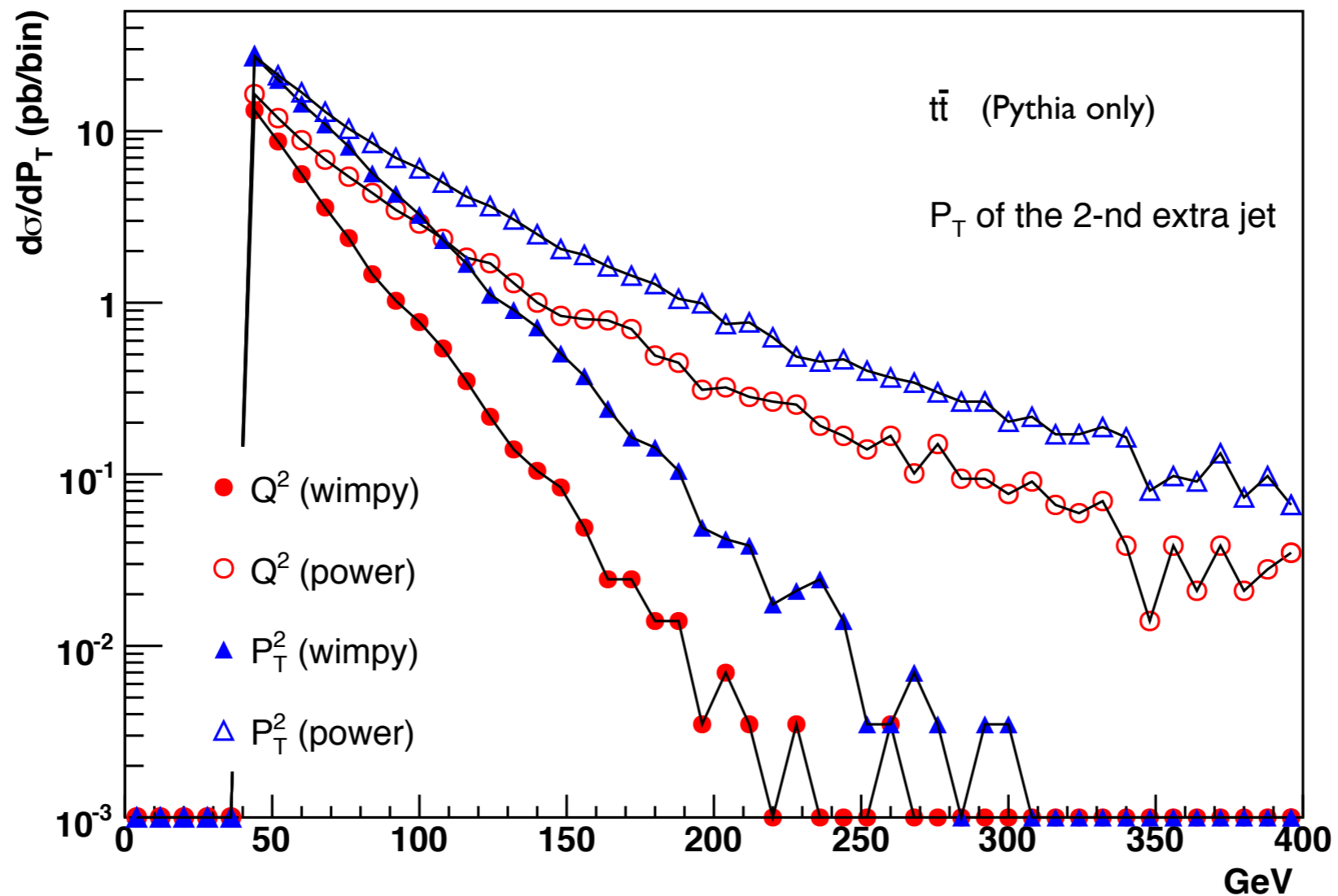


Matching/Merging

Olivier Mattelaer
IPPP/Durham

In the soft-collinear approximation of Parton Shower MCs, parameters are used to tune the result \Rightarrow Large variation in results (small prediction power)

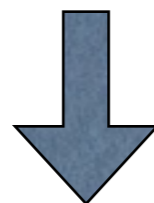




Matrix Elements vs. Parton Showers

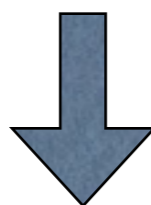


ME



1. Fixed order calculation
2. Computationally expensive
3. Limited number of particles
4. Valid when partons are **hard and well separated**
5. Quantum interference correct
6. Needed for multi-jet description

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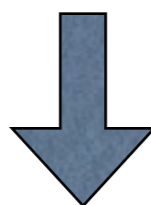
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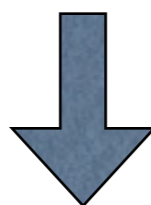
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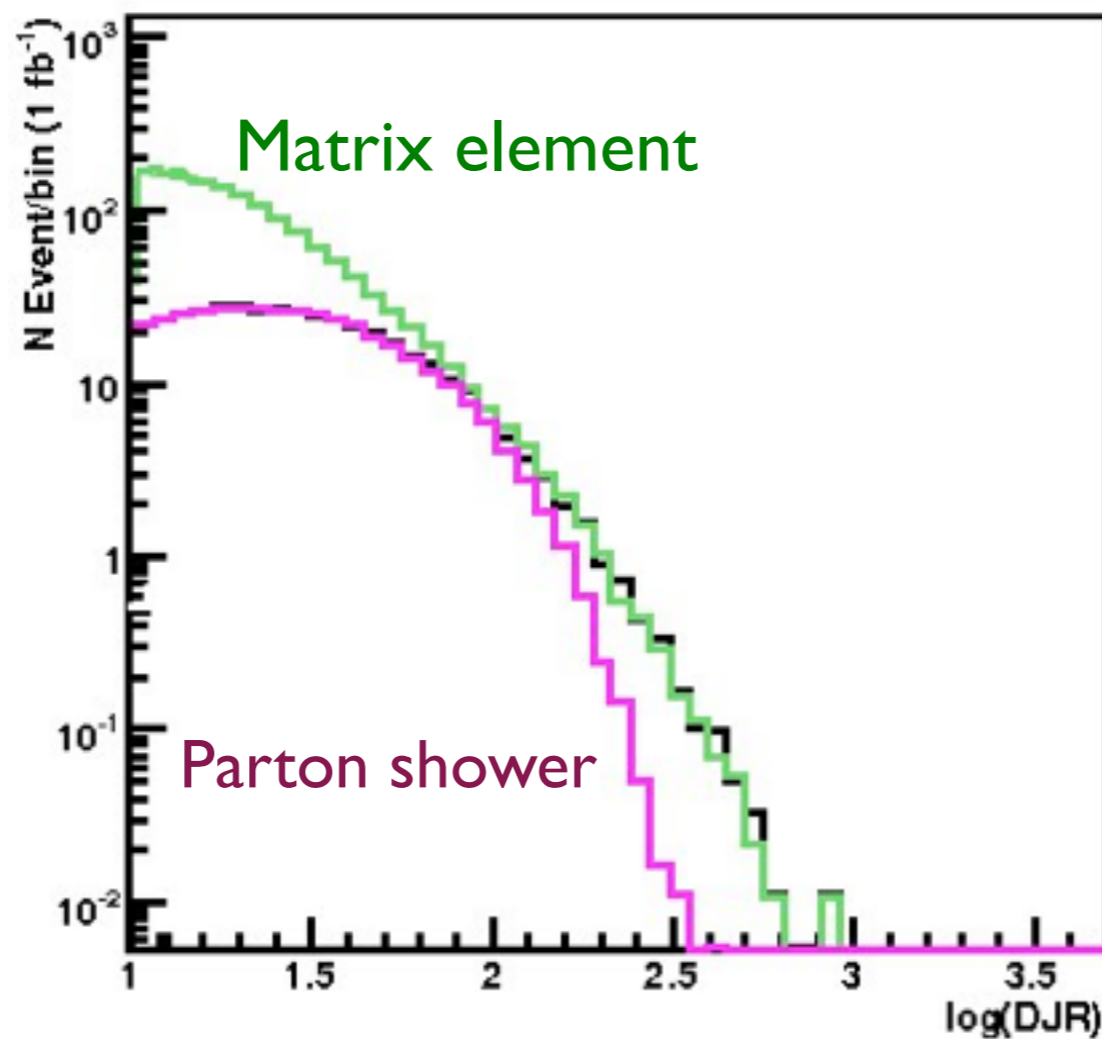
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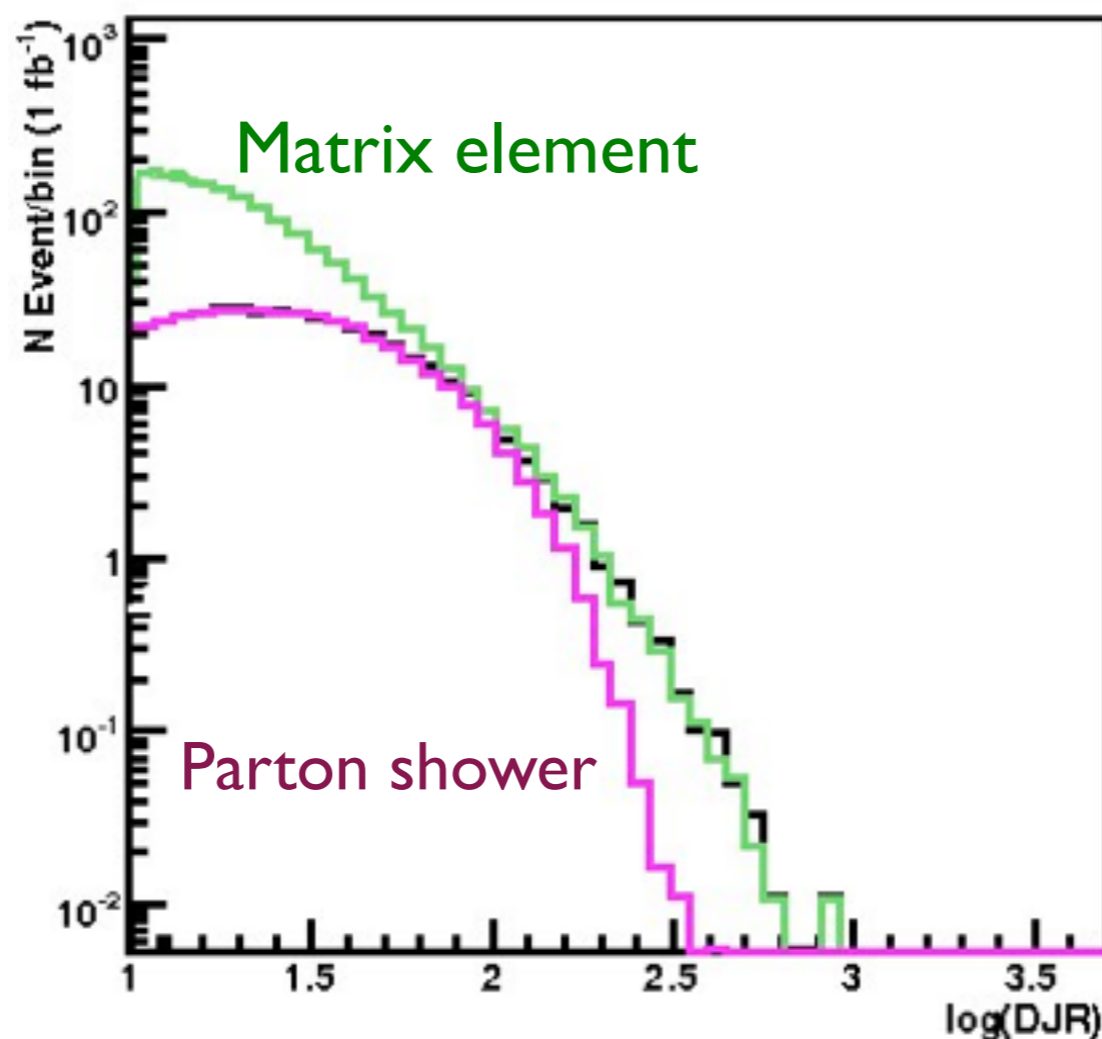
Approaches are complementary: merge them!

Difficulty: avoid double counting, ensure smooth distributions



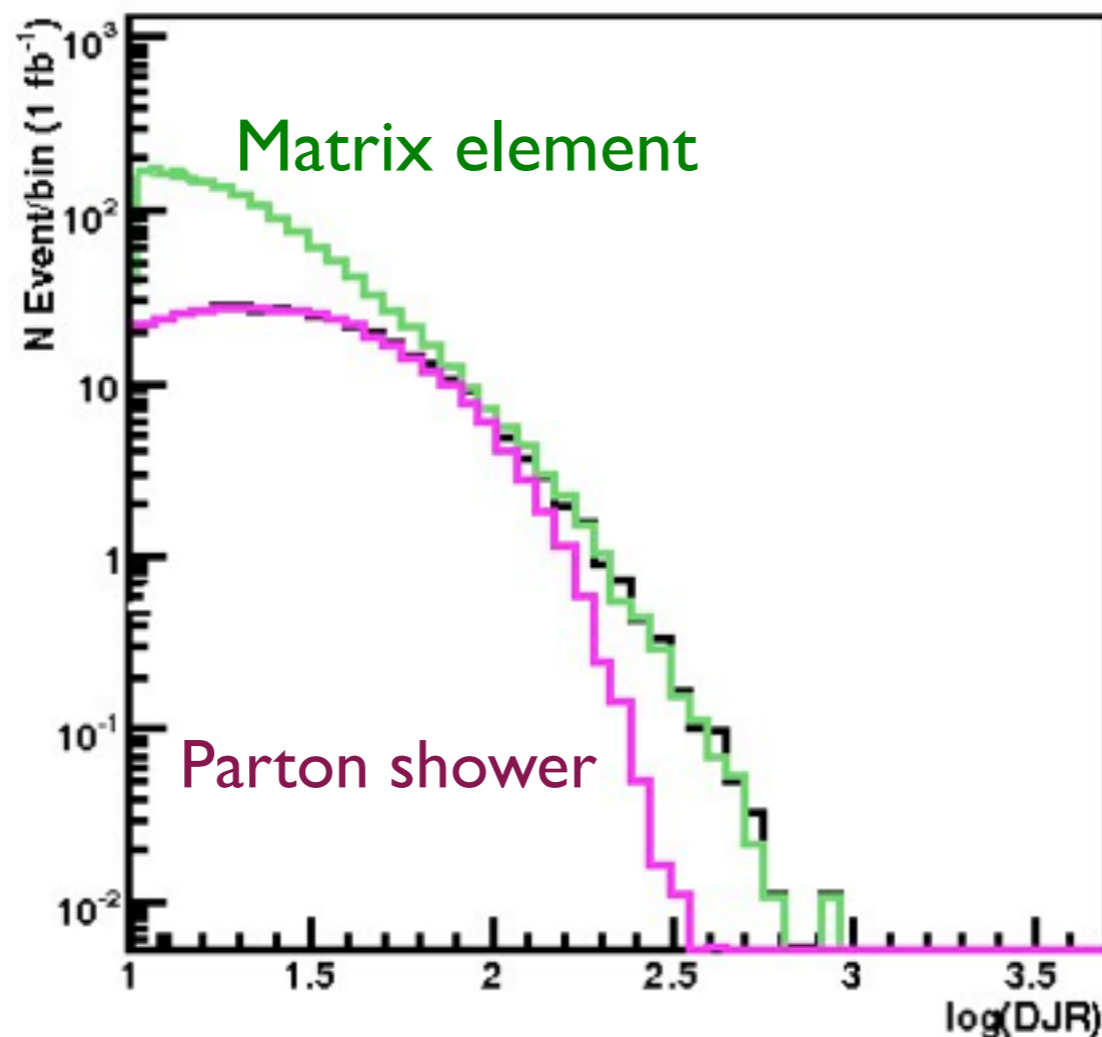
2nd QCD radiation jet in top pair production at the LHC, using MadGraph + Pythia

- Regularization of matrix element divergence



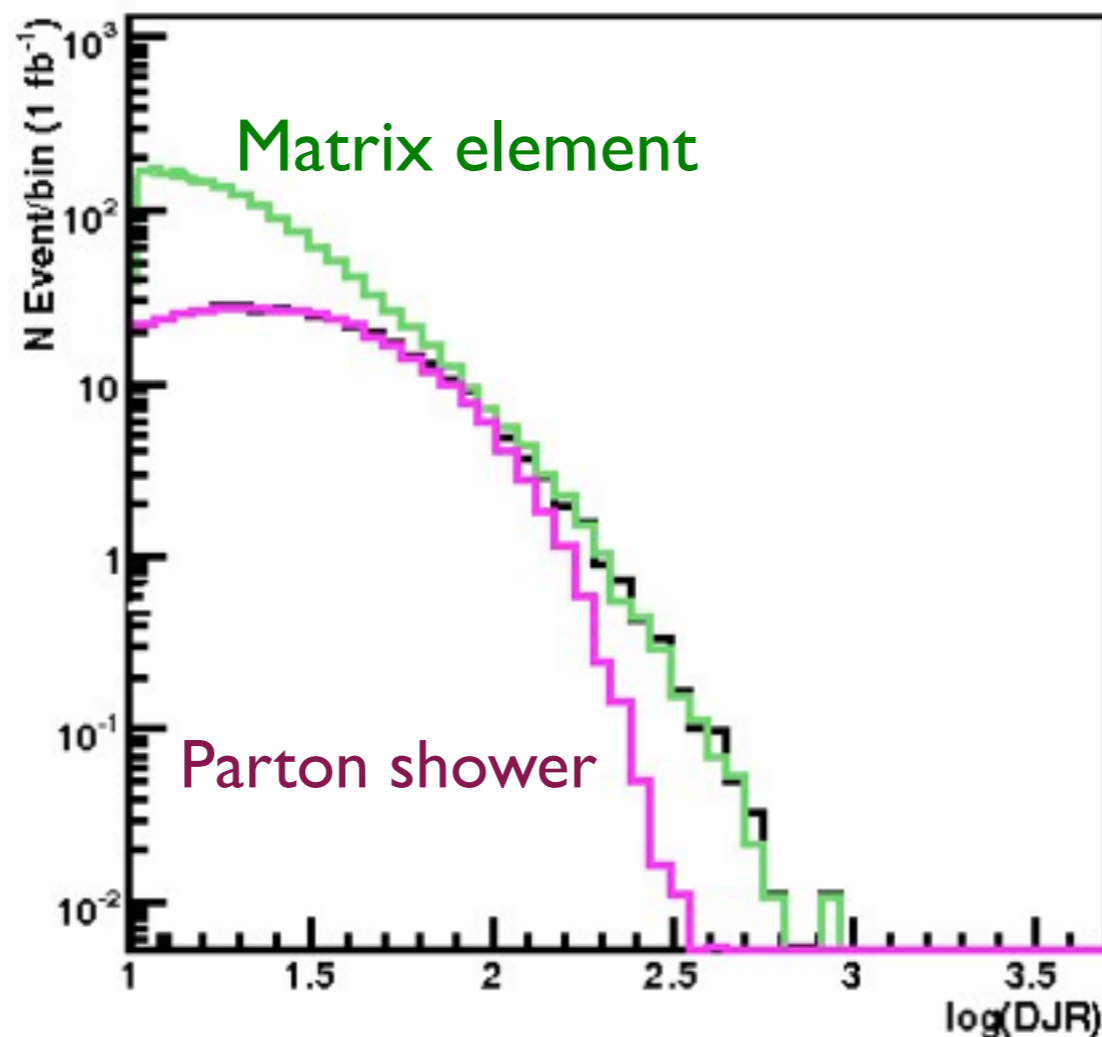
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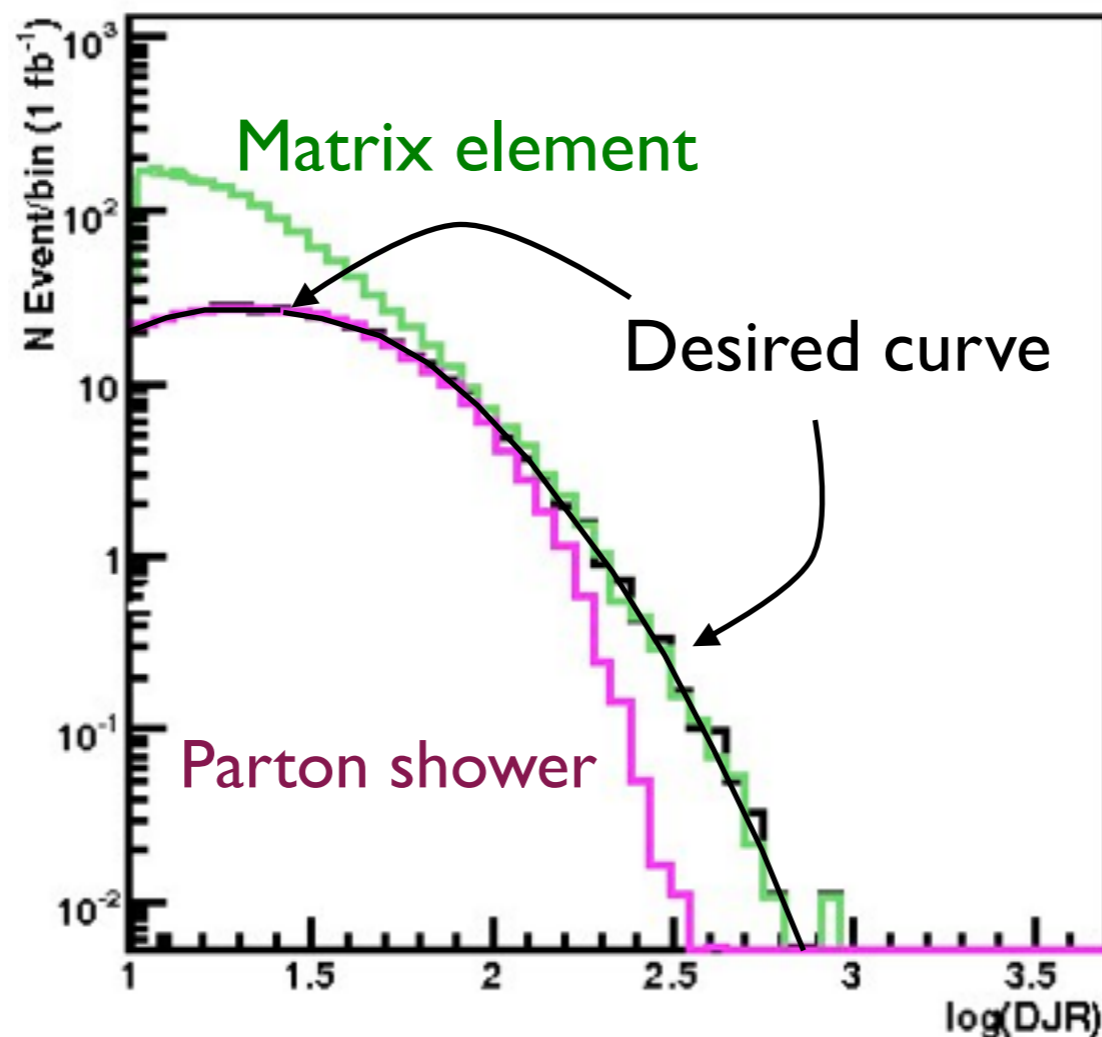
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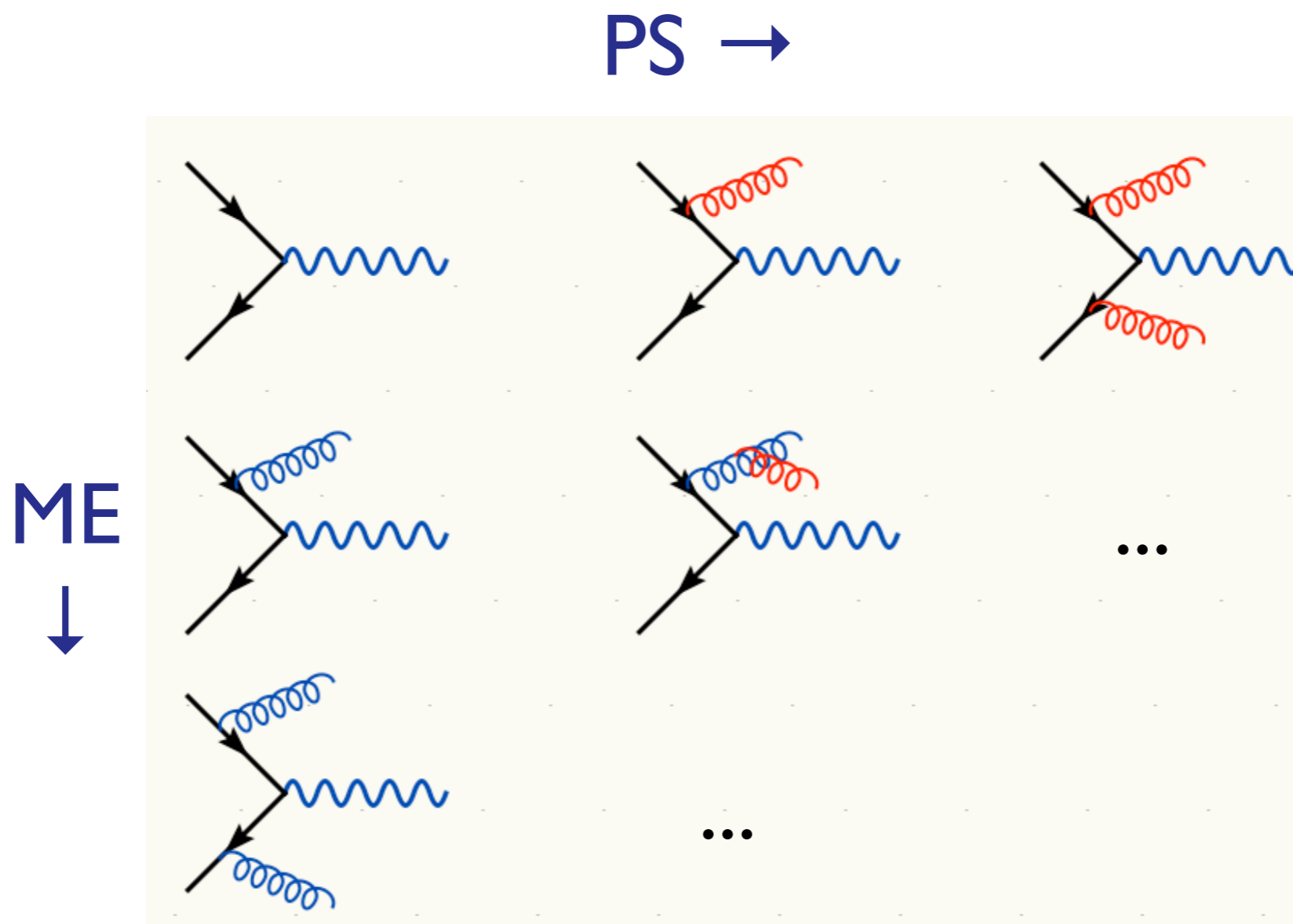
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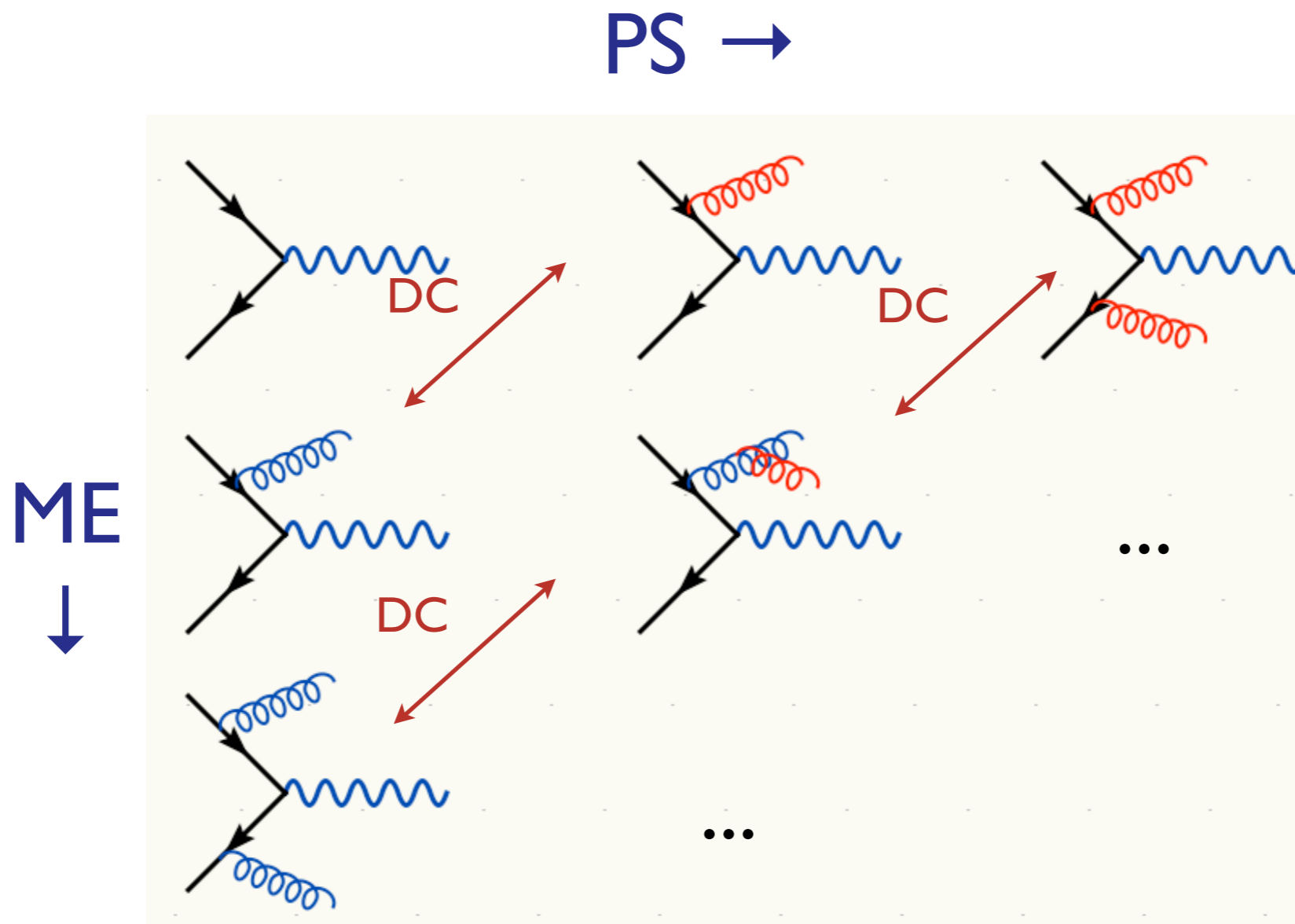


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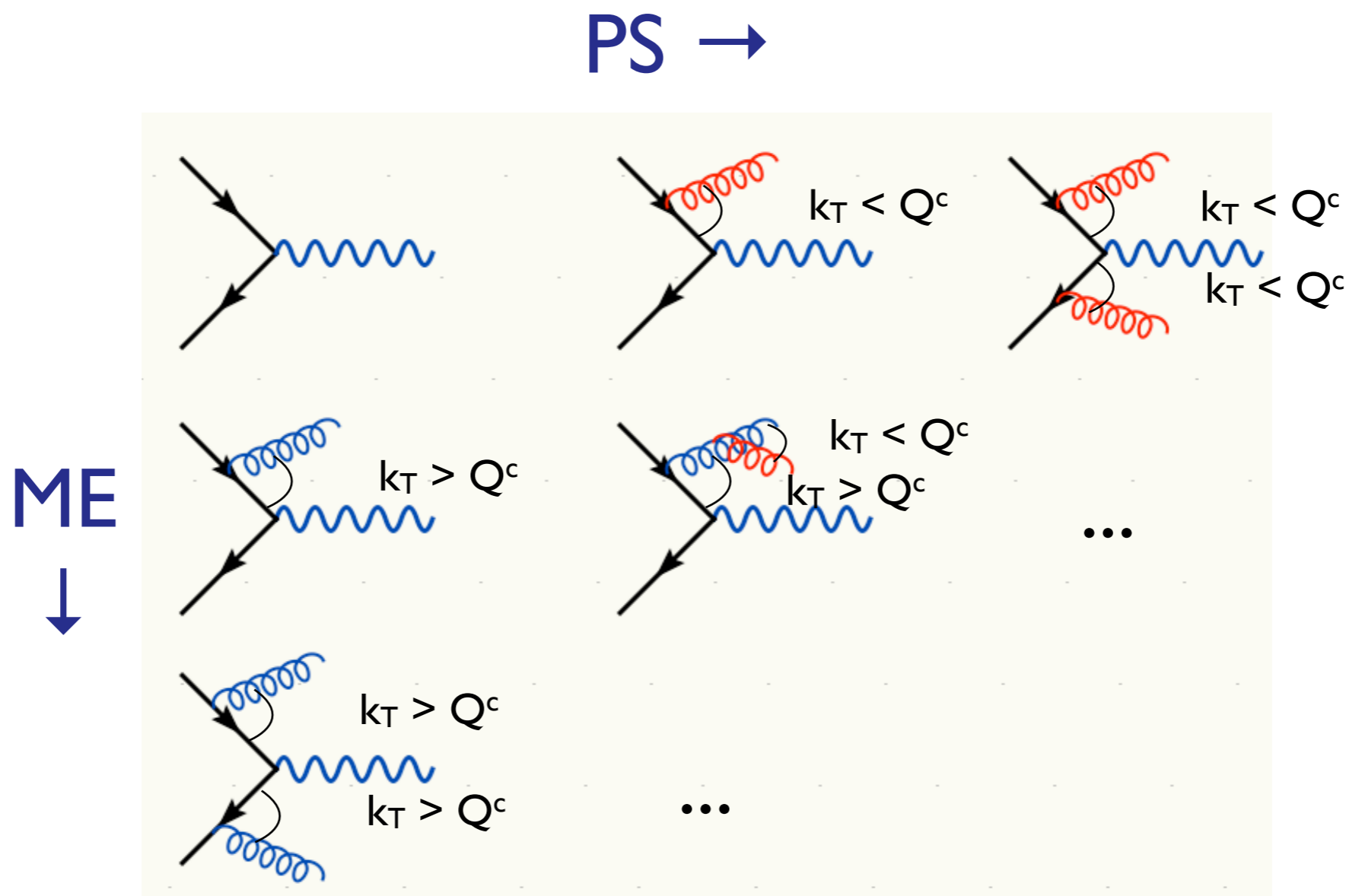
[Mangano]
 [Catani, Krauss, Kuhn, Webber]
 [Lönnblad]



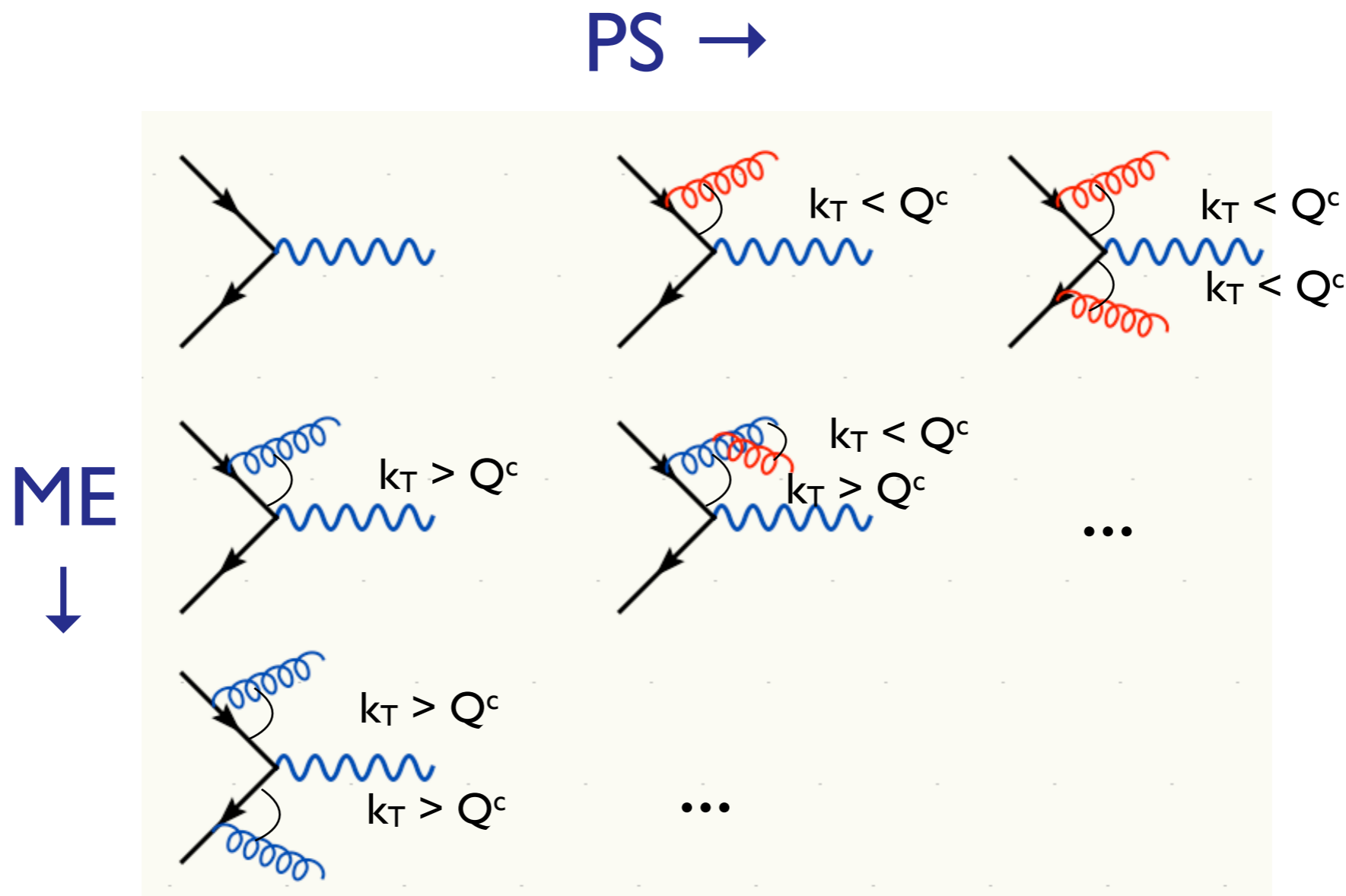
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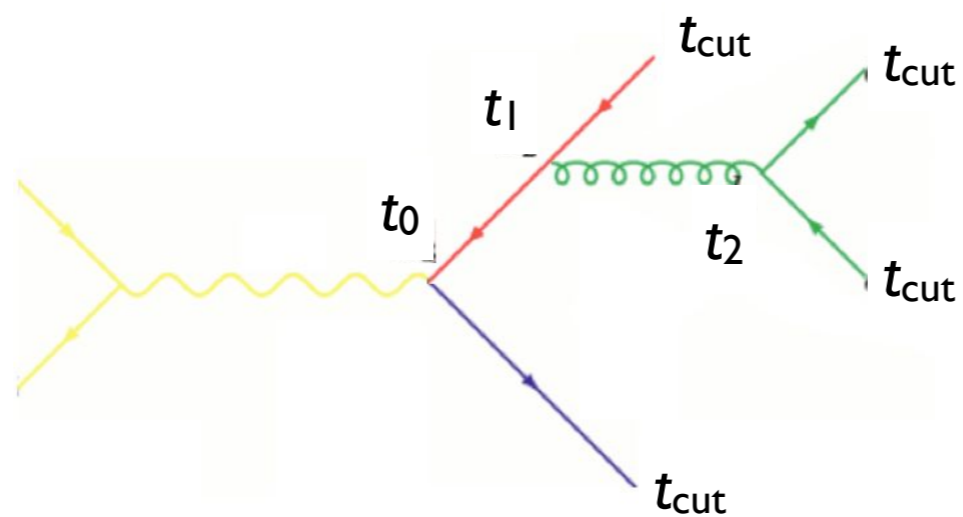


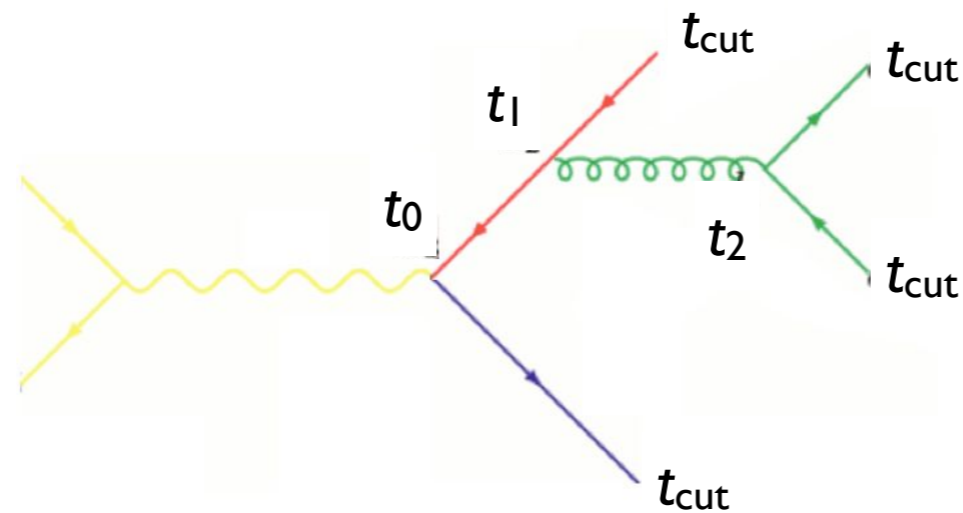
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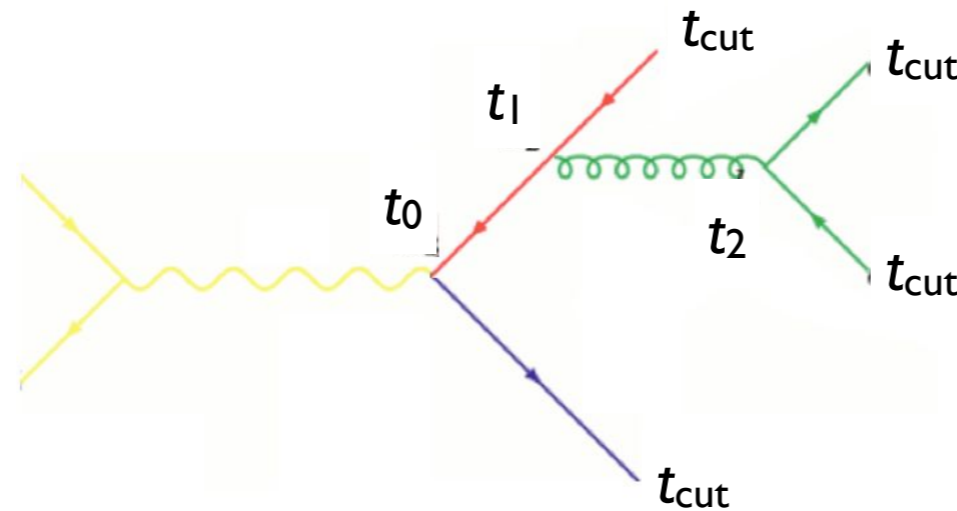
Double counting between ME and PS easily avoided using phase space cut between the two: PS below cutoff, ME above cutoff.

- So double counting problem easily solved, but what about getting smooth distributions that are independent of the precise value of Q^c ?
- Below cutoff, distribution is given by PS
 - need to make ME look like PS near cutoff
- Let's take another look at the PS!



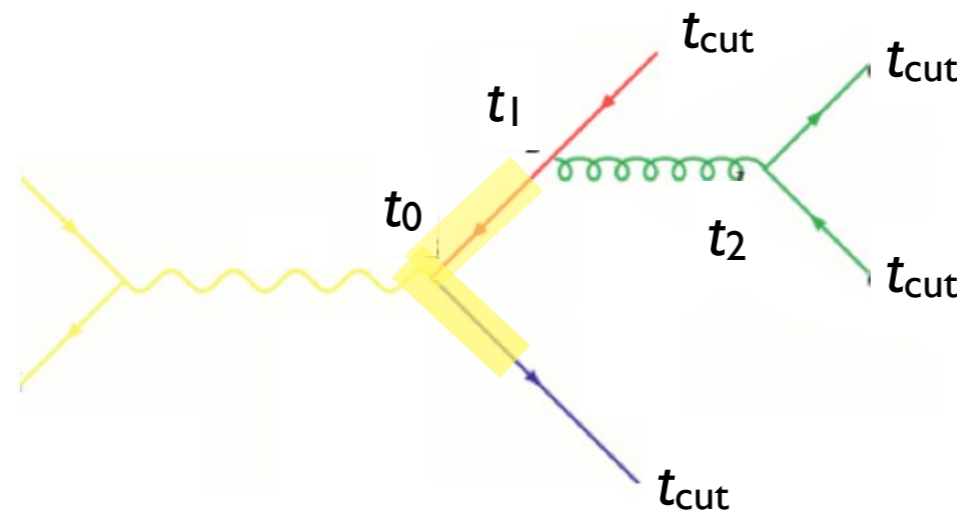


- How does the PS generate the configuration above?



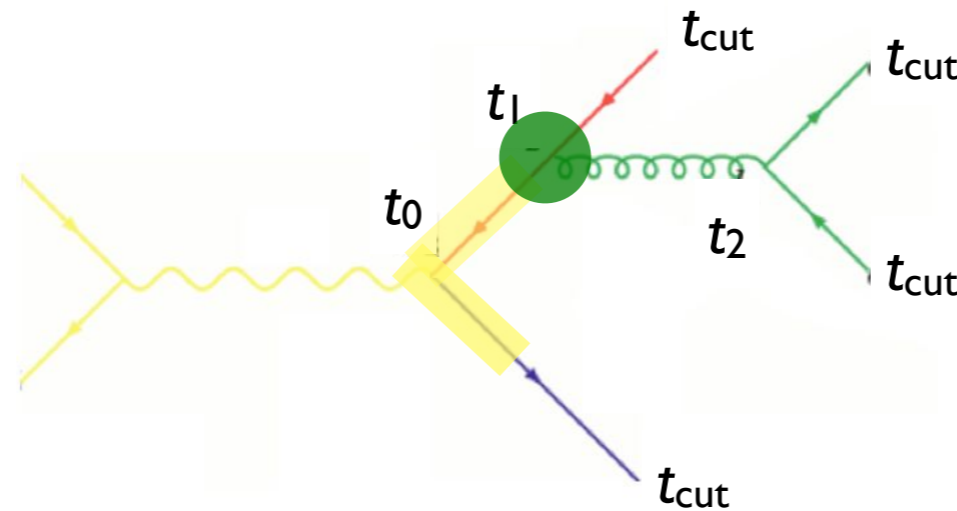
- How does the PS generate the configuration above?
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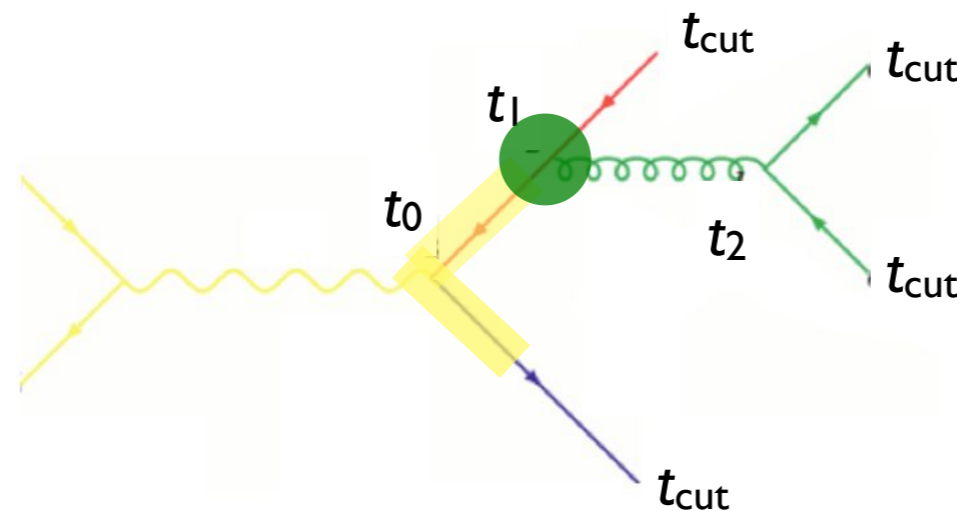
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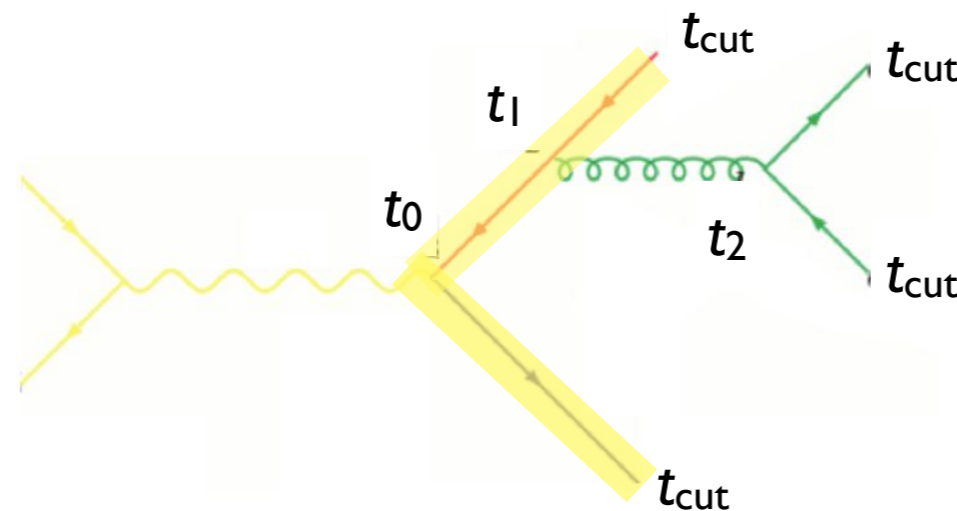


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and for the whole tree

$$(\Delta_q(t_{\text{cut}}, t_0))^2 \Delta_g(t_2, t_1) (\Delta_q(t_{\text{cut}}, t_2))^2 \frac{\alpha_s(t_1)}{2\pi} P_{gq}(z) \frac{\alpha_s(t_2)}{2\pi} P_{qg}(z')$$

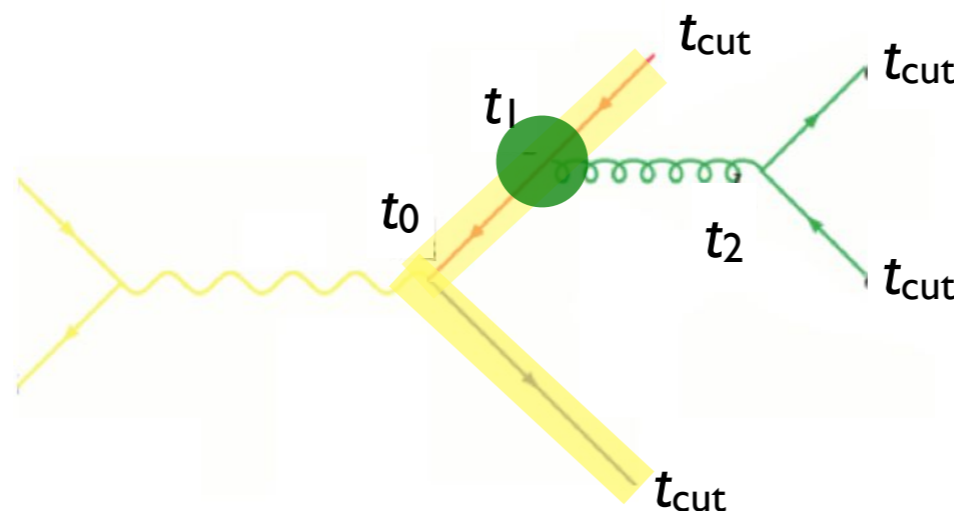


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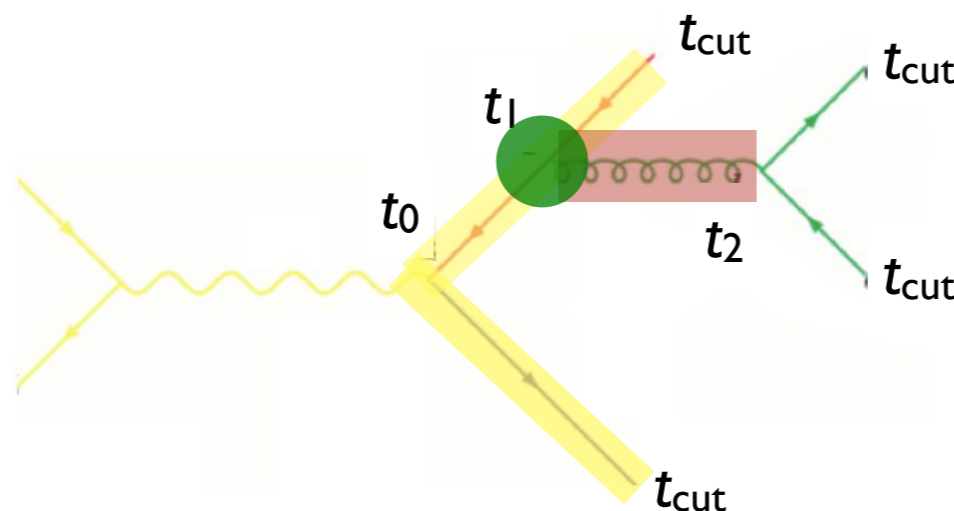


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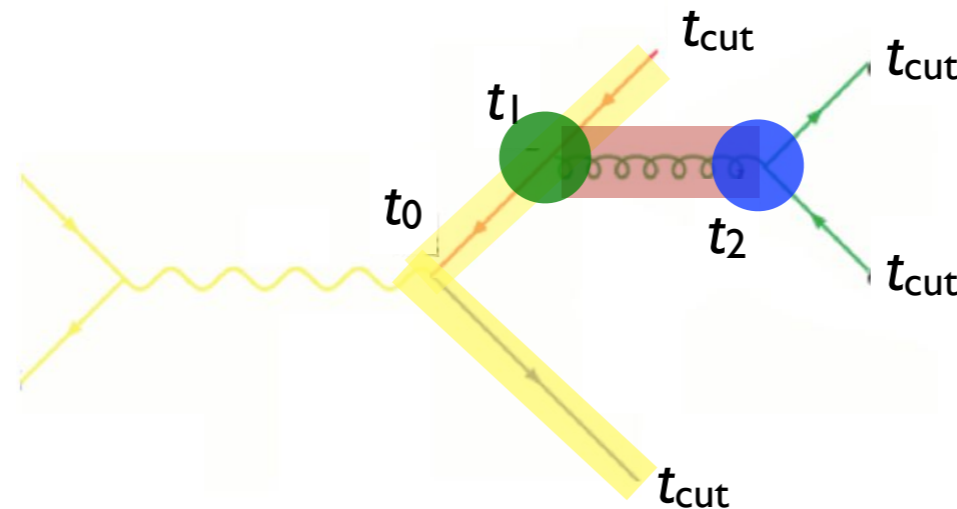


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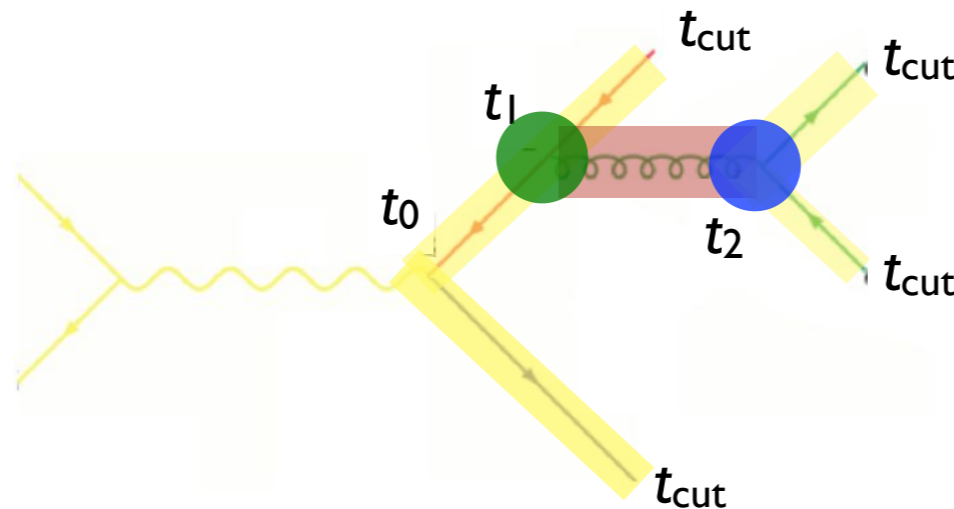


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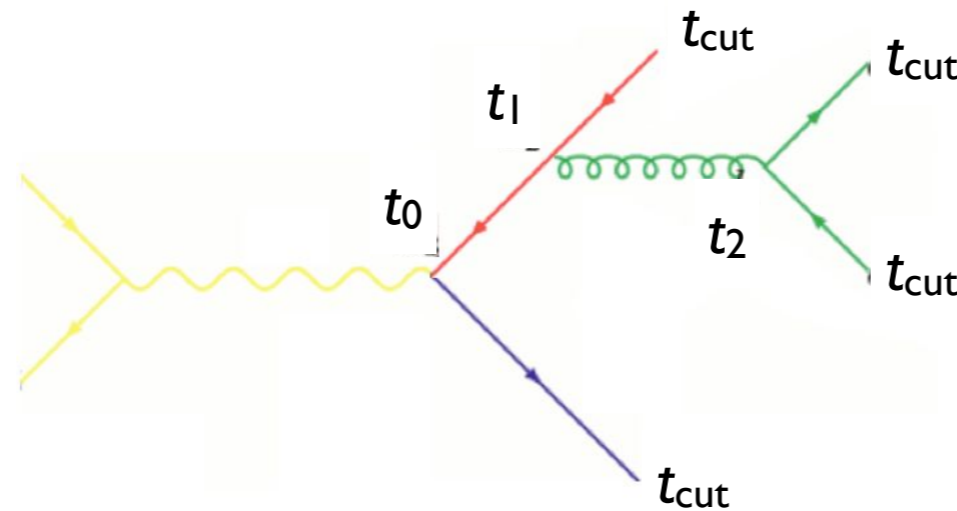


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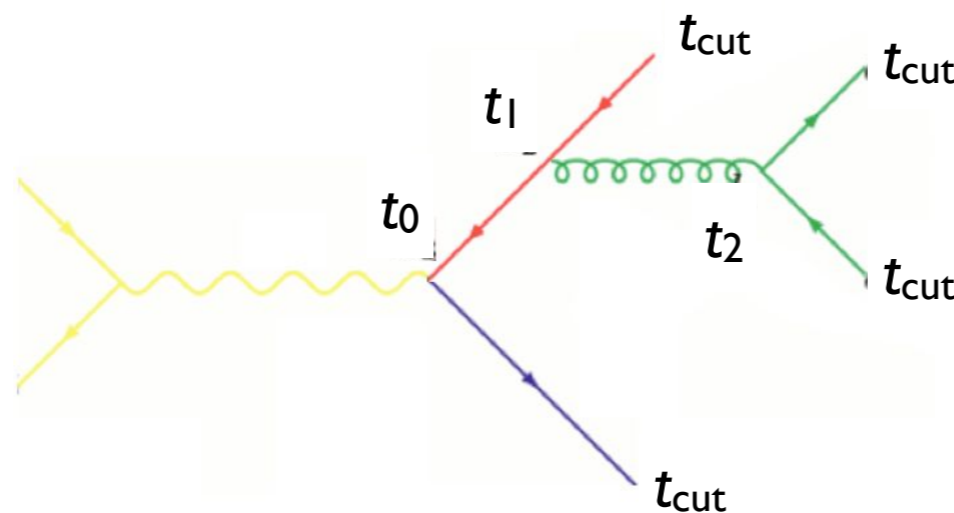
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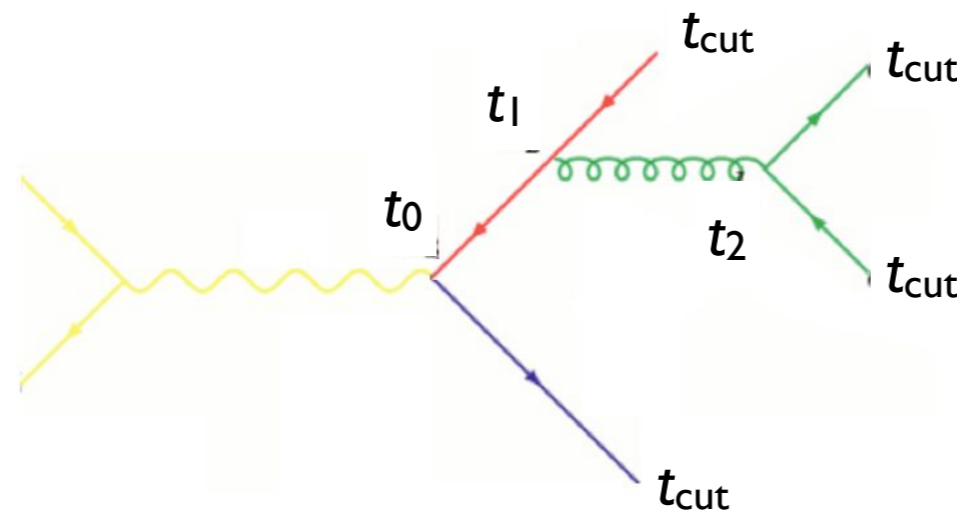


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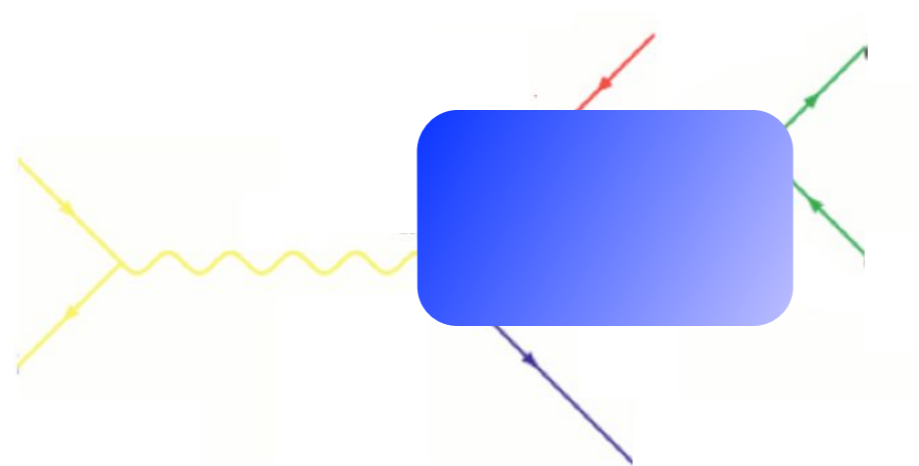
Corresponds to the matrix element
 BUT with α_s evaluated at the scale of each splitting



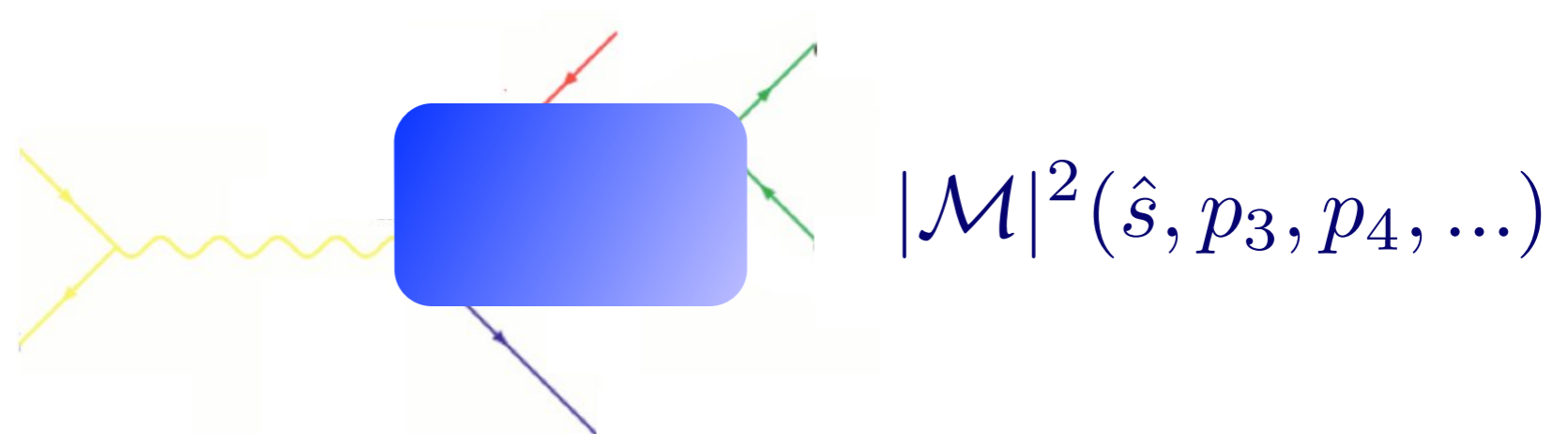
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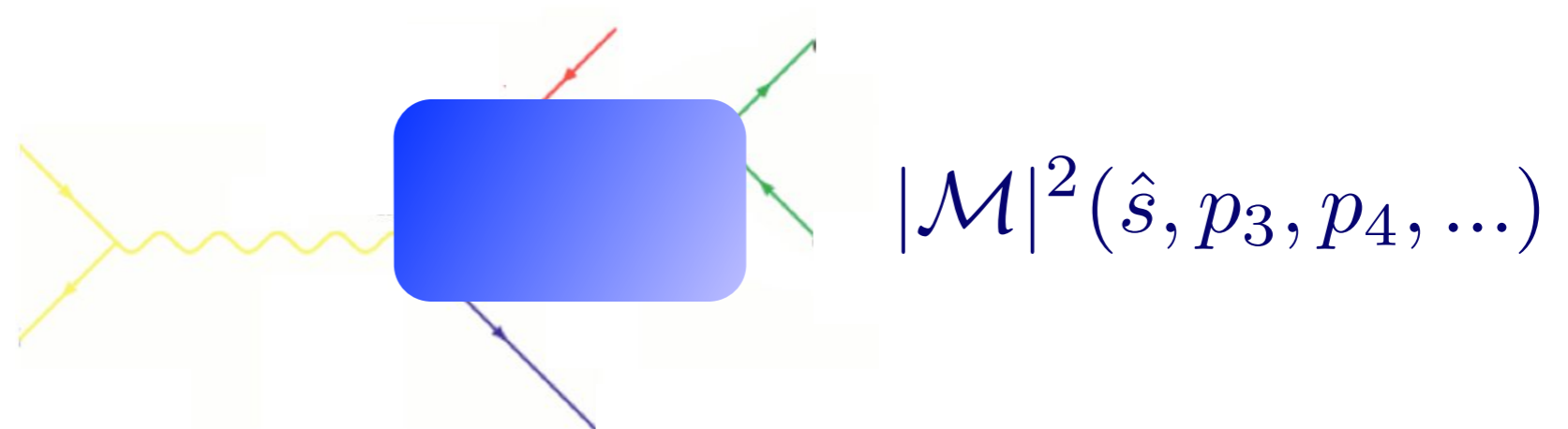
Sudakov suppression due to disallowing additional radiation
 above the scale t_{cut}



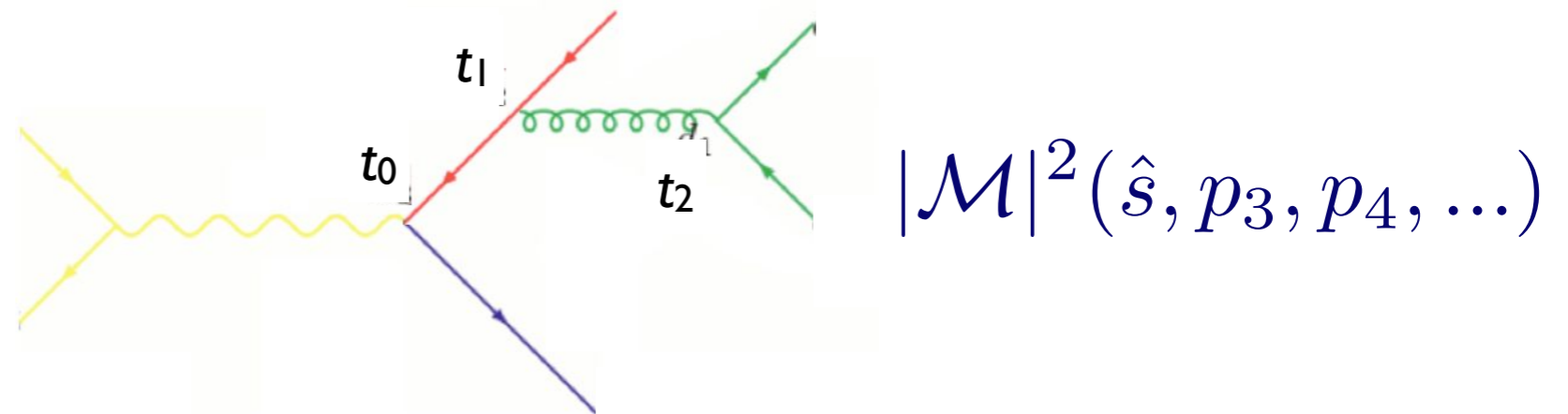
$$|\mathcal{M}|^2(\hat{s}, p_3, p_4, \dots)$$



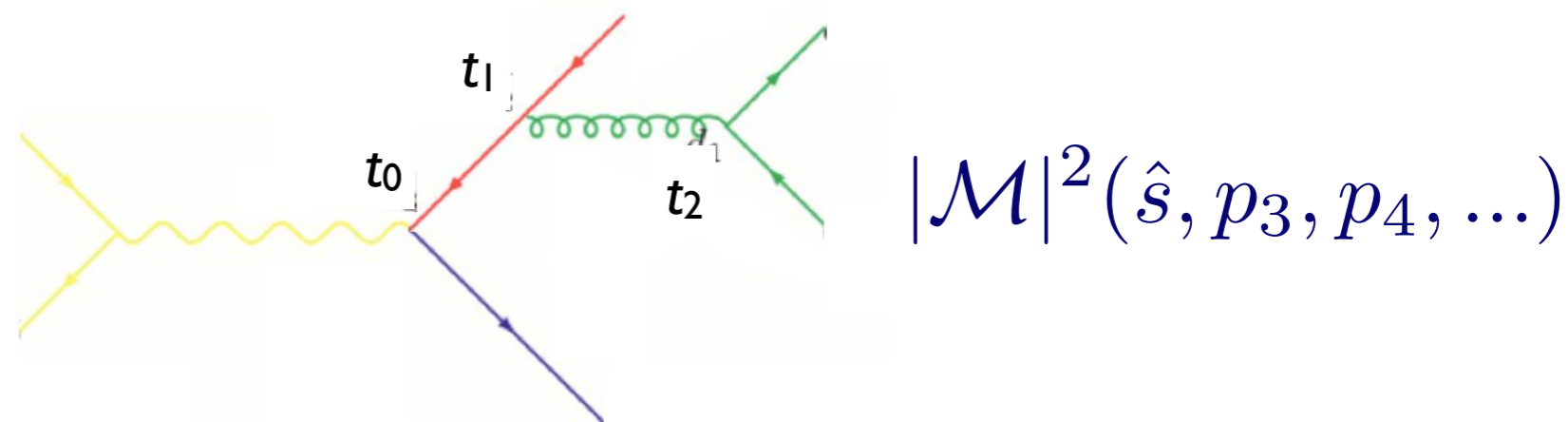
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 - I. Cluster the event using some clustering algorithm
 - this gives us a corresponding “parton shower history”

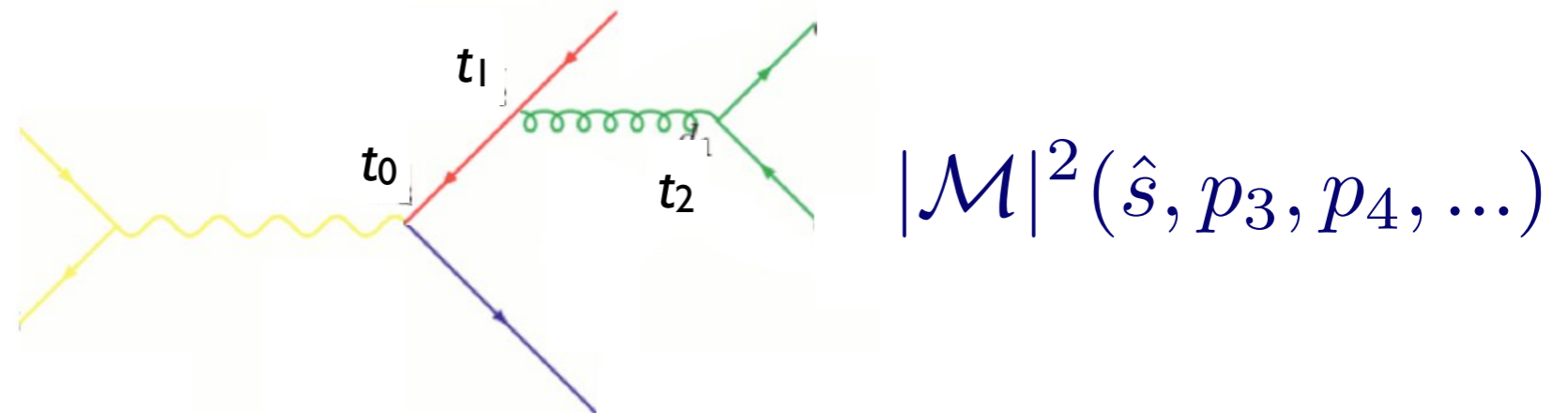


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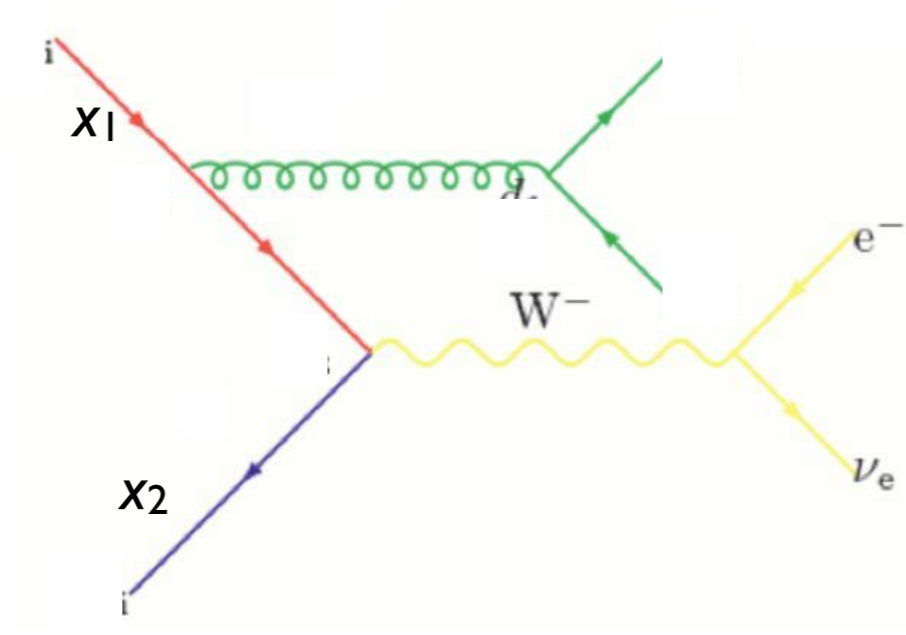


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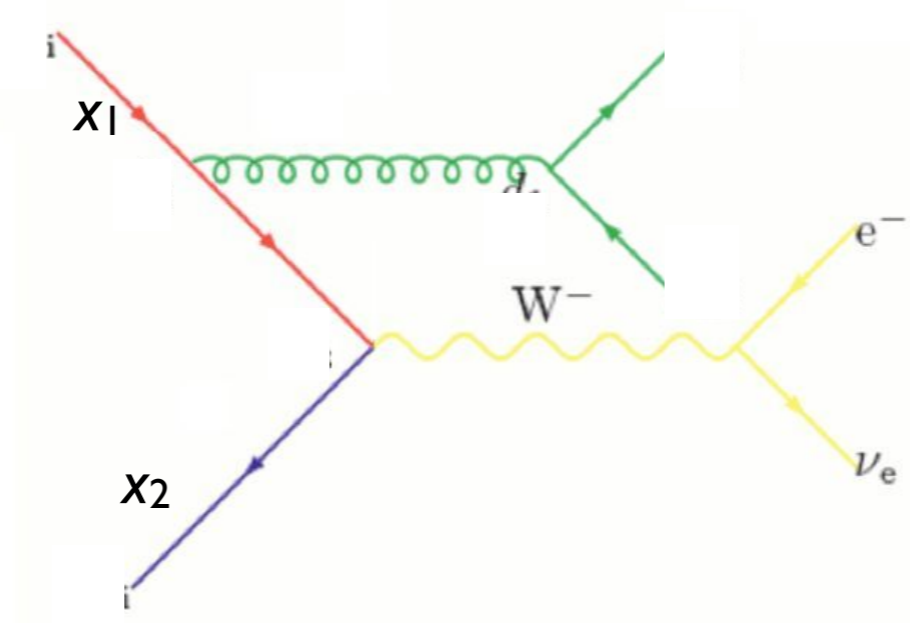
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3. Use some algorithm to apply the equivalent Sudakov suppression $(\Delta_q(t_{\text{cut}}, t_0))^2 \Delta_g(t_2, t_1) (\Delta_q(\text{cut}, t_2))^2$



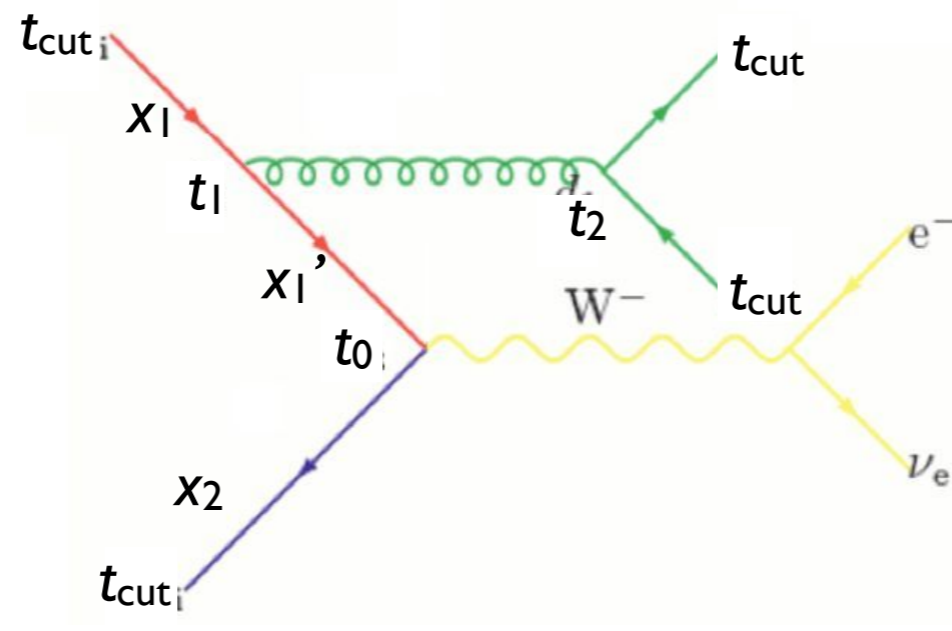
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- what happens for initial state radiation?



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- Let's do the same exercise as before:

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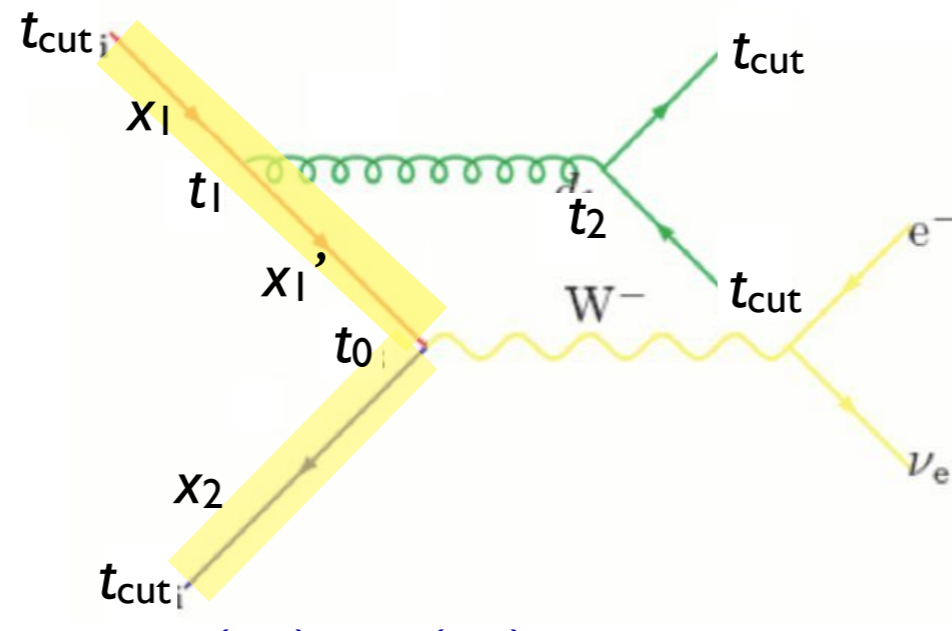
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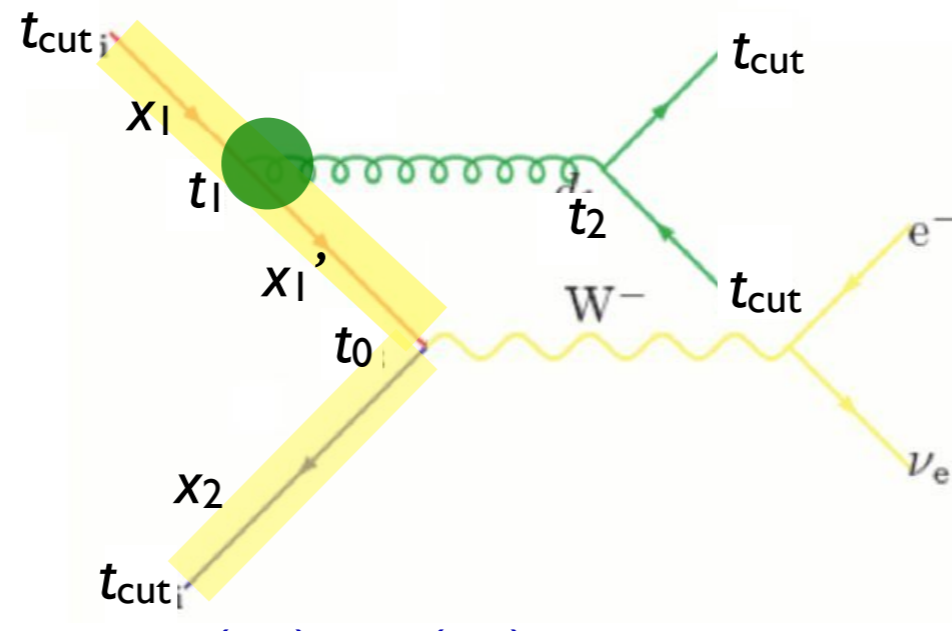
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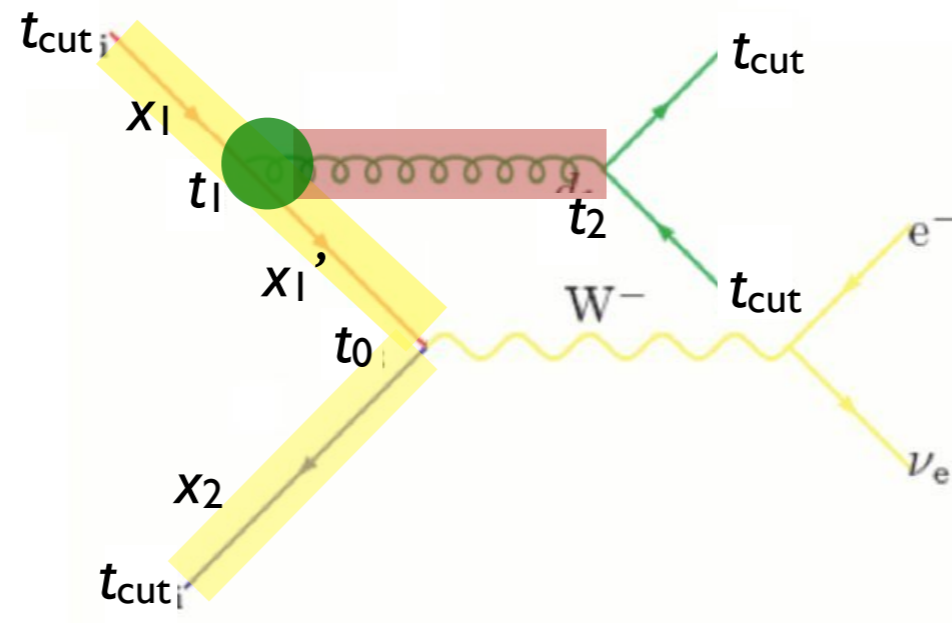
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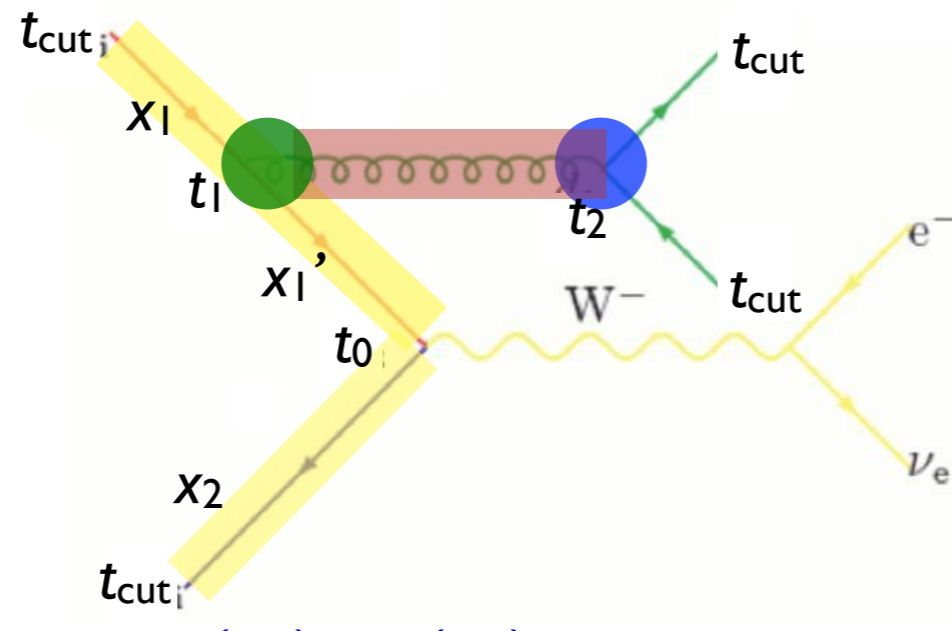
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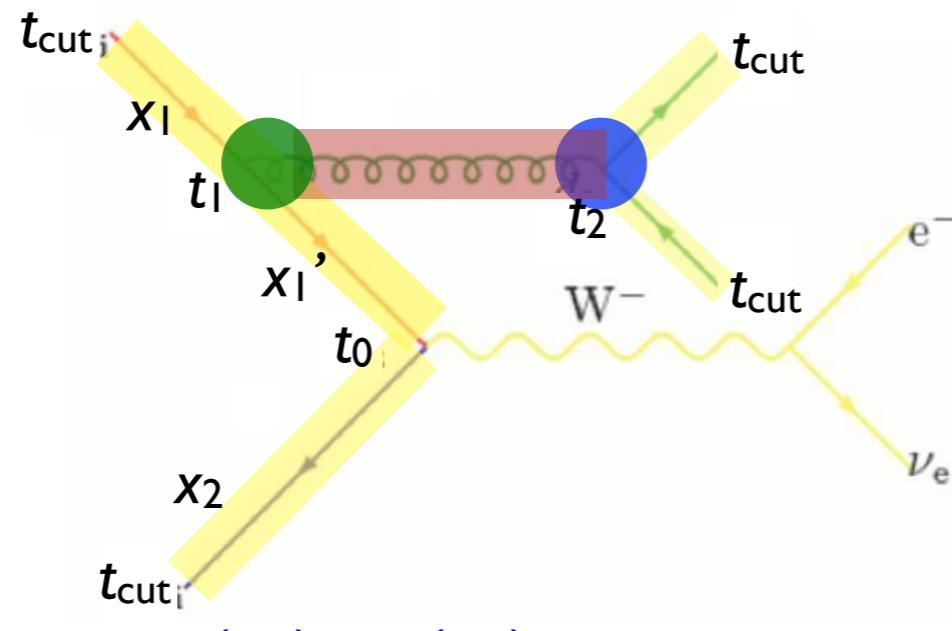
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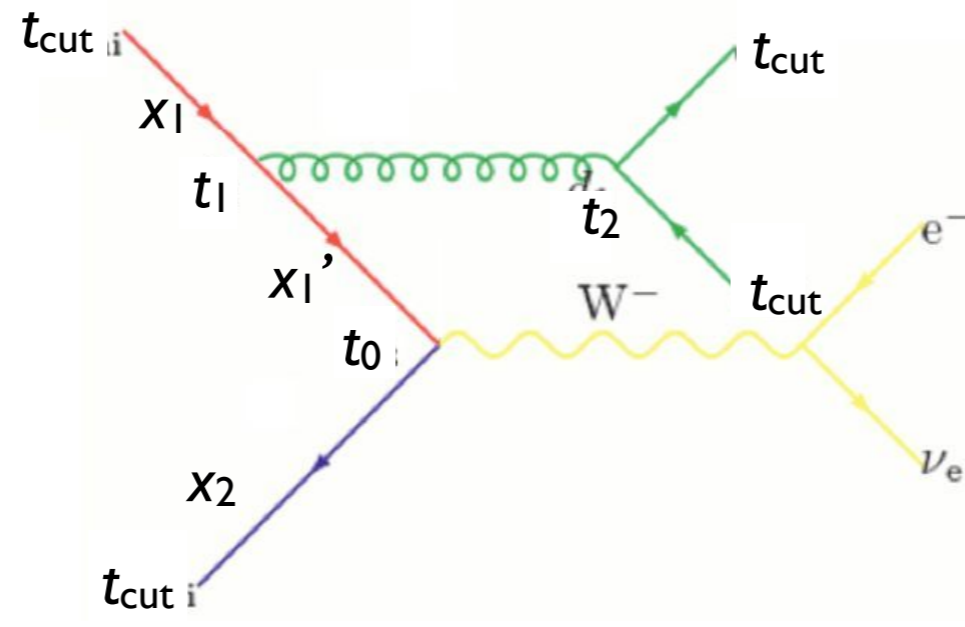


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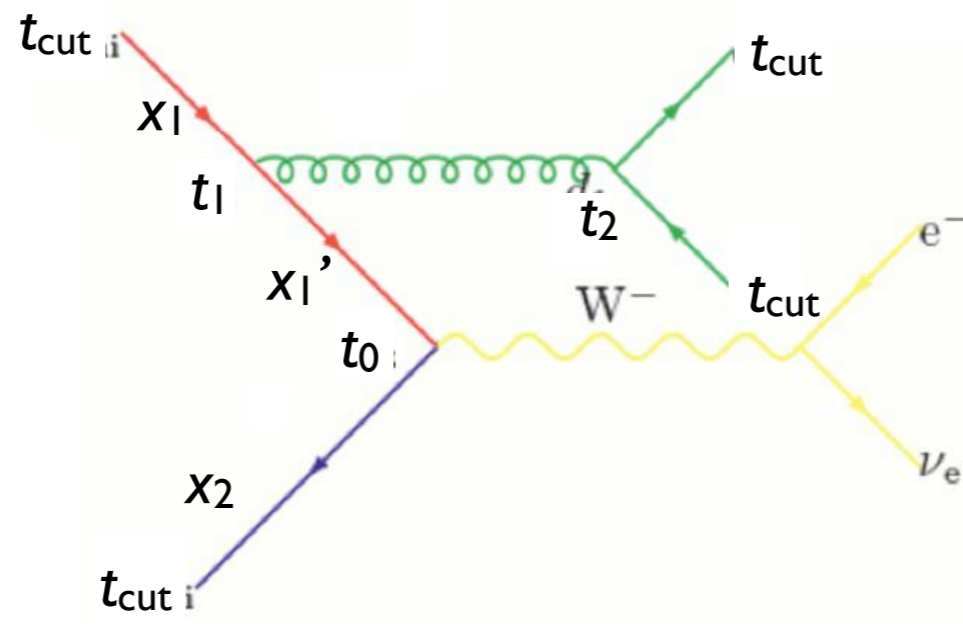


$$\begin{aligned}
 & (\Delta_{Iq}(t_{\text{cut}}, t_0))^2 \Delta_g(t_2, t_1) (\Delta_q(t_{\text{cut}}, t_2))^2 \frac{\alpha_s(t_1)}{2\pi} \frac{P_{gq}(z)}{z} \frac{f_q(x_1, t_1)}{f_q(x'_1, t_1)} \frac{\alpha_s(t_2)}{2\pi} P_{qg}(z') \\
 & \quad \times \hat{\sigma}_{q\bar{q} \rightarrow e\nu}(\hat{s}, \dots) f_q(x'_1, t_0) f_{\bar{q}}(x_2, t_0)
 \end{aligned}$$



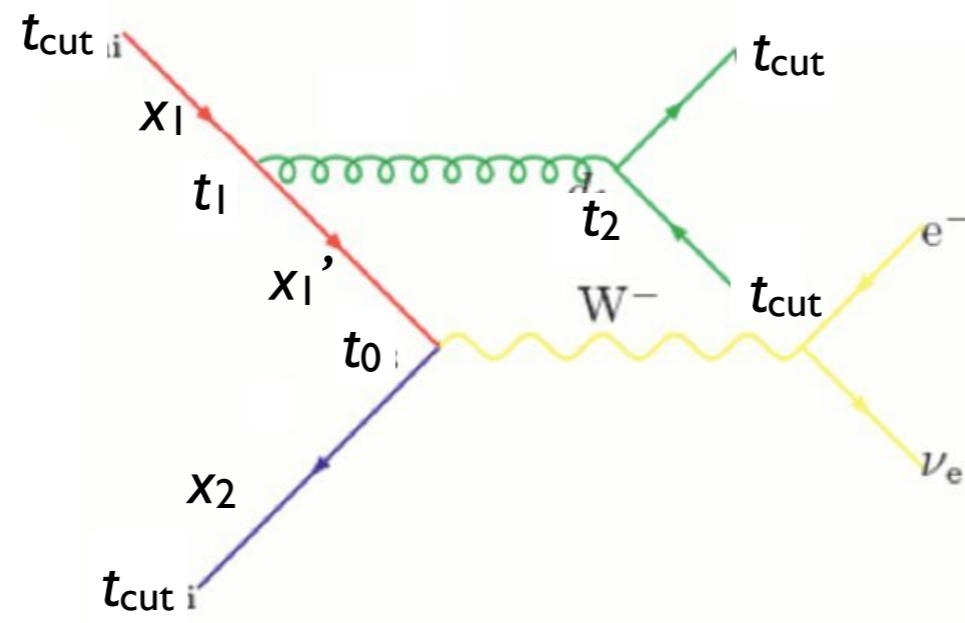
$$(\Delta_{Iq}(t_{\text{cut}}, t_0))^2 \Delta_g(t_2, t_1) (\Delta_q(t_{\text{cut}}, t_2))^2 \frac{\alpha_s(t_1)}{2\pi} \frac{P_{gq}(z)}{z} \frac{f_q(x_1, t_1)}{f_q(x'_1, t_1)} \frac{\alpha_s(t_2)}{2\pi} P_{qg}(z') \\
 \times \hat{\sigma}_{q\bar{q} \rightarrow e\nu}(\hat{s}, \dots) f_q(x'_1, t_0) f_{\bar{q}}(x_2, t_0)$$

ME with α_s evaluated at the scale of each splitting



$$(\Delta_{Iq}(t_{\text{cut}}, t_0))^2 \Delta_g(t_2, t_1) (\Delta_q(t_{\text{cut}}, t_2))^2 \frac{\alpha_s(t_1)}{2\pi} \frac{P_{gq}(z)}{z} \frac{f_q(x_1, t_1)}{f_q(x'_1, t_1)} \frac{\alpha_s(t_2)}{2\pi} P_{qg}(z') \\
 \times \hat{\sigma}_{q\bar{q} \rightarrow e\nu}(\hat{s}, \dots) f_q(x'_1, t_0) f_{\bar{q}}(x_2, t_0)$$

ME with α_s evaluated at the scale of each splitting
 PDF reweighting

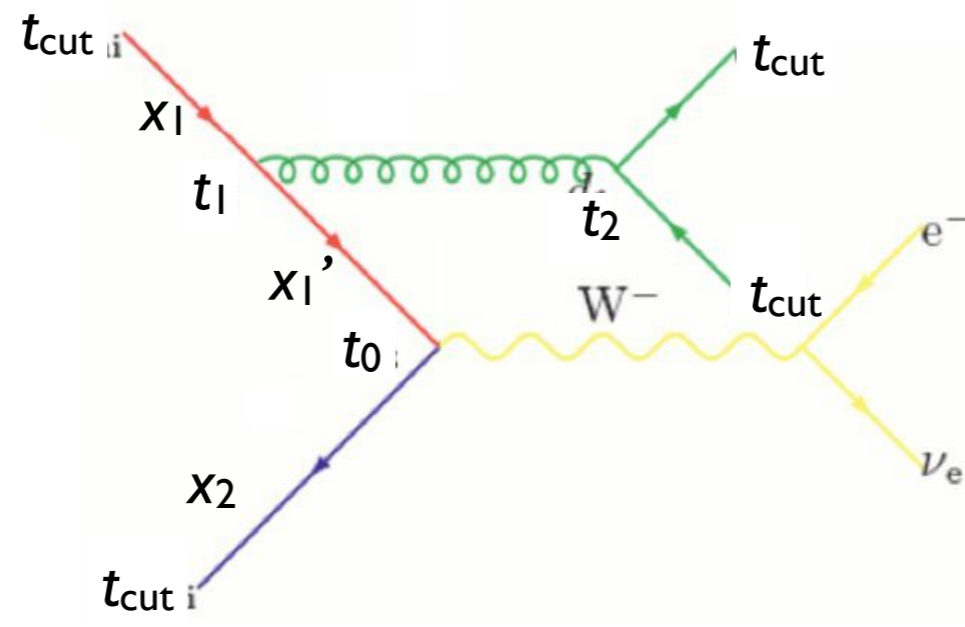


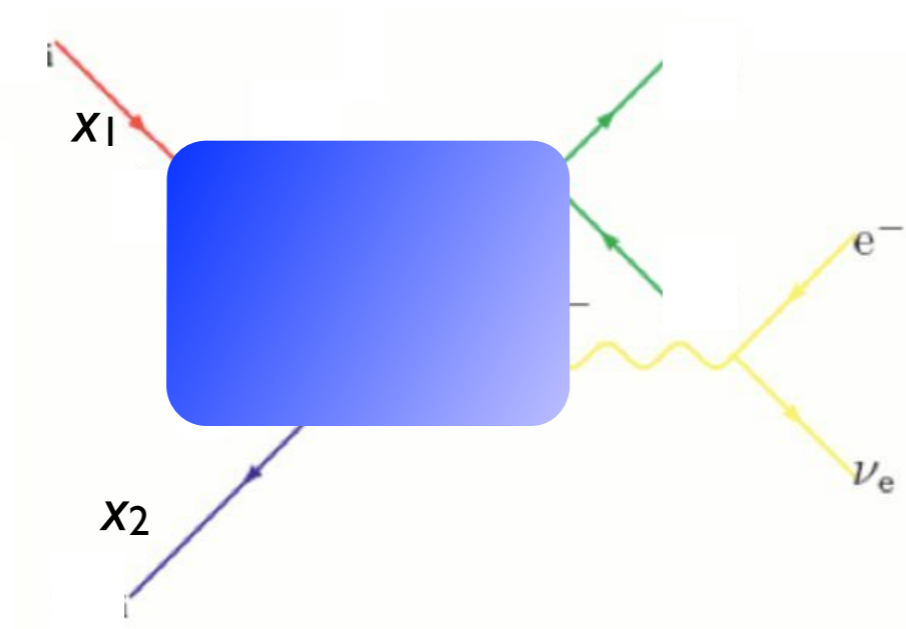
$$(\Delta_{Iq}(t_{\text{cut}}, t_0))^2 \Delta_g(t_2, t_1) (\Delta_q(t_{\text{cut}}, t_2))^2 \frac{\alpha_s(t_1)}{2\pi} \frac{P_{gq}(z)}{z} \frac{f_q(x_1, t_1)}{f_q(x'_1, t_1)} \frac{\alpha_s(t_2)}{2\pi} P_{qg}(z') \\
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ME with α_s evaluated at the scale of each splitting

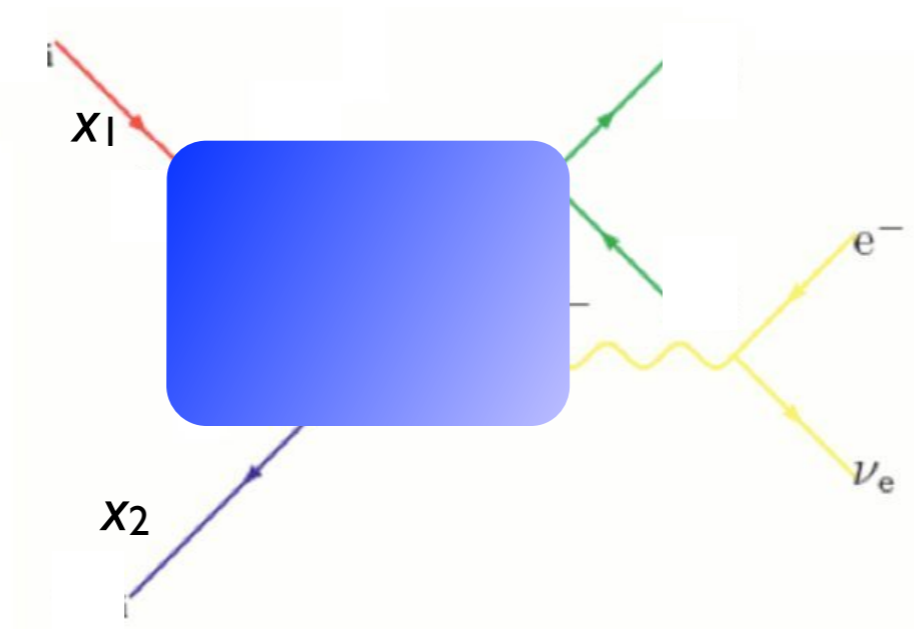
PDF reweighting

Sudakov suppression due to non-branching above scale t_{cut}

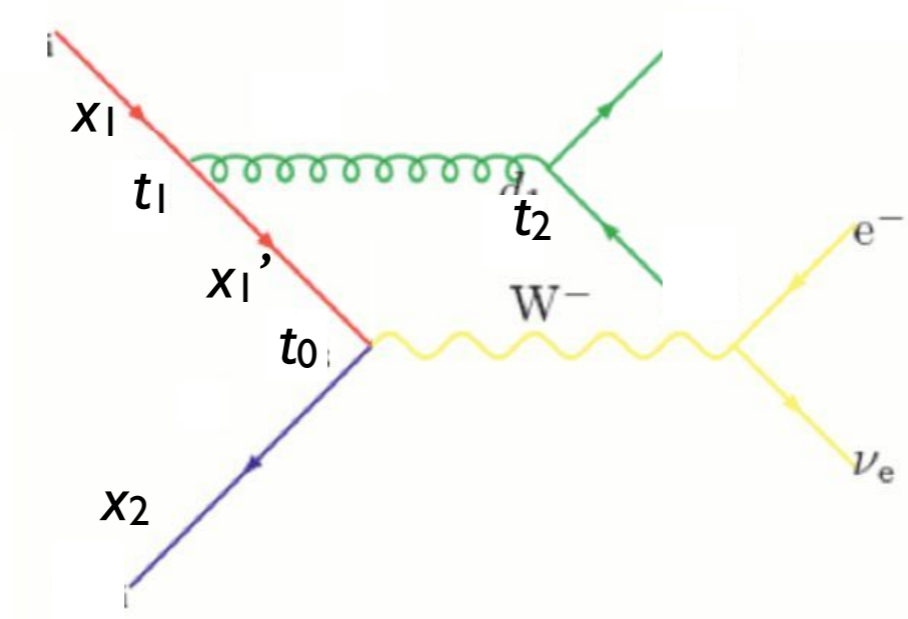




- Again, use a clustering scheme to get a parton shower history

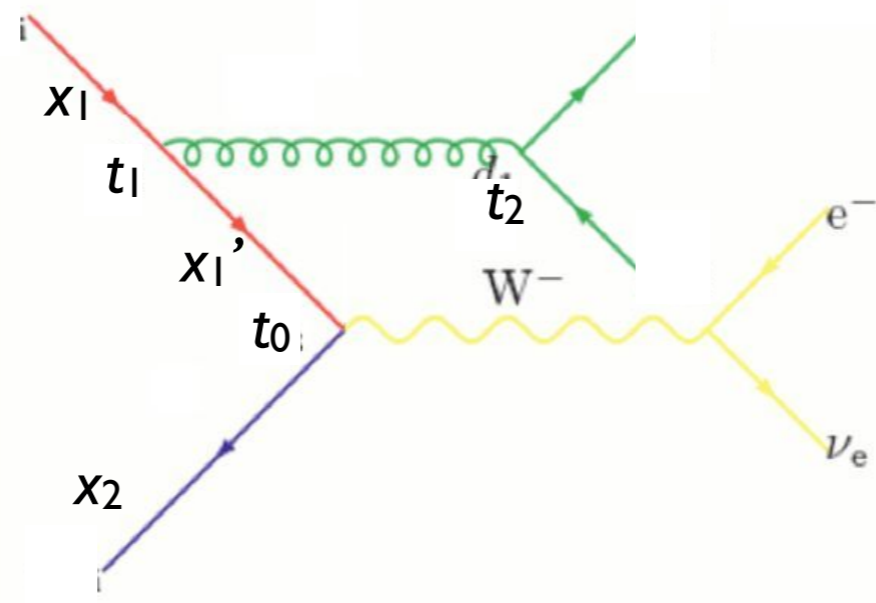


- Again, use a clustering scheme to get a parton shower history



- Again, use a clustering scheme to get a parton shower history
- Now, reweight both due to α_s and PDF

$$|\mathcal{M}|^2 \rightarrow |\mathcal{M}|^2 \frac{\alpha_s(t_1)}{\alpha_s(t_0)} \frac{\alpha_s(t_2)}{\alpha_s(t_0)} \frac{f_q(x'_1, t_0)}{f_q(x'_1, t_1)}$$

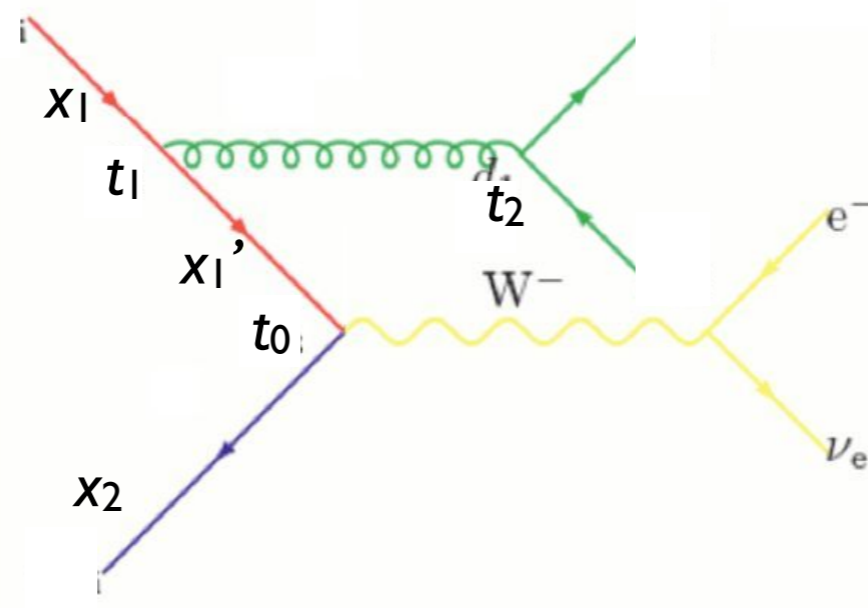


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- Remember to use first clustering scale on each side for PDF scale:

$$\mathcal{P}_{\text{event}} = \hat{\sigma}(x_1, x_2, p_3, p_4, \dots) f_q(x_1, t_1) f_{\bar{q}}(x_2, t_0)$$



- We still haven't specified how to apply the Sudakov reweighting to the matrix element
- Three general schemes available in the literature:
 - ➔ CKKW scheme [Catani,Krauss,Kuhn,Webber 2001; Krauss 2002]
 - ➔ Lönnblad scheme (or CKKW-L) [Lönnblad 2002]
 - ➔ MLM scheme [Mangano *unpublished* 2002; Mangano et al. 2007]

[Catani, Krauss, Kuhn, Webber 2001]
[Krauss 2002]

[Catani, Krauss, Kuhn, Webber 2001]
[Krauss 2002]

- Apply the required Sudakov suppression

$$(\Delta_{Iq}(t_{\text{cut}}, t_0))^2 \Delta_g(t_2, t_1) (\Delta_q(t_{\text{cut}}, t_2))^2$$

analytically, using the best available (NLL) Sudakovs.

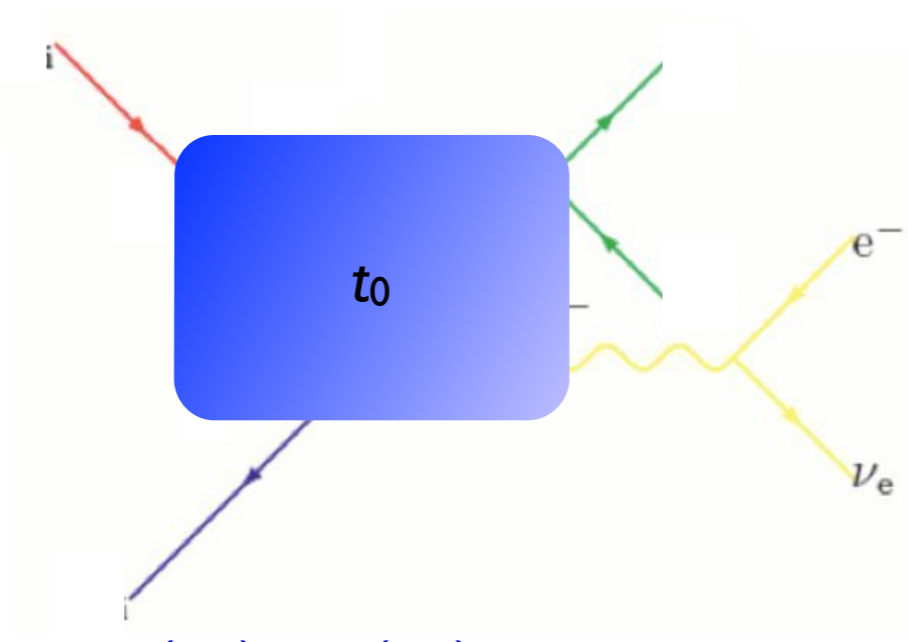
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analytically, using the best available (NLL) Sudakovs.

- Perform “truncated showering”: Run the parton shower starting at t_0 , but forbid any showers above the cutoff scale t_{cut} .



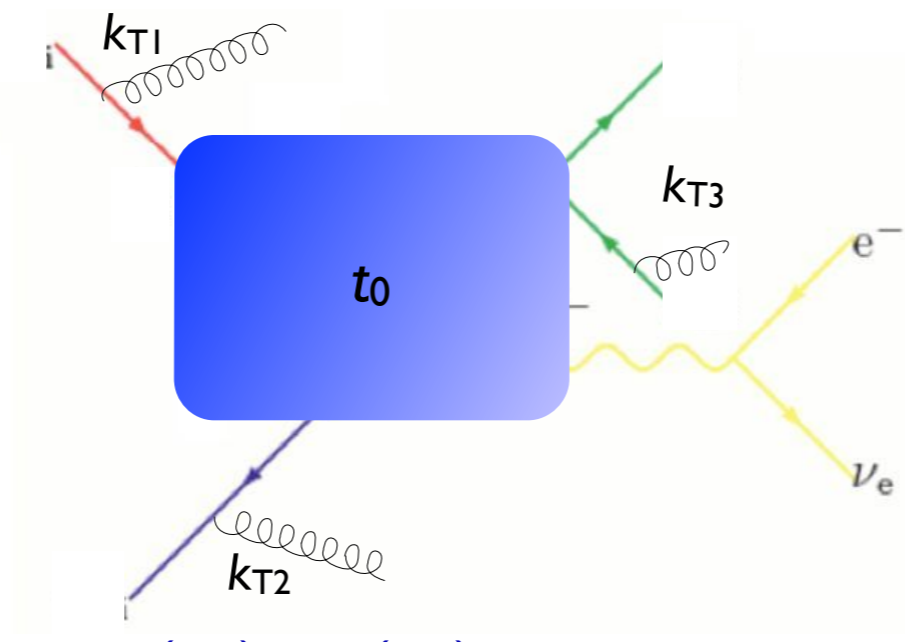
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- Apply the required Sudakov suppression

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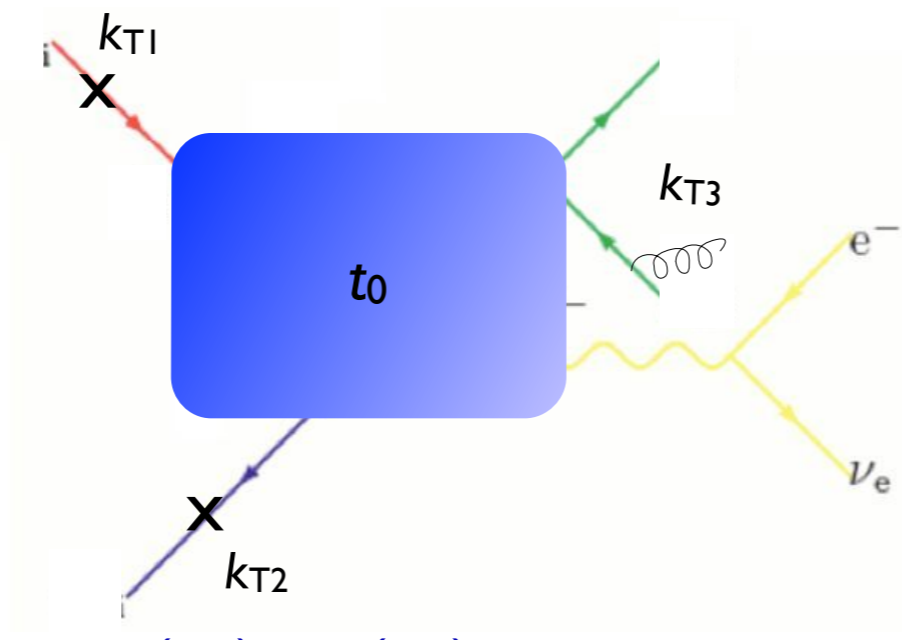
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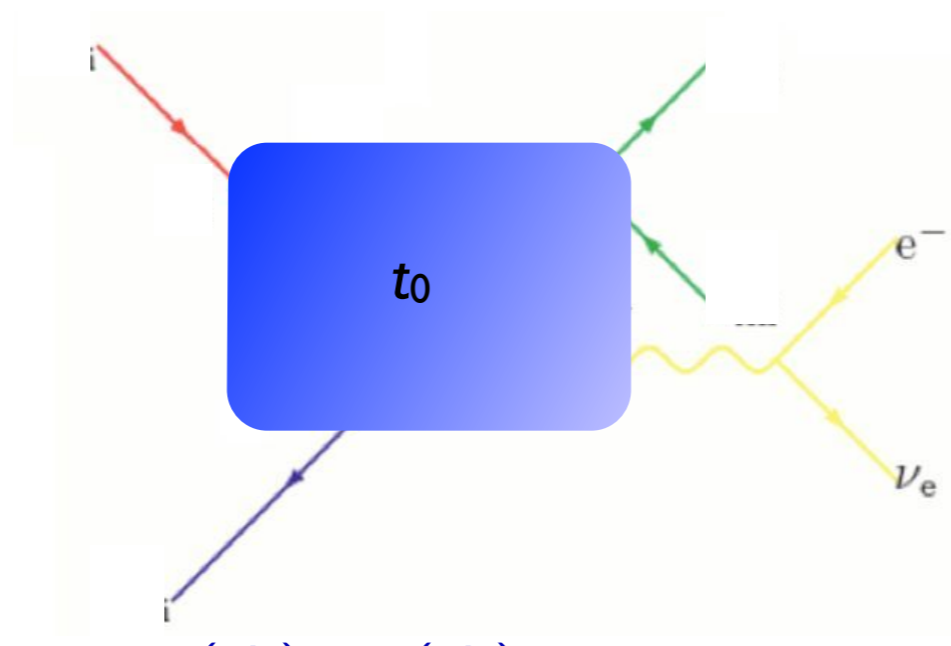
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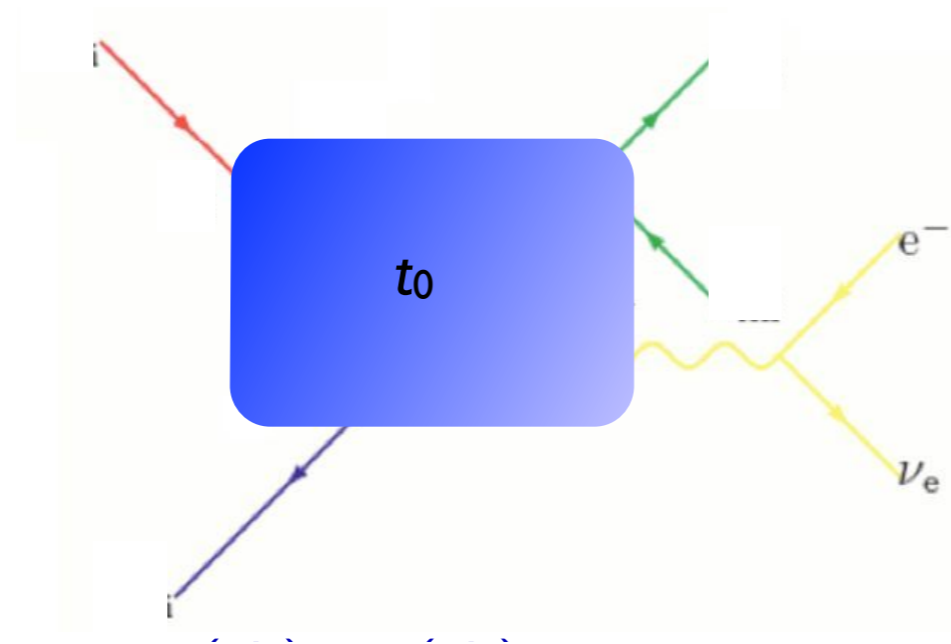
analytically, using the best available (NLL) Sudakovs.

- Perform “truncated showering”: Run the parton shower starting at t_0 , but forbid any showers above the cutoff scale t_{cut} .
- ✓ Best theoretical treatment of matrix element
 - Requires dedicated PS implementation
 - Mismatch between analytical Sudakov and (non-NLL) shower
- Implemented in Sherpa (v. 1.1)

[Lönnblad 2002]
[Hoeche et al. 2009]

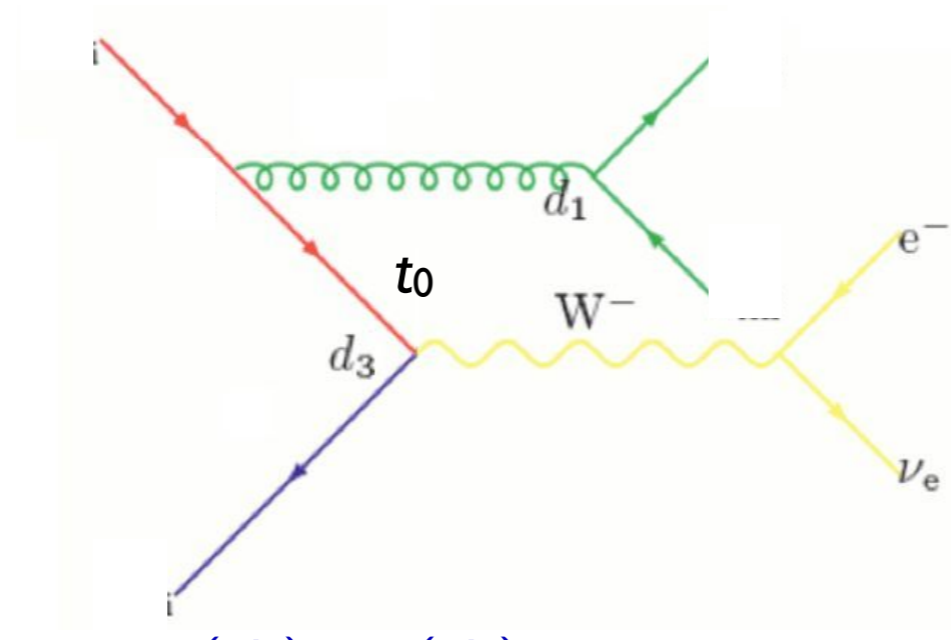


[Lönblad 2002]
[Hoeche et al. 2009]



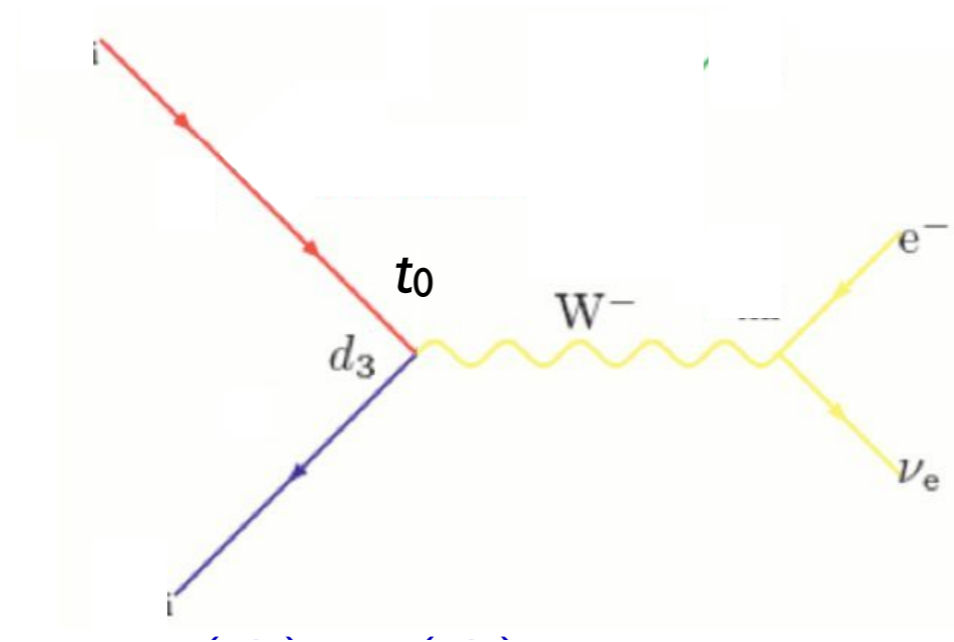
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[Hoeche et al. 2009]



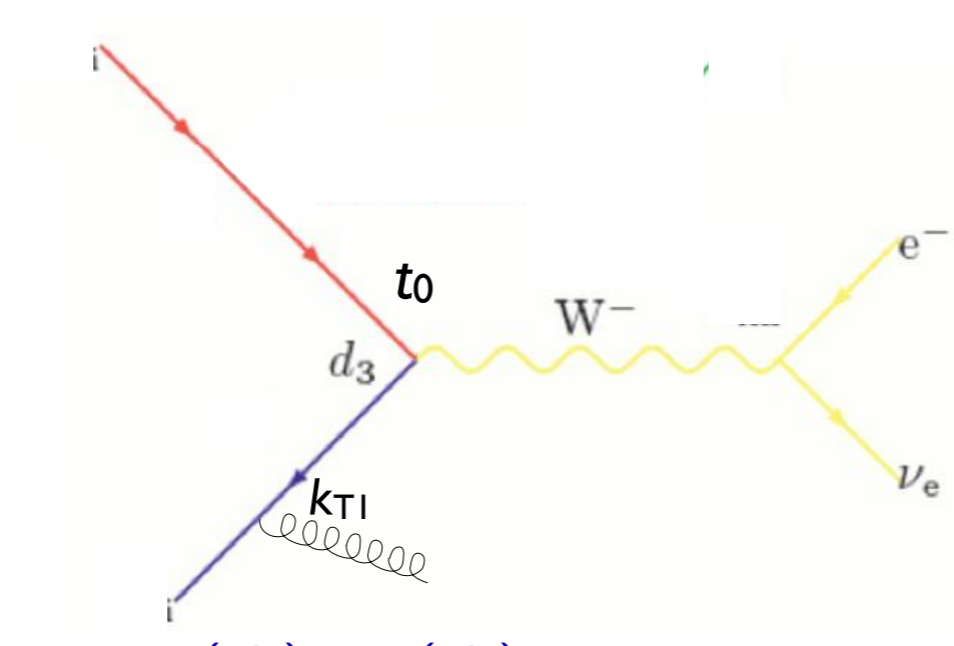
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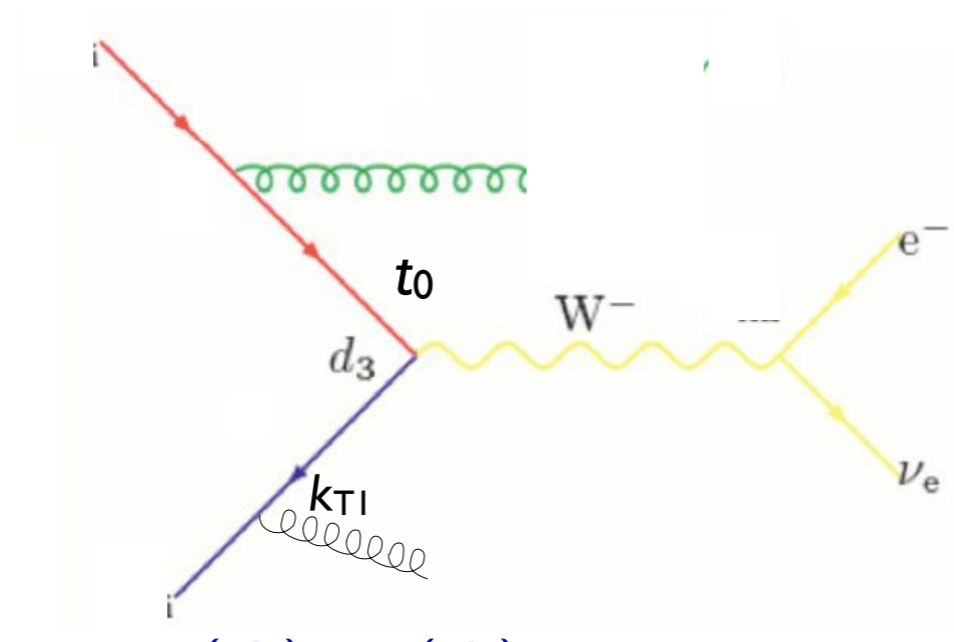
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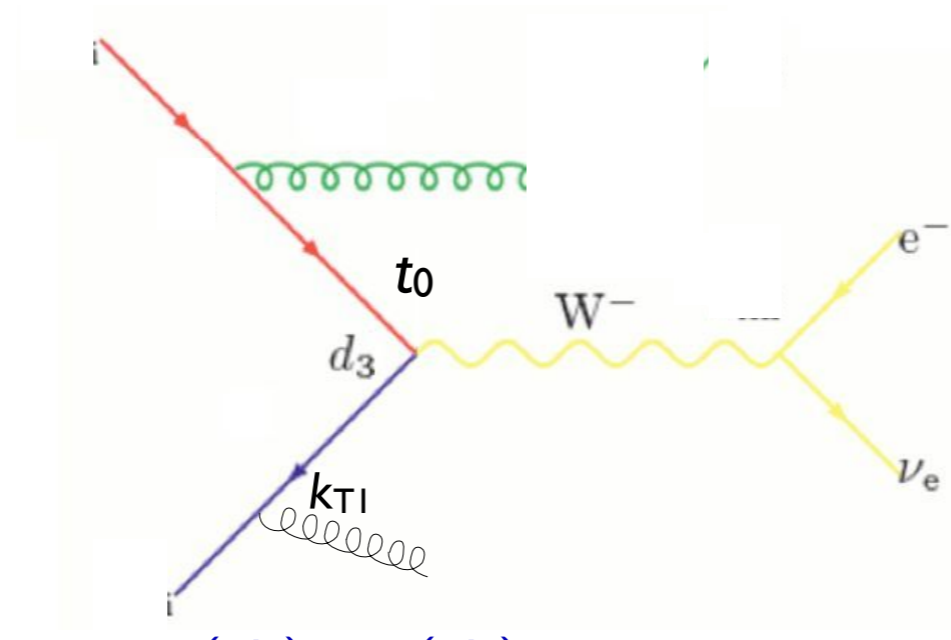
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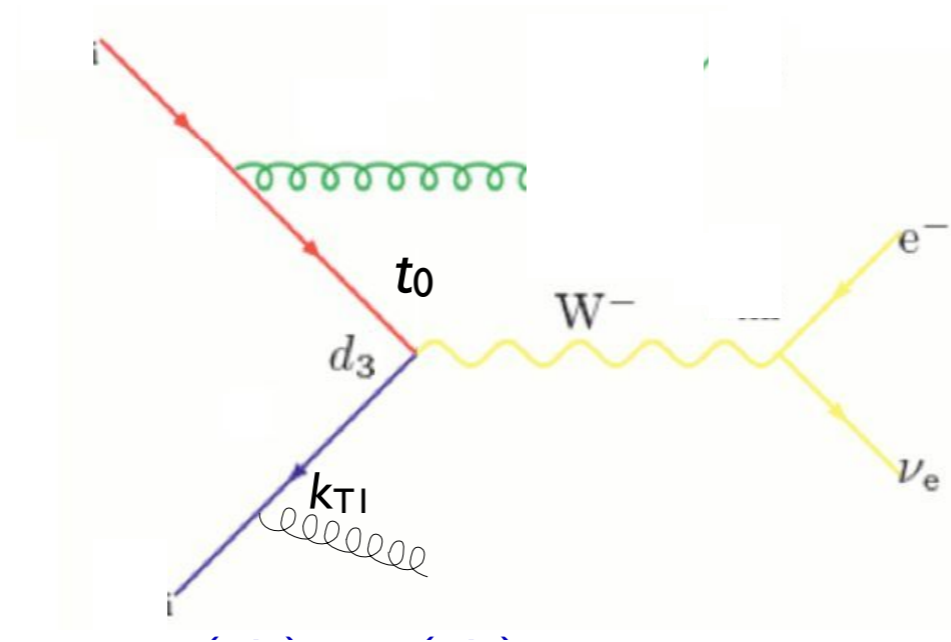
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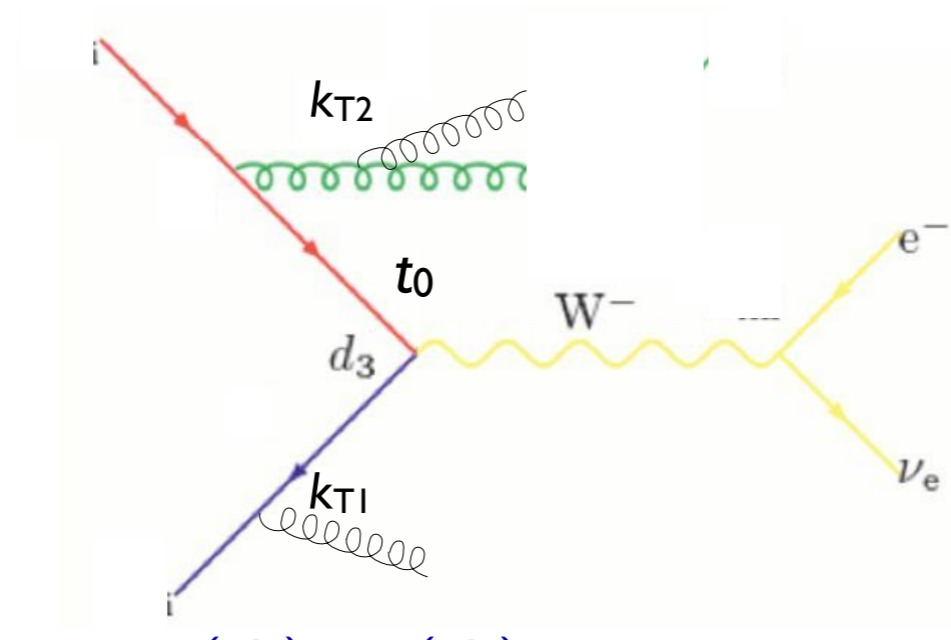
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[Lönblad 2002]
[Hoeche et al. 2009]



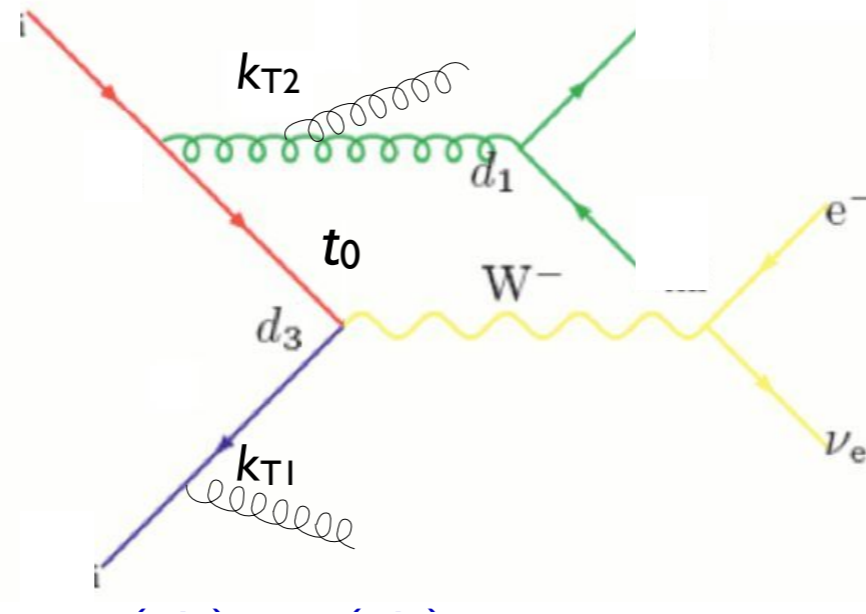
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[Lönblad 2002]
[Hoeche et al. 2009]



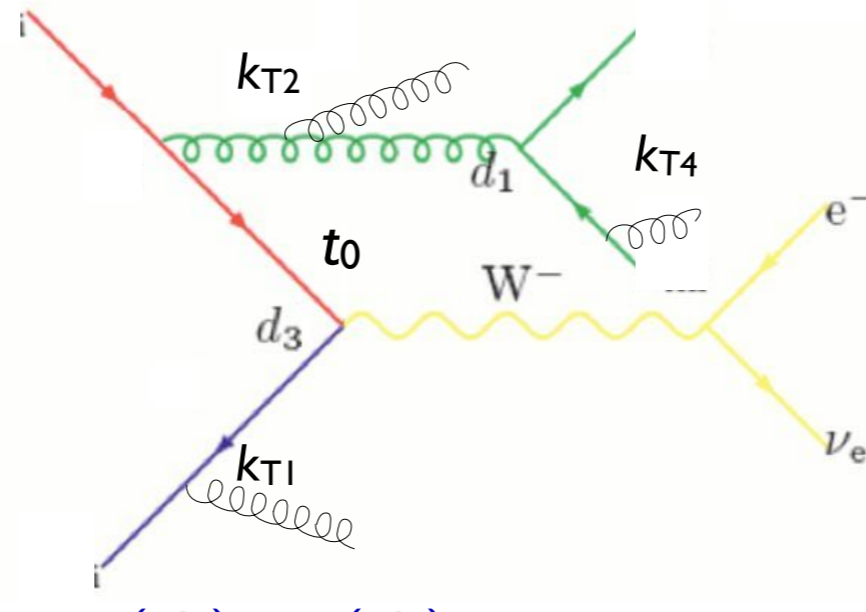
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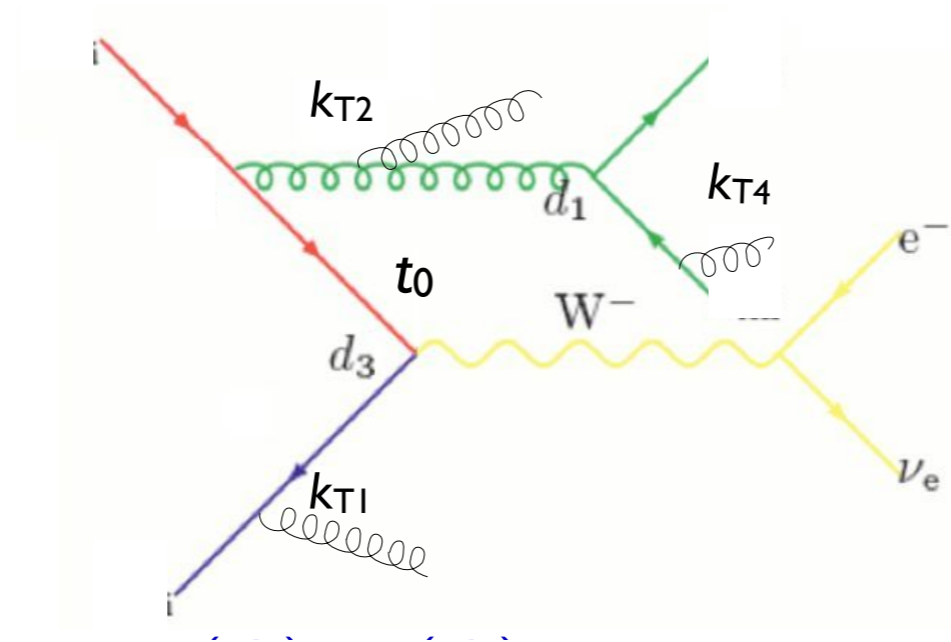
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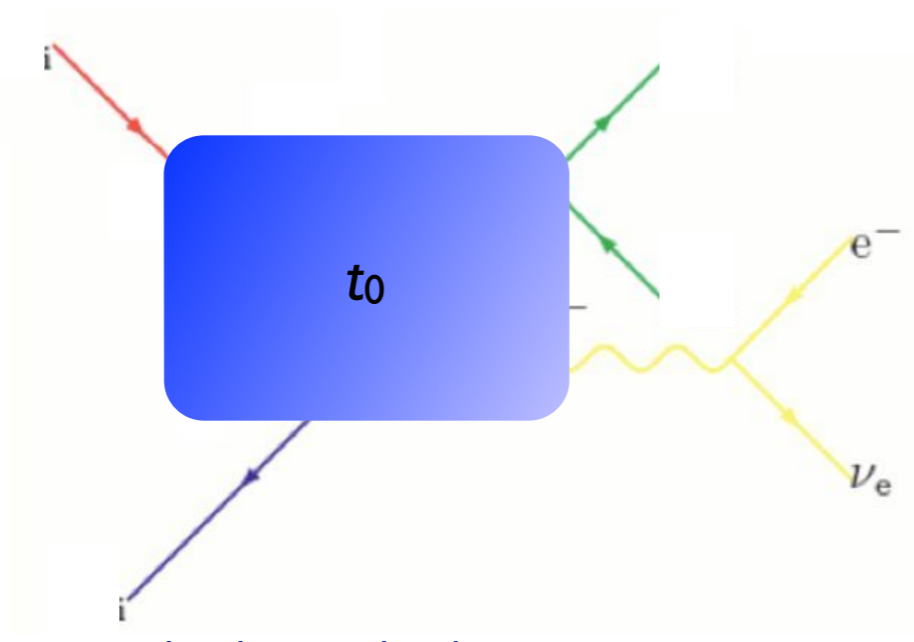
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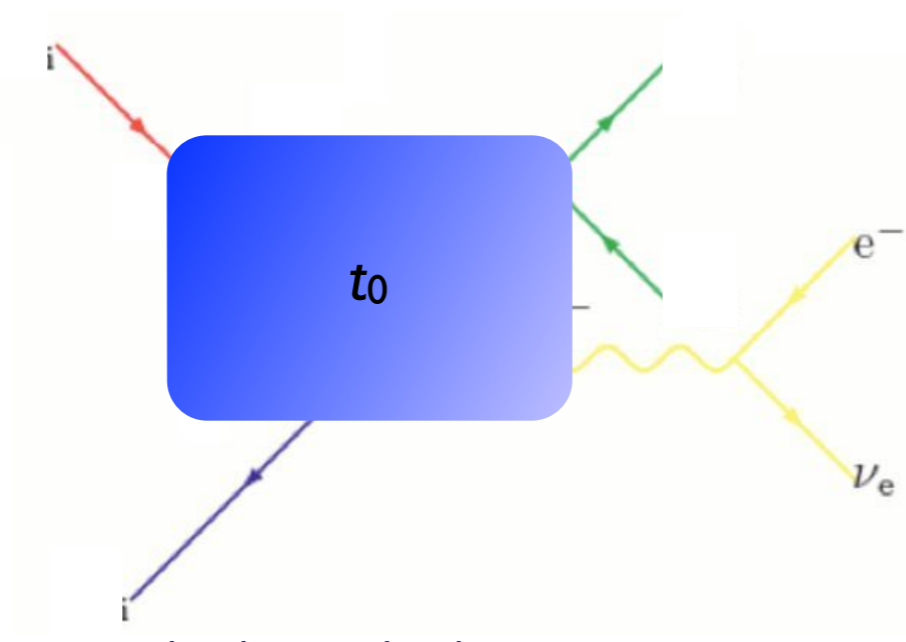
- ✓ Automatic agreement between Sudakov and shower
- Requires dedicated PS implementation
 - ➔ Need multiple implementations to compare between showers
- Implemented in Ariadne, Sherpa (v. 1.2), and Pythia 8

[M.L. Mangano, ~2002, 2007]
[J.A. et al 2007, 2008]



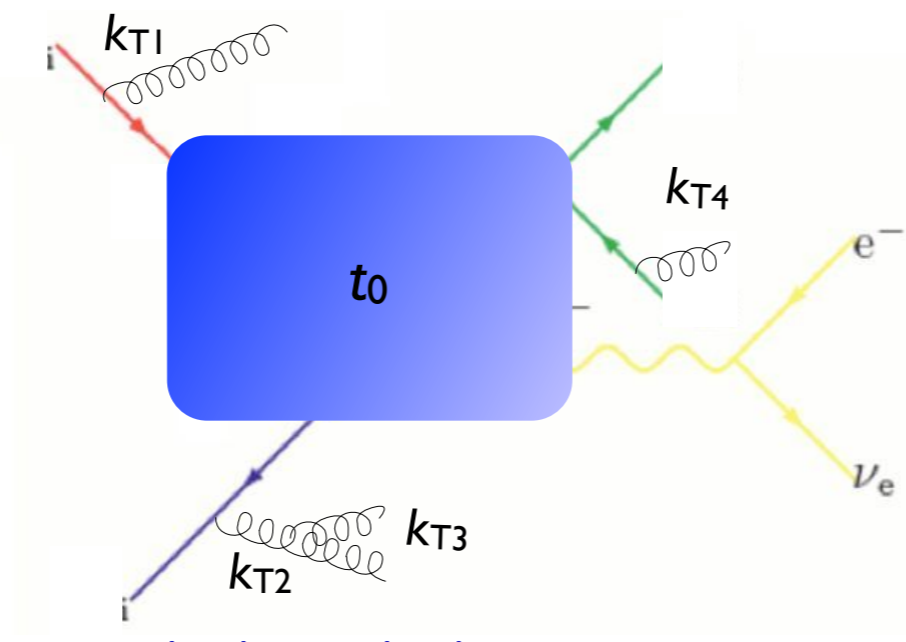
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- The simplest way to do the Sudakov suppression is to run the shower on the event, starting from t_0 !



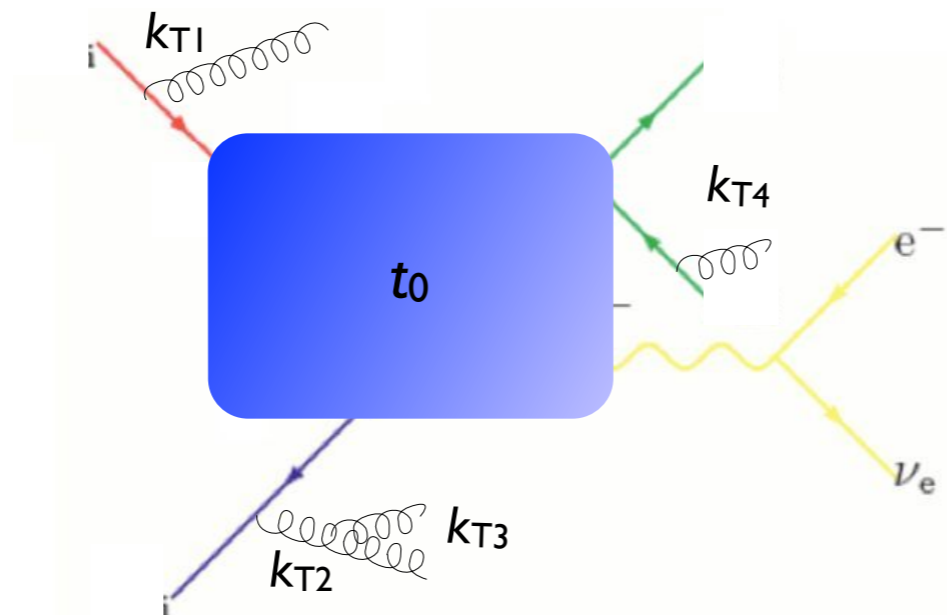
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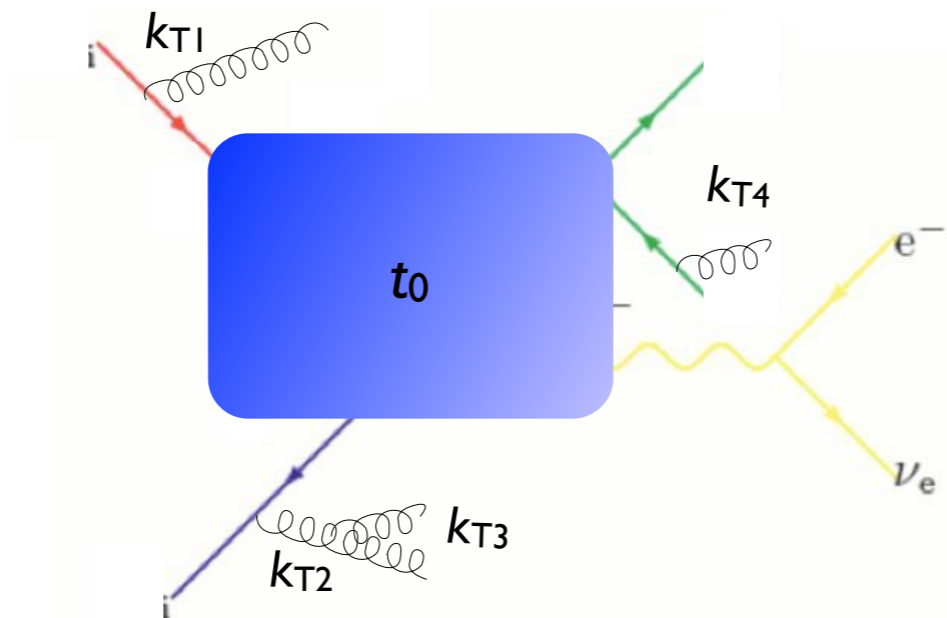
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[M.L. Mangano, ~2002, 2007]
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- The simplest way to do the Sudakov suppression is to run the shower on the event, starting from t_0 !



- Perform jet clustering after PS - if hardest jet $k_{T1} > t_{cut}$ or there are jets not matched to partons, reject the event
- The resulting Sudakov suppression from the procedure is

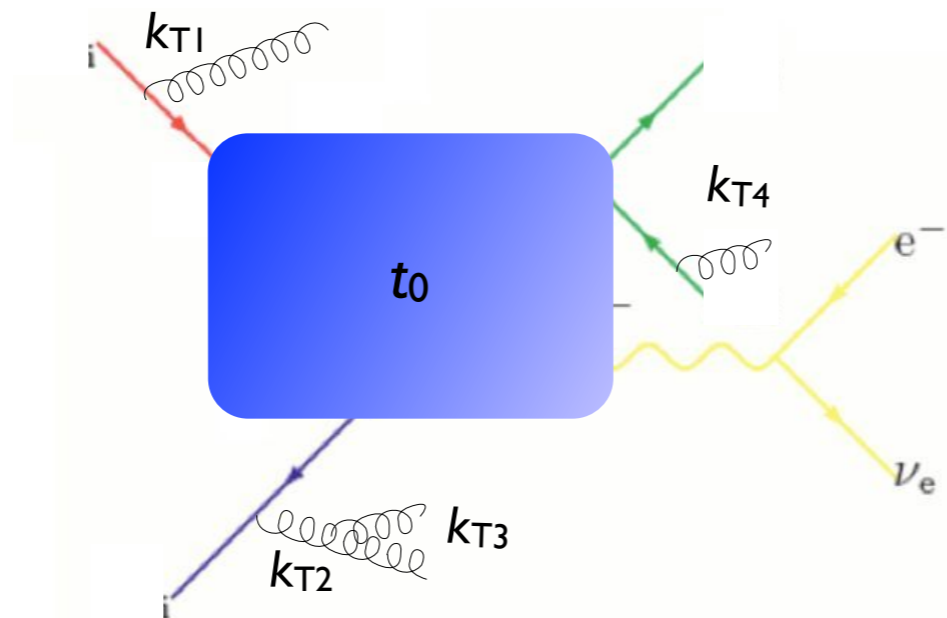
$$(\Delta_{Iq}(t_{cut}, t_0))^2 (\Delta_q(t_{cut}, t_0))^2$$

which turns out to be a good enough approximation of the correct expression

$$(\Delta_{Iq}(t_{cut}, t_0))^2 \Delta_g(t_2, t_1) (\Delta_q(t_{cut}, t_2))^2$$

[M.L. Mangano, ~2002, 2007]
[J.A. et al 2007, 2008]

- The simplest way to do the Sudakov suppression is to run the shower on the event, starting from t_0 !



- ✓ Simplest available scheme
- ✓ Allows matching with any shower, without modification
- ➔ Sudakov suppression not exact, minor mismatch with shower
- Implemented in AlpGen, HELAC, MadGraph+Pythia 6

- In the previous, assumed we can simulate all parton multiplicities by the ME
 - In practice, we can only do limited number of final-state partons with matrix element (up to 4-5 or so)
 - For the highest jet multiplicity that we generate with the matrix element, we need to allow additional jets above the matching scale t_{cut} , since we will otherwise not get a jet-inclusive description – but still can't allow PS radiation harder than the ME partons
- ➔ Need to replace t_{cut} by the clustering scale for the softest ME parton for the highest multiplicity

- We have a number of choices to make in the above procedure. The most important are:
 1. The clustering scheme used to determine the parton shower history of the ME event
 2. What to use for the scale Q^2 (factorization scale)
 3. How to divide the phase space between parton showers and matrix elements

1. The clustering scheme used inside MadGraph and Sherpa to determine the parton shower history is the Durham k_T scheme. For e^+e^- :

$$k_{Tij}^2 = 2 \min(E_i^2, E_j^2)(1 - \cos \theta_{ij})$$

and for hadron collisions, the minimum of:

$$k_{Tibeam} = m_i^2 + p_{Ti}^2 = (E_i + p_{zi})(E_i - p_{zi})$$

and

$$k_{Tij}^2 = \min(p_{Ti}^2, p_{Tj}^2) R_{ij}$$

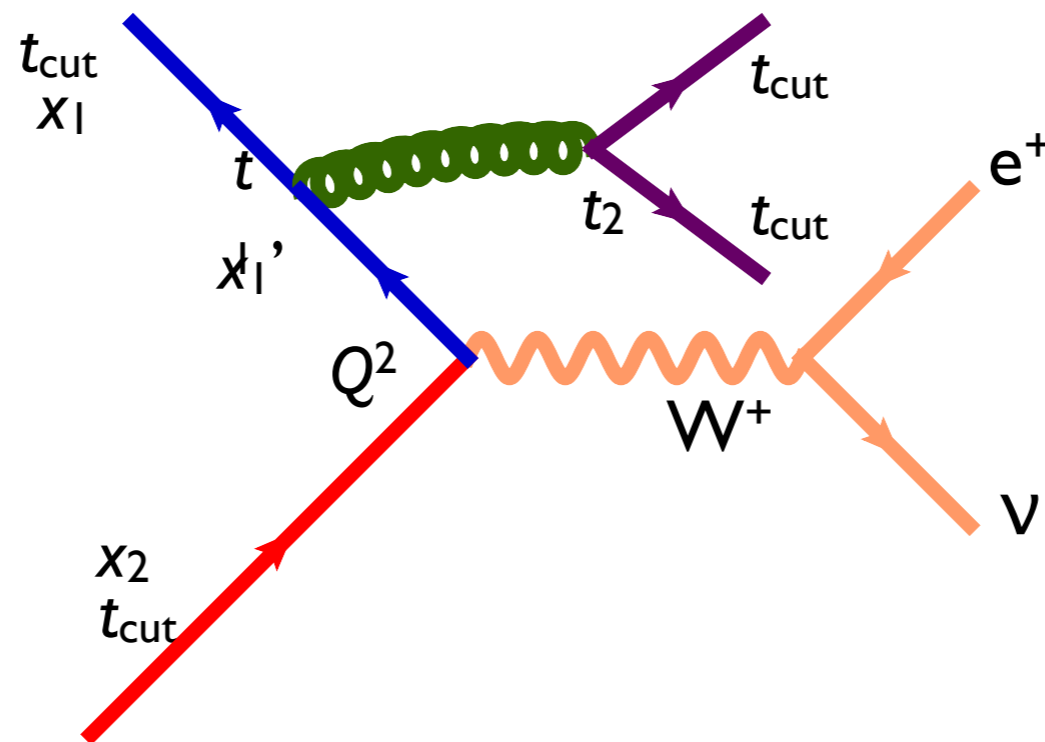
with $R_{ij} = 2[\cosh(y_i - y_j) - \cos(\phi_i - \phi_j)] \simeq (\Delta y)^2 + (\Delta \phi)^2$

Find the smallest k_{Tij} (or k_{Tibeam}), combine partons i and j (or i and the beam), and continue until you reach a $2 \rightarrow 2$ (or $2 \rightarrow 1$) scattering.

2. In AlpGen a more naive cone algorithm is used.

- Cannot use the standard k_T clustering:
 - MadGraph and Sherpa only allow clustering according to valid diagrams in the process. This means that, e.g., two quarks or quark-antiquark of different flavor are never clustered, and the clustering always gives a physically allowed parton shower history.
 - If there is an on-shell propagator in the diagram (e.g. a top quark), only clustering according to diagrams with this propagator is allowed.

2. The clustering provides a convenient choice for factorization scale Q^2 :



Cluster back to the $2 \rightarrow 2$ (here $qq \rightarrow W^-g$) system, and use the W boson transverse mass in that system.

3. How to divide the phase space between PS and ME:
This is where the schemes really differ:

AlpGen: MLM Cone

MadGraph: MLM Cone, k_T or shower- k_T

Sherpa: CKKW

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Sherpa: CKKW

- a. **Cone jet MLM scheme** (better suited for angular ordered showers, i.e. herwig, but works for all showers):
- Use cuts in p_T (p_T^{ME}) and ΔR between partons in ME
 - Cluster events after parton shower using a cone jet algorithm with the same ΔR and $p_T^{match} > p_T^{ME}$
 - Keep event if all jets are matched to ME partons (i.e., all ME partons are within ΔR of a jet)

3. How to divide the phase space between PS and ME:
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AlpGen: MLM Cone

MadGraph: MLM Cone, k_T or shower- k_T

Sherpa: CKKW

- b. **k_T -jet MLM scheme** (better suited for k_T ordered showers, i.e. pythia, but works for all showers):
- Use cut in the Durham k_T in ME
 - Cluster events after parton shower using the same k_T clustering algorithm into k_T jets with $k_T^{\text{match}} > k_T^{\text{ME}}$
 - Keep event if all jets are matched to ME partons (i.e., all partons are within k_T^{match} to a jet)

3. How to divide the phase space between PS and ME:
This is where the schemes really differ:

AlpGen: MLM Cone

MadGraph: MLM Cone, k_T or shower- k_T

Sherpa: CKKW

c. **Shower- k_T scheme:**

- Use cut in the Durham k_T in ME
- After parton shower, get information from the PS generator about the k_T^{PS} of the hardest shower emission
- Keep event if $k_T^{\text{PS}} < k_T^{\text{match}}$

3. How to divide the phase space between PS and ME:
This is where the schemes really differ:

AlpGen: MLM Cone

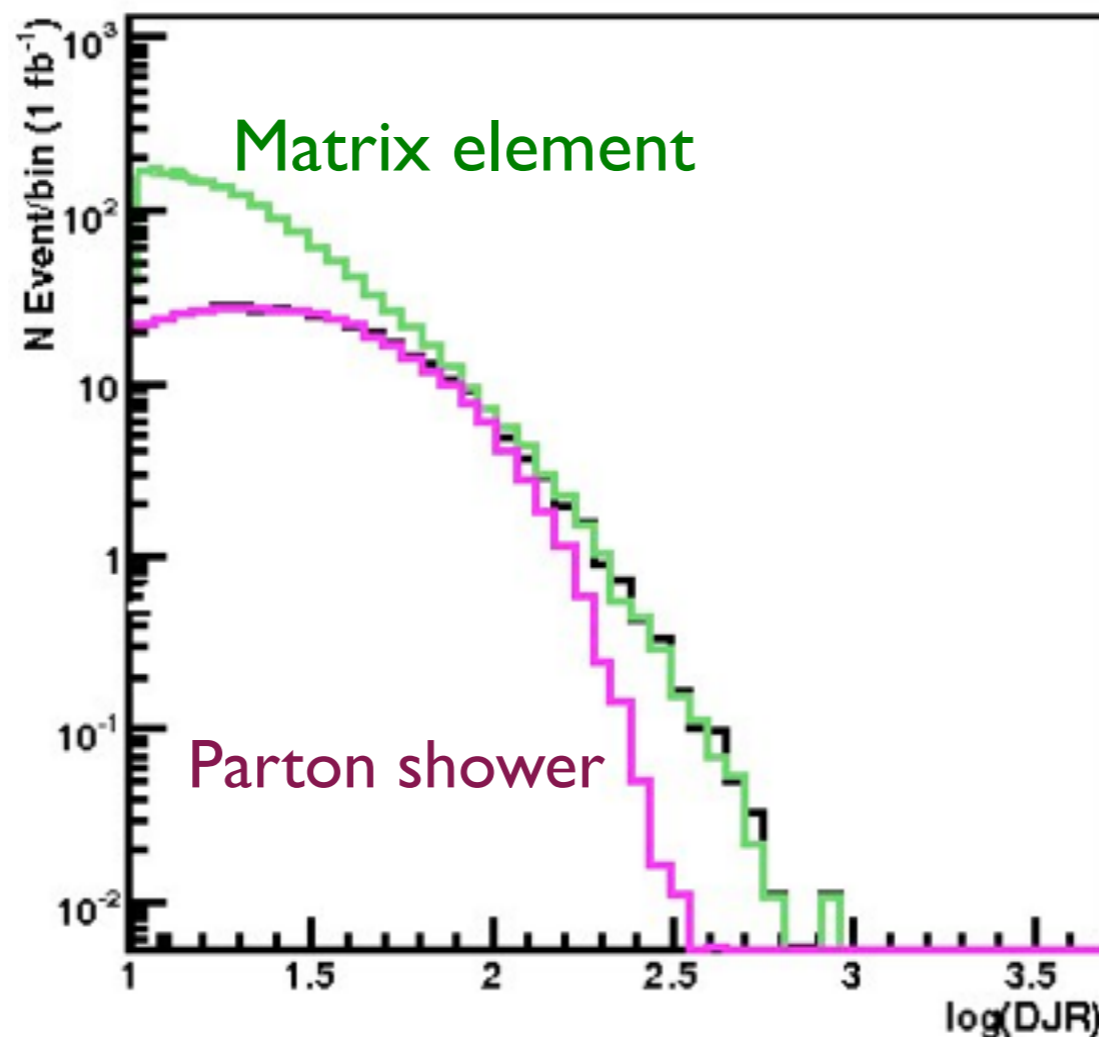
MadGraph: MLM Cone, k_T or shower- k_T

Sherpa: CKKW

- d. **CKKW Scheme** (Need special veto'ed shower):

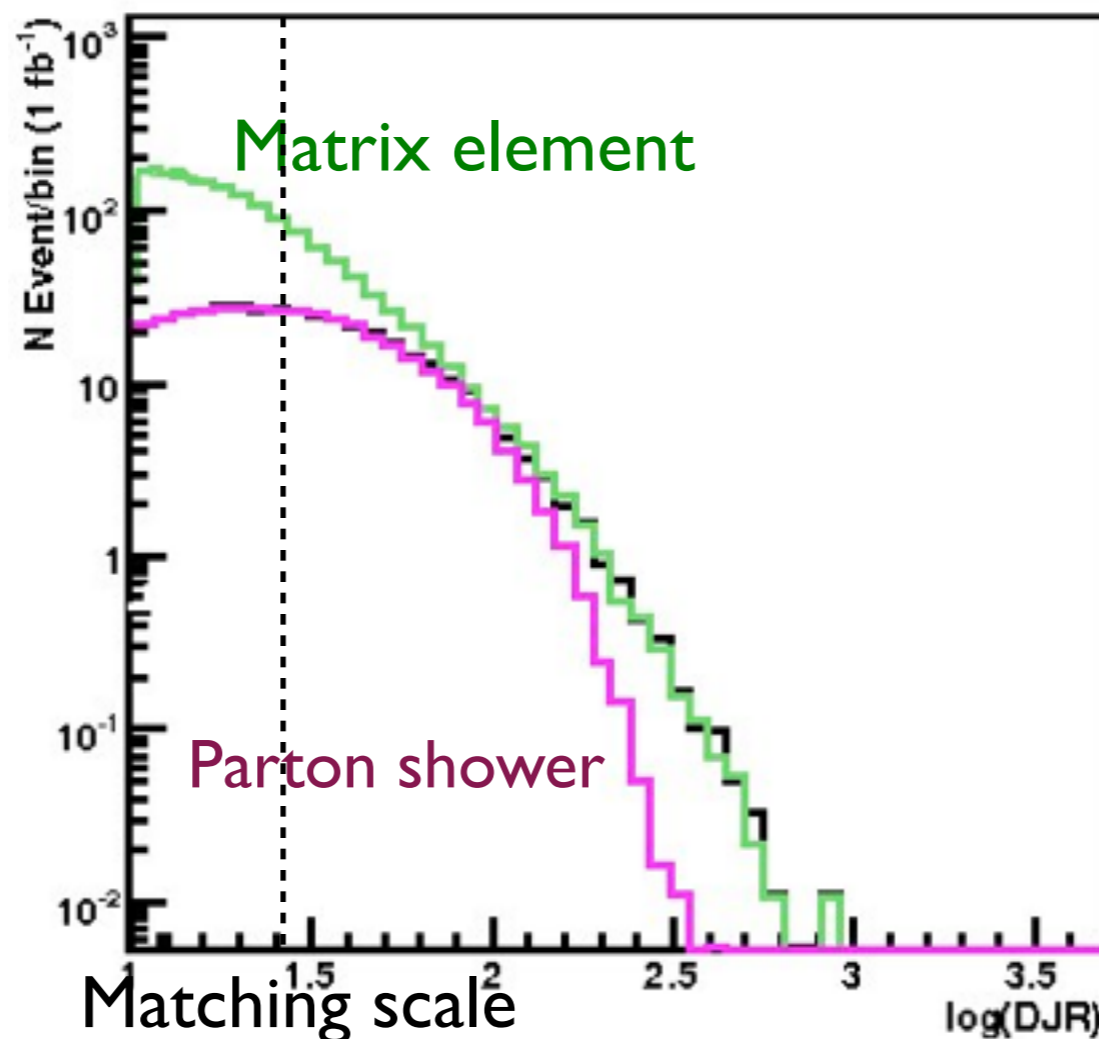
- Use cut in the Durham k_T in ME (k_T^{match})
- Because the Durham k_T is not the same as the evolution parameter of the shower, we might miss contributions, therefore
- Start the shower at the original scale, and after each emission, check the value of t_i :
- if $t_i > k_T^{\text{match}}$ veto that emission, i.e. continue the shower as if that emission never happened

- Regularization of matrix element divergence
- Correction of the parton shower for large momenta
- Smooth jet distributions



2nd QCD radiation jet in top pair production at the LHC, using MadGraph + Pythia

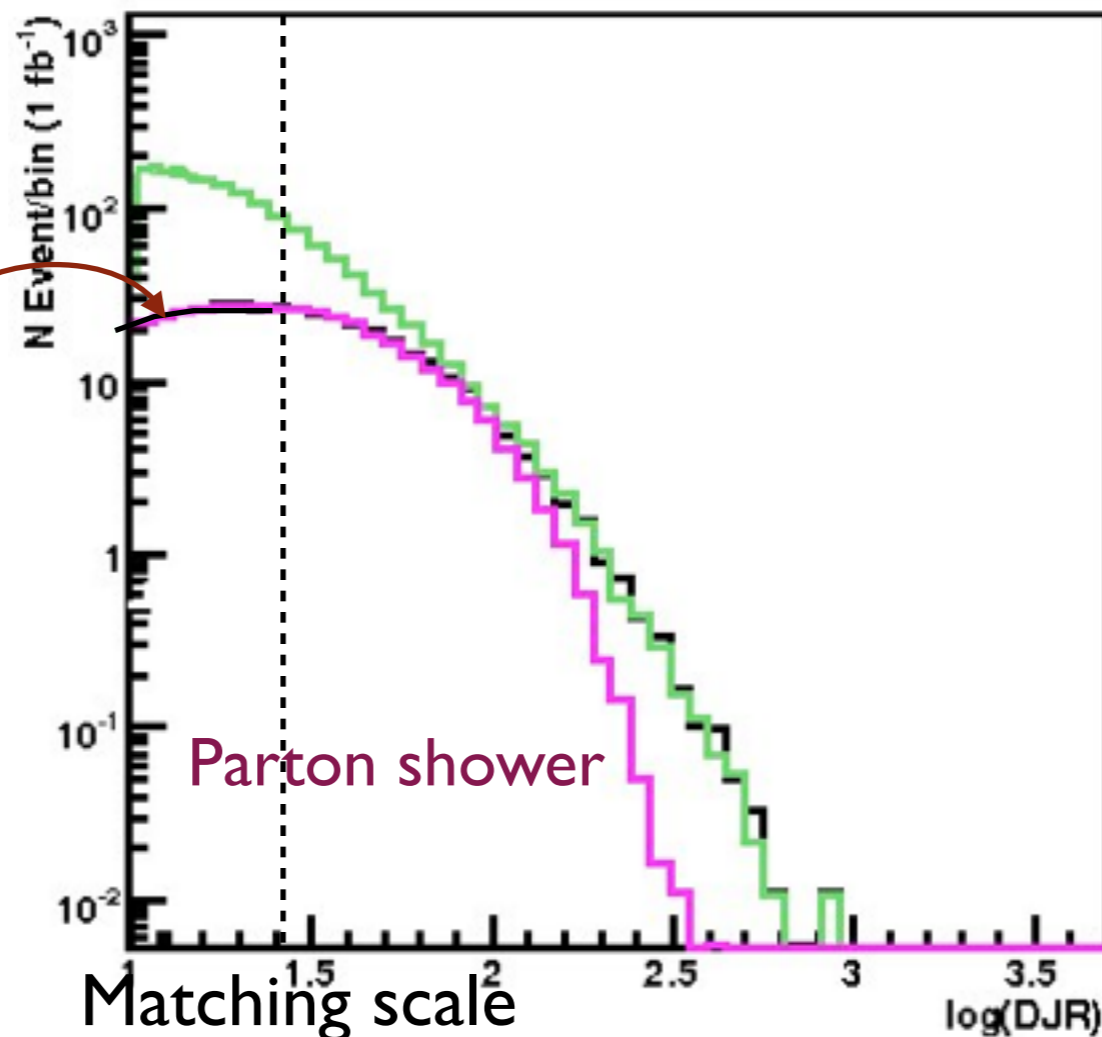
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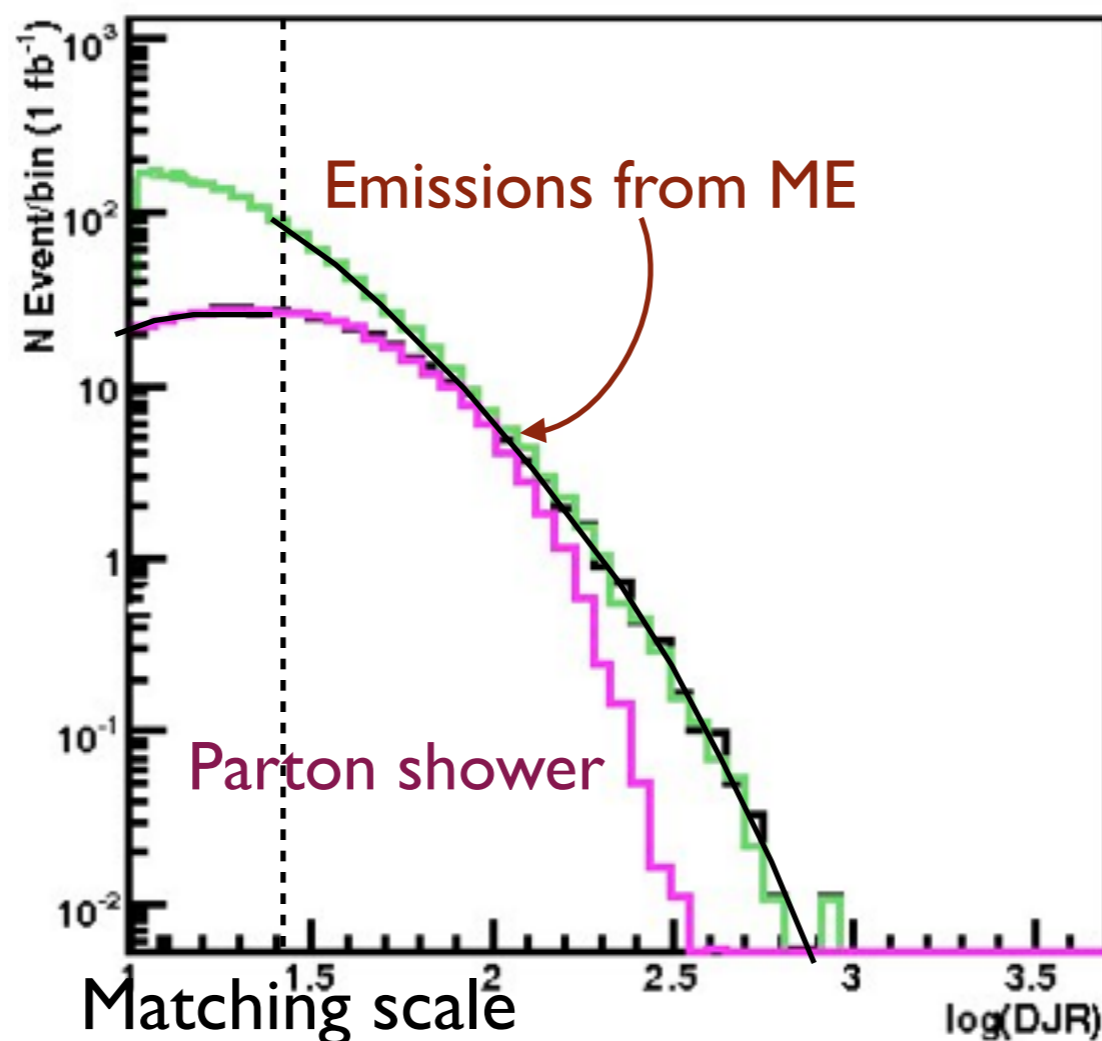
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Emissions from PS



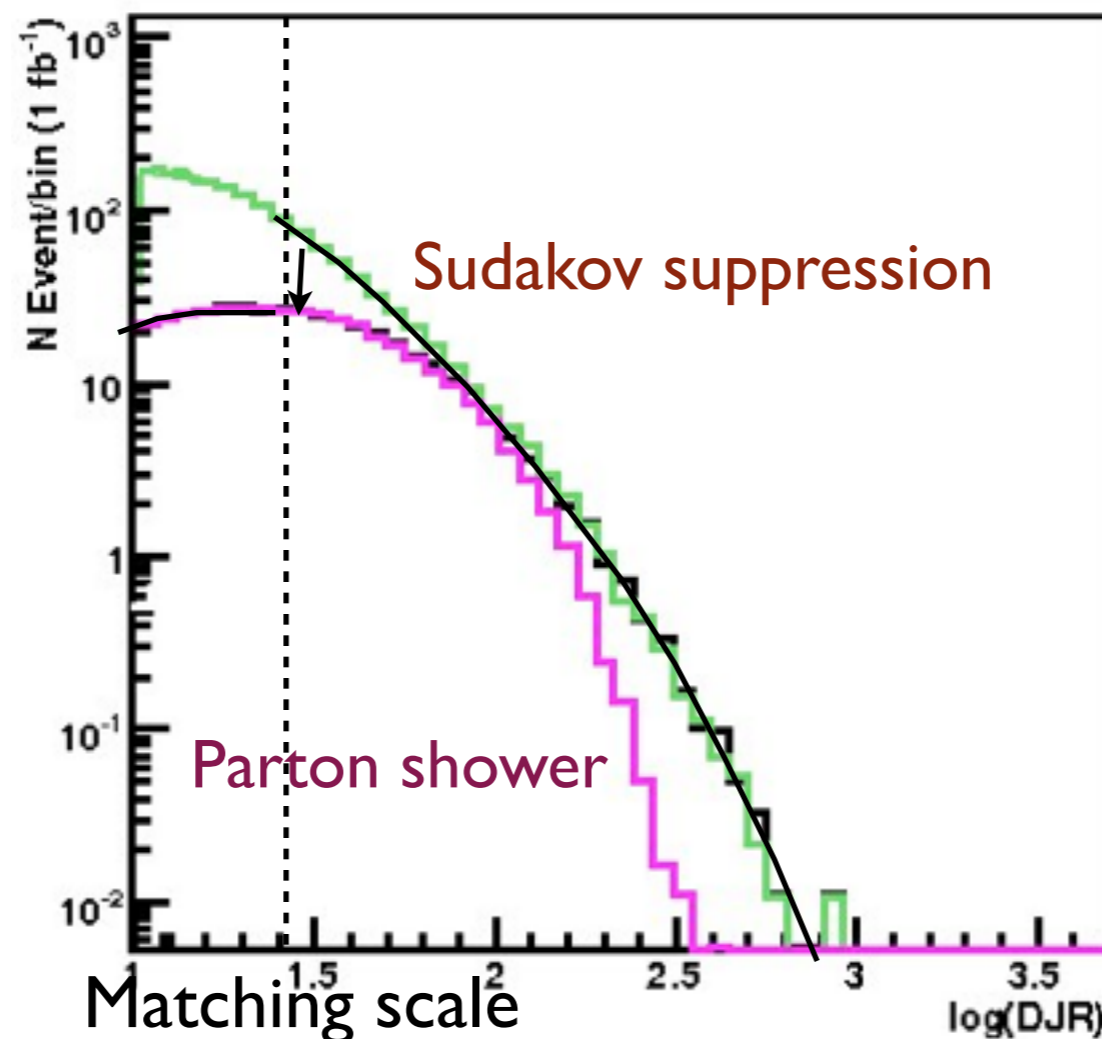
2nd QCD radiation jet in top pair production at the LHC, using MadGraph + Pythia

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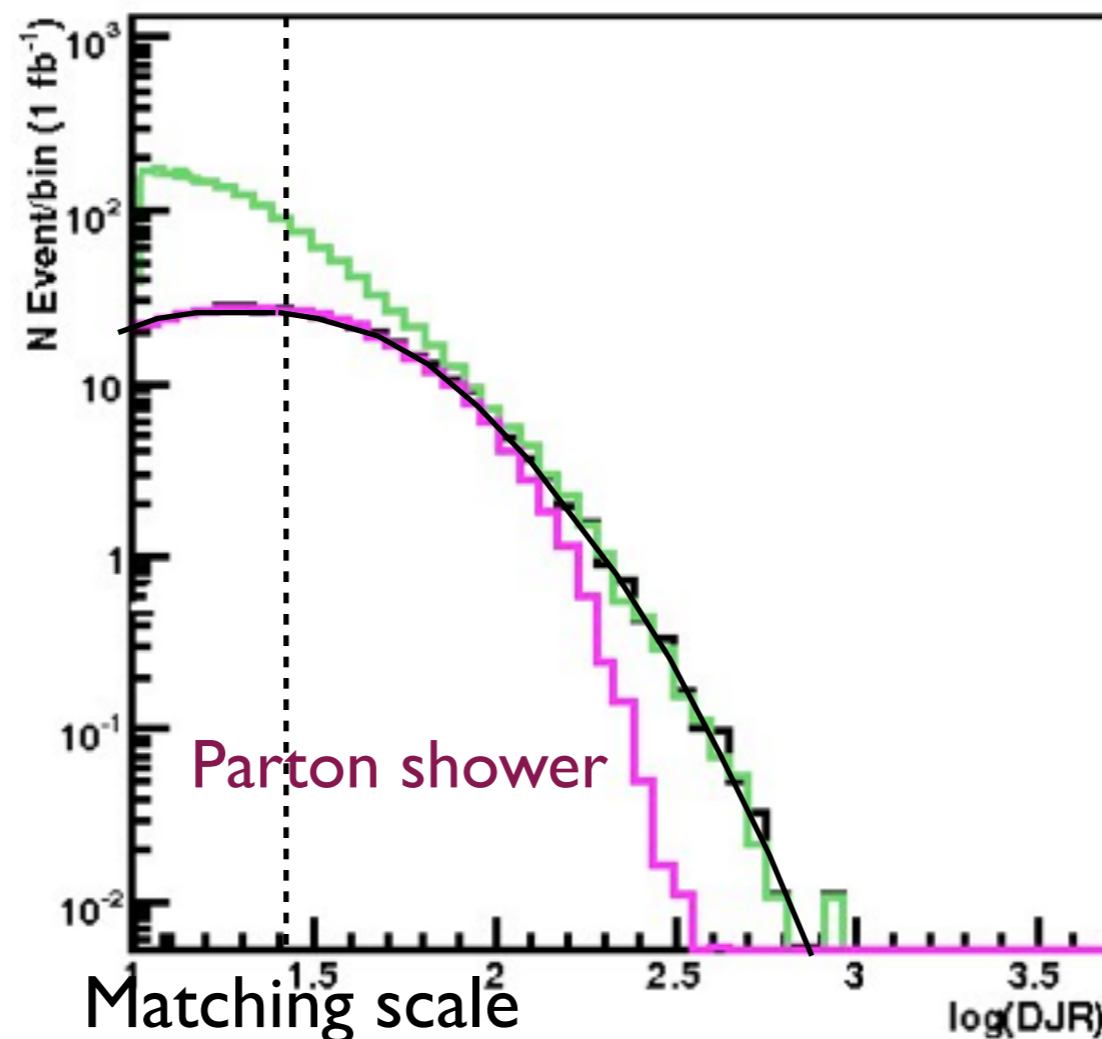
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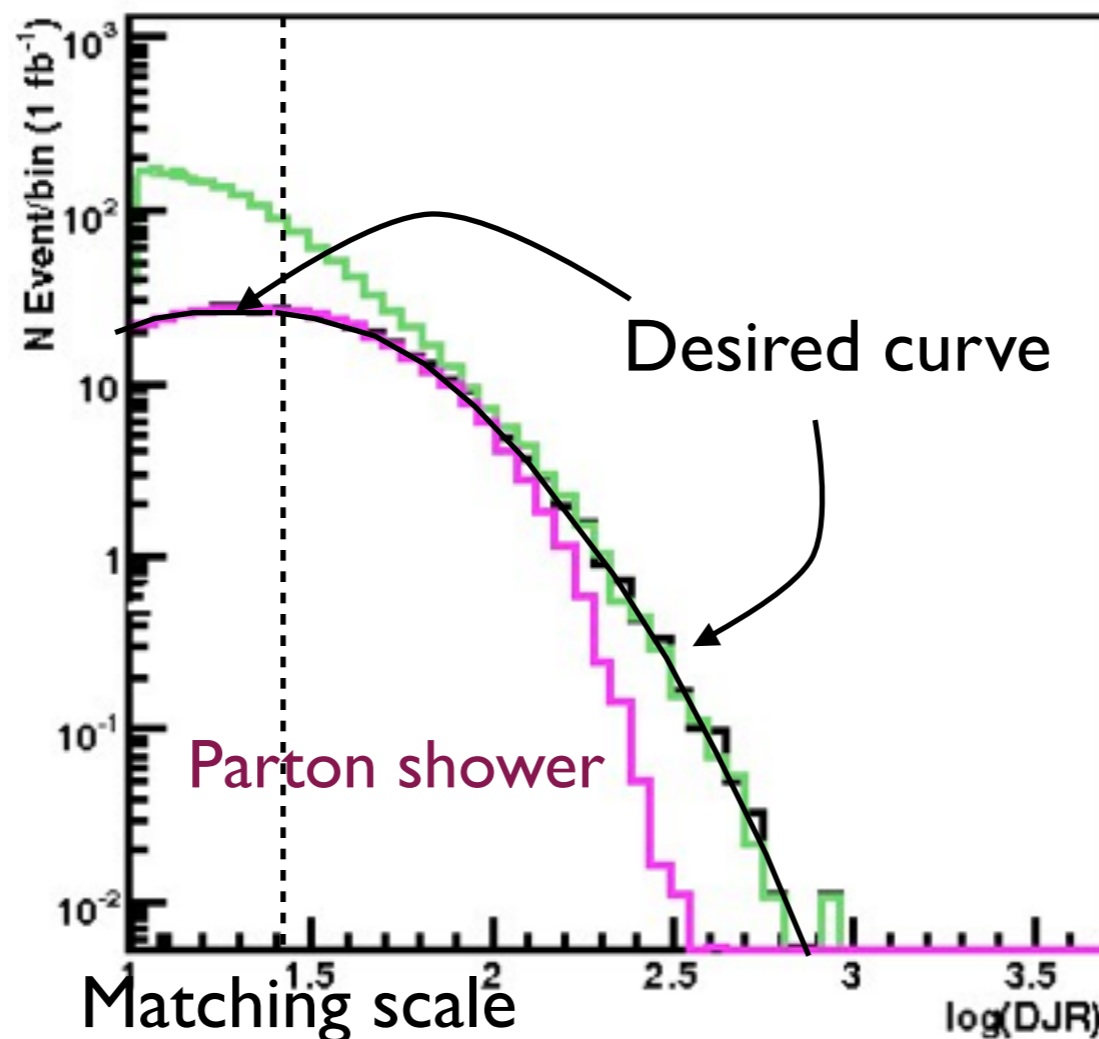
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- Correction of the parton shower for large momenta
- Smooth jet distributions



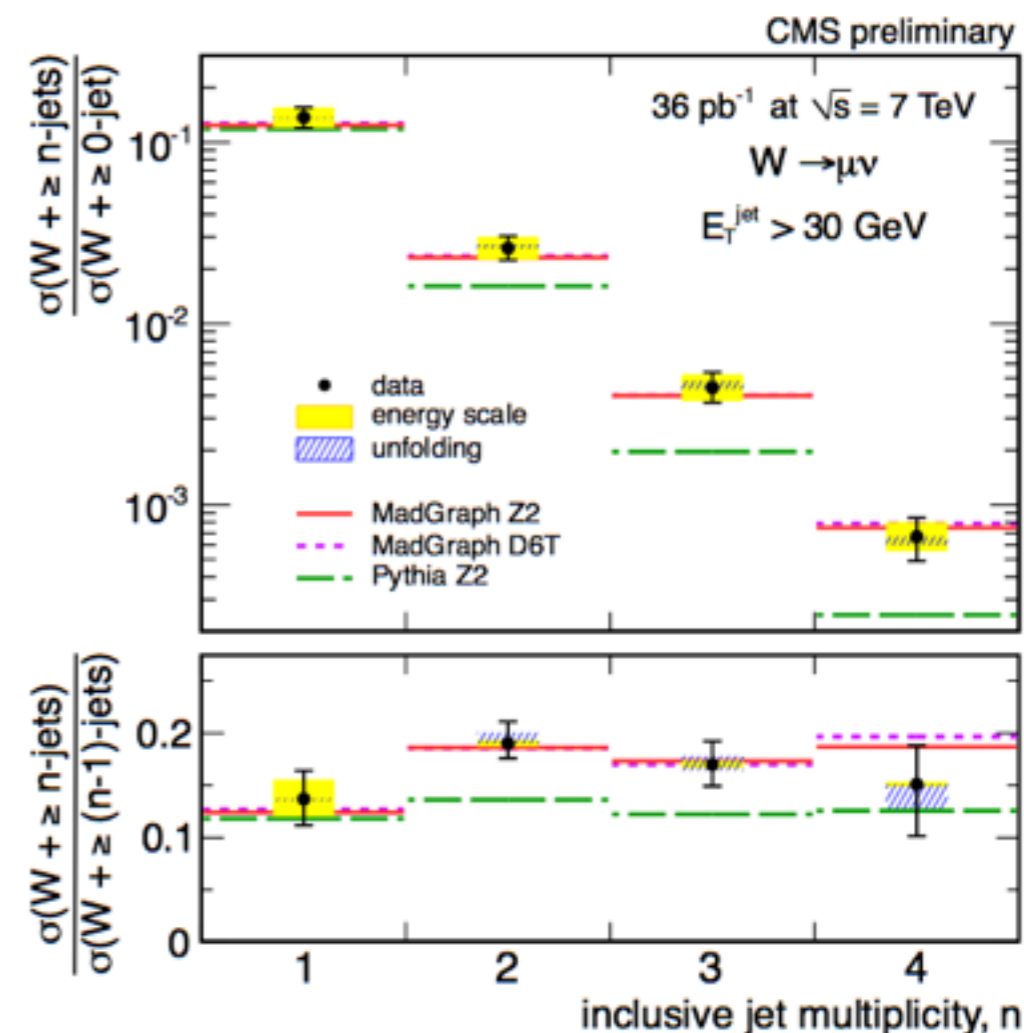
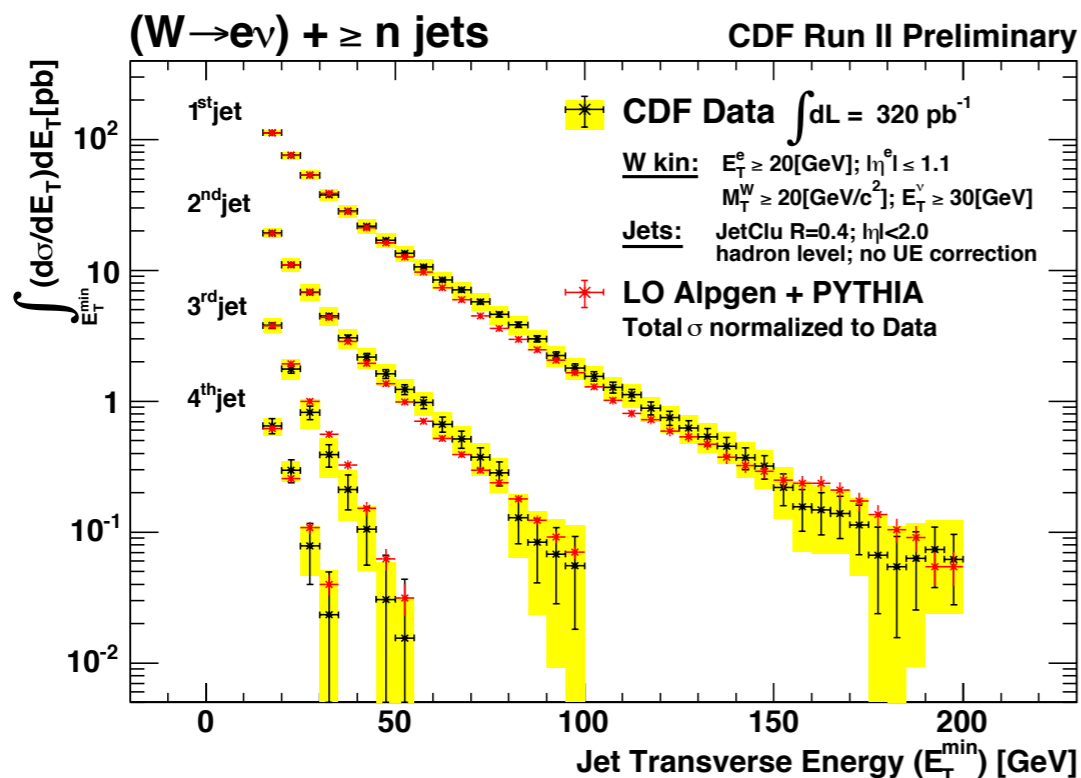
2nd QCD radiation jet in top pair production at the LHC, using MadGraph + Pythia

- Regularization of matrix element divergence
- Correction of the parton shower for large momenta
- Smooth jet distributions



2nd QCD radiation jet in top pair production at the LHC, using MadGraph + Pythia

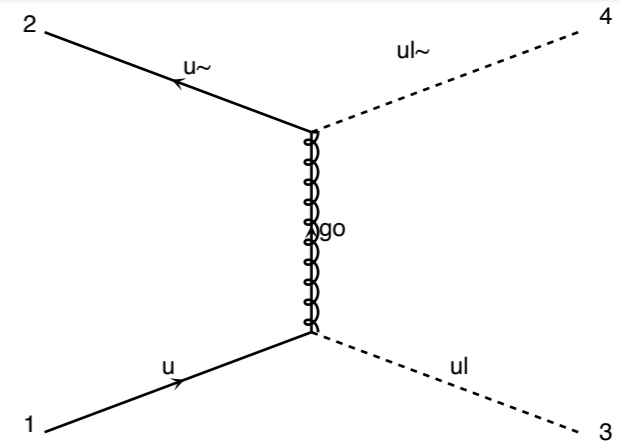
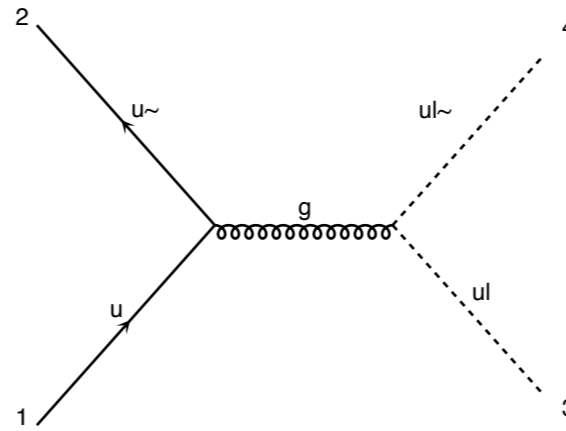
1. Generate ME events (with different parton multiplicities) using parton-level cuts ($p_T^{\text{ME}}/\Delta R$ or k_T^{ME})
2. Cluster each event and reweight α_s and PDFs based on the scales in the clustering vertices
3. Apply Sudakov factors to account for the required non-radiation above clustering cutoff scale and generate parton shower emissions below clustering cutoff:
 - a. (CKKW) Analytical Sudakovs + truncated showers
 - b. (CKKW-L) Sudakovs from truncated showers
 - c. (MLM) Sudakovs from reclustered shower emissions



- Very good agreement at Tevatron (left) and LHC (right)
- Matched samples obtained via different matching schemes (MLM and CKKW) consistent within the expected uncertainties.
- Pure parton shower (Pythia) doesn't describe the data beyond 1st jet.

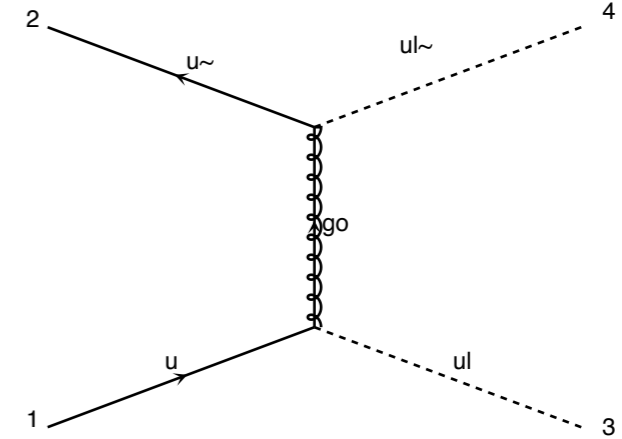
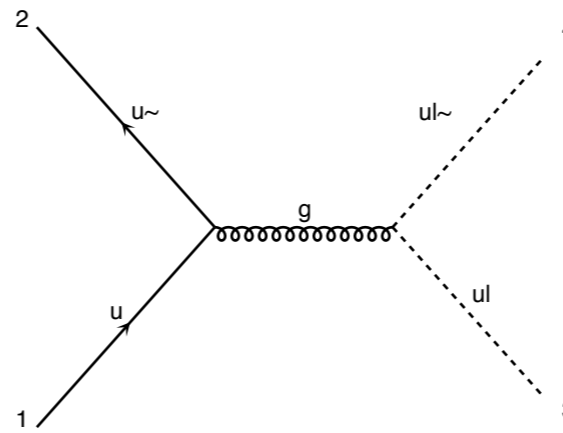
- Double counting of decay with jets
 - ➔ Especially for BSM

Squark pair production

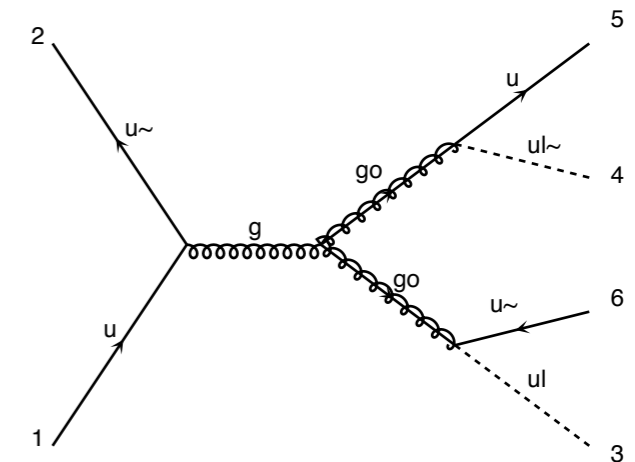
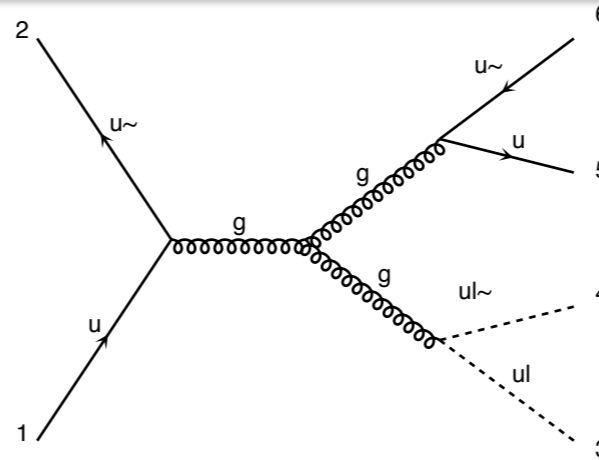


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Squark pair production

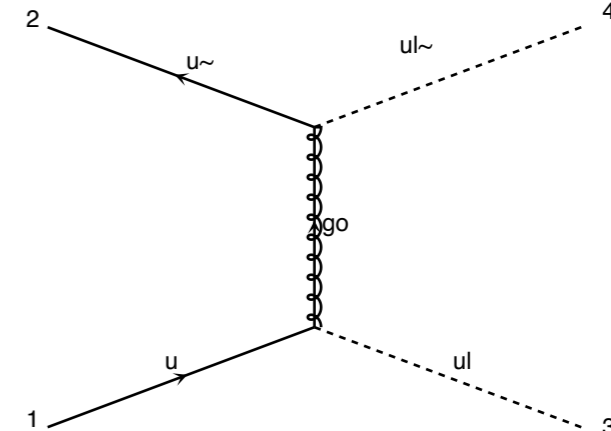
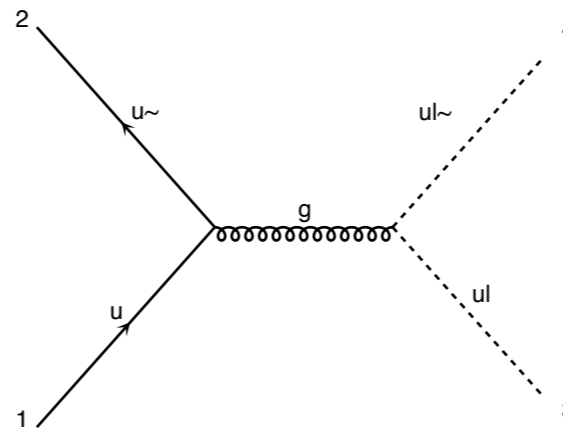


- 2j
- Presence of new production mechanism

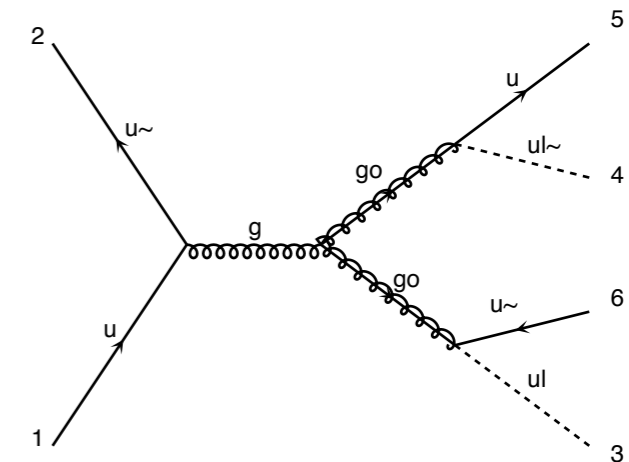
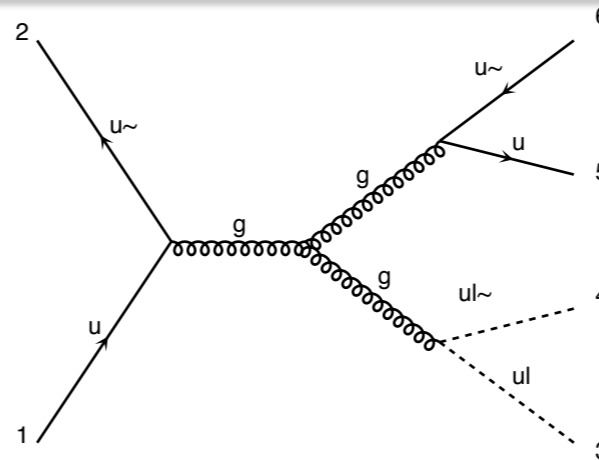


- Double counting of decay with jets
 - ➔ Especially for BSM

Squark pair production



- 2j
- Presence of new production mechanism



- Those sample should be consider as gluino production if gluino is on shell

- This has been solved in recent versions of MadGraph 5 by the new “\$” syntax

```
mg5> import model mssm
```

```
mg5> generate p p > dr dr~ j j $ go
```

- This has been solved in recent versions of MadGraph 5 by the new “\$” syntax
`mg5> import model mssm`
`mg5> generate p p > dr dr~ j j $ go`
- This removes any on-shell gluinos from the event generation (where on-shell is defined as $m \pm n \cdot \Gamma$ with n set by `bwcutoff` in the `run_card.dat`)

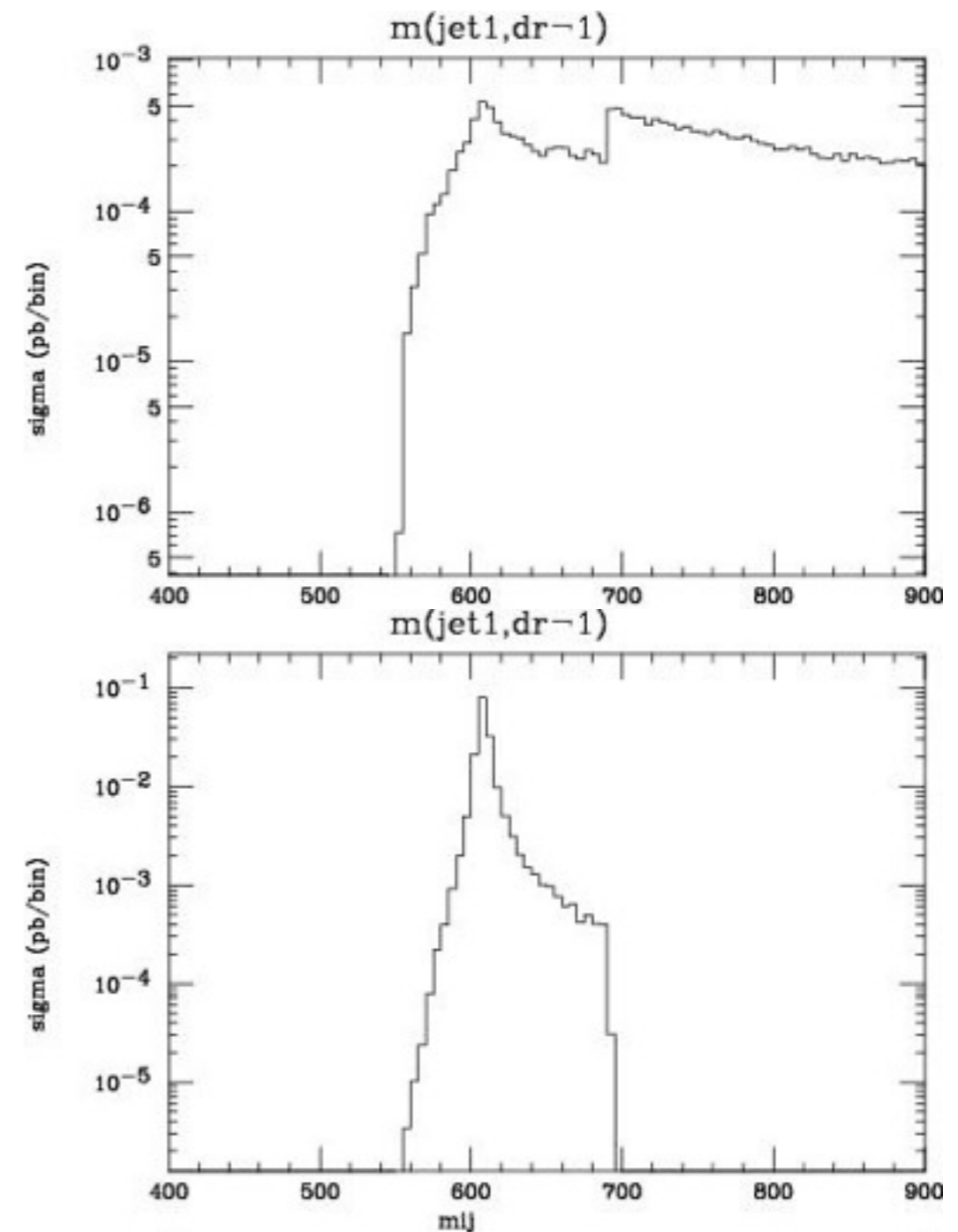
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mg5> import model mssm  
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- This removes any on-shell gluinos from the event generation (where on-shell is defined as $m \pm n \cdot \Gamma$ with n set by `bwcutoff` in the `run_card.dat`)
- The corresponding region is exactly filled if you run gluino production with gluinos decaying to `dr j` (using the same `bwcutoff`).

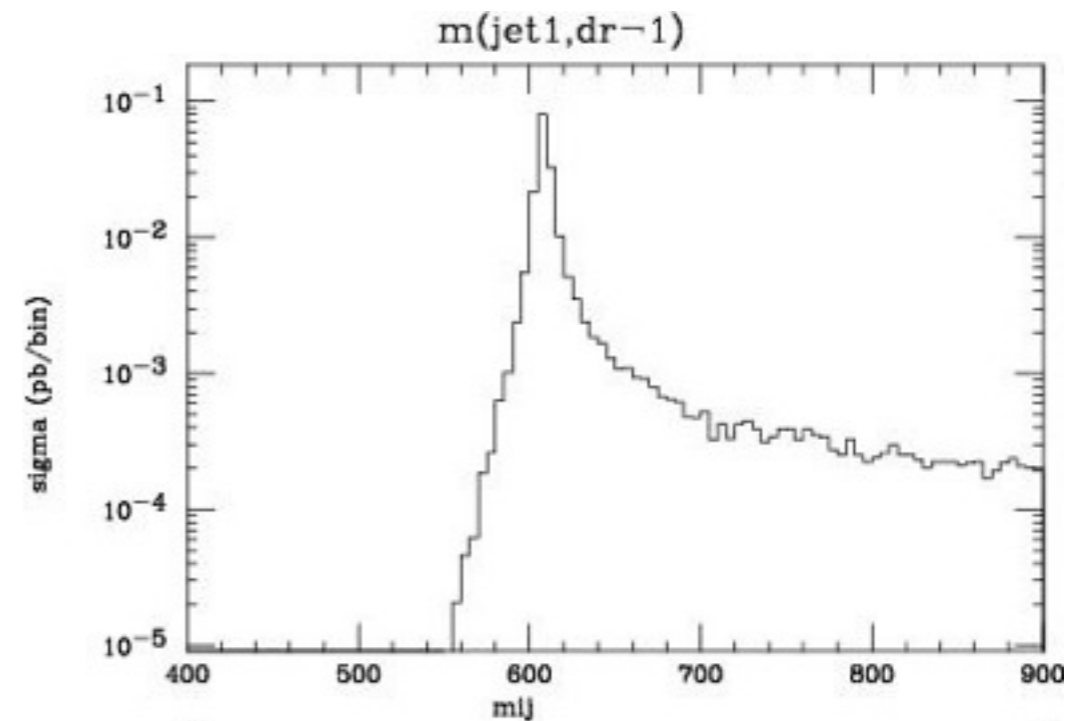
Invariant mass distributions
of d_r squark and d quark

$pp \rightarrow d_r d_r^* d g$

$pp \rightarrow d_r g, g \rightarrow d_r^* d$

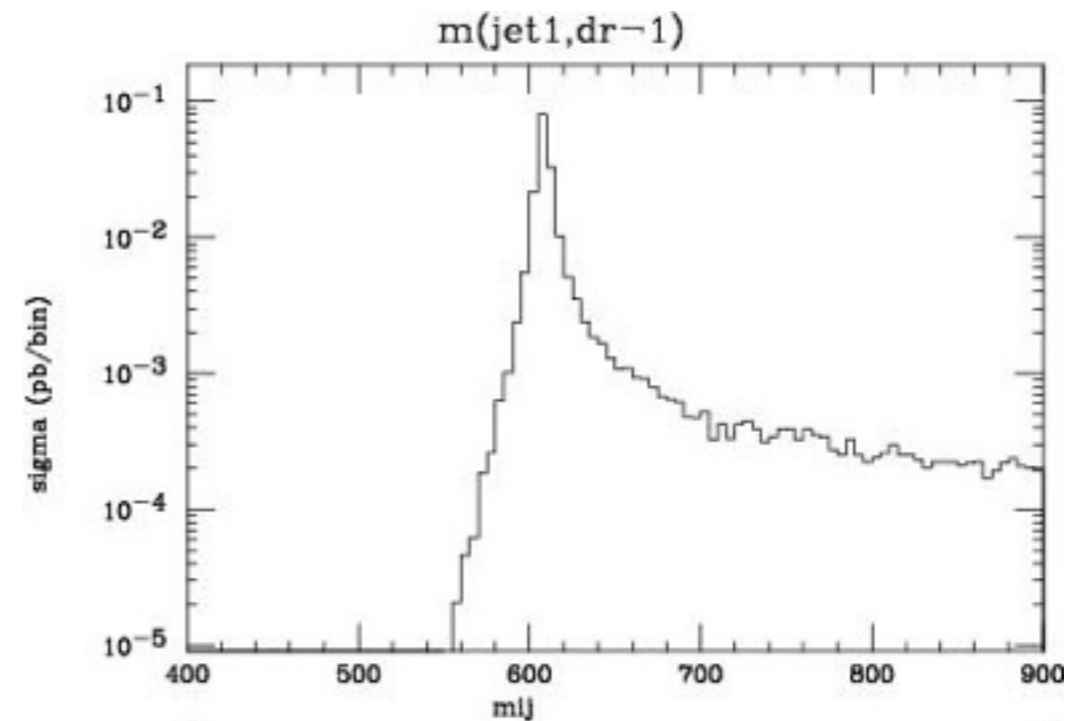


Invariant mass distributions
of d_r squark and d quark

$$\begin{aligned}
 & p p \rightarrow d_r \tilde{d}_r \rightarrow d g \\
 & + \\
 & p p \rightarrow d_r g, g \rightarrow \tilde{d}_r d
 \end{aligned}$$


Invariant mass distributions
of d_r squark and d quark

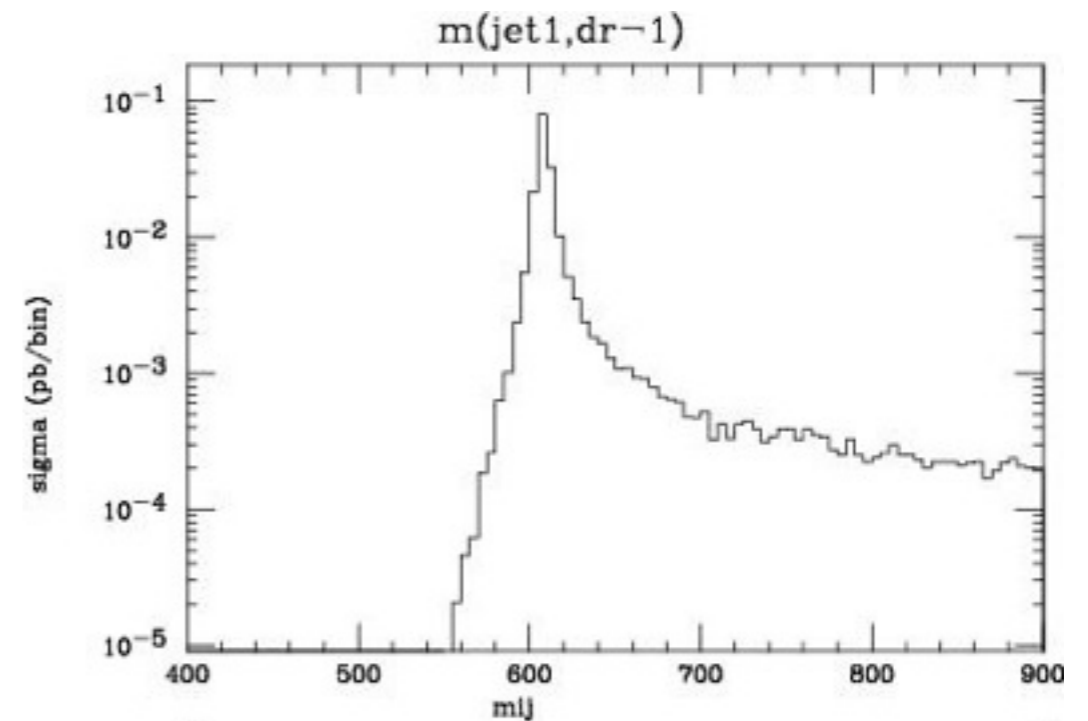
$$\begin{aligned}
 & p p \rightarrow d_r \bar{d}_r \sim d \bar{d} \text{ } g_0 \\
 & \quad \quad \quad + \\
 & p p \rightarrow d_r g_0, g_0 \rightarrow d_r \bar{d}_r
 \end{aligned}$$



Double counting between samples completely removed!

Invariant mass distributions
of d_r squark and d quark

$$\begin{aligned}
 & p p \rightarrow d_r \bar{d}_r \sim d \bar{d} \text{ } \text{\$ } g o \\
 & \quad \quad \quad + \\
 & p p \rightarrow d_r g o, g o \rightarrow d_r \bar{d}_r \text{ } d
 \end{aligned}$$



Double counting between samples completely removed!

Jet correctly handle by the matching/merging

Example: Simulation of $pp \rightarrow W$ with 0, 1, 2 jets (comfortable on a laptop)

```
mg5> generate p p > w+, w+ > l+ vl @0
mg5> add process p p > w+ j, w+ > l+ vl @1
mg5> add process p p > w+ j j, w+ > l+ vl @2
mg5> output
```

In run_card.dat:

...

1 = ickkw

...

0 = ptj

...

15 = xqcut

Matching on

No cone matching

k_T matching scale

Matching automatically done when run through
MadEvent and Pythia!

- By default, k_T -MLM matching is run if $xqcut > 0$, with the matching scale $QCUT = \max(xqcut * 1.4, xqcut + 10)$
- For shower- k_T , by default $QCUT = xqcut$
- If you want to change the Pythia setting for matching scale or switch to shower- k_T matching:

```
In pythia_card.dat:
```

```
...
```

```
! This sets the matching scale, needs to be > xqcut
```

```
QCUT = 30
```

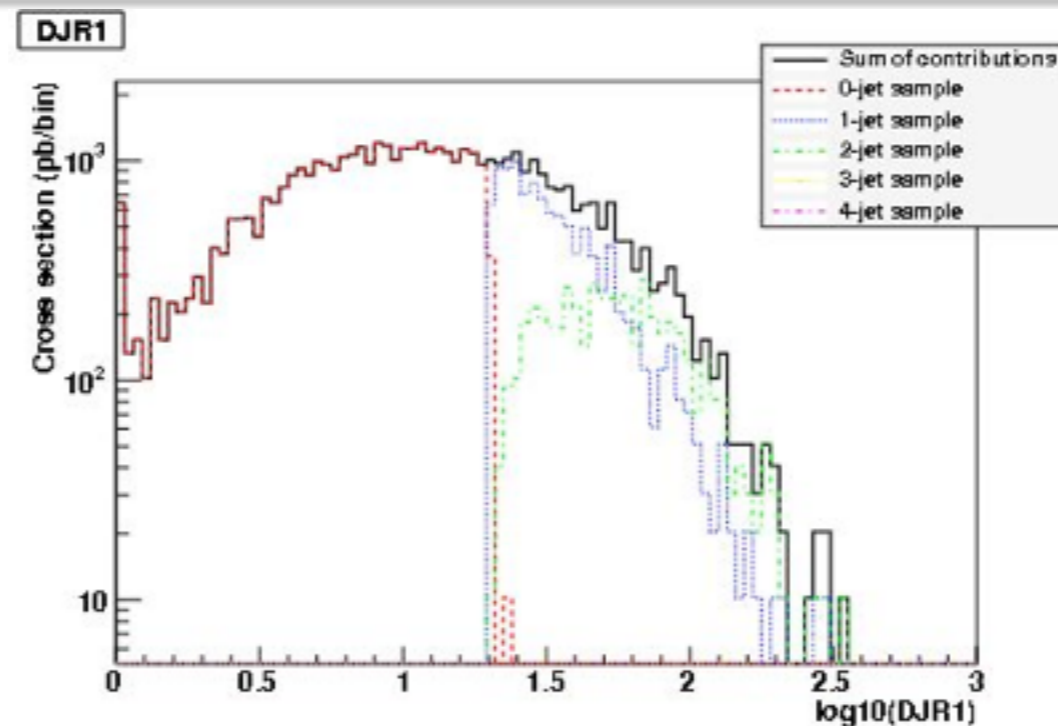
```
! This switches from  $k_T$ -MLM to shower- $k_T$  matching
```

```
! Note that  $MSTP(81) \geq 20$  needed (pT-ordered shower)
```

```
SHOWERKT = T
```

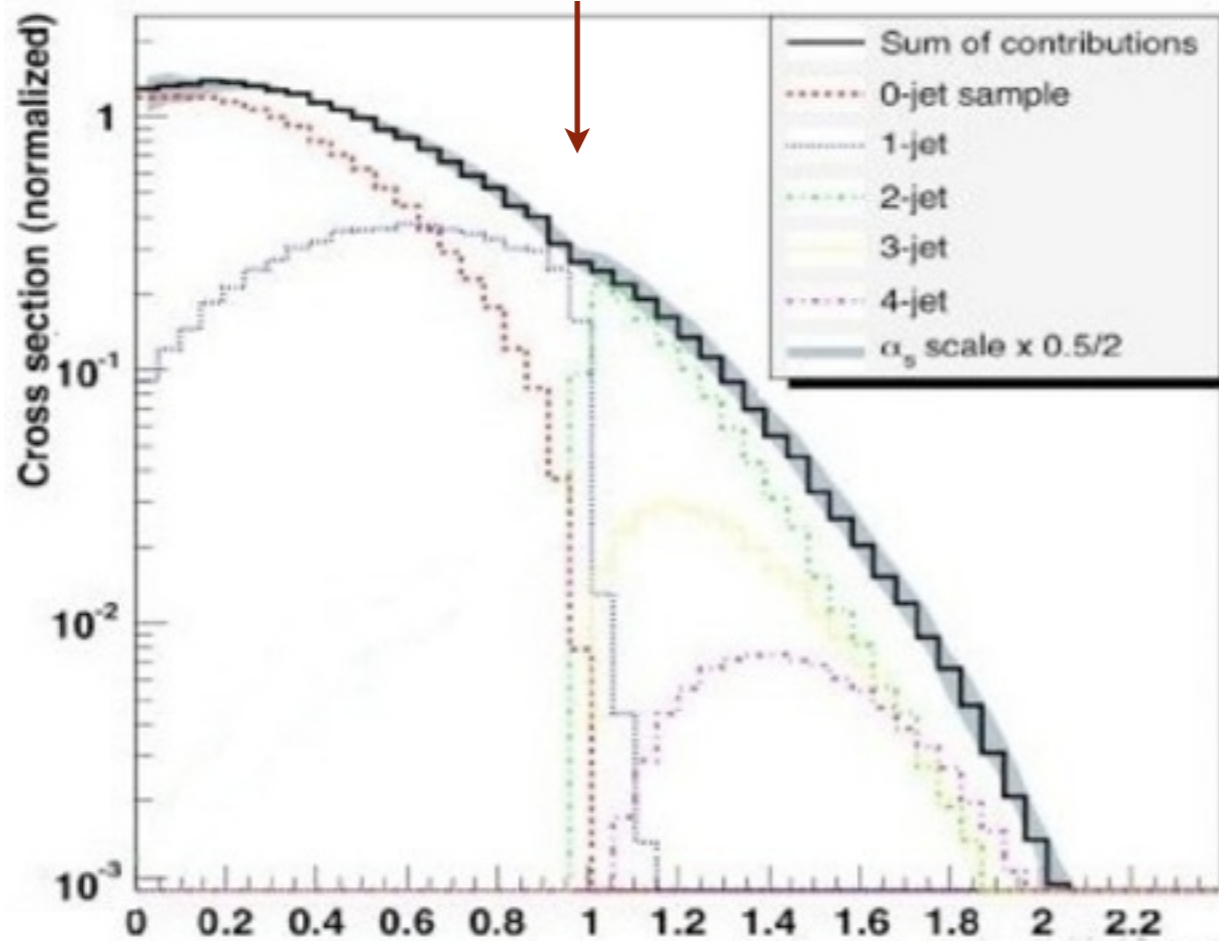
- The matching scale (QCUT) should typically be chosen around $1/6$ - $1/2$ x hard scale (so x_{qcut} correspondingly lower)
- The matched cross section (for $X+0, 1, \dots$ jets) should be close to the unmatched cross section for the 0-jet sample (found on the process HTML page)
- The differential jet rate plots should be smooth
- When QCUT is varied (within the region of validity), the matched cross section or differential jet rates should not vary significantly

- These are the clustering scales in the kt-jet clustering scheme
- DJR1: pT of the last remaining jet
- DJR2: The **minimum** between the pT of the second to last remaining jet **and** the kt between the last two jets.
- Only radiative jets (not those from decay) should enter those plots.

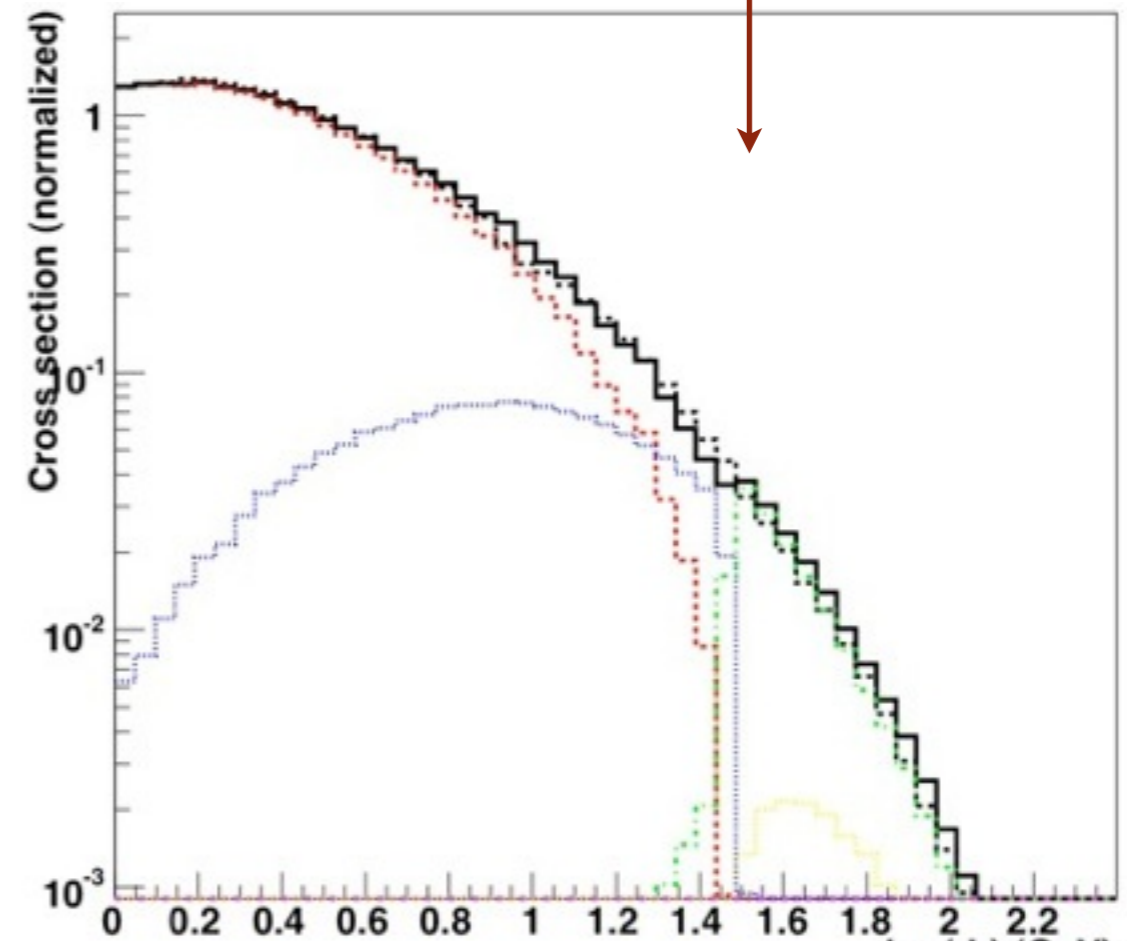


W+jets production at the Tevatron for MadGraph+Pythia (k_T -jet MLM scheme, q^2 -ordered Pythia showers)

$Q^{\text{match}} = 10 \text{ GeV}$

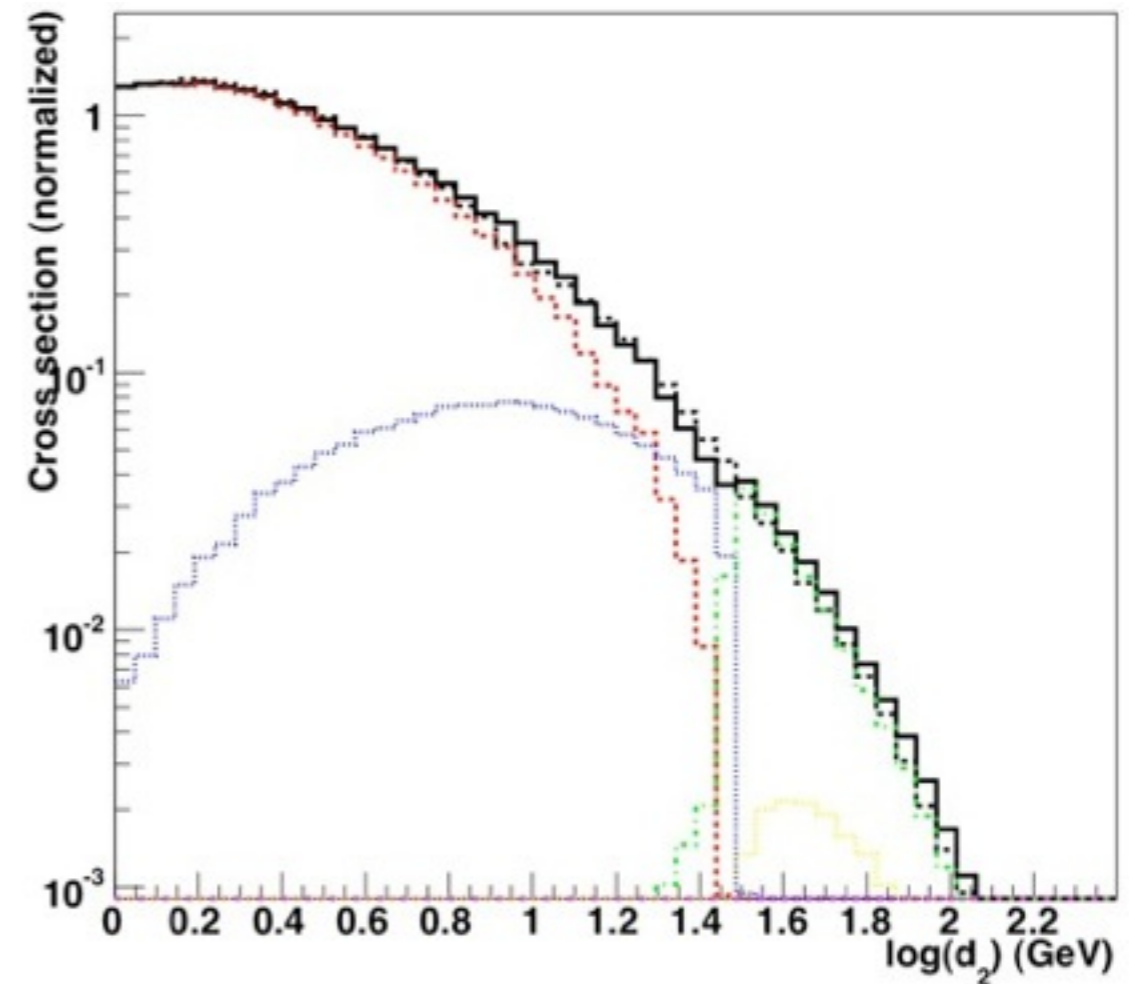
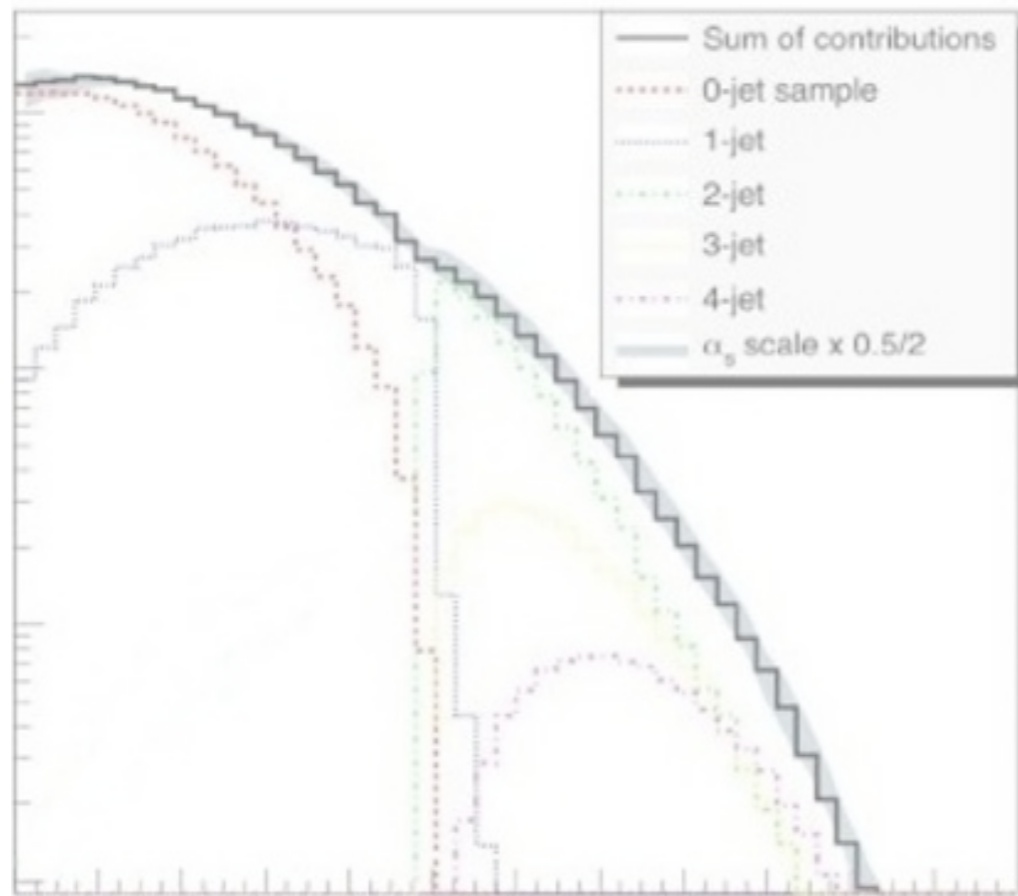


$Q^{\text{match}} = 30 \text{ GeV}$

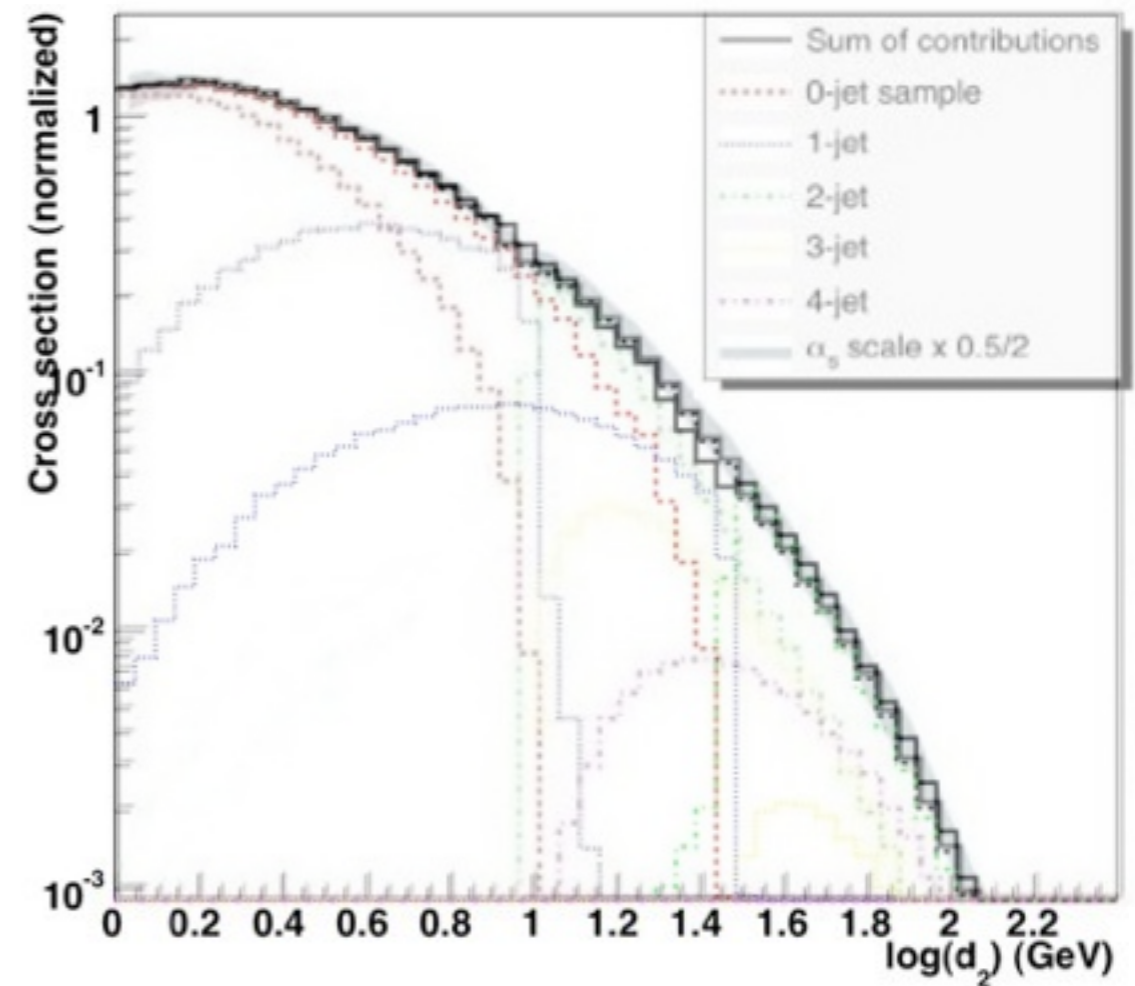


log(Differential jet rate for $1 \rightarrow 2$ radiated jets $\sim p_T(2\text{nd jet})$)

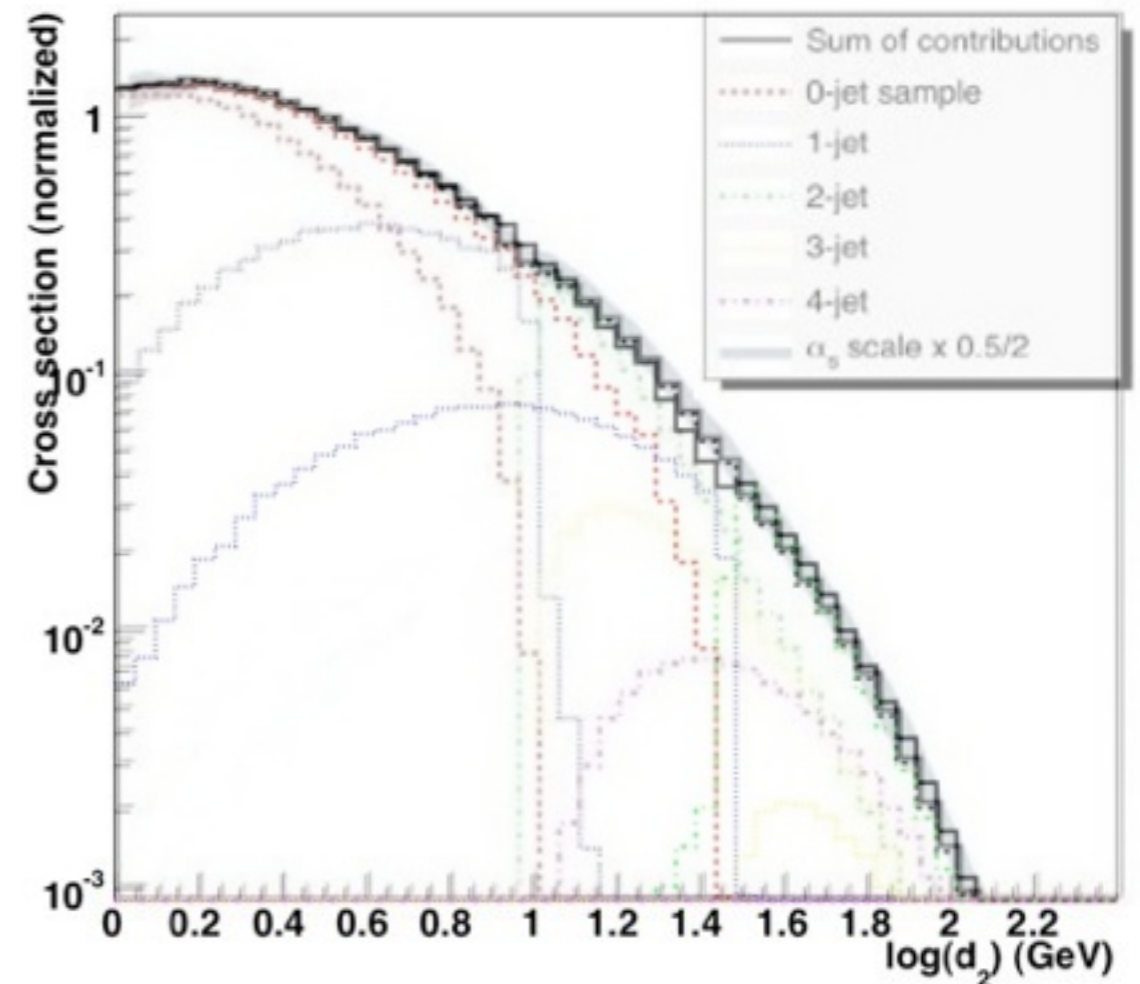
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(k_T -jet MLM scheme, q^2 -ordered Pythia showers)



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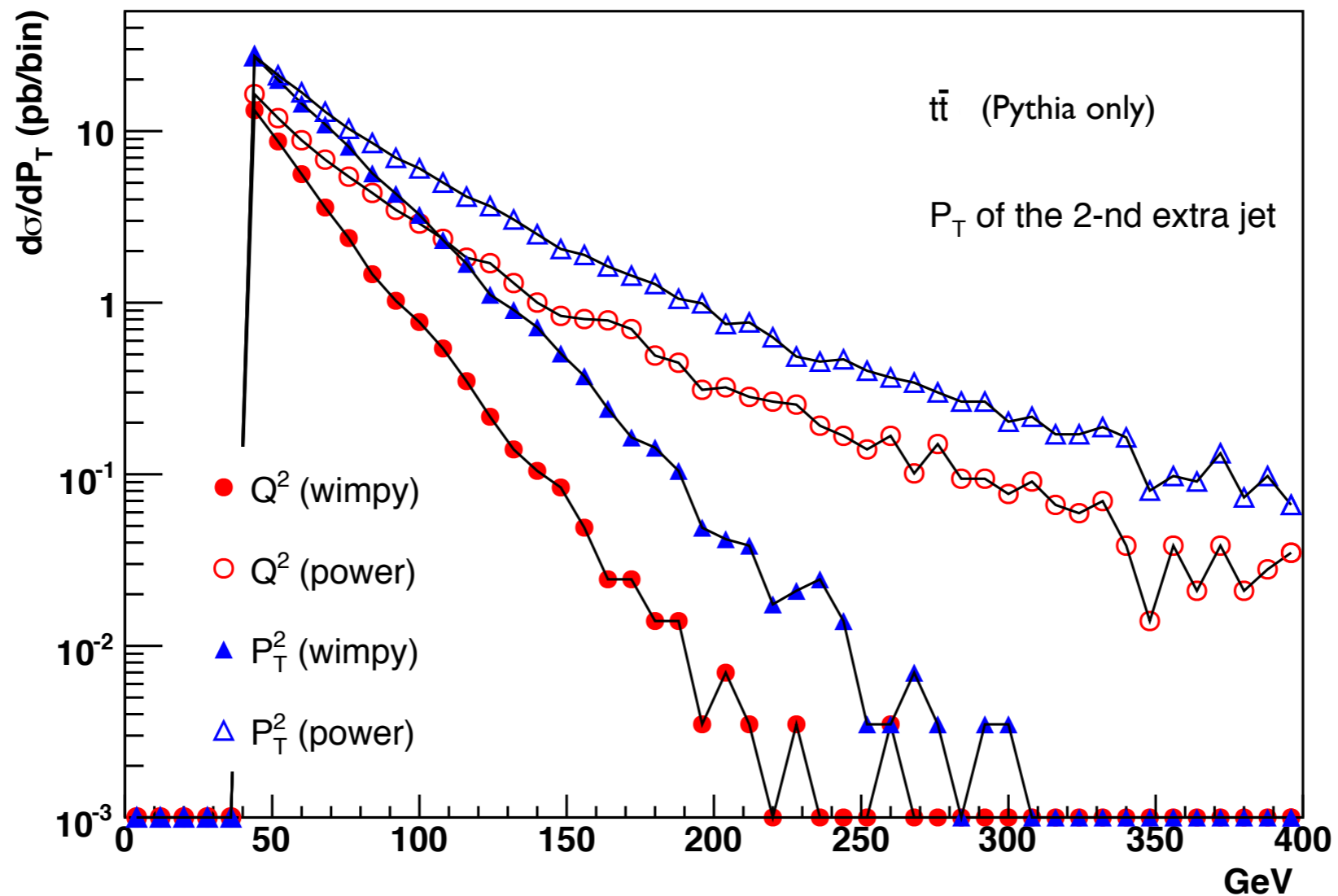
W+jets production at the Tevatron for MadGraph+Pythia
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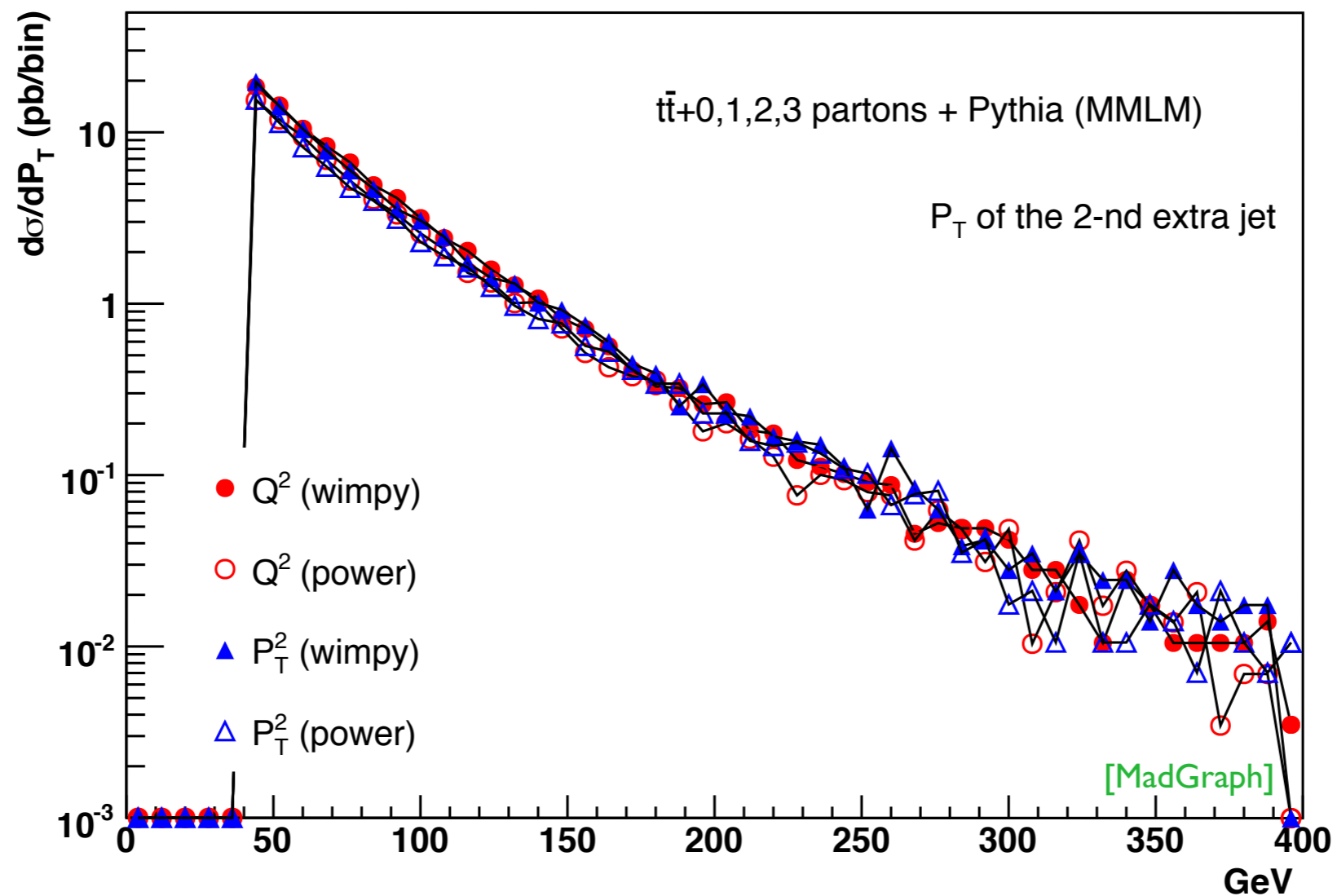
Jet distributions smooth, and stable when we vary the matching scale!

- Time for live Demonstration!

In the soft-collinear approximation of Parton Shower MCs, parameters are used to tune the result \Rightarrow Large variation in results (small prediction power)



In a matched sample these differences are irrelevant since the behavior at high p_T is dominated by the matrix element.



- Despite the apparent enormous complexity of simulation of complete collider events, nature has kindly allowed us to factorize the simulation into separate steps
- The Monte Carlo method allows us to step-by-step simulate hard scattering, parton shower, particle decays, hadronization, and underlying event
- Jet matching between matrix elements and parton showers gives crucial improvement of simulation of background as well as signal processes
- Running matching with MadGraph + Pythia is very easy, but the results should always be checked for consistency