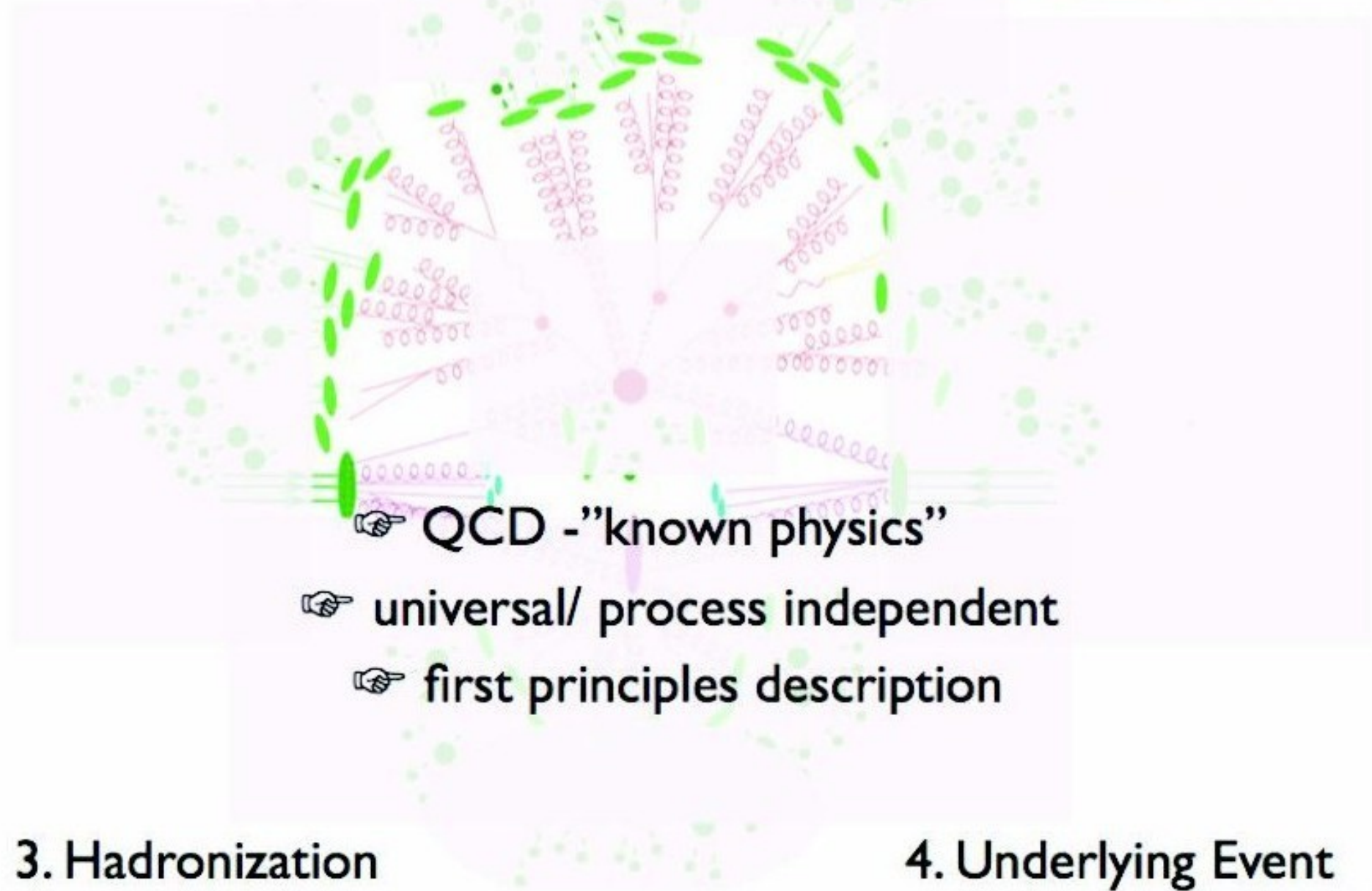
☞ it can be systematically improved

# 3. Hadronization

# 4. Underlying Event

# I. High- $Q^2$ Scattering

# 2. Parton Shower



☞ QCD - "known physics"

☞ universal/ process independent

☞ first principles description

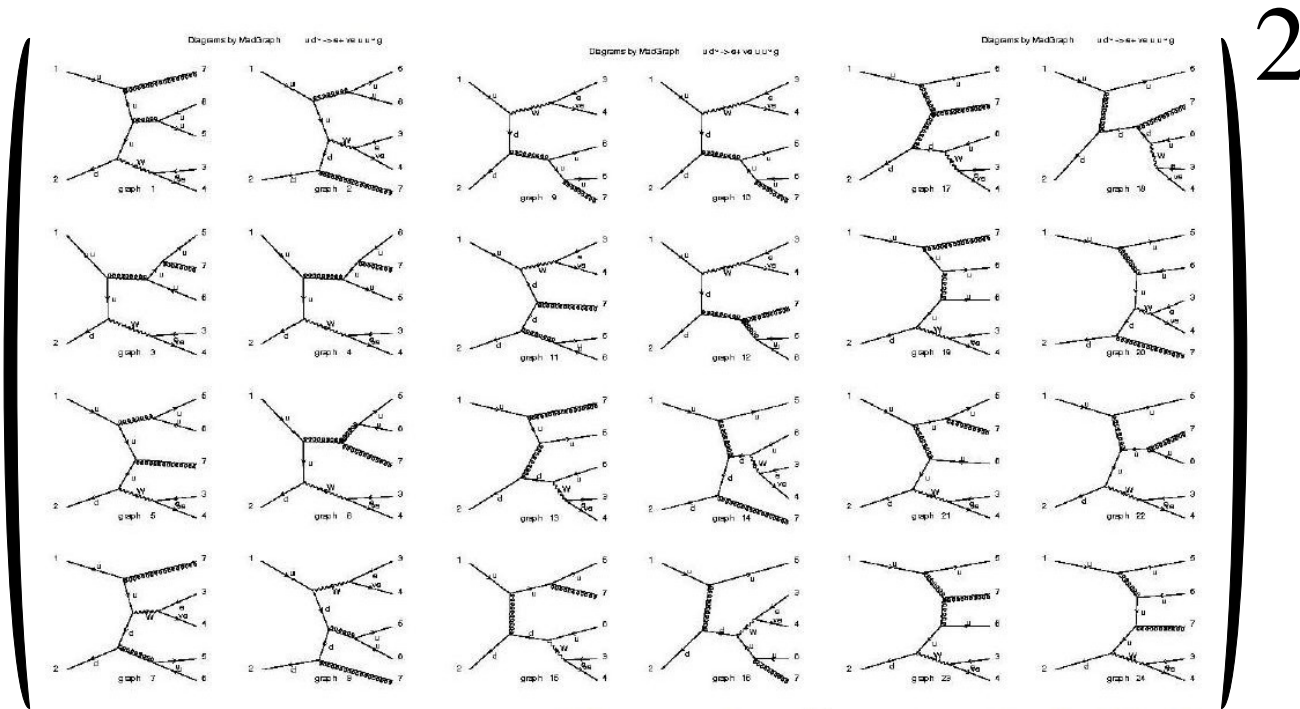
# 3. Hadronization

# 4. Underlying Event

# How to describe QCD emissions

- Matrix Element Generators
  - Parton Showers

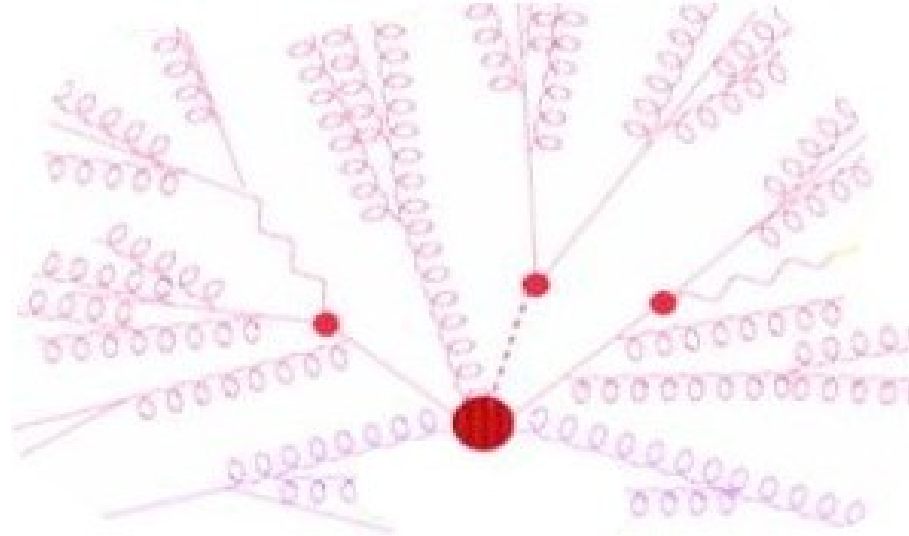
# Matrix Element Generators



Diagrams for  $u\bar{d} \rightarrow e^+\nu_e u\bar{u}g$  by MadGraph

- Use full Matrix Element information
- Different methods in different generators
  - MadGraph, Grace, Sherpa: Feynman diags + Helicity ampls
  - AlpGen, Helac: Recursion relations
  - Whizard: 1POW/DAG optimization

# Parton Showers



- Based on soft-collinear approximation
- Markov process to perform subsequent QCD emissions (due to factorization)
- Strict ordering of emissions in ordering variable
  - $Q^2$  (Pythia  $< 6.3$ ),  $p_T$  (Pythia  $> 6.3$ , Ariadne)
  - $\theta E$  (Herwig)



# Jet matching/merging

## Matrix elements

- 1 Fixed order calculation
- 2 Computationally expensive
- 3 Limited number of particles
- 4 Valid when partons are hard and well separated
- 5 Quantum interference correct
- 6 Needed for multi-jet description

## Parton showers

- 1 Resums logs to all orders
- 2 Computationally cheap
- 3 No limit on particle multiplicity
- 4 Valid when partons are collinear and/or soft
- 5 Partial quantum interference through angular ordering
- 6 Needed for hadronization/detector simulation

Complementary approximations

Need to combine without double counting

# Matching schemes

## The simple idea behind matching

- Use **matrix element description** for well separated jets, and **parton showers** for collinear jets
  - Phase-space cutoff to separate regions
- ⇒ No double-counting between jet multiplicities

## Difficulties

- Get smooth transition between regions
- No/small dependence from precise cutoff
- No/small dependence from largest multiplicity sample

## How to accomplish this

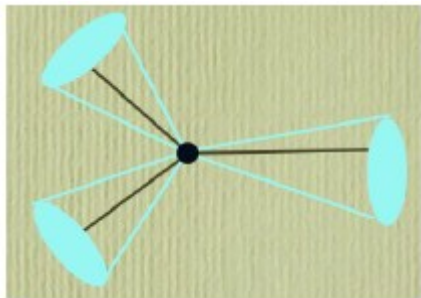
- CKKW scheme (Catani, Krauss, Kuhn, Webber)
- Lönnblad scheme
- MLM scheme

# MLM matching

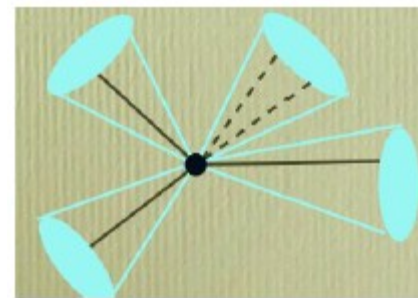
J.A. et al. [arXiv:0706.2569],  
cf. M.L. Mangano [2002, Alpgen home page]

Use shower hardness to separate ME/PS

- 1 Generate multiparton event with cut on jet  $k_T$
- 2 Cluster event and use  $k_T^2$  for  $\alpha_s$  scale
- 3 Shower event (using Pythia) starting from hard scale
- 4 Collect showered partons in  $k_T$  jets with  $k_{T\text{cut}} > k_{T\text{min}}$
- 5 Keep event only if each jet matched to one parton
- 6 For highest multiplicity sample, allow extra jets softer than  $k_{T\text{min}}$

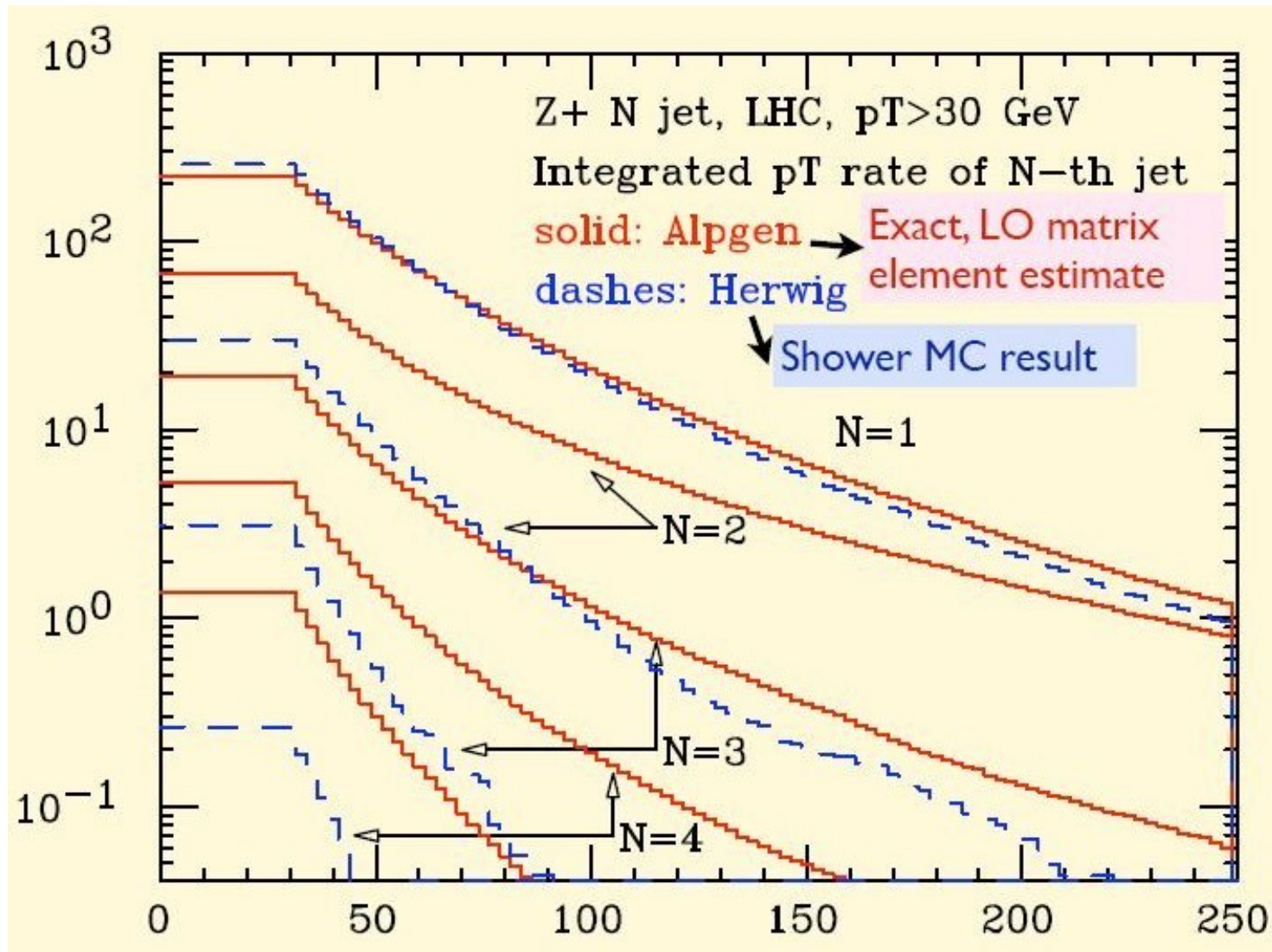


Keep



Discard unless highest multiplicity

# Impact of matching on Z+jets



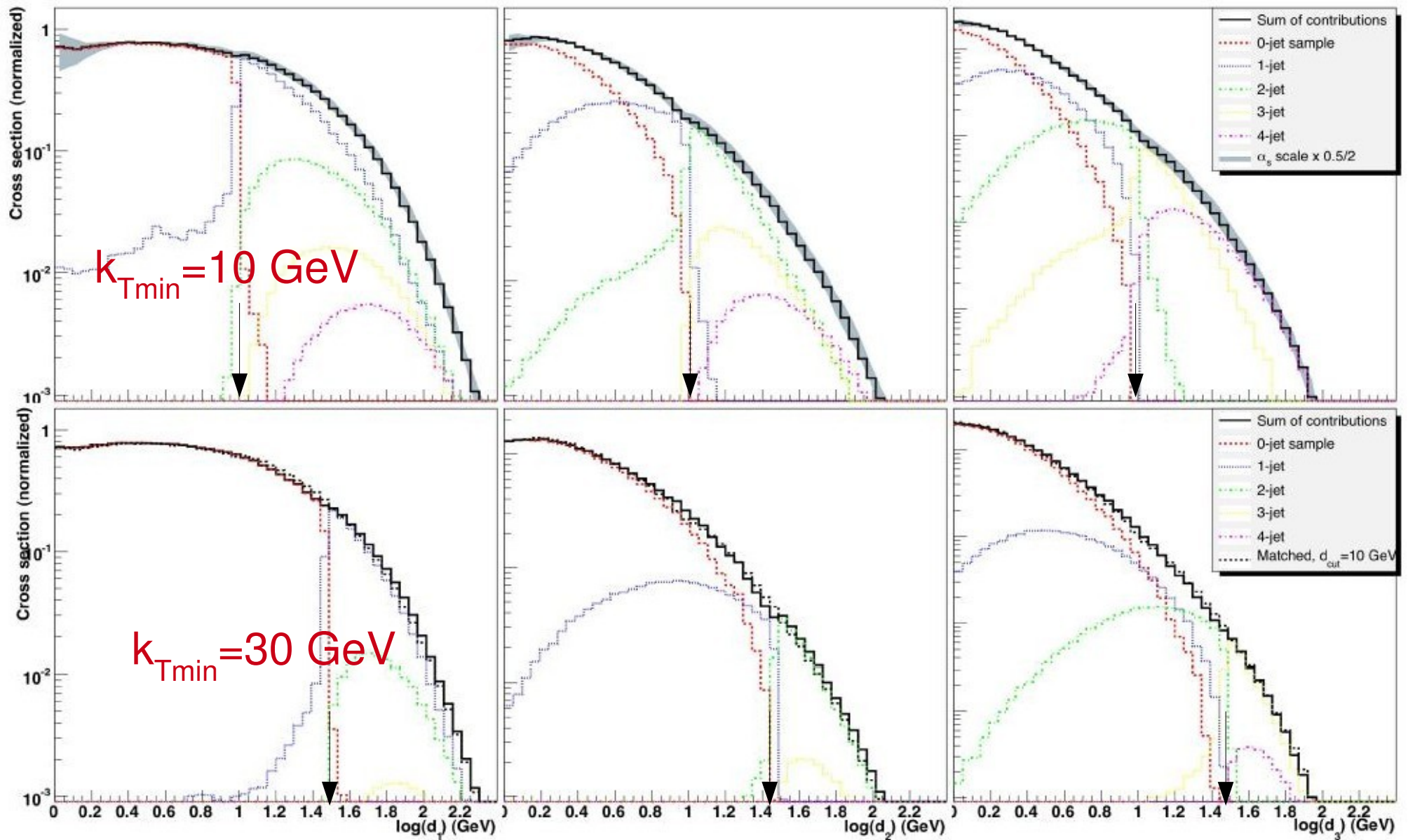
Matching vital to get multijet backgrounds right!

# Matching in MG/ME+Pythia

[MG/ME = MadGraph/MadEvent, an automatized Matrix Element generator and event generator, see arXiv:0706.2334]

- $k_T$  and cone jet MLM schemes
- New “shower  $k_T$ ” scheme
- Both  $Q^2$  and  $p_T$ -ordered Pythia showers
- Extensively validated in  $V$ +jets [arXiv:0706.2569],  $VV$ +jets,  $t$  pair+jets,  $H$ +jets and inclusive jets
- Only generator that allows matching in BSM processes (e.g. gluino/squark production)

# Smoothness of matching

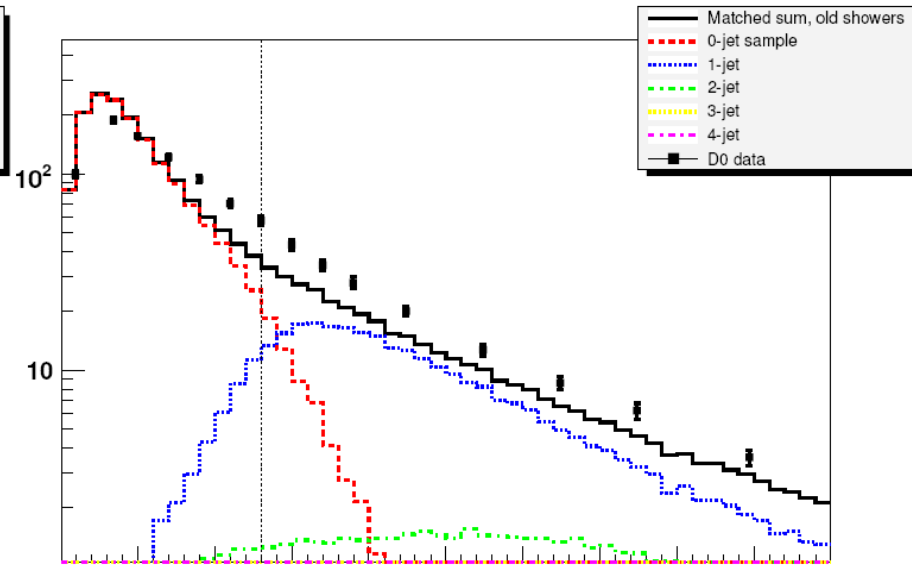
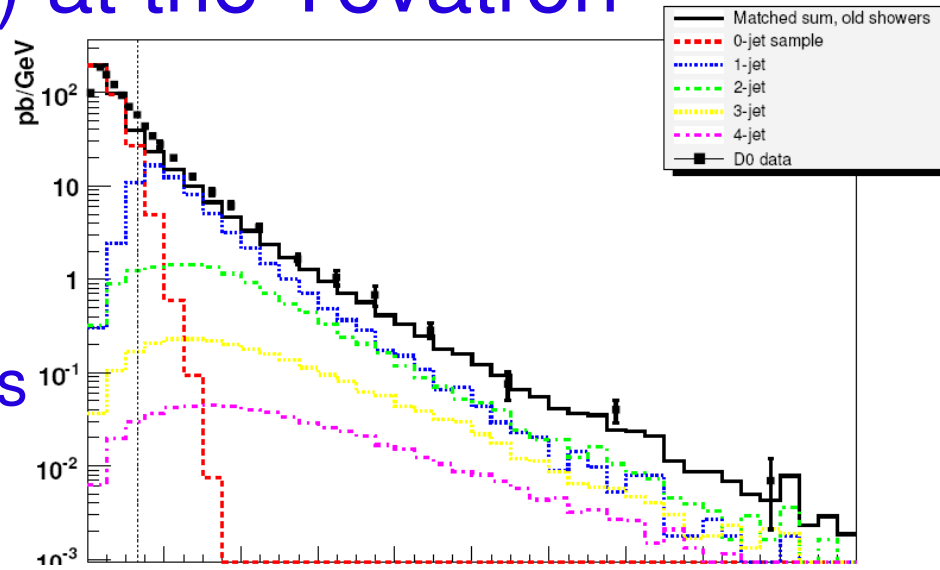


Diff. jet rates by MadEvent+Pythia in  $W$ +jets prod. at the Tevatron

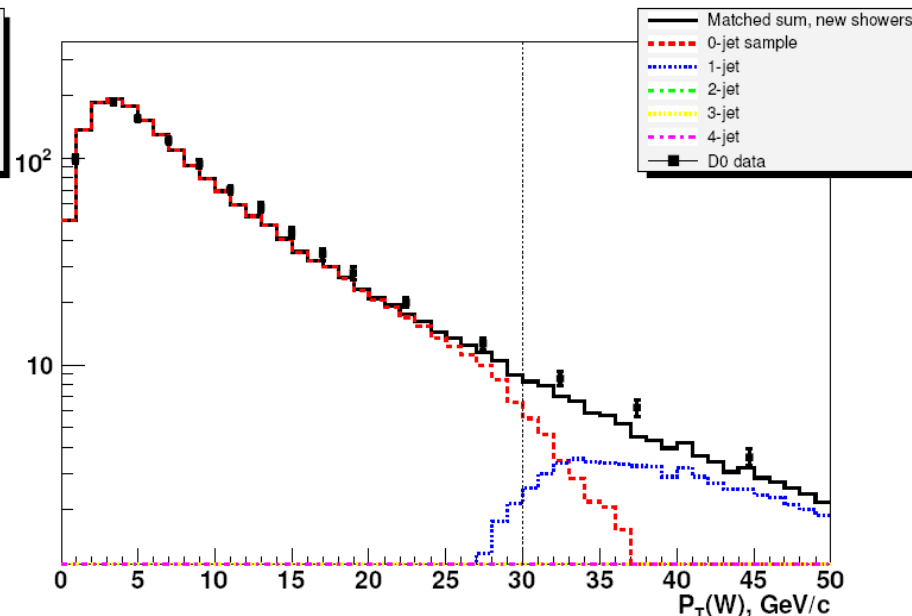
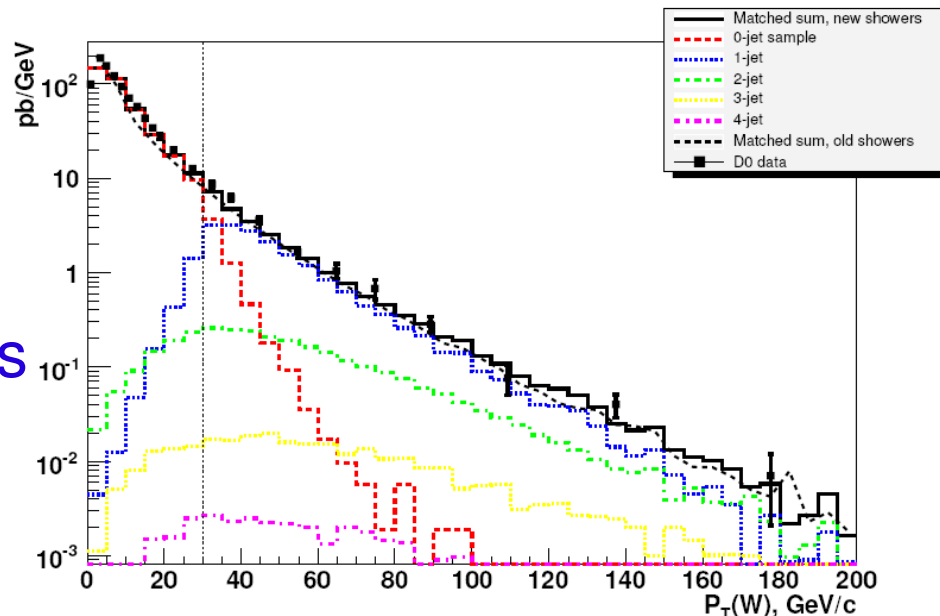
# Comparison with Tevatron Data

## $p_T(W)$ at the Tevatron

“Old”  
Pythia  
showers



“New”  
Pythia  
showers



Tail given by ME, head by pure showers

# Matching in gluino/squark production

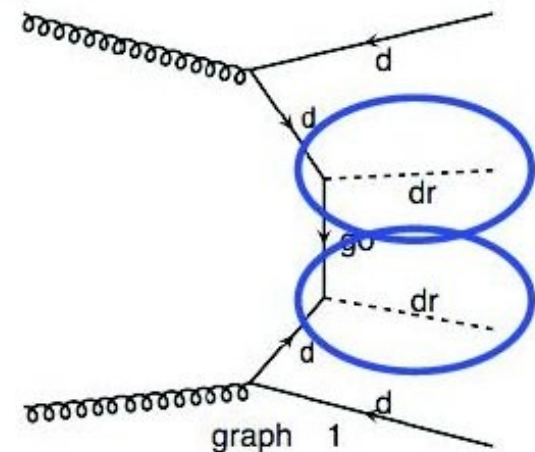
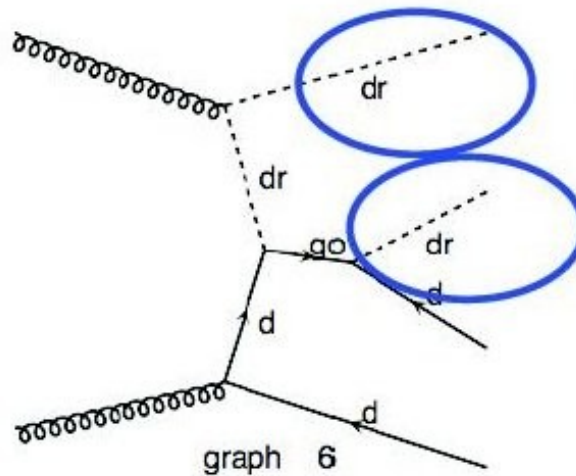
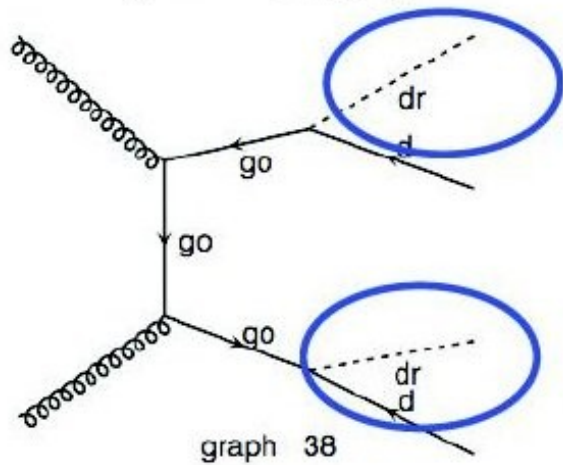
- We know that matching is vital for jet production in SM backgrounds
- But is it relevant for heavy QCD particle production?
  - Very hard jets from decays
  - Parton showers more accurate for larger masses
- Turns out there are many cases where **matching has a large impact!**
- Most important at hadron colliders: ISR matching



# Double counting

- Special difficulty in SUSY matching – double counting between squark and gluino production

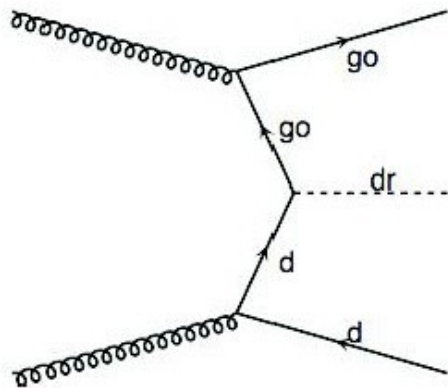
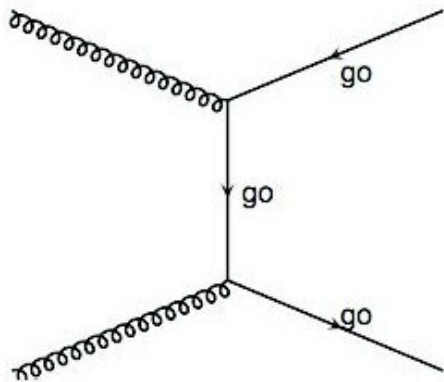
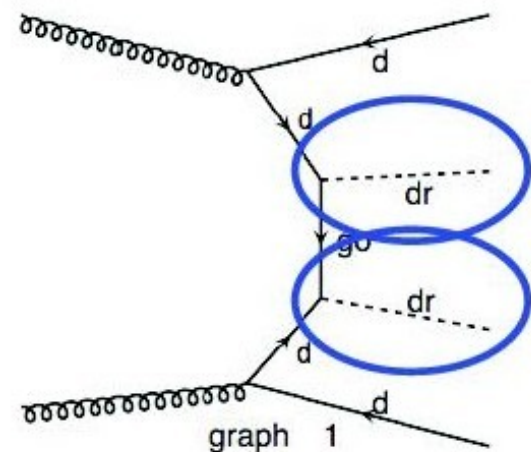
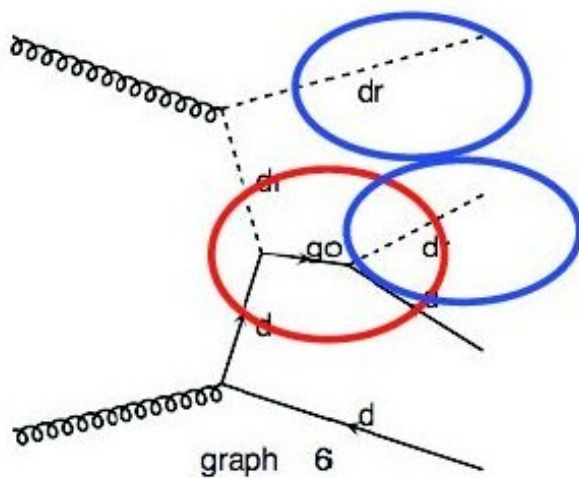
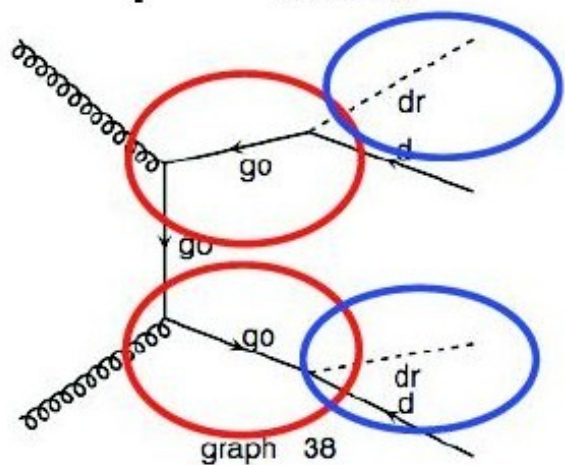
Example:  $\tilde{q}\tilde{q}jj$



# Double counting

- Special difficulty in SUSY matching – double counting between squark and gluino production

Example:  $\tilde{q}\tilde{q}jj$



Double counted  
with on-shell  
gluino prod with  
 $\tilde{g} \rightarrow \tilde{d}_R + q$

# Double counting

- Solved by keeping track of on-shell resonances in the production event files

```
<event>
6 0 0.7992762E-04 0.9118800E+02 0.7816531E-02 0.1300000E+00
 21 -1 0 0 502 503 0.00000000000E+00 0.00000000000E+00 0.38916243784E+03 0.38916243784E+03 0.00000000000E+00 0. 1.
 1 -1 0 0 501 0 0.00000000000E+00 0.00000000000E+00 -0.16355197391E+04 0.16355197391E+04 0.00000000000E+00 0. 1.
1000021 2 1 2 501 503 -0.22162854802E+03 -0.24366260777E+03 -0.12022753376E+04 0.13861620323E+04 0.60620830799E+03 0. 0.
-1 1 3 3 0 503 0.18372150189E+02 0.27121177112E+02 -0.34707630298E+02 0.47725399437E+02 0.00000000000E+00 0. -1.
2000001 1 3 3 501 0 -0.24000069821E+03 -0.27078378488E+03 -0.11675677073E+04 0.13384366329E+04 0.54522846200E+03 0. -1.
2000001 1 1 2 502 0 0.22162854802E+03 0.24366260777E+03 -0.44081963594E+02 0.63852014456E+03 0.54522846200E+03 0. -1.
</event>
```

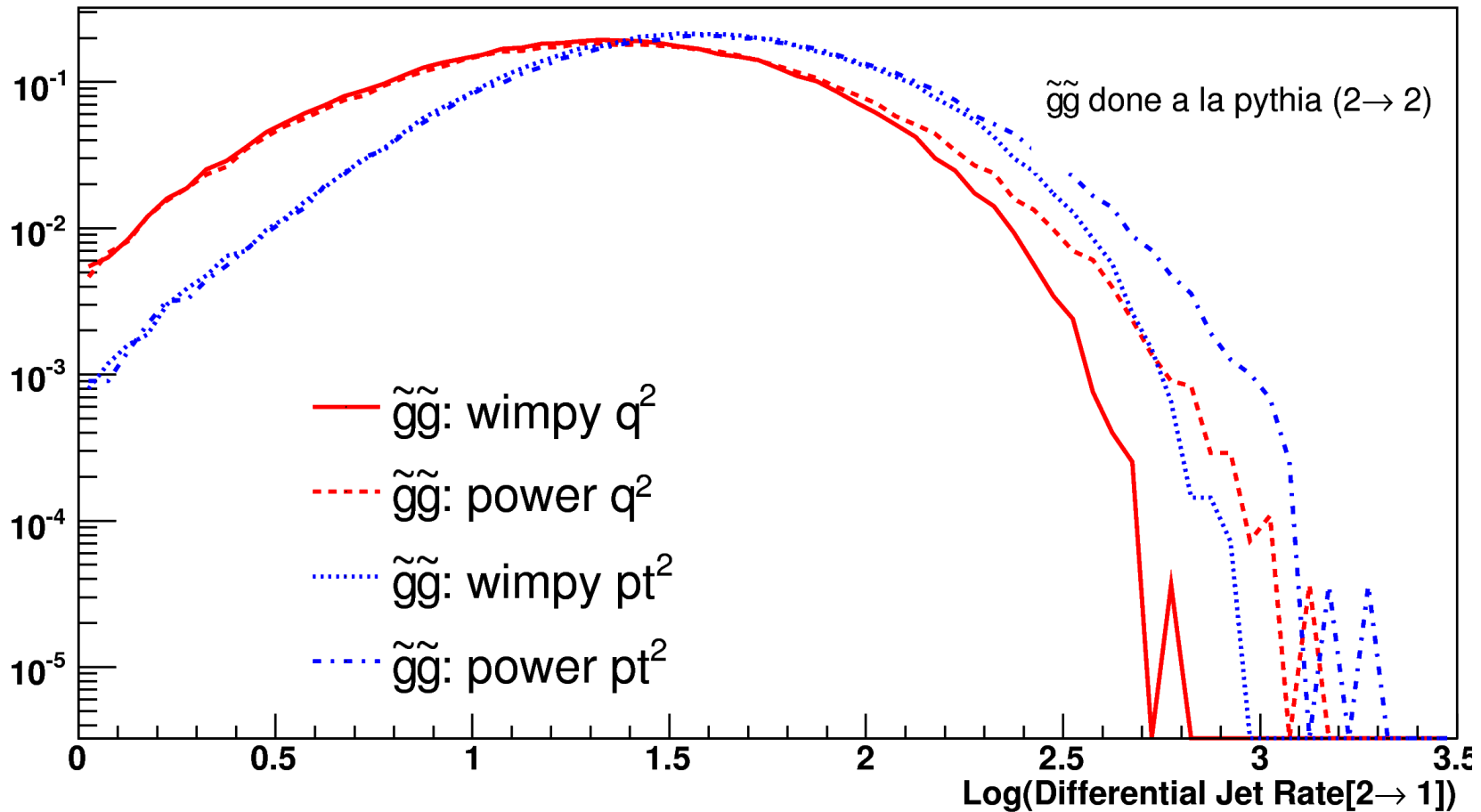
- Allows to remove double-counted events after production
- Double-check – perform generation without resonant diagrams (gauge-inv. only in NWA!)
  - Excellent agreement

# Shower parameter dependence

- Shower “tweakable”
  - Strength for fitting data (after-the-fact)
  - Weakness for predictivity
- Most important parameters used here:
  - Type of shower ( $Q^2$  or  $p_T$ -ordered)
  - Shower starting scale
    - Factorization scale - “wimpy”
    - Total energy of collider - “power”
- Matching quite insensitive to shower parameters

# Shower parameter dependence

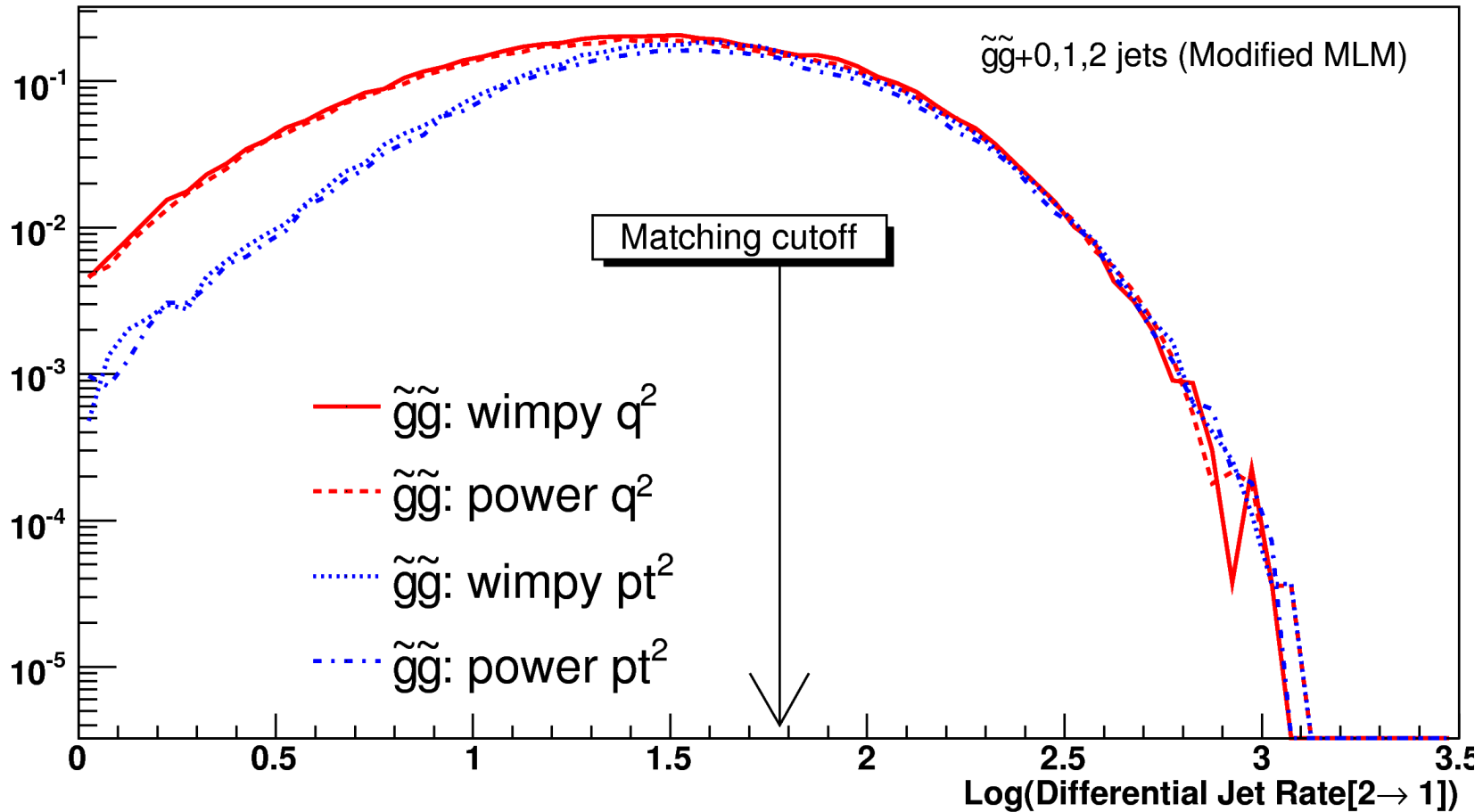
## ISR jets for different Pythia shower params



600 GeV gluino pair production at the LHC

# Shower parameter dependence

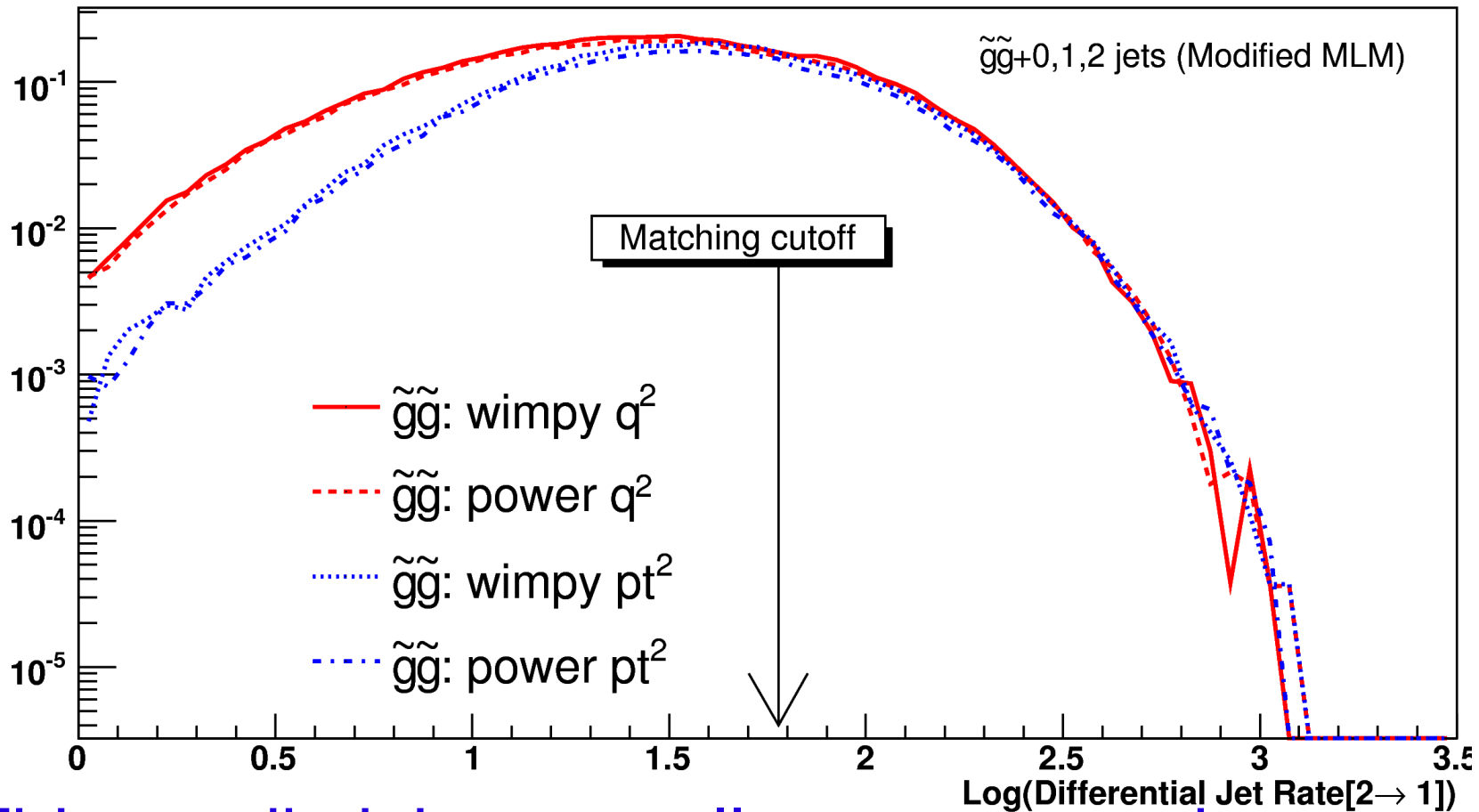
## ISR jets after matching with MG/ME



600 GeV gluino pair production at the LHC

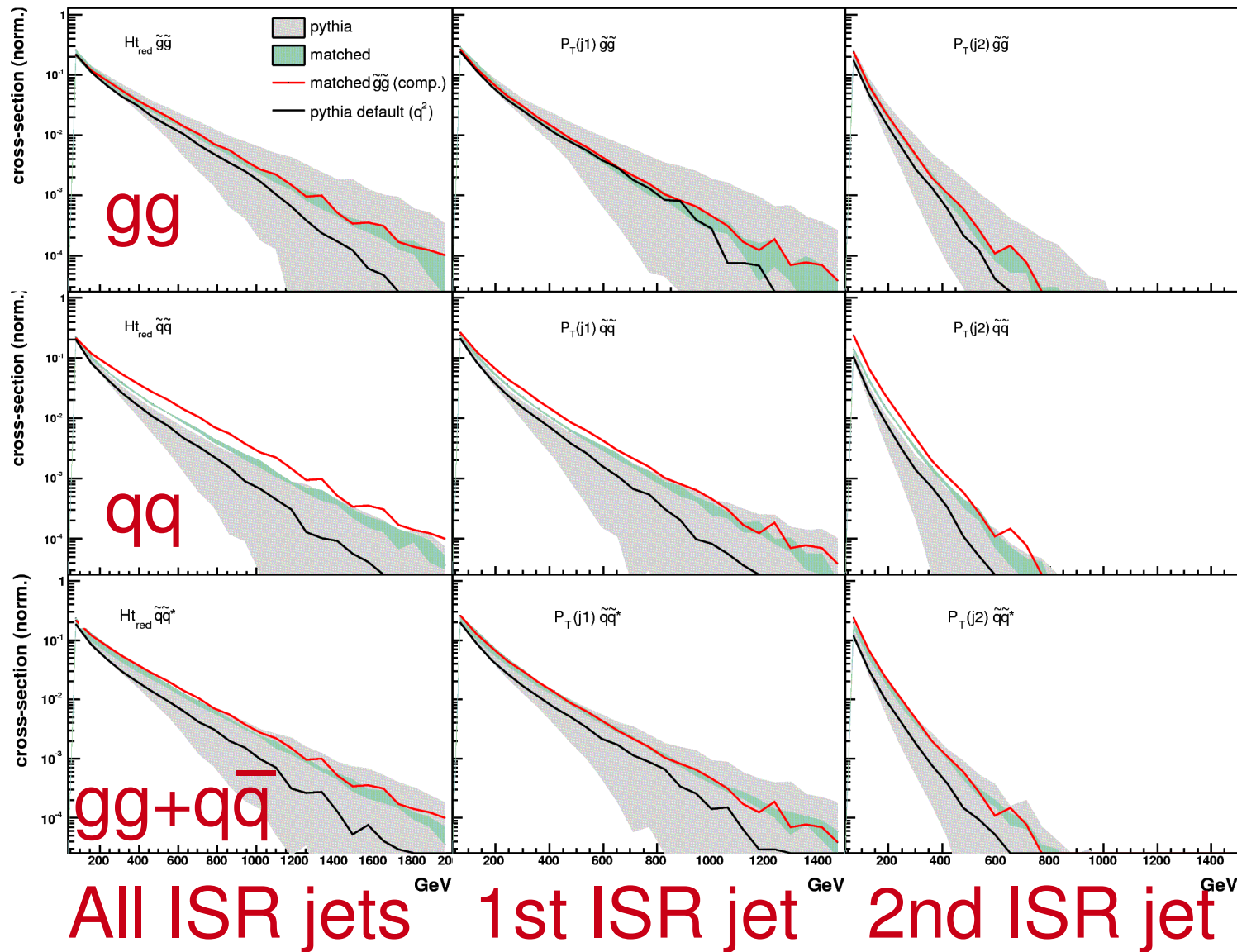
# Shower parameter dependence

## ISR jets after matching with MG/ME



This predictivity now allows us to make proper investigations of dependence on different factors

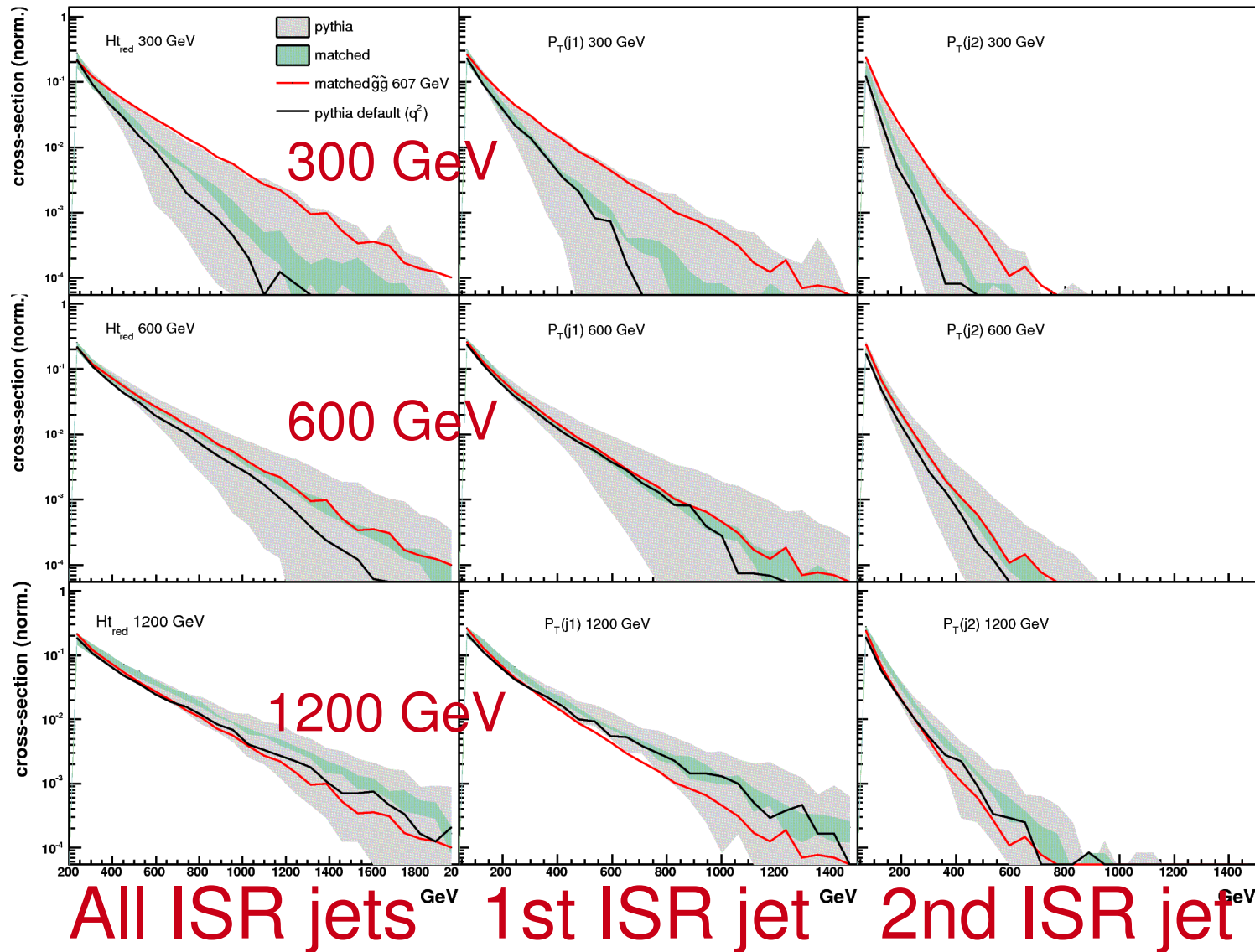
# Dependence on the initial state: $gg$ , $qq$ , $q\bar{q}$



No single shower tune for all initial states!



# Dependence on the produced particle mass



Can ISR help determine overall mass scale?

# Impact of matching with decay

Matching obviously important for radiated jets, but what about including decay jets from the squarks/gluinos?

Still important in many scenarios!

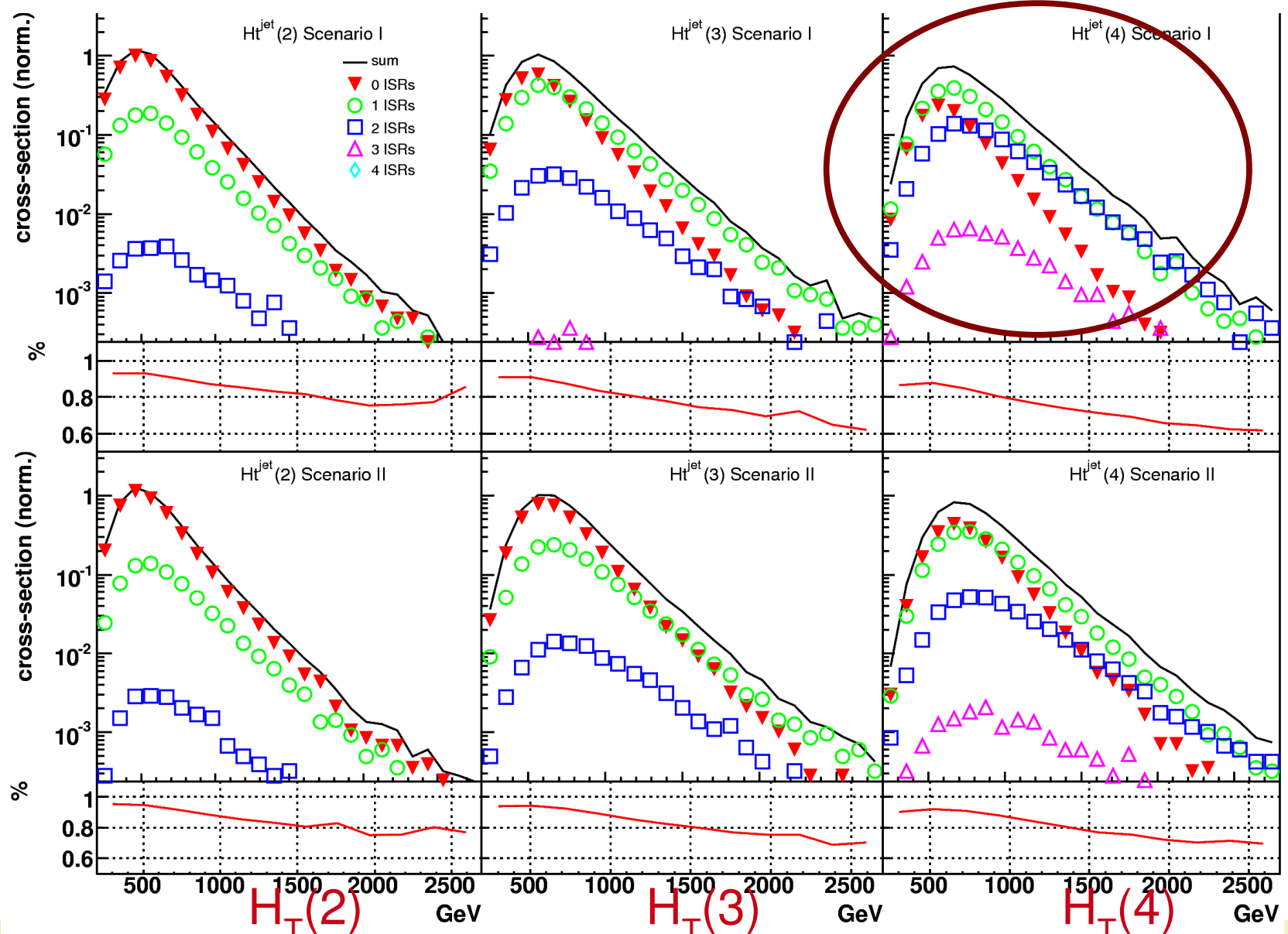
Example: Small mass splitting between gluino and squark – gives gluino decay with 2 hard + 2 soft jets

# Impact of matching with decay

## 600 GeV gluino pair production

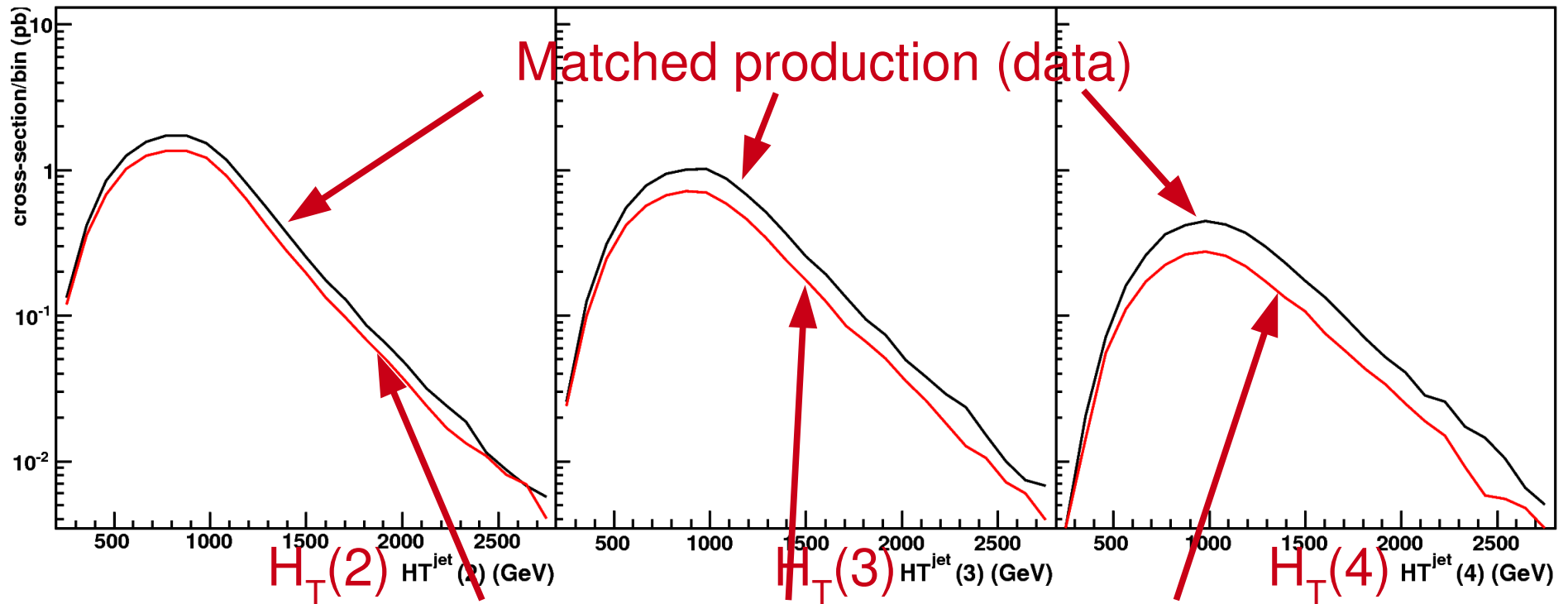
$M_g - M_q =$   
50 GeV

squarks  
heavy –  
3-body  
~g decay



# Risk for misinterpretation

Scenario: s<sub>q</sub>s<sub>q</sub> production only (gluinos heavy)

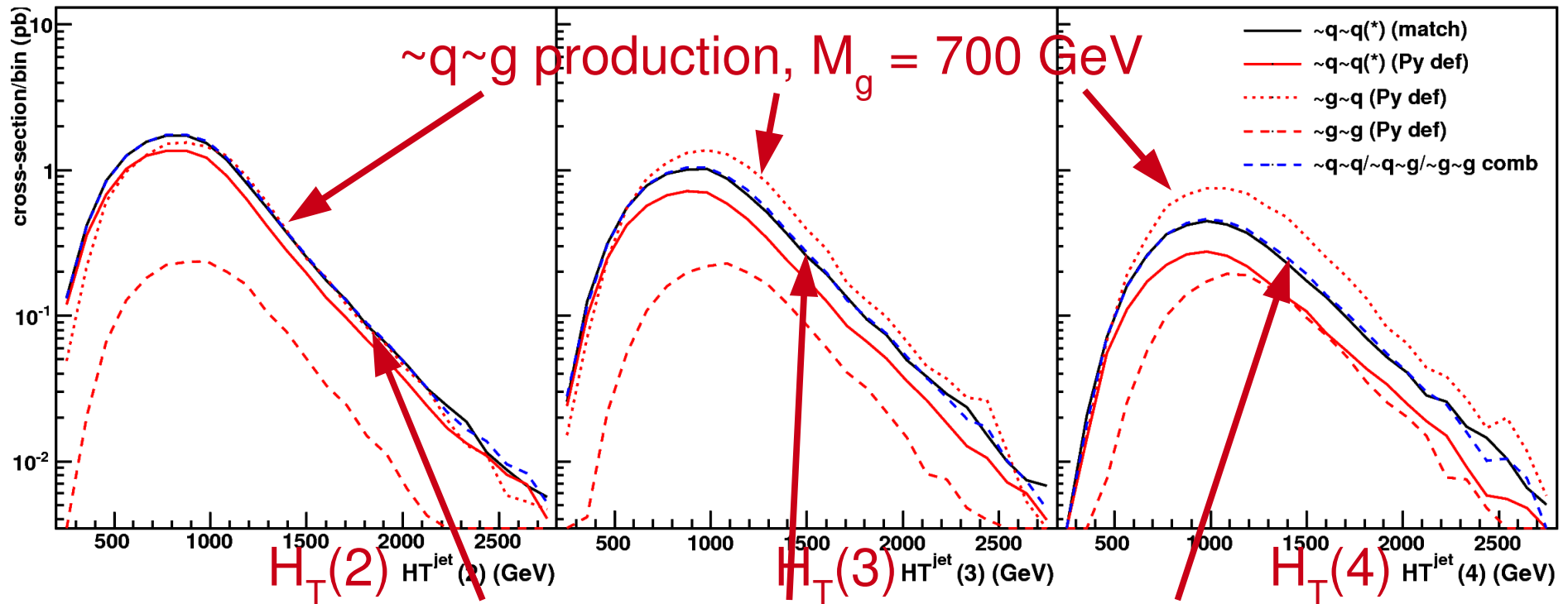


Unmatched production (simulation, Pythia default)

Looks like we're missing gluino component!

# Risk for misinterpretation

Scenario:  $\tilde{q}\tilde{q}$  production only (gluinos heavy)



Unmatched  $\tilde{q}\tilde{q} + 25\%\tilde{q}\tilde{g}$  (fits “data”)

Easy misinterpretation:  $\tilde{q}\tilde{g}$  component!

# Light gluinos at the Tevatron

JA, Le, Lisanti, Wacker [arXiv:0803.0019,  
arXiv:0809.3264]

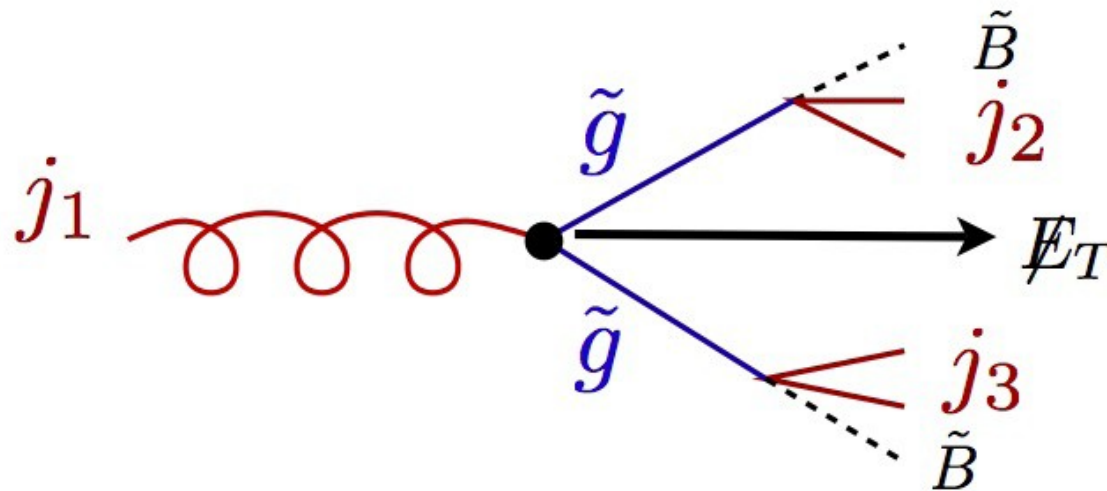
- Searches at the Tevatron have always been done in the mSUGRA framework
- mSUGRA (and mGMSB) “special” scenarios: fixed mass ratio  $m_g:m_W:m_B \sim 6:2:1$
- Not representative for general MSSM (or other BSM models!)
- Study of projected exclusion region at Tevatron with free ratio  $m_g:m_B$  (arXiv:0803.0019)

# Light gluinos at the Tevatron

JA, Le, Lisanti, Wacker [arXiv:0803.0019,  
arXiv:0809.3264]

Special difficulty when decay products are soft  
(nearly degenerate masses):

- No (small) missing transverse energy in decay
- Need recoil against jets to get  $\cancel{E}_T$  signature



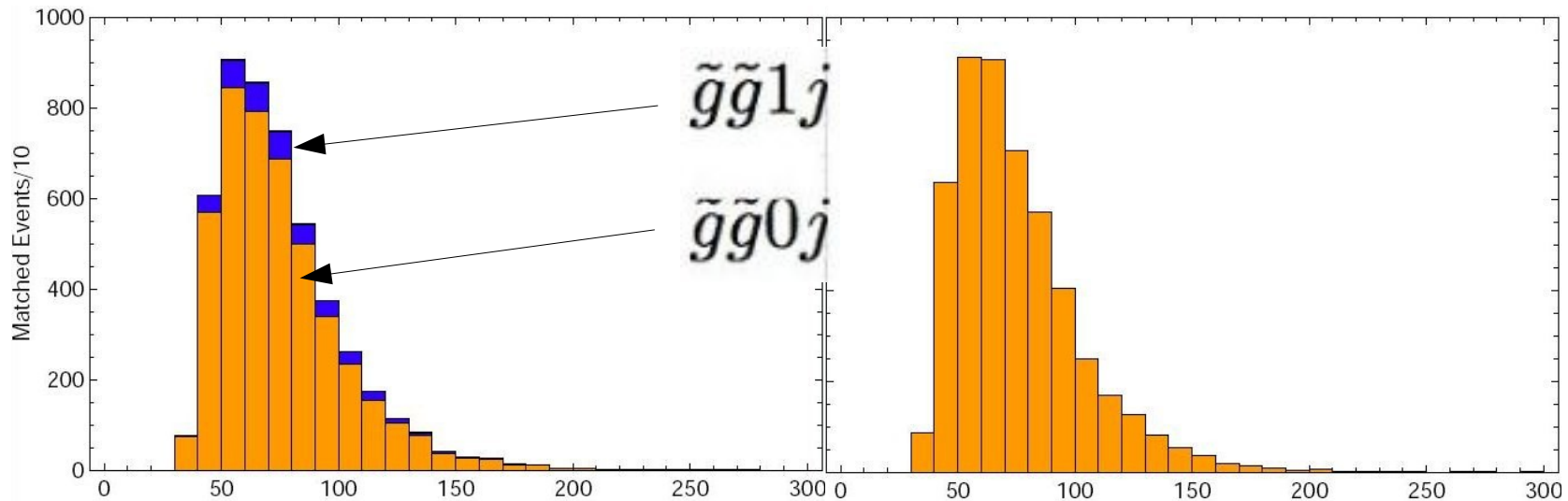
# Light gluinos at the Tevatron

Matched

Unmatched

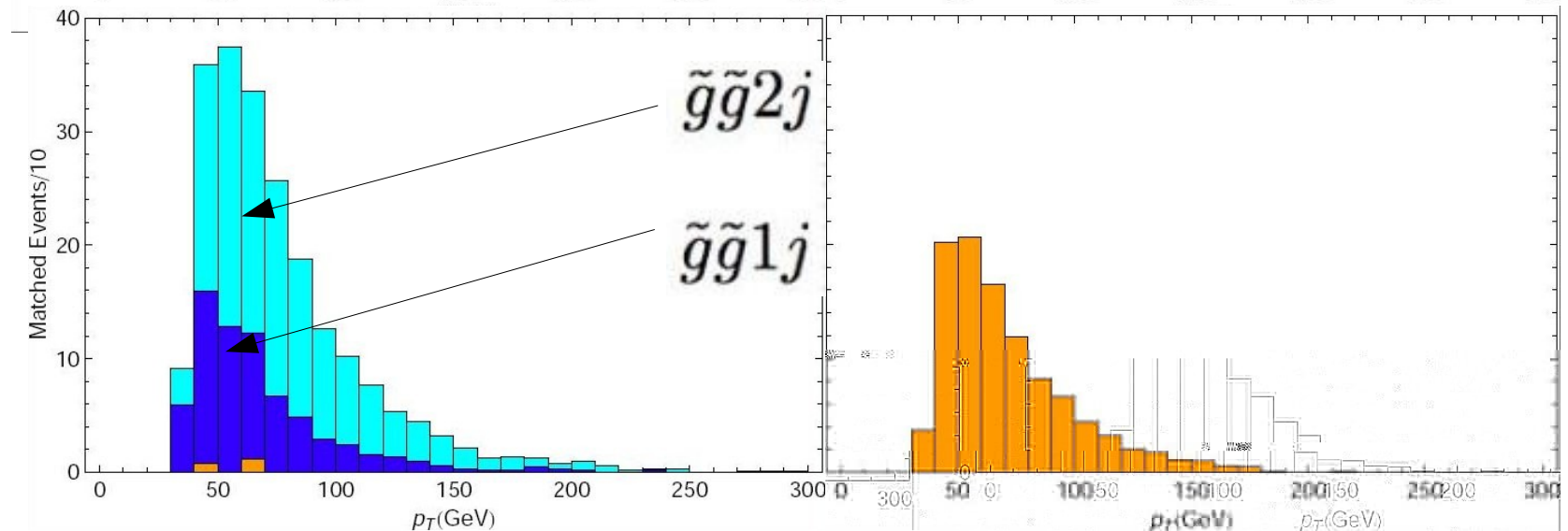
$M_g = 150$  GeV

$M_B = 40$  GeV



$M_g = 150$  GeV

$M_B = 130$  GeV

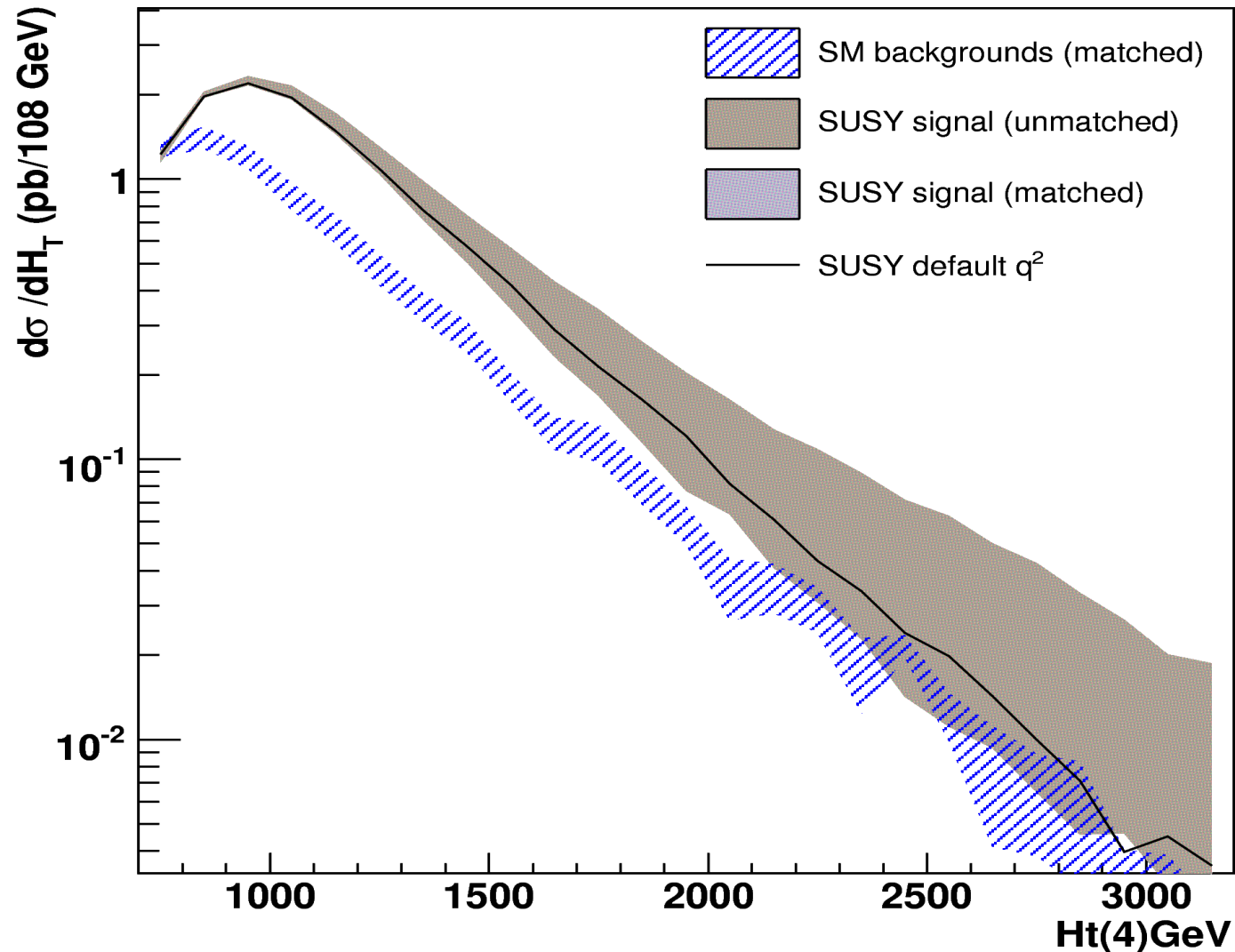


Tevatron, after 2-jet and missing  $E_T$  cuts



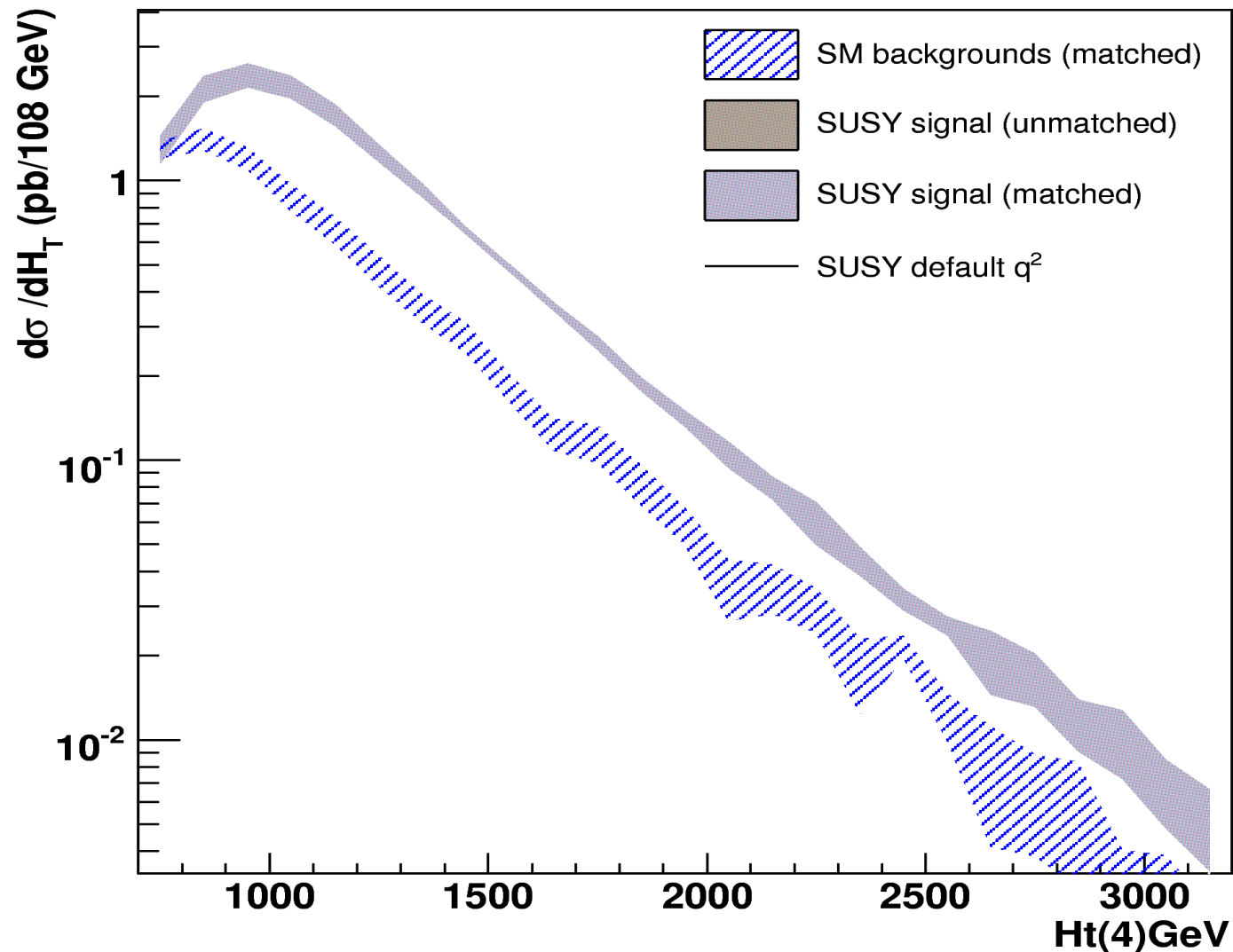
# Impact on collective exp. searches

“SPS1a” SUSY production at the LHC - unmatched



# Impact on collective exp. searches

“SPS1a” SUSY production at the LHC - matched



# Conclusions

- The LHC is a hadronic collider – busy and complicated QCD environment
- Any search for new physics must take QCD radiation effects into account
  - Extra jet production from ISR QCD emissions
  - Boost of the central production system
- Many scenarios where jet matching between matrix elements and parton showers is crucial to get a good description of the signal as well as SM backgrounds

# Backup slides

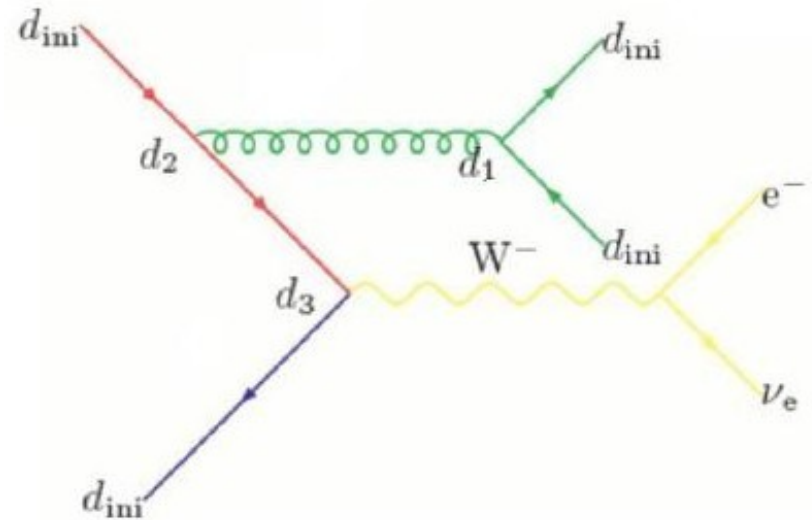
# CKKW matching

## Imitate parton shower procedure for matrix elements

- 1 Choose a cutoff (jet resolution) scale  $d_{\text{ini}}$
- 2 Generate multiparton event with  $d_{\text{min}} = d_{\text{ini}}$  and factorization scale  $d_{\text{ini}}$
- 3 Cluster event with  $k_T$  algorithm to find “parton shower history”
- 4 Use  $d_i \simeq k_T^2$  in each vertex as scale for  $\alpha_s$
- 5 Weight event with NLL Sudakov factor  $\Delta(d_j, d_{\text{ini}})/\Delta(d_i, d_{\text{ini}})$  for each parton line between vertices  $i$  and  $j$  ( $d_j$  can be  $d_{\text{ini}}$ )
- 6 Shower event, allowing only emissions with  $k_T < d_{\text{ini}}$  (“vetoed showers”)
- 7 For highest multiplicity sample, use  $\min(d_i)$  of event as  $d_{\text{ini}}$

Boost-invariant  $k_T$  measure:

$$\begin{cases} d_{iB} = p_{T,i}^2 \\ d_{ij} = \min(p_{T,i}^2, p_{T,j}^2) F_{ij} \\ F_{ij} = 2 \{ \cosh(\eta_i - \eta_j) - \cos(\phi_i - \phi_j) \} \end{cases}$$



- For final-state showers: Combination of NLL Sudakov factors and vetoed NLL showers **guarantees independence of  $d_{ini}$  to NLL order**
- For initial-state showers: No proof but **works ok**
- Problem in practice: No NLL shower implementation! (Sherpa uses Pythia-like showers)

# More about matching in MG

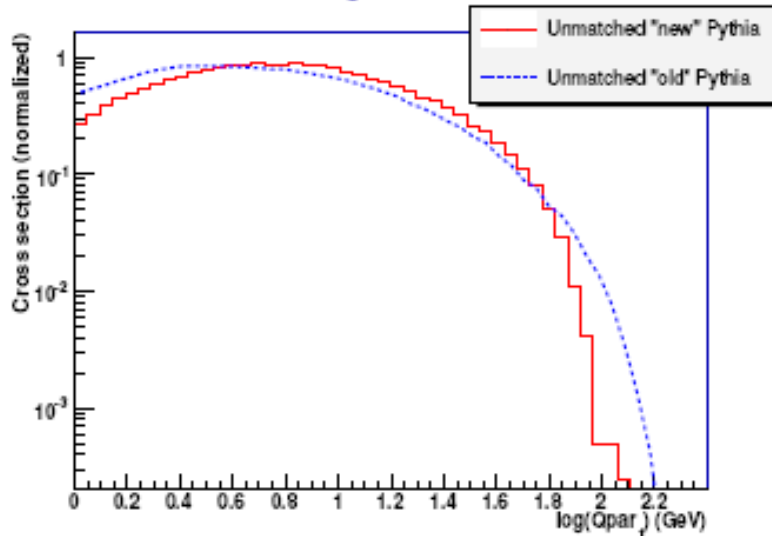
## Shower kT scheme

- Keep/reject event based on  $k_T$  of hardest shower emission (as reported by Pythia)
- Highest multiplicity treatment as in CKKW, use min  $d_{\text{parton}}$  as cutoff
- No jet clustering
- No need of “fiducial region”, can use  $k_T^{\text{match}} = d_{\text{cut}}^{\text{ME}}$
- Need similar kT definitions in ME and PS (only “new”,  $p_T$ -ordered showers at present)

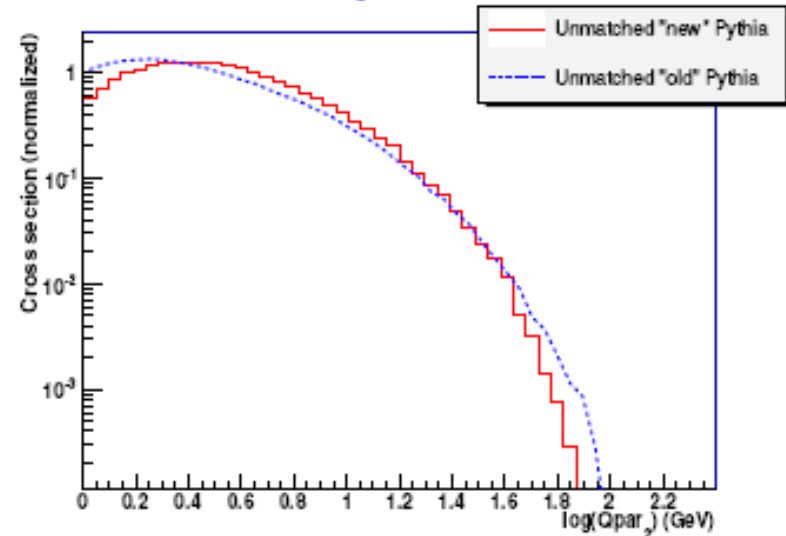
# Comparisons between old and new Pythia showers

Differential jet rates in  $W$  production at the Tevatron

$0 \rightarrow 1$  jet rate



$1 \rightarrow 2$  jet rate



$p_T(W)$  in  $W$  production at the Tevatron

