

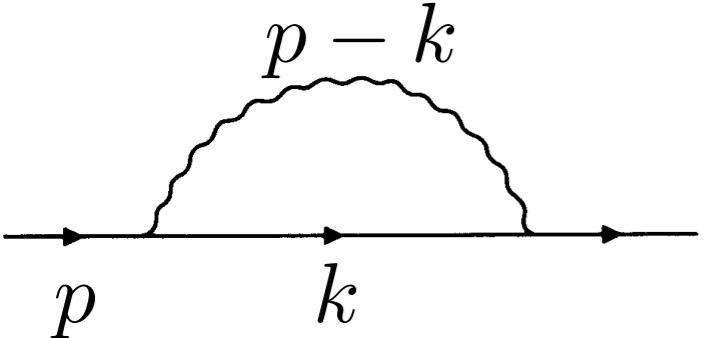
Next-to-Leading Order with FeynRules

Celine Degrande (IPPP, Durham University)
Lund 2014

Plan

- Renormalization
- KLN Theorem
- Rational Terms
- FeynRules at NLO

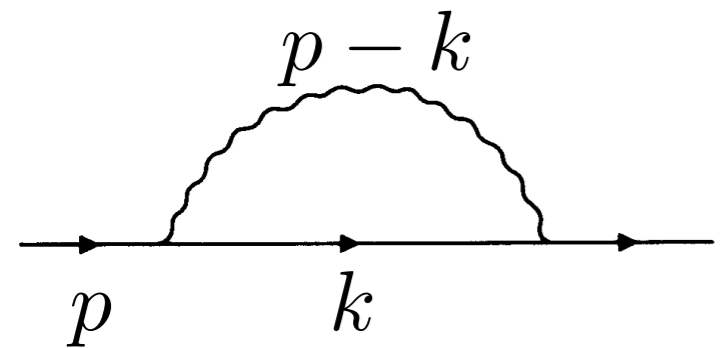
Electron self-energy



The diagram shows a horizontal line representing an electron with momentum p entering from the left and exiting to the right. A loop is formed by a wavy line (photon) and a solid line (electron). The loop starts at a vertex on the electron line, goes up and right, then down and right, then left, and finally up and left back to the starting vertex. The momentum of the incoming electron is labeled p at the left end. The momentum of the loop electron is labeled k at the bottom vertex. The momentum of the loop photon is labeled $p - k$ at the top vertex.

$$= (-ie)^2 \int \frac{d^4 k}{(2\pi)^4} \gamma^\mu \frac{i(\not{k} + m)}{k^2 - m^2 + i\epsilon} \gamma^\nu \frac{-i\eta_{\mu\nu}}{(p - k)^2 + i\epsilon}$$

Electron self-energy



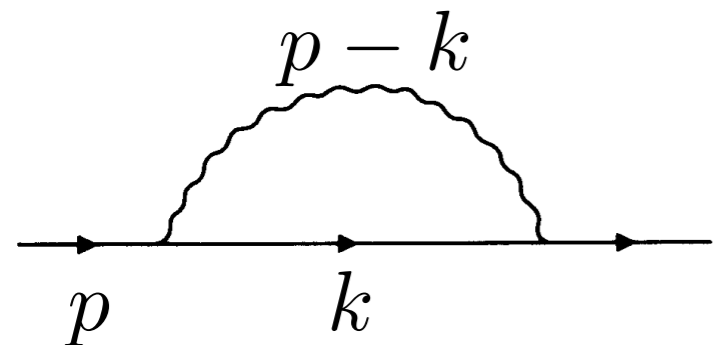
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Feynman parameter :

$$\frac{1}{AB} = \int_0^1 dx \frac{1}{(Ax + B(1-x))}$$

$$= -e^2 \int_0^1 dx \int \frac{d^4 l}{(2\pi)^4} \frac{\gamma^\mu (\not{l} + \cancel{x\not{p}} + m) \gamma_\mu}{(l^2 - \Delta + i\epsilon)^2} \left\{ \begin{array}{l} l \equiv k - xp \\ \Delta = (1-x)(m^2 - xp^2) \end{array} \right.$$

Electron self-energy



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Feynman parameter :

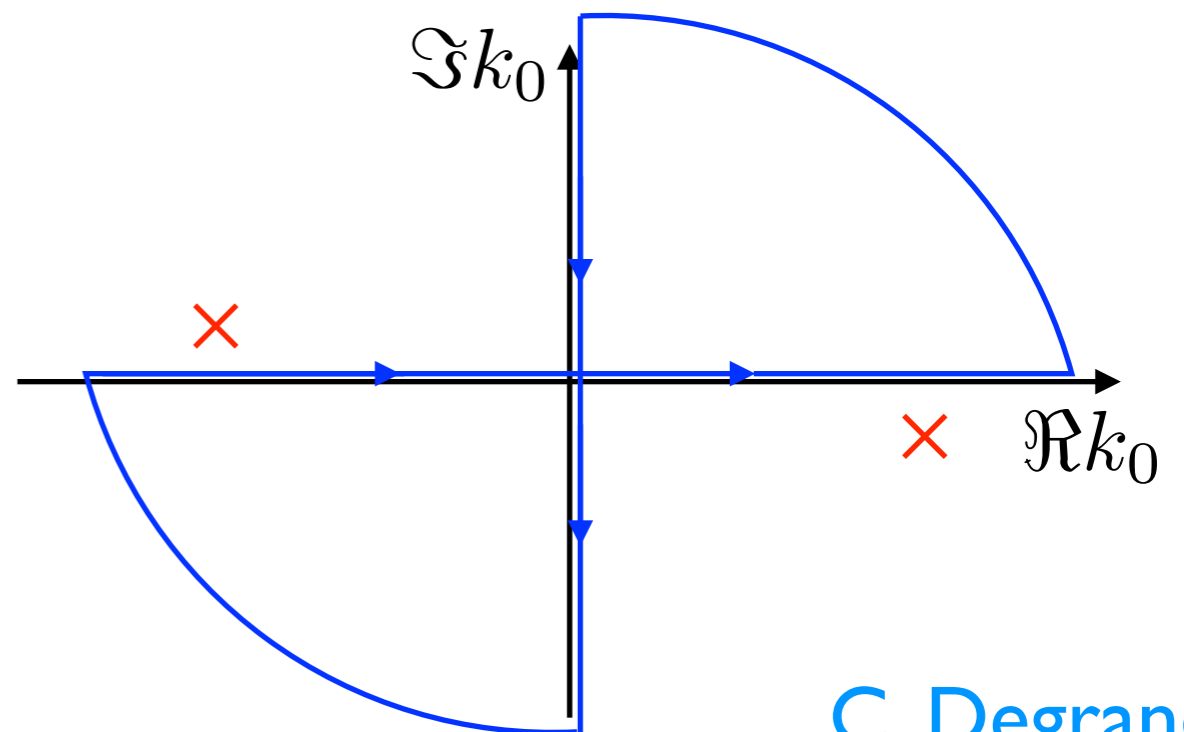
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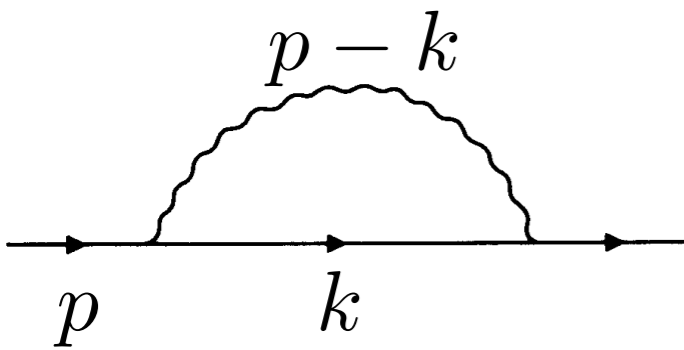
Minkowski to Euclidian space

$$\int_{-\infty}^{+\infty} dl_0 = - \int_{i\infty}^{-i\infty} dl_0 = i \int_{-\infty}^{+\infty} dl_0^E$$

$$l_0^E \equiv -il_0, \quad 1^E \equiv 1$$



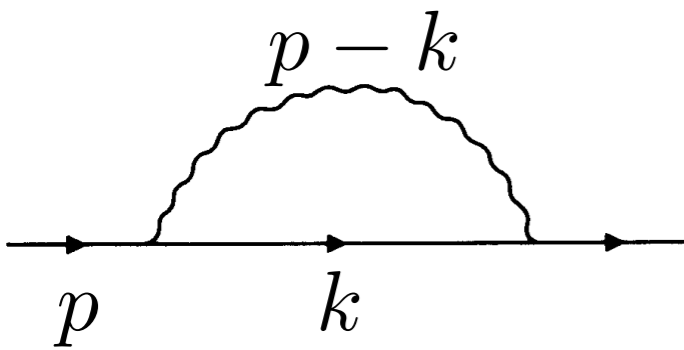
Electron self-energy



$$= -ie^2 \int_0^1 dx \int \frac{d^4 l^E}{(2\pi)^4} \frac{\gamma^\mu (\cancel{l}^E + x \cancel{p} + m) \gamma_\mu}{((l^E)^2 + \Delta - i\epsilon)^2}$$

$$l^2 = -l_0^2 - \mathbf{l}^2 = -l^E{}^2$$

Electron self-energy



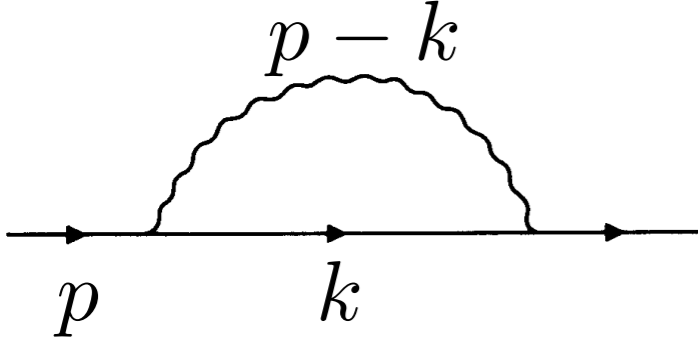
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Dimensional regularization

$$= -ie^2 \int_0^1 dx \gamma^\mu (x\cancel{p} + m) \gamma_\mu \int \frac{d^d l^E}{(2\pi)^d} \frac{1}{((l^E)^2 + \Delta)^2}$$

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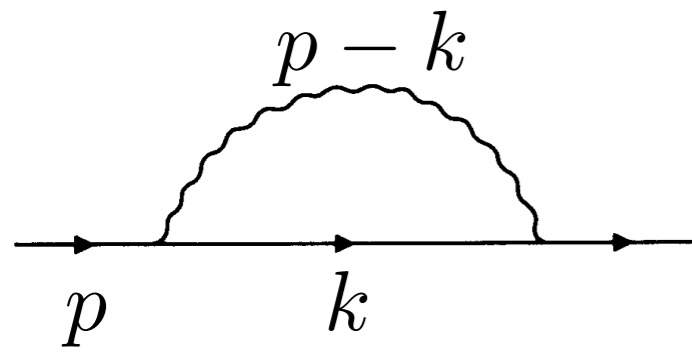
0

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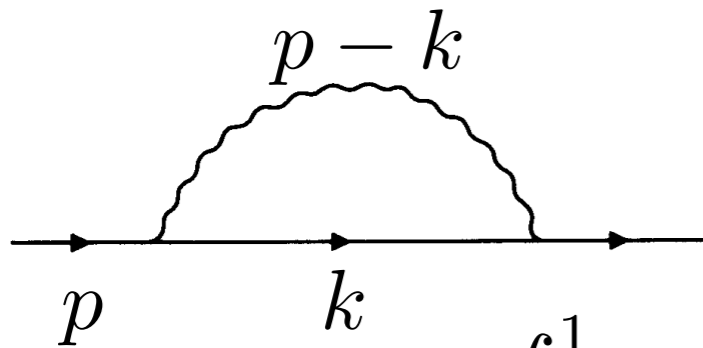
$$= -ie^2 \int_0^1 dx \gamma^\mu (x\cancel{p} + m) \gamma_\mu \int \frac{d^d l^E}{(2\pi)^d} \frac{1}{((l^E)^2 + \Delta)^2}$$

$$\int \frac{d\Omega_d}{(2\pi)^d} \frac{1}{2} \int_0^\infty dl^{E^2} \frac{(l^{E^2})^{d/2-1}}{(l^{E^2} + \Delta)^2} = \frac{1}{(4\pi)^{d/2} \Gamma(d/2)} \left(\frac{1}{\Delta}\right)^{2-d/2} \int_0^1 dz z^{1-d/2} (1-z)^{d/2-1}$$

$$z \equiv \frac{\Delta}{l^{E^2} + \Delta}$$

Electron self-energy

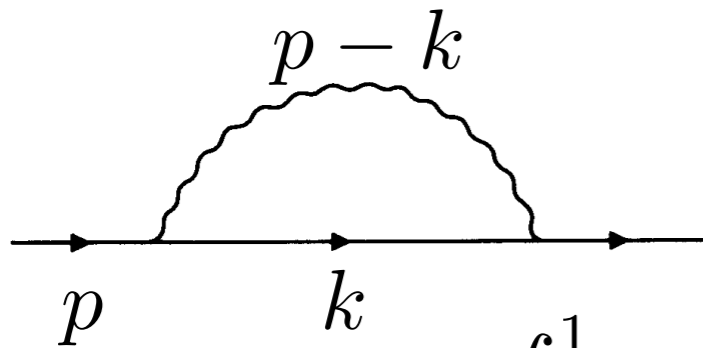
$$B(a, b) \equiv \frac{\Gamma(a)\Gamma(b)}{\Gamma(a+b)} \equiv \int dz z^{a-1} (1-z)^{b-1}$$



$$= -ie^2 \int_0^1 dx \gamma^\mu (\cancel{x}p + m) \gamma_\mu \frac{1}{(4\pi)^{d/2}} \frac{\Gamma(2-d/2)}{\Gamma(2)} \left(\frac{1}{\Delta}\right)^{2-d/2}$$

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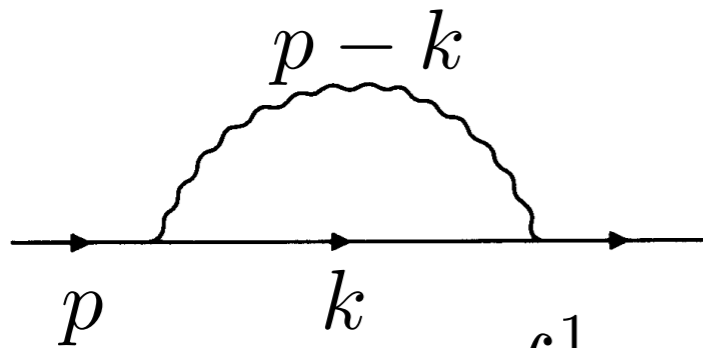
$$\frac{1}{(4\pi)^{d/2}} \frac{\Gamma(2-d/2)}{\Gamma(2)} \left(\frac{1}{\Delta}\right)^{2-d/2}$$

↓ $d \equiv 4 - 2\epsilon$

$$\frac{1}{(4\pi)^2} \left(\frac{1}{\epsilon} - \gamma + \log(4\pi) - \log \Delta \right)$$

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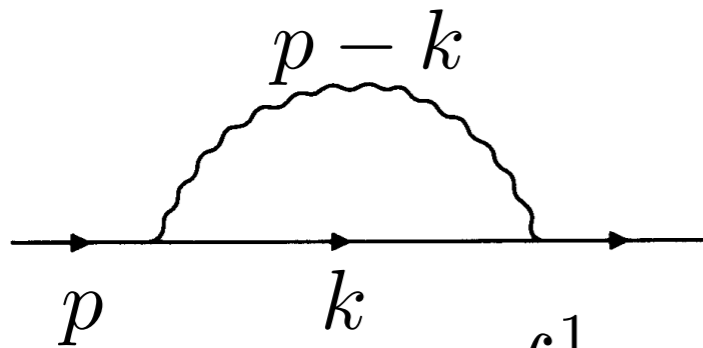
$$\frac{1}{(4\pi)^2}$$

$$\left(\frac{1}{\epsilon} - \gamma + \log(4\pi) - \log \Delta\right)$$

$$\frac{1}{\bar{\epsilon}}$$

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$$\frac{1}{\bar{\epsilon}}$$

$$= -i \frac{e^2}{(4\pi)^2} \frac{1}{\bar{\epsilon}} (-\cancel{p} + 4m) + \text{finite}$$

Renormalization

$$x_0 \rightarrow x + \delta x,$$

$$\phi_0 \rightarrow \left(1 + \frac{1}{2}\delta Z_{\phi\phi}\right)\phi + \sum_{\chi} \frac{1}{2}\delta Z_{\phi\chi}\chi,$$

$$\mathcal{L}_{Dirac} \rightarrow \bar{\psi} (\not{\partial} - m) (1 + \delta Z_{\psi\psi}) \psi - \delta m \bar{\psi}\psi$$

Renormalization

$$x_0 \rightarrow x + \delta x \quad \text{real constants}$$

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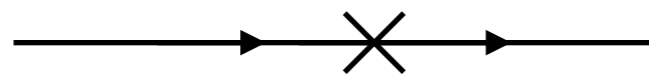
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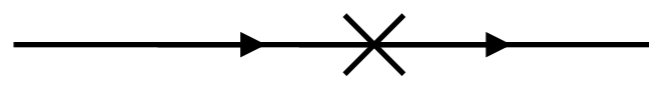
$$i (\not{p} - m) \delta Z_{\psi\psi} - i \delta m$$

Renormalization

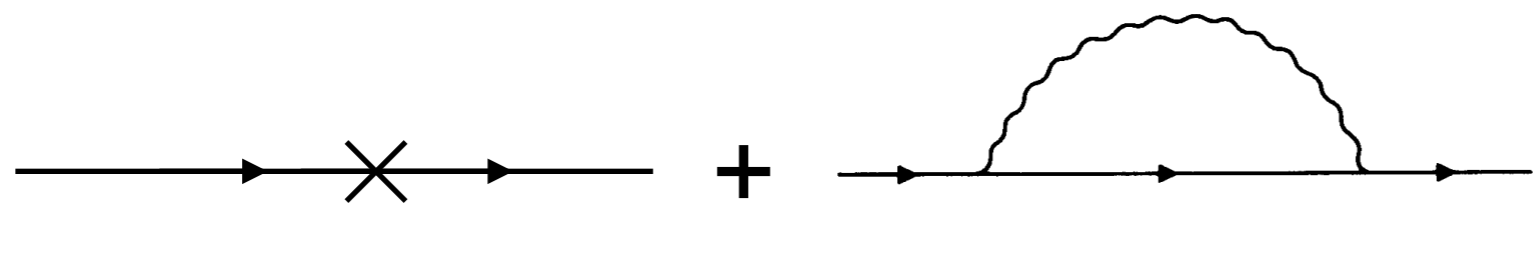
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$$= 0 \frac{1}{\bar{\epsilon}} + \text{finite}$$

$$\delta Z_{\psi\psi} = -\frac{e^2}{(4\pi)^2} \frac{1}{\bar{\epsilon}} + \dots$$

$$\delta m = -\frac{3e^2}{(4\pi)^2} \frac{1}{\bar{\epsilon}} + \dots$$

Renormalization

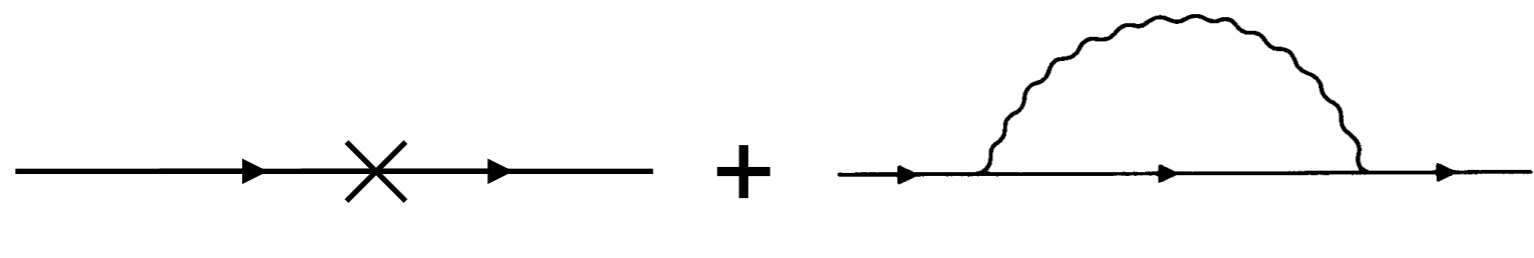
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**Arbitrary:
Renormalization
scheme**

Operator dimension

- Fermion fields : 3/2
- Boson fields : 1
- derivatives : 1

Dimension 4



$$\mathcal{L} = \frac{1}{2} (\partial_\mu \phi \partial^\mu \phi - m^2 \phi^2)$$
$$\mathcal{L} = \bar{\psi}(x) [i \not{\partial} - m] \psi(x)$$

$$D_\mu = \partial_\mu - ieQA_\mu$$

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Dimension 4


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Definition : dimension of the operator is the sum of the dimensions of its fields and derivatives

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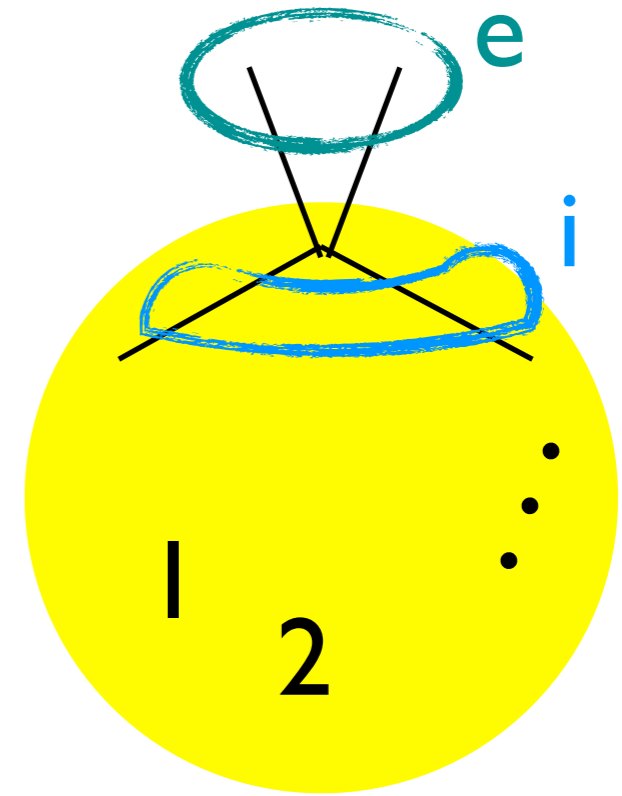
$d=2$
 $D_\mu = \partial_\mu - ieQA_\mu$

Definition : dimension of the operator is the sum of the dimensions of its fields and derivatives

Exclude the dimension of the coefficient!

Divergent amplitudes

Each vertex: p^{d-e-i}



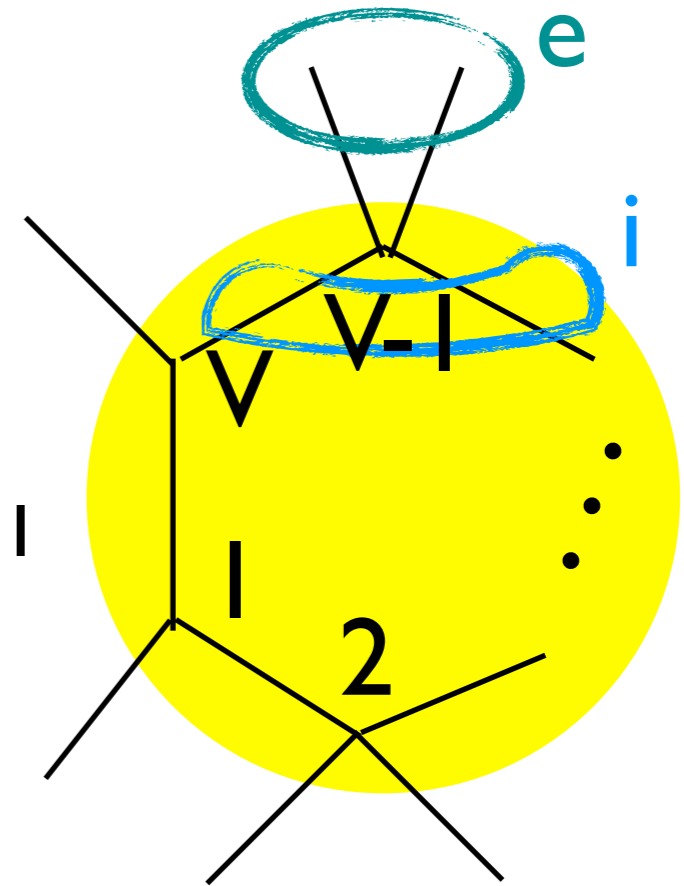
Divergent amplitudes

Each vertex: p^{d-e-i}

One loop: V propagators

Each internal fermion propagator: p^{-1}

Each internal boson propagator: p^{-2}



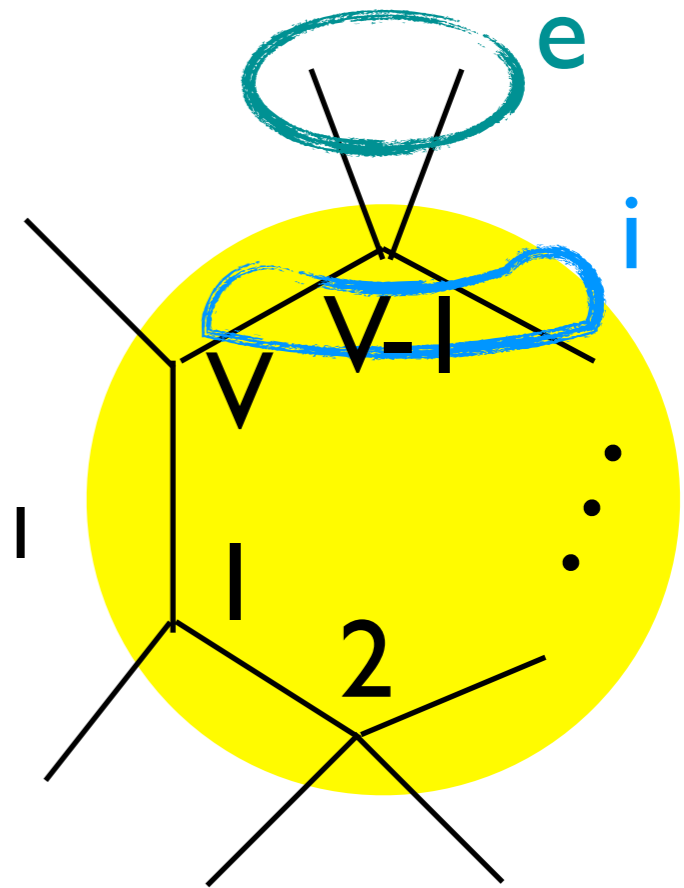
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$$I = \sum i = 3F + 2B \quad (F/B \text{ \# of internal fermions/bosons})$$

$$F + B = V$$

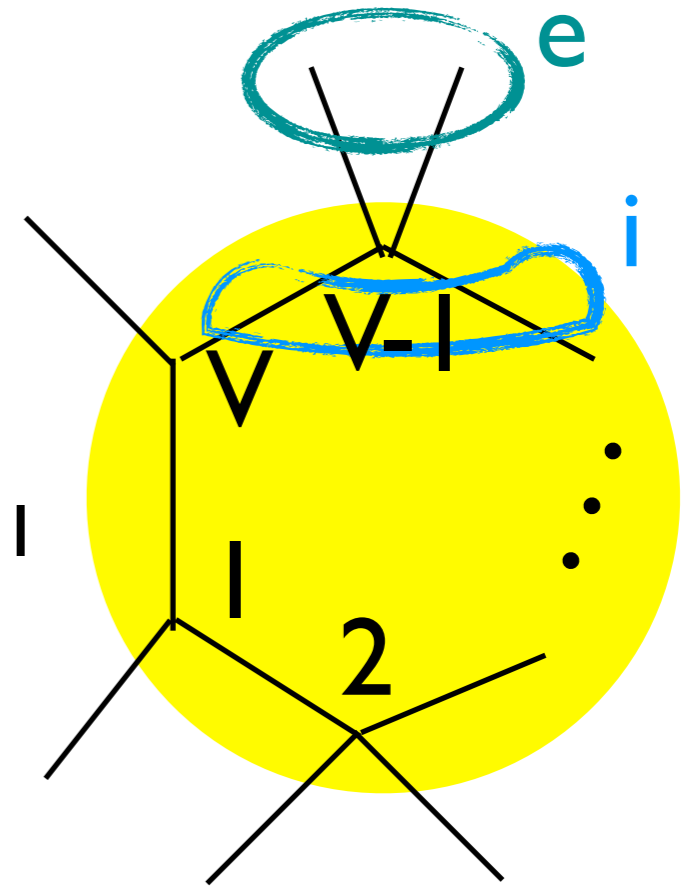
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$$F + B = V + (L - 1)$$

Each extra loop: one extra propagator

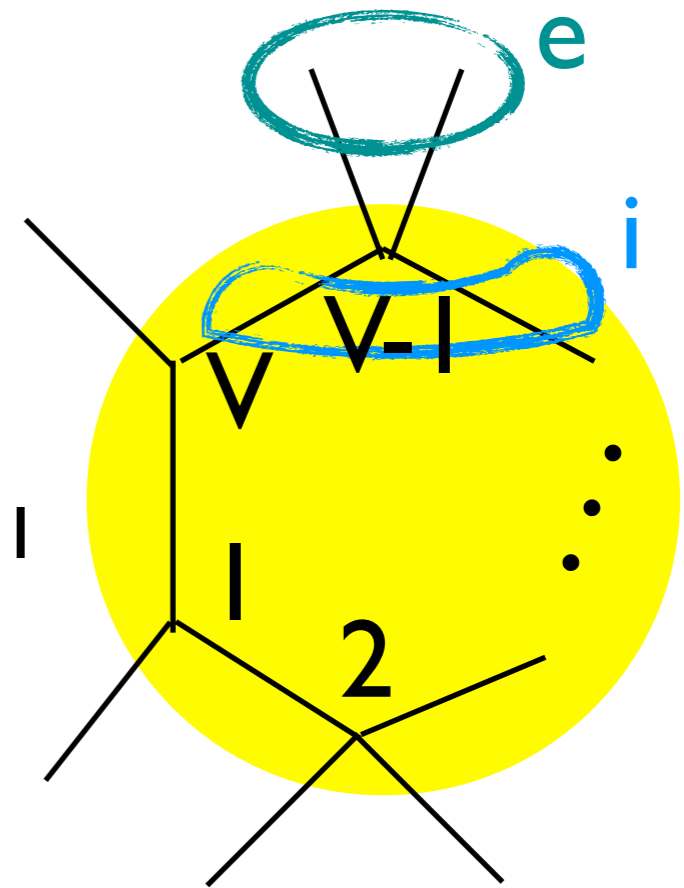
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Each extra loop: one extra propagator

$$A(D, E, L, V) \sim \left(\int d^4 p \right)^L p^{D-E-I-2B-F} \sim \left(\int d^4 p \right)^L p^{D-E-4V-4(L-1)}$$

Renormalization

$$A(D, E, L, V) \sim \left(\int d^4 p \right)^L p^{D-E-4V-4(L-1)}$$

Renormalizable model if $d \leq 4$ for all operators

$D \leq 4V$: $A \sim \Lambda^{4-E}$ diverges only if $E \leq 4$

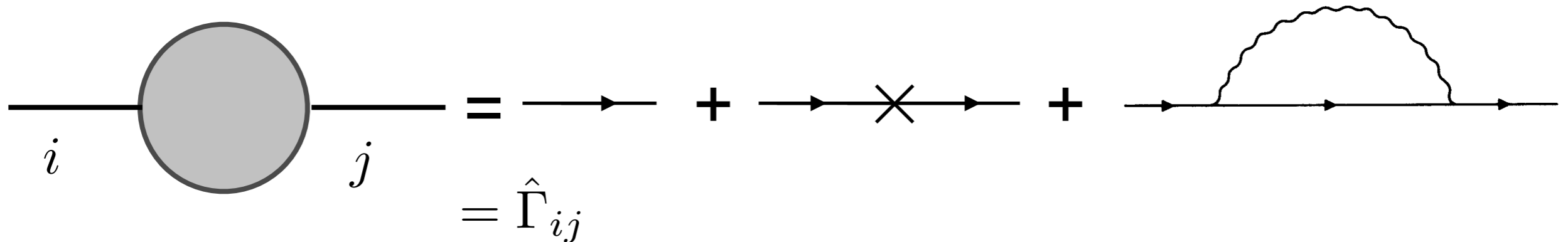
- All the divergence are absorbed by the fields and external parameters redefinitions
- Yang-mills theories
- With spontaneous symmetry breaking

On-shell scheme

Renormalized mass = Physical mass

Two-point function vanishes on-shell (No loop corrections on the external legs)

Cancel the mixings



$$\tilde{\mathcal{R}} \hat{\Gamma}_{ij}^f(p) u_j(p) \Big|_{p^2 = m_{f,j}^2} = 0,$$

$$\lim_{p^2 \rightarrow m_{f,i}^2} \frac{\not{p} + m_{f,i}}{p^2 - m_{f,i}^2} \tilde{\mathcal{R}} \hat{\Gamma}_{ii}^f(p) u_i(p) = i u_i(p),$$

Renormalization conditions

$$\tilde{\mathfrak{R}} \bar{u}_i(p') \hat{\Gamma}_{ij}^f(p') \Big|_{p'^2=m_{f,i}^2} = 0,$$

$$\lim_{p'^2 \rightarrow m_{f,i}^2} \bar{u}_i(p') \tilde{\mathfrak{R}} \hat{\Gamma}_{ii}^f(p') \frac{\not{p}' + m_{f,i}}{p'^2 - m_{f,i}^2} = i\bar{u}_i(p').$$

$$\hat{\Gamma}_{ij} = i\delta_{ij} (\not{p} - m_i) + i [f_{ij}^L(p^2) \not{p}\gamma_- + f_{ij}^R(p^2) \not{p}\gamma_+ + f_{ij}^{SL}(p^2) \gamma_- + f_{ij}^{SR}(p^2) \gamma_+]$$

$$\tilde{\mathfrak{R}} [f_{ij}^L(p^2) m_i + f_{ij}^{SR}(p^2)] \Big|_{p^2=m_i^2} = 0$$

$$\tilde{\mathfrak{R}} [f_{ij}^R(p^2) m_i + f_{ij}^{SL}(p^2)] \Big|_{p^2=m_i^2} = 0$$

$$\tilde{\mathfrak{R}} \left[2m_i \frac{\partial}{\partial p^2} [(f_{ii}^L(p^2) + f_{ii}^R(p^2)) m_i + f_{ii}^{SL}(p^2) + f_{ii}^{SR}(p^2)] + f_{ii}^L(p^2) + f_{ii}^R(p^2) \right] \Big|_{p^2=m_i^2} = 0$$

Similar for the vectors and scalars

Electron renormalization

$$f^L(p^2) = f^R(p^2) = \delta Z_{\psi\psi} - e^2 \int_0^1 dx (2-d)x \frac{1}{(4\pi)^{d/2}} \frac{\Gamma(\epsilon)}{\Gamma(2)} \left(\frac{1}{\Delta}\right)^\epsilon$$

$$f^{SL}(p^2) = f^{SR}(p^2) = -m\delta Z_{\psi\psi} - \delta m - e^2 \int_0^1 dx \frac{m d}{(4\pi)^{d/2}} \frac{\Gamma(\epsilon)}{\Gamma(2)} \left(\frac{1}{\Delta}\right)^\epsilon$$

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Gamma algebra in d dimension!

Electron renormalization

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$$\delta m = -\frac{e^2}{(4\pi)^2} \left(\frac{3}{\bar{\epsilon}} + 4 - 6 \log \left(\frac{m}{\mu} \right) \right)$$

Eqs 1 and 2

Electron renormalization

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Eqs 1 and 2

$$\Delta = (1-x)(m^2 - xp^2)$$

$$\left. \frac{\partial \Delta^{-\epsilon}}{\partial p^2} \right|_{p^2=m^2} = \epsilon x (1-x) \Delta^{-\epsilon-1} \Big|_{p^2=m^2} = \epsilon x (1-x)^{-1-2\epsilon} m^{-2-2\epsilon}$$

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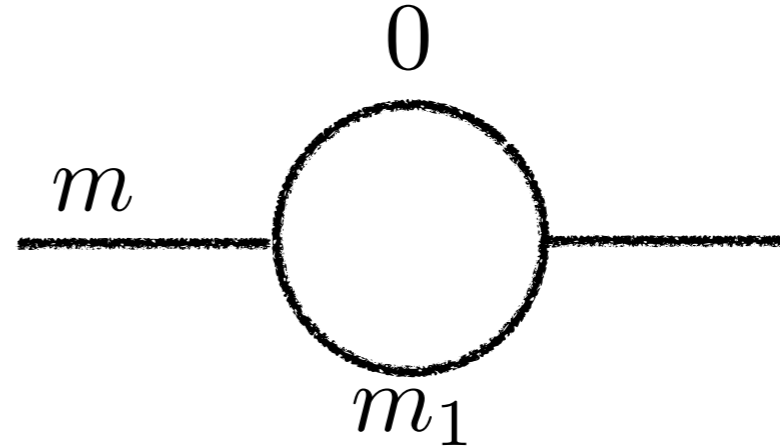
$$\left. \frac{\partial \Delta^{-\epsilon}}{\partial p^2} \right|_{p^2=m^2} = \epsilon x (1-x) \Delta^{-\epsilon-1} \Big|_{p^2=m^2} = \epsilon x (1-x)^{-1-2\epsilon} m^{-2-2\epsilon}$$

After integration over x

$$\frac{\epsilon \Gamma(2) \Gamma(-2\epsilon)}{\Gamma(2-2\epsilon)}$$

$$\delta Z_{\psi\psi} = -\frac{e^2}{(4\pi)^2} \left(\frac{1}{\bar{\epsilon}_{UV}} + \frac{2}{\bar{\epsilon}_{IR}} + 4 - 6 \log \left(\frac{m}{\mu} \right) \right)$$

Real/Complex masses



Real masses

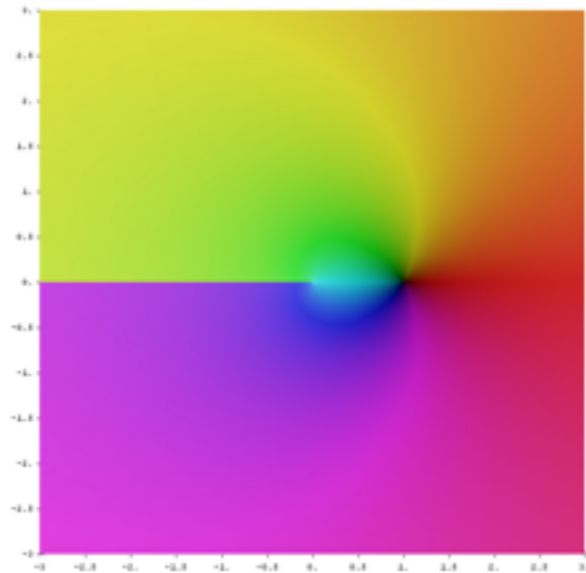
$$m_1^2 > m^2$$

$$\Re (\log [m_1^2 - p^2]) \Big|_{p^2=m^2}$$

$$m_1^2 < m^2$$

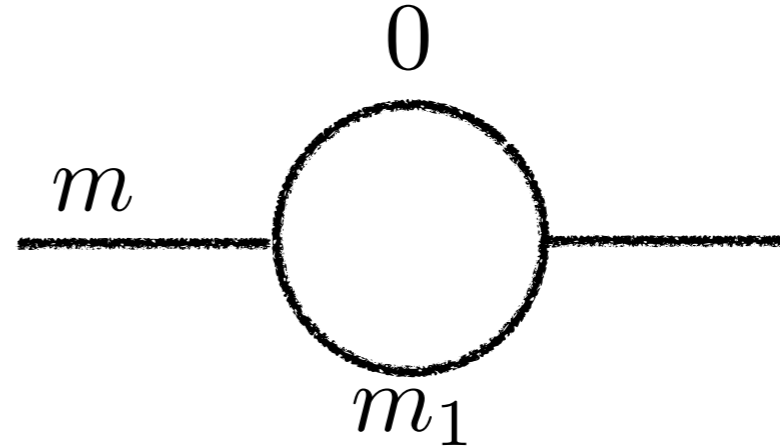
$$\Re (\log [p^2 - m_1^2] + i\pi) \Big|_{p^2=m^2}$$

Mass corrections are complex if the particle can decay



Epsilon prescription of the propagator

Real/Complex masses

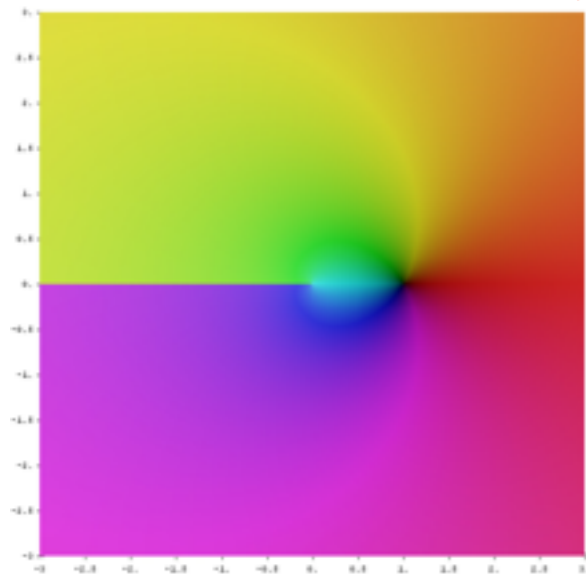


Real masses

$$m_1^2 > m^2 \quad \cancel{\Re}(\log [m_1^2 - p^2]) \Big|_{p^2=m^2}$$

$$m_1^2 < m^2 \quad \cancel{\Re}(\log [p^2 - m_1^2] + i\pi) \Big|_{p^2=m^2}$$

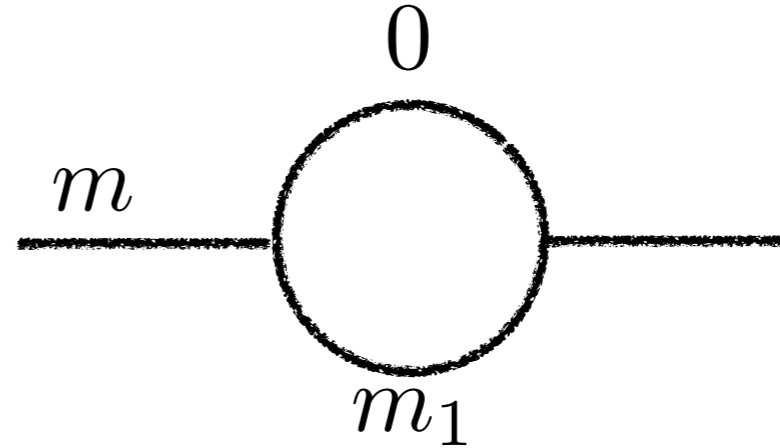
Mass corrections are complex if the particle can decay



$$m \in \mathbb{R}$$

Epsilon prescription of the propagator

Real/Complex masses

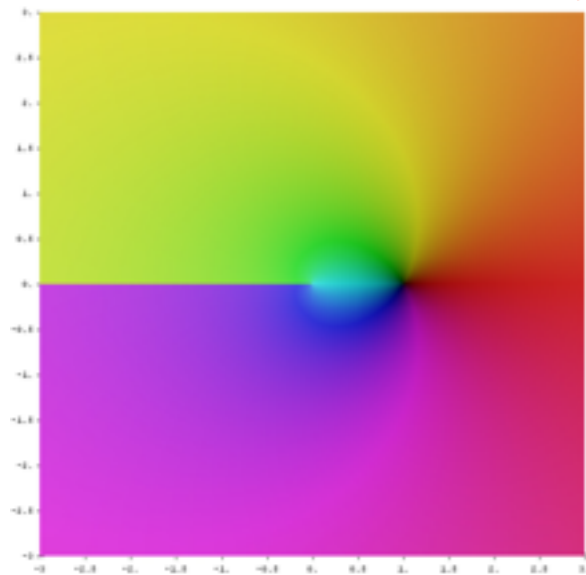


Real masses

$$m_1^2 > m^2 \quad \cancel{\Re}(\log [m_1^2 - p^2]) \Big|_{p^2=m^2}$$

$$m_1^2 < m^2 \quad \cancel{\Re}(\log [p^2 - m_1^2] + i\pi) \Big|_{p^2=m^2}$$

Mass corrections are complex if the particle can decay



$$m \in \mathbb{R}$$

Epsilon prescription of the propagator

Complex masses

$$\log [m_1^2 - p^2] \Big|_{p^2=m^2}$$

$$m^2 \rightarrow m^2 - im\Gamma$$

Complex mass scheme

$$\left. \begin{array}{l} \phi_0 \rightarrow (1 + \frac{1}{2}\delta Z_{\phi\phi})\phi \\ \phi_0^\dagger \rightarrow (1 + \frac{1}{2}\delta Z_{\phi\phi}^*)\phi^\dagger \end{array} \right\} \Rightarrow \partial^\mu \phi_0 \partial_\mu \phi_0^\dagger \rightarrow (1 + \Re \delta Z_{\phi\phi}) \partial^\mu \phi \partial_\mu \phi^\dagger$$

Complex mass scheme

$$\left. \begin{array}{l} \phi_0 \rightarrow (1 + \frac{1}{2}\delta Z_{\phi\phi})\phi \\ \phi_0^\dagger \rightarrow (1 + \frac{1}{2}\delta Z_{\phi\phi}^*)\phi^\dagger \end{array} \right\} \Rightarrow \partial^\mu \phi_0 \partial_\mu \phi_0^\dagger \rightarrow (1 + \Re \delta Z_{\phi\phi}) \partial^\mu \phi \partial_\mu \phi^\dagger$$

$$\left. \begin{array}{l} \phi_0 \rightarrow (1 + \frac{1}{2}\delta Z_{\phi\phi})\phi \\ \phi_0^\dagger \rightarrow (1 + \frac{1}{2}\delta Z_{\phi\phi})\phi^\dagger \end{array} \right\} \Rightarrow \partial^\mu \phi_0 \partial_\mu \phi_0^\dagger \rightarrow (1 + \delta Z_{\phi\phi}) \partial^\mu \phi \partial_\mu \phi^\dagger$$

Complex mass scheme

~~$$\left. \begin{aligned} \phi_0 &\rightarrow (1 + \frac{1}{2}\delta Z_{\phi\phi})\phi \\ \phi_0^\dagger &\rightarrow (1 + \frac{1}{2}\delta Z_{\phi\phi}^*)\phi^\dagger \end{aligned} \right\} \Rightarrow \partial^\mu \phi_0 \partial_\mu \phi_0^\dagger \rightarrow (1 + \Re \delta Z_{\phi\phi}) \partial^\mu \phi \partial_\mu \phi^\dagger$$~~

$$\left. \begin{aligned} \phi_0 &\rightarrow (1 + \frac{1}{2}\delta Z_{\phi\phi})\phi \\ \phi_0^\dagger &\rightarrow (1 + \frac{1}{2}\delta Z_{\phi\phi})\phi^\dagger \end{aligned} \right\} \Rightarrow \partial^\mu \phi_0 \partial_\mu \phi_0^\dagger \rightarrow (1 + \delta Z_{\phi\phi}) \partial^\mu \phi \partial_\mu \phi^\dagger$$

$$m \rightarrow m + \delta m$$

Now with complex masses

Complex mass scheme

$$\left. \begin{aligned} \phi_0 &\rightarrow (1 + \frac{1}{2}\delta Z_{\phi\phi})\phi \\ \phi_0^\dagger &\rightarrow (1 + \frac{1}{2}\delta Z_{\phi\phi}^*)\phi^\dagger \end{aligned} \right\} \Rightarrow \partial^\mu \phi_0 \partial_\mu \phi_0^\dagger \rightarrow (1 + \Re \delta Z_{\phi\phi}) \partial^\mu \phi \partial_\mu \phi^\dagger$$

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$$m \rightarrow m + \delta m$$

Now with complex masses

Hermitian

$$\mathcal{L}_0 = \mathcal{L} + \delta\mathcal{L}$$

Not hermitian

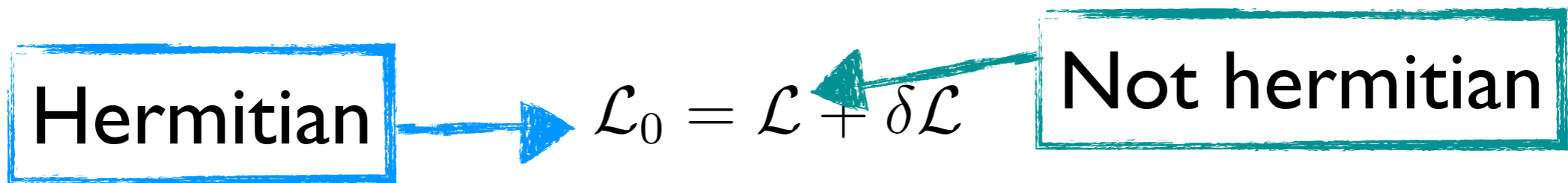
Complex mass scheme

$$\left. \begin{aligned} \phi_0 &\rightarrow (1 + \frac{1}{2}\delta Z_{\phi\phi})\phi \\ \phi_0^\dagger &\rightarrow (1 + \frac{1}{2}\delta Z_{\phi\phi}^*)\phi^\dagger \end{aligned} \right\} \Rightarrow \partial^\mu \phi_0 \partial_\mu \phi_0^\dagger \rightarrow (1 + \Re \delta Z_{\phi\phi}) \partial^\mu \phi \partial_\mu \phi^\dagger$$

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$$m \rightarrow m + \delta m$$

Now with complex masses



$$\cancel{\Re} [f_{ij}^L(p^2) m_i + f_{ij}^{SR}(p^2)] \Big|_{p^2=m_i^2} = 0$$

$$\cancel{\Re} [f_{ij}^R(p^2) m_i + f_{ij}^{SL}(p^2)] \Big|_{p^2=m_i^2} = 0$$

$$\cancel{\Re} \left[2m_i \frac{\partial}{\partial p^2} [(f_{ii}^L(p^2) + f_{ii}^R(p^2)) m_i + f_{ii}^{SL}(p^2) + f_{ii}^{SR}(p^2)] + f_{ii}^L(p^2) + f_{ii}^R(p^2) \right] \Big|_{p^2=m_i^2} = 0$$

Similar for the vectors and scalars

Renormalization conditions

Zero momentum scheme available for the gauge couplings

$$\Gamma_{FFV}^\mu(p_1, p_2) = igT^a \delta_{f_1, f_2} \left[\gamma^\mu \left(\frac{\delta g}{g} + \frac{1}{2} \delta Z_{VV} + \frac{1}{2} \delta Z_{FF}^R + \frac{1}{2} \delta Z_{FF}^L + \frac{g'_V}{2g} \delta Z_{V'V} \right) \right. \\ \left. + \gamma^\mu \gamma_5 \left(\frac{1}{2} \delta Z_{FF}^R - \frac{1}{2} \delta Z_{FF}^L + \frac{g'_A}{2g} \delta Z_{V'V} \right) \right. \\ \left. + \left(\gamma^\mu h^V(k^2) + \gamma^\mu \gamma_5 h^A(k^2) + \frac{(p_1 - p_2)^\mu}{2m} h^S(k^2) + \frac{k_\mu}{2m} h^P(k^2) \right) \right]$$



$$\frac{\delta g}{g} + \frac{1}{2} \delta Z_{VV} + \frac{1}{2} \delta Z_{FF}^R + \frac{1}{2} \delta Z_{FF}^L + \frac{g'_V}{2g} \delta Z_{V'V} + h^V(0) + h^S(0) = 0 \\ \frac{1}{2} \delta Z_{FF}^R - \frac{1}{2} \delta Z_{FF}^L + \frac{g'_A}{2g} \delta Z_{V'V} + h^A(0) = 0.$$

By gauge invariance



$$\frac{\delta g}{g} + \frac{1}{2} \delta Z_{VV} + \frac{g'_V}{2g} \delta Z_{V'V} + \frac{g'_A}{2g} \delta Z_{V'V} = 0$$

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$$\frac{\delta g}{g} + \frac{1}{2} \delta Z_{VV} + \frac{1}{2} \delta Z_{FF}^R + \frac{1}{2} \delta Z_{FF}^L + \frac{g'_V}{2g} \delta Z_{V'V} + h^V(0) + h^S(0) = 0 \\ \frac{1}{2} \delta Z_{FF}^R - \frac{1}{2} \delta Z_{FF}^L + \frac{g'_A}{2g} \delta Z_{V'V} + h^A(0) = 0.$$

By gauge invariance

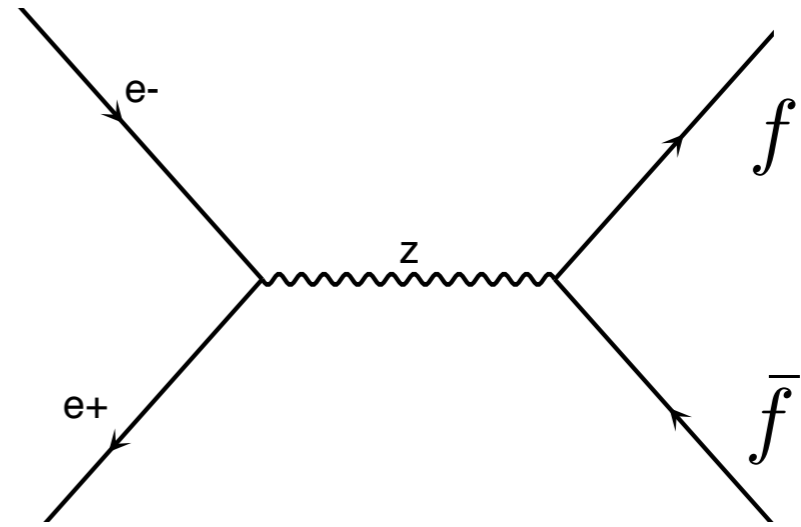
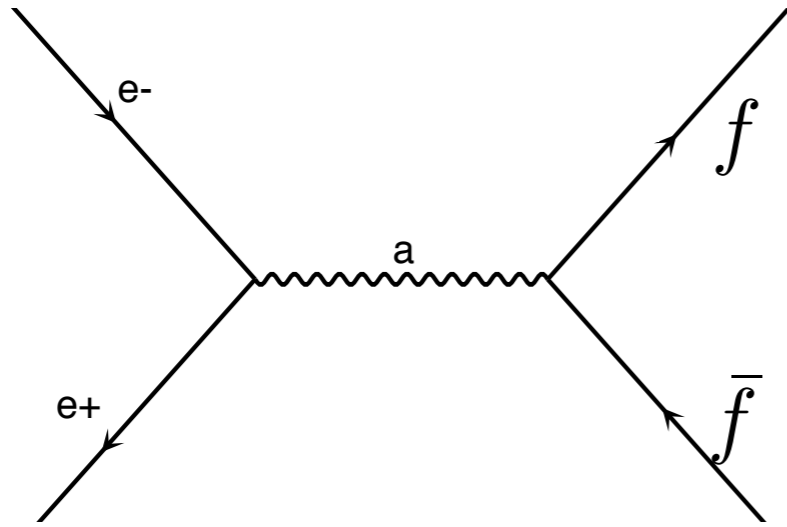
$$\frac{\delta g}{g} + \frac{1}{2} \delta Z_{VV} + \frac{g'_V}{2g} \delta Z_{V'V} + \frac{g'_A}{2g} \delta Z_{V'V} = 0$$

Only from
two-point
functions

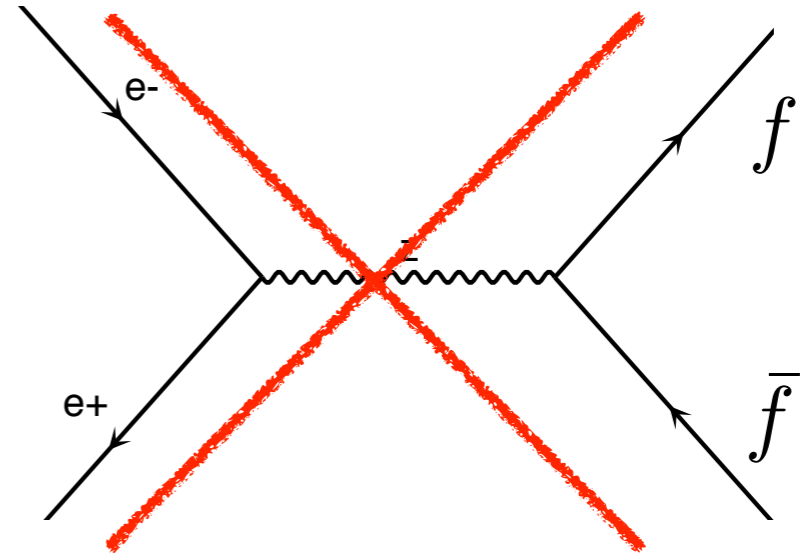
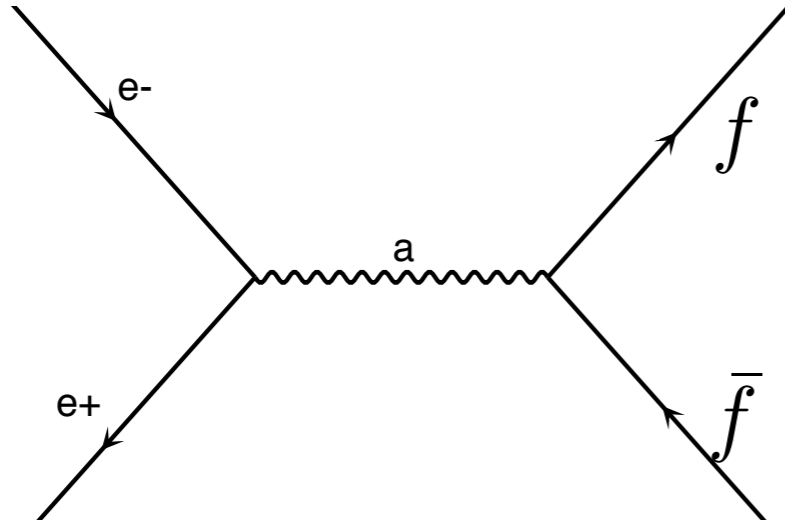
Plan

- Renormalization
- KLN Theorem
- Rational Terms
- FeynRules at NLO

Infrared divergences

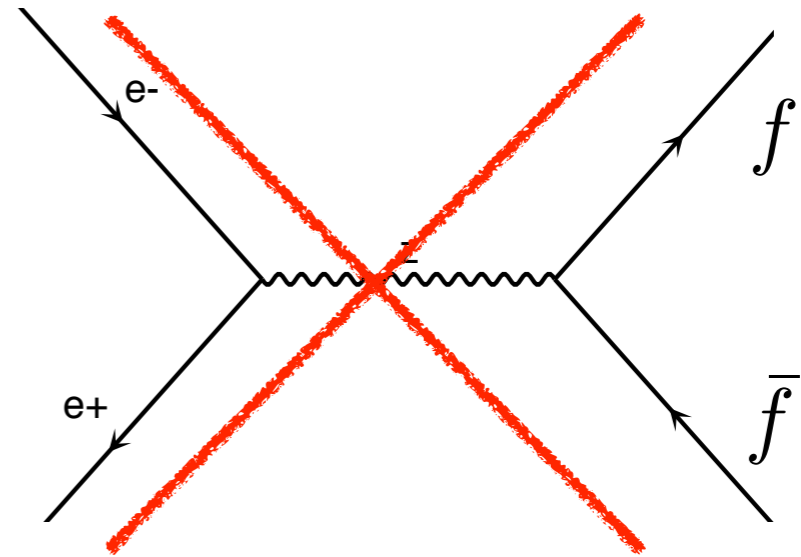
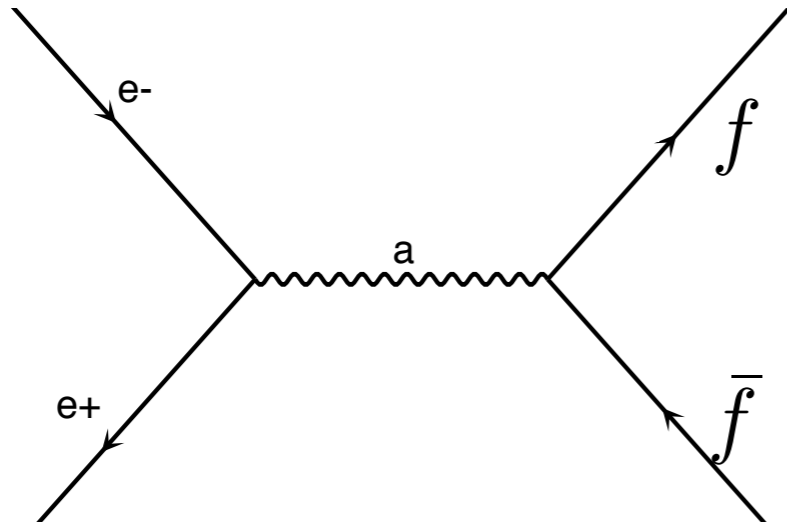


Infrared divergences



Well below the Z mass

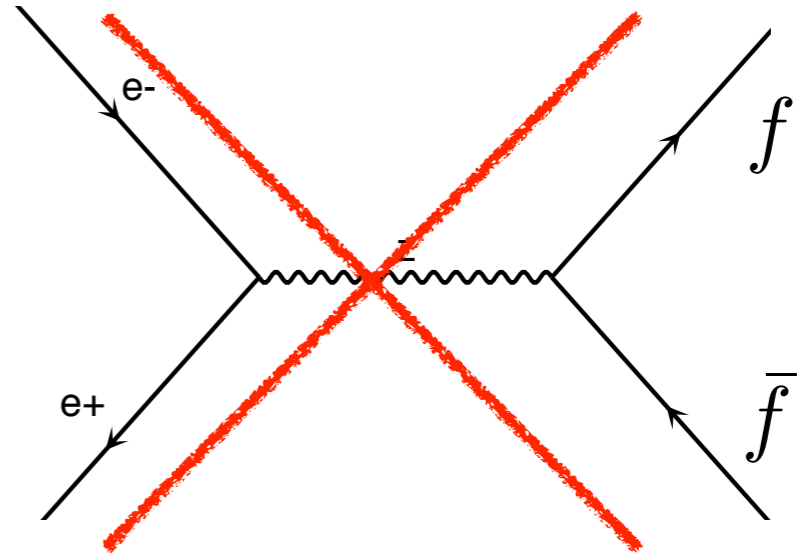
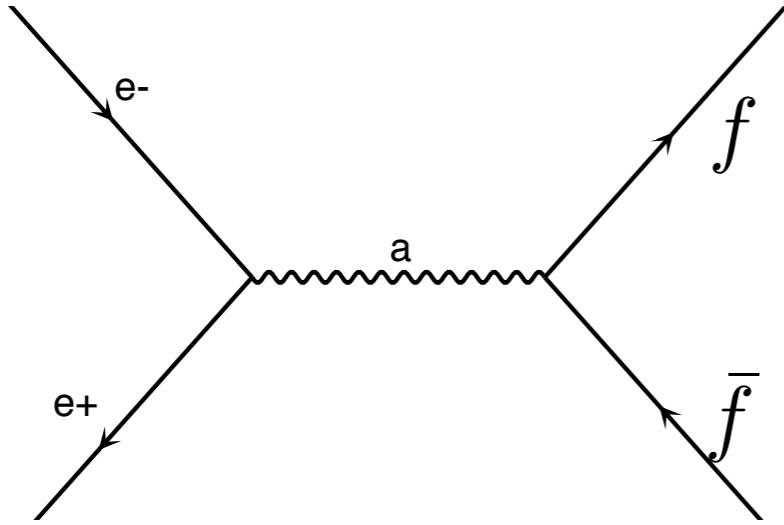
Infrared divergences



Well below the Z mass

$$\sigma_0 = \frac{4\pi\alpha^2}{3s} Q_f^2$$

Infrared divergences



Well below the Z mass

$$\sigma_0 = \frac{4\pi\alpha^2}{3s} Q_f^2$$

$$R \equiv \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)} = \frac{\sum_q \sigma(e^+e^- \rightarrow q\bar{q})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)} = 3 \sum_q Q_q^2$$

KLN illustration

Virtual corrections

$$\sigma^V = \sigma_0 \frac{2\alpha_S}{\pi} H(\epsilon) \left[-\frac{2}{\epsilon^2} + \frac{3}{\epsilon} - 8 + \pi^2 + \mathcal{O}(\epsilon) \right]$$

where
$$H(\epsilon) = \frac{3(1 + \epsilon)}{(3 + 2\epsilon)\Gamma(2 + 2\epsilon)(4\pi)^{2\epsilon}} = 1 + \mathcal{O}(\epsilon)$$

Real corrections

$$\sigma^{q\bar{q}g} = \sigma_0 \frac{2\alpha_S}{\pi} H(\epsilon) \left[\frac{2}{\epsilon^2} - \frac{3}{\epsilon} + \frac{19}{2} - \pi^2 + \mathcal{O}(\epsilon) \right]$$

$$R = 3 \sum_q Q_q^2 \left\{ 1 + \frac{\alpha_S}{\pi} + \mathcal{O}(\alpha_S^2) \right\}$$

Pole cancellation provides a check

Plan

- Renormalization
- KLN Theorem
- Rational Terms
- FeynRules at NLO

Loop computation

$$\begin{aligned} \mathcal{A}^{1-loop} &= \sum_i d_i \text{Box}_i + \sum_i c_i \text{Triangle}_i + \sum_i b_i \text{Bubble}_i \\ &+ \sum_i a_i \text{Tadpole}_i + R \end{aligned}$$

- Box, Triangle, Bubble and Tadpole are known scalar integrals
- Loop computation = find the coefficients
 - Unitarity
 - Multiple cuts
 - Tensor reduction (OPP)

R₂

$$\bar{A}(\bar{q}) = \frac{1}{(2\pi)^4} \int d^d \bar{q} \frac{\bar{N}(\bar{q})}{\bar{D}_0 \bar{D}_1 \dots \bar{D}_{m-1}}, \quad \bar{D}_i = (\bar{q} + p_i)^2 - m_i^2$$

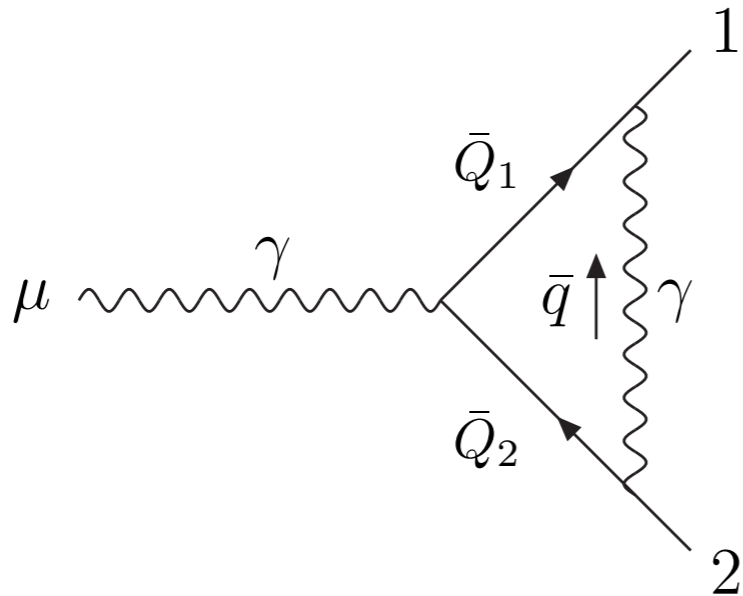
$$\bar{N}(\bar{q}) = N(q) + \tilde{N}(\tilde{q}, q, \epsilon)$$

d 4 ε

$$R_2 \equiv \lim_{\epsilon \rightarrow 0} \frac{1}{(2\pi)^4} \int d^d \bar{q} \frac{\tilde{N}(\tilde{q}, q, \epsilon)}{\bar{D}_0 \bar{D}_1 \dots \bar{D}_{m-1}}$$

Finite set of vertices that can be computed once
for all

R₂ example



$$\bar{Q}_1 = \bar{q} + p_1 = Q_1 + \tilde{q}$$

$$\bar{Q}_2 = \bar{q} + p_2 = Q_2 + \tilde{q}$$

$$\bar{D}_0 = \bar{q}^2$$

$$\bar{D}_1 = (\bar{q} + p_1)^2$$

$$\bar{D}_2 = (\bar{q} + p_2)^2$$

't Hooft Veltman
scheme

$$\bar{\eta}^{\bar{\mu}\bar{\nu}} \bar{\eta}_{\bar{\mu}\bar{\nu}} = d,$$

$$\bar{\gamma}^{\bar{\mu}} \bar{\gamma}_{\bar{\mu}} = d \mathbb{1},$$

$$\begin{aligned} \bar{N}(\bar{q}) &\equiv e^3 \left\{ \bar{\gamma}_{\bar{\beta}} (\bar{Q}_1 + m_e) \gamma_{\mu} (\bar{Q}_2 + m_e) \bar{\gamma}^{\bar{\beta}} \right\} \\ &= e^3 \left\{ \gamma_{\beta} (Q_1 + m_e) \gamma_{\mu} (Q_2 + m_e) \gamma^{\beta} \right. \\ &\quad \left. - \epsilon (Q_1 - m_e) \gamma_{\mu} (Q_2 - m_e) + \epsilon \tilde{q}^2 \gamma_{\mu} - \tilde{q}^2 \gamma_{\beta} \gamma_{\mu} \gamma^{\beta} \right\} \end{aligned}$$

$$R_2 = -\frac{ie^3}{8\pi^2} \gamma_{\mu}$$

$$\int d^n \bar{q} \frac{\tilde{q}^2}{\bar{D}_0 \bar{D}_1 \bar{D}_2} = -\frac{i\pi^2}{2} + \mathcal{O}(\epsilon),$$

$$\int d^n \bar{q} \frac{q_{\mu} q_{\nu}}{\bar{D}_0 \bar{D}_1 \bar{D}_2} = -\frac{i\pi^2}{2\epsilon} g_{\mu\nu} + \mathcal{O}(1)$$

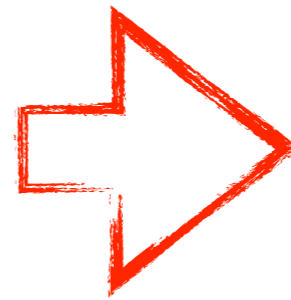
R₁

Due to the ϵ dimensional parts of the denominators

Like for the 4 dimensional part but with a different set of integrals

$$\int d^n \bar{q} \frac{\tilde{q}^2}{\bar{D}_i \bar{D}_j} = -\frac{i\pi^2}{2} \left[m_i^2 + m_j^2 - \frac{(p_i - p_j)^2}{3} \right] + \mathcal{O}(\epsilon),$$
$$\int d^n \bar{q} \frac{\tilde{q}^2}{\bar{D}_i \bar{D}_j \bar{D}_k} = -\frac{i\pi^2}{2} + \mathcal{O}(\epsilon),$$
$$\int d^n \bar{q} \frac{\tilde{q}^4}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_l} = -\frac{i\pi^2}{6} + \mathcal{O}(\epsilon).$$

Only $R = R_1 + R_2$ is gauge invariant



Check

Plan

- Renormalization
- KLN Theorem
- Rational Terms
- FeynRules at NLO

FR@NLO

- Goal : Automate the one-loop computation for BSM models
- Required ingredients :
 - Tree-level vertices
 - R2 vertices (OPP)
 - UV counterterm vertices
- Solution : UFO at NLO

FR@NLO

- Goal : Automate the one-loop computation for BSM models
- Required ingredients :
 - Tree-level vertices Done(FeynRules)
 - R2 vertices (OPP)
 - UV counterterm vertices
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FR@NLO

- Goal : Automate the one-loop computation for BSM models
- Required ingredients :

- Tree-level vertices

Done(FeynRules)

- R2 vertices (OPP)

- UV counterterm vertices

Missing

- Solution : UFO at NLO

FR@NLO

- Goal : Automate the one-loop computation for BSM models
- Required ingredients :

- Tree-level vertices

Done(FeynRules)

- R2 vertices (OPP)

- UV counterterm vertices

Missing

- Solution : UFO at NLO

UV counterterm vertices

External parameters

$$\begin{aligned}x_0 &\rightarrow x + \delta x, \\ \phi_0 &\rightarrow \left(1 + \frac{1}{2}\delta Z_{\phi\phi}\right)\phi + \sum_{\chi} \frac{1}{2}\delta Z_{\phi\chi}\chi.\end{aligned}$$

Same for the conjugate field

One renormalization constant for each fermion chirality

Internal parameters are renormalised by replacing the external parameters in their expressions

$$\mathcal{L}_0 = \mathcal{L} + \delta\mathcal{L} \leftarrow \text{vertices after solving the reno. cond.}$$

How does it work?

FeynRules

Renormalize the Lagrangian



model.mod
model.gen

FeynArts

Write the amplitudes

NLOCT.m

Compute the NLO vertices

How does it work?

FeynRules
Renormalize the Lagrangian

model.mod
model.gen

model.nlo

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Compute the NLO vertices

How does it work?

FeynRules
Renormalize the Lagrangian

model.mod
model.gen

FeynArts
Write the amplitudes

NLOCT.m
Compute the NLO vertices

model.nlo



How does it work?

FeynRules :

...

```
Lren = OnShellRenormalization[ LSM , QCDOnly ->True];  
WriteFeynArtsOutput[ Lren , Output -> "SMrenoL",  
GenericFile -> False]
```

How does it work?

FeynRules :

...

```
Lren = OnShellRenormalization[ LSM , QCDOOnly ->True];  
WriteFeynArtsOutput[ Lren , Output -> "SMrenoL",  
GenericFile -> False]
```

FeynArts / NLOCT :

```
WriteCT[ "SMrenoL/SMrenoL" , "Lorentz", Output->  
"SMQCDreno", QCDOonly -> True]
```

How does it work?

FeynRules :

...

```
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GenericFile -> False]
```

FeynArts / NLOCT :

```
WriteCT[ "SMrenoL/SMrenoL" , "Lorentz", Output->  
"SMQCDreno", QCDOonly -> True]
```

How does it work?

FeynRules :

...

```
Lren = OnShellRenormalization[ LSM , QCDDOnly -> True];  
WriteFeynArtsOutput[ Lren , Output -> "SMrenoL",  
GenericFile -> False]
```

FeynArts / NLOCT :

```
WriteCT[ "SMrenoL/SMrenoL" , "Lorentz", Output->  
"SMQCDreno", QCDDOnly -> True]
```

How does it work?

FeynRules :

...

```
Lren = OnShellRenormalization[ LSM , QCDOOnly ->True];  
WriteFeynArtsOutput[ Lren , Output -> "SMrenoL",  
GenericFile -> False]
```

FeynArts / NLOCT :

```
WriteCT[ "SMrenoL/SMrenoL" , "Lorentz", Output->  
"SMQCDreno", QCDOOnly -> True]
```

FeynRules :

...

```
Get["SMQCDreno.nlo"];  
WriteUFO[ LSM , UVCounterterms -> UV$vertlist ,  
R2Vertices -> R2$vertlist]
```

How does it work?

FeynRules :

...

```
Lren = OnShellRenormalization[ LSM , QCDOOnly ->True];  
WriteFeynArtsOutput[ Lren , Output -> "SMrenoL",  
GenericFile -> False]
```

FeynArts / NLOCT :

```
WriteCTI "SMrenoL/SMrenoL" , "Lorentz", Output->  
"SMQCDreno", QCDOOnly -> True]
```

FeynRules :

...

```
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WriteUFO[ LSM , UVCounterterms -> UV$vertlist ,  
R2Vertices -> R2$vertlist]
```

model.nlo

Model information (FR+FeynArts model/generic files)

```
R2$vertlist = {  
  {{{anti[u], 1}, {u, 2}}, ((-I/12)*gs^2*  
  IndexDelta[Index[Colour, Ext[1]], Index[Colour, Ext[2]]]*IPL[{u, G}]*  
  (TensDot[SlashedP[2], ProjM][Index[Spin, Ext[1]], Index[Spin, Ext[2]]] +  
  TensDot[SlashedP[2], ProjP][Index[Spin, Ext[1]], Index[Spin, Ext[2]]]))/Pi^2},  
  ...  
}
```

~FeynRules syntaxe

UV\$vertlist (ϵ is FR\$Eps)

```
FR$InteractionOrderPerturbativeExpansion = {{QCD, 1}, {QED, 0}};
```

NLOCT\$assumptions


QCDOnly

WriteCT[... , Assumptions->{...}]

UFO@NLO



- CT_vertices.py

```
V_1 = CTVertex(name = 'V_1',  
              type = 'R2',  
              particles = [ P.g, P.g, P.g ],  
              color = [ 'f(1,2,3)' ],  
              lorentz = [ L.VVV2 ],  
              loop_particles = [ [ [P.b], [P.c], [P.d], [P.s], [P.t], [P.u] ], [ [P.g] ] ],  
              couplings = {(0,0,0):C.R2GC_273_53,(0,0,1):C.R2GC_273_54})
```



- CT_couplings.py

```
UVGC_271_34 = Coupling(name = 'UVGC_271_34',  
                      value = {-1:'( 0 if MB else -(complex(0,1)*G**2)/(24.*cmath.pi**2) ) +  
                                (complex(0,1)*G**2)/(24.*cmath.pi**2)', 0:'( -(complex(0,1)*G**2*reglog(MB/MU_R))/  
                                (12.*cmath.pi**2) if MB else 0 )'},  
                      order = {'QCD':2})
```



- In coupling_order.py

```
QCD = CouplingOrder(name = 'QCD',  
                   expansion_order = 99,  
                   hierarchy = 1,  
                   perturbative_expansion = 1)
```

```
QED = CouplingOrder(name = 'QED',  
                   expansion_order = 99,  
                   hierarchy = 2)
```

Restrictions/Assumptions

- Renormalizable Lagrangian, maximum dimension of the operators is 4
- Feynman Gauge
- $\{\gamma_\mu, \gamma_5\} = 0$
- 't Hooft-Veltman scheme
- On-shell/complex mass scheme for the masses and wave functions
- \overline{MS} by default for everything else (zero-momentum possible for fermion gauge boson interaction)

NLOCT

- Amplitudes from FeynArts (discard irrelevant diagrams like ghost boxes)
- Compute terms at the generic level

$$\vec{c} \cdot \vec{L} = \sum_i c_i L_i$$

- Feynman parameters
- Remove terms with an odd or too low rank
- Gather loop momentum

$$q^\mu q^\nu q^\rho q^\sigma \rightarrow q^4 \frac{1}{d(d+2)} (\eta^{\mu\nu} \eta^{\rho\sigma} + \eta^{\mu\rho} \eta^{\nu\sigma} + \eta^{\mu\sigma} \eta^{\rho\nu})$$

$$q^\mu q^\nu \rightarrow q^2 \frac{1}{d} \eta^{\mu\nu}.$$

NLOCT

- Replace momentum integrals

$$\begin{array}{l}
 \int d^d q \frac{\epsilon}{q^2 - m^2} \Big|_{R_2} = i\pi^2 m^2, \\
 \int d^d q \frac{\epsilon}{(q^2 - \Delta)^2} \Big|_{R_2} = i\pi^2, \\
 \int d^d q \frac{q^2 (a\epsilon + b)}{(q^2 - \Delta)^2} \Big|_{R_2} = i\pi^2 (2a - b)\Delta, \\
 \int d^d q \frac{q^2 (a\epsilon + b)}{(q^2 - \Delta)^3} \Big|_{R_2} = i\pi^2 \left(a - \frac{1}{2}b \right), \\
 \int d^d q \frac{q^4 (a\epsilon + b)}{(q^2 - \Delta)^4} \Big|_{R_2} = i\pi^2 \left(a - \frac{5}{6}b \right), \\
 \mu^{2\epsilon} \int d^d q \frac{a\epsilon + b}{q^2 - m^2} \Big|_{UV} = i\pi^2 m^2 \left(\frac{b}{\bar{\epsilon}} + a + b - b \log \left(\frac{m^2}{\mu^2} \right) \right), \\
 \mu^{2\epsilon} \int d^d q \frac{a\epsilon + b}{(q^2 - \Delta)^2} \Big|_{UV} = i\pi^2 (a\epsilon + b) \left(\frac{1}{\bar{\epsilon}} - \log \left(\frac{\Delta}{\mu^2} \right) \right), \\
 \mu^{2\epsilon} \int d^d q \frac{q^2 (a\epsilon + b)}{(q^2 - \Delta)^2} \Big|_{UV} = i\pi^2 (2a\epsilon + b\epsilon + 2b) \left(\frac{1}{\bar{\epsilon}} - \log \left(\frac{\Delta}{\mu^2} \right) \right) \Delta, \\
 \mu^{2\epsilon} \int d^d q \frac{q^2 (a\epsilon + b)}{(q^2 - \Delta)^3} \Big|_{UV} = i\pi^2 \frac{b}{\bar{\epsilon}}, \\
 \mu^{2\epsilon} \int d^d q \frac{q^4 (a\epsilon + b)}{(q^2 - \Delta)^4} \Big|_{UV} = i\pi^2 \frac{b}{\bar{\epsilon}},
 \end{array}$$

- Integrate over the Feynman parameters (but for the two-point UV finite terms)
- Replace masses and couplings by their values for each field insertion

NLOCT

- Replace momentum integrals

$$\begin{array}{l}
 \int d^d q \frac{\epsilon}{q^2 - m^2} \Big|_{R_2} = i\pi^2 m^2, \\
 \int d^d q \frac{\epsilon}{(q^2 - \Delta)^2} \Big|_{R_2} = i\pi^2, \\
 \int d^d q \frac{q^2 (a\epsilon + b)}{(q^2 - \Delta)^2} \Big|_{R_2} = i\pi^2 (2a - b)\Delta, \\
 \int d^d q \frac{q^2 (a\epsilon + b)}{(q^2 - \Delta)^3} \Big|_{R_2} = i\pi^2 \left(a - \frac{1}{2}b \right), \\
 \int d^d q \frac{q^4 (a\epsilon + b)}{(q^2 - \Delta)^4} \Big|_{R_2} = i\pi^2 \left(a - \frac{5}{6}b \right), \\
 \mu^{2\epsilon} \int d^d q \frac{a\epsilon + b}{q^2 - m^2} \Big|_{UV} = i\pi^2 m^2 \left(\frac{b}{\bar{\epsilon}} + a + b - b \log \left(\frac{m^2}{\mu^2} \right) \right), \\
 \mu^{2\epsilon} \int d^d q \frac{a\epsilon + b}{(q^2 - \Delta)^2} \Big|_{UV} = i\pi^2 (a\epsilon + b) \left(\frac{1}{\bar{\epsilon}} - \log \left(\frac{\Delta}{\mu^2} \right) \right), \\
 \mu^{2\epsilon} \int d^d q \frac{q^2 (a\epsilon + b)}{(q^2 - \Delta)^2} \Big|_{UV} = i\pi^2 (2a\epsilon + b\epsilon + 2b) \left(\frac{1}{\bar{\epsilon}} - \log \left(\frac{\Delta}{\mu^2} \right) \right) \Delta, \\
 \mu^{2\epsilon} \int d^d q \frac{q^2 (a\epsilon + b)}{(q^2 - \Delta)^3} \Big|_{UV} = i\pi^2 \frac{b}{\bar{\epsilon}}, \\
 \mu^{2\epsilon} \int d^d q \frac{q^4 (a\epsilon + b)}{(q^2 - \Delta)^4} \Big|_{UV} = i\pi^2 \frac{b}{\bar{\epsilon}},
 \end{array}$$

- Integrate over the Feynman parameters (but for the two-point UV finite terms)
- Replace masses and couplings by their values for each field insertion

NLOCT

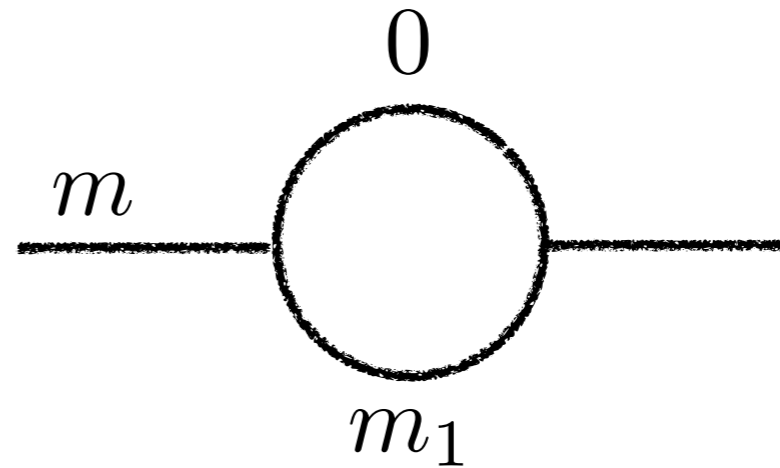
- Perform the color algebra for triplets and octets
- Write the renormalization conditions (fix p^2) **End R_2**
- Do the integration over the feynman parameters for the UV-finite parts

$$b_0(p^2, m_1, m_2) \equiv \int_0^1 dx \log \left(\frac{p^2(x-1)x + x(m_1^2 - m_2^2) + m_2^2 - i\epsilon_p}{\mu^2} \right)$$

$$b_0(0, 0, 0) = \frac{1}{\bar{\epsilon}}$$

- Solve the renormalization conditions
- Replace the counterterms by their values in the CT vertices

Real/Complex masses



Real masses

$$m \in \mathbb{R} \quad m_1^2 < m^2 \quad \cancel{\Re}(\log [p^2 - m_1^2] + \cancel{i\pi}) \Big|_{p^2=m^2}$$

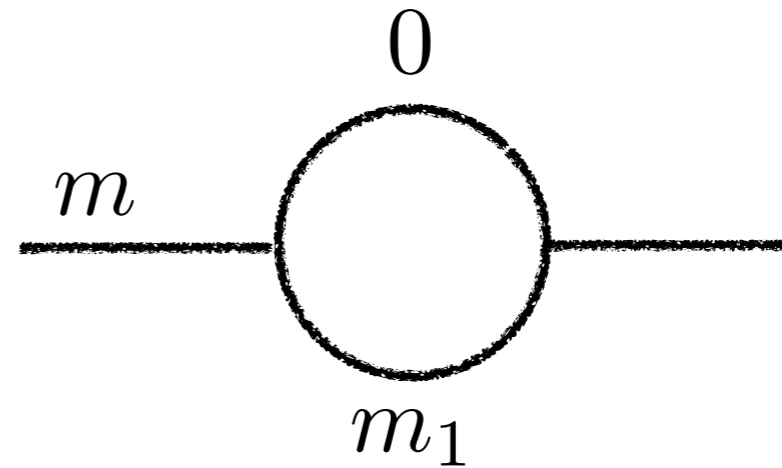
$$m_1^2 > m^2 \quad \cancel{\Re}(\log [m_1^2 - p^2]) \Big|_{p^2=m^2}$$

Complex masses

$$\log [m_1^2 - p^2] \Big|_{p^2=m^2}$$

All cases are kept unless the users put some assumptions

Real/Complex masses



Real masses

$$m \in \mathbb{R} \quad m_1^2 < m^2 \quad \cancel{\Re}(\log [p^2 - m_1^2] + \cancel{i\pi}) \Big|_{p^2=m^2}$$
$$m_1^2 > m^2 \quad \cancel{\Re}(\log [m_1^2 - p^2]) \Big|_{p^2=m^2}$$

Complex masses

$$\log [m_1^2 - p^2] \Big|_{p^2=m^2} \quad \text{Faster!}$$

All cases are kept unless the users put some assumptions

Renormalization options

```
FR$LoopSwitches = {{Gf, MW}};
```

Switch internal masses with an external parameter
appearing in its expression

Renormalization options

```
FR$LoopSwitches = {{Gf, MW}};
```

Switch into *Before calling OnShellRenormalization* external parameter
appearing in its caption

Renormalization options

```
FR$LoopSwitches = {{Gf, MW}};
```

Switch into *Before calling OnShellRenormalization* external parameter appearing in its definition

OnShellRenormalization options

QCDOnly : Only the coloured fields and their masses and the couplings with QCD if True

FlavorMixing : Forbid all the mixing or allow only some of them

Exclude4ScalarCT : No CT for the 4 scalars vertices (but keep the 4 scalars TL)

Simplify2Point : Put the quadratic part of the Lagrangian in canonical form (Avoided if False)

WriteCT options

`WriteCT [<model>, <genericfile>, options]`

OnShellRenormalization options

`QCDOnly` : Only QCD corrections

`Assumptions` : Mass spectrum for the UV counterterms

`Exclude4ScalarCT` : No computation of the CT for the 4 scalars vertices (but keep the 4 scalars TL)

`ZeroMom` : {coupling, vertex} use zero momentum for coupling on vertex

`ComplexMass` : complex mass scheme if True

R2 : Validation

- tested* on the SM (QCD:P. Draggiotis et al. +QED:M.V. Garzelli et al)
- tested* on MSSM (QCD:H.-S. Shao, Y.-J. Zhang) : test the Majorana

*Analytic comparison of the expressions

UV Validation

- SM QCD : tested* (W. Beenakker, S. Dittmaier, M. Kramer, B. Plumper)
- SM EW : tested* (expressions given by H.-S. Shao from A. Denner)

*Analytic comparison of the expressions

Tests in event generators

- aMC@NLO
- The SM QCD has been tested by V. Hirschi (Comparison with the built-in version)
- The MSSM QCD and SM EW are tested by H.-S. Shao and V. Hirschi
- 2HDM QCD is currently tested ($p p \rightarrow S, H^\pm t$)
 - gauge invariance
 - pole cancelation

SM tests

=== Finite ===

Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d d~ > w+ w- g	-1.2565695610e+01	-1.2565705416e+01	-1.2565696276e+01	3.9018817097e-07	Pass

=== Born ===

Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d d~ > w+ w- g	1.8518318521e-06	1.8518318521e-06	1.8518318521e-06	8.0617231411e-15	Pass

=== Single pole ===

Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d d~ > w+ w- g	-1.9397426502e+01	-1.9397426502e+01	-1.9397426504e+01	5.5894073017e-11	Pass

=== Double pole ===

Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d d~ > w+ w- g	-5.6666666667e+00	-5.6666666667e+00	-5.6666666667e+00	3.0015206007e-14	Pass

=== Summary ===

1/1 passed, 0/1 failed=== Finite ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
d~ d > a g g	-5.3971186943e+01	-5.3971193753e+01	-5.3971189940e+01	6.3091071914e-08	Pass

=== Born ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
d~ d > a g g	6.4168774056e-05	6.4168764370e-05	6.4168764370e-05	7.5467680882e-08	Pass

=== Single pole ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
d~ d > a g g	-3.7439549398e+01	-3.7439549398e+01	-3.7439549397e+01	6.8122965983e-12	Pass

=== Double pole ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
d~ d > a g g	-8.6666666667e+00	-8.6666666667e+00	-8.6666666667e+00	2.2443585452e-14	Pass

=== Summary ===

1/1 passed, 0/1 failed=== Finite ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
d~ d > z g g	-5.3769573669e+01	-5.3769573347e+01	-5.3769566412e+01	6.7475496780e-08	Pass

SM tests

=== Born ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
d~ d > z g g	3.1531233900e-04	3.1531235770e-04	3.1531235770e-04	2.9654886777e-08	Pass

=== Single pole ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
d~ d > z g g	-3.7464897007e+01	-3.7464897007e+01	-3.7464897007e+01	4.2333025503e-12	Pass

=== Double pole ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
d~ d > z g g	-8.6666666667e+00	-8.6666666667e+00	-8.6666666667e+00	2.1316282073e-14	Pass

=== Summary ===

l/l passed, 0/l failed=== Finite ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
d~ d > z z g	-5.9990384275e+00	-5.9990511729e+00	-5.9990379587e+00	1.1013604745e-06	Pass

=== Born ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
d~ d > z z g	2.2616997126e-06	2.2617000449e-06	2.2617000449e-06	7.3450366526e-08	Pass

=== Single pole ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
d~ d > z z g	-1.5469587040e+01	-1.5469587040e+01	-1.5469587040e+01	1.5226666708e-11	Pass

=== Double pole ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
d~ d > z z g	-5.6666666667e+00	-5.6666666667e+00	-5.6666666667e+00	2.6645352591e-15	Pass

=== Summary ===

l/l passed, 0/l failed=== Finite ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
g g > h t t~	2.9740187004e+01	2.9740187005e+01	2.9740187036e+01	5.3265970697e-10	Pass

SM tests

=== Born ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
g g > h t t~	1.1079653971e-07	1.1079653974e-07	1.1079653974e-07	1.3190849004e-10	Pass

=== Single pole ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
g g > h t t~	-7.0825709000e+00	-7.0825709000e+00	-7.0825709000e+00	5.0901237085e-13	Pass

=== Double pole ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
g g > h t t~	-6.0000000000e+00	-6.0000000000e+00	-6.0000000000e+00	1.7023419711e-15	Pass

=== Summary ===

l/l passed, 0/l failed=== Finite ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
g g > z t t~	3.6409017466e+01	3.6409021125e+01	3.6409021117e+01	5.0242920154e-08	Pass

=== Born ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
g g > z t t~	7.0723041711e-07	7.0723046101e-07	7.0723046101e-07	3.1039274206e-08	Pass

=== Single pole ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
g g > z t t~	-7.1948086812e+00	-7.1948086773e+00	-7.1948086773e+00	2.7349789963e-10	Pass

=== Double pole ===

Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
g g > z t t~	-6.0000000000e+00	-6.0000000000e+00	-6.0000000000e+00	2.5165055225e-15	Pass

=== Summary ===

l/l passed, 0/l failed=== Finite ===

Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d d~ > w+ w- g	-1.2565695610e+01	-1.2565705416e+01	-1.2565696276e+01	3.9018817097e-07	Pass

SM tests

=== Born ===

Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d d~ > w+ w- g	1.8518318521e-06	1.8518318521e-06	1.8518318521e-06	8.0617231411e-15	Pass

=== Single pole ===

Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d d~ > w+ w- g	-1.9397426502e+01	-1.9397426502e+01	-1.9397426504e+01	5.5894073017e-11	Pass

=== Double pole ===

Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d d~ > w+ w- g	-5.6666666667e+00	-5.6666666667e+00	-5.6666666667e+00	3.0015206007e-14	Pass

=== Summary ===

1/1 passed, 0/1 failed=== Finite ===

Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d~ d > a g g	-1.1504816412e+01	-1.1504816557e+01	-1.1504815497e+01	4.6089385415e-08	Pass

=== Born ===

Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d~ d > a g g	2.3138920858e-06	2.3138920858e-06	2.3138920858e-06	4.3012538015e-15	Pass

=== Single pole ===

Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d~ d > a g g	-2.8637049838e+01	-2.8637049838e+01	-2.8637049838e+01	1.5718407645e-13	Pass

=== Double pole ===

Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d~ d > a g g	-8.6666666667e+00	-8.6666666667e+00	-8.6666666667e+00	1.7421961310e-15	Pass

=== Summary ===

1/1 passed, 0/1 failed=== Finite ===

Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d~ d > z g g	-1.0306105482e+01	-1.0306105654e+01	-1.0306102645e+01	1.4600800434e-07	Pass

+2/3

Summary/Prospect

- Automatic BSM@NLO
 - renormalizable
 - Feynman gauge



Summary/Prospect

- Automatic BSM@NLO
 - renormalizable
 - Feynman gauge
- Next version
 - EFT
 - Any gauge
 - other renormalization scheme (EW)



Summary/Prospect

- Automatic BSM@NLO
 - renormalizable
 - Feynman gauge
- Next version
 - EFT
 - Any gauge
 - other renormalization scheme (EW)
- With the help of the FeynRules and Madgraph_aMC@NLO teams

