

Monte-Carlo Generation

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IPPP/Durham

Topic

- Collider Physics



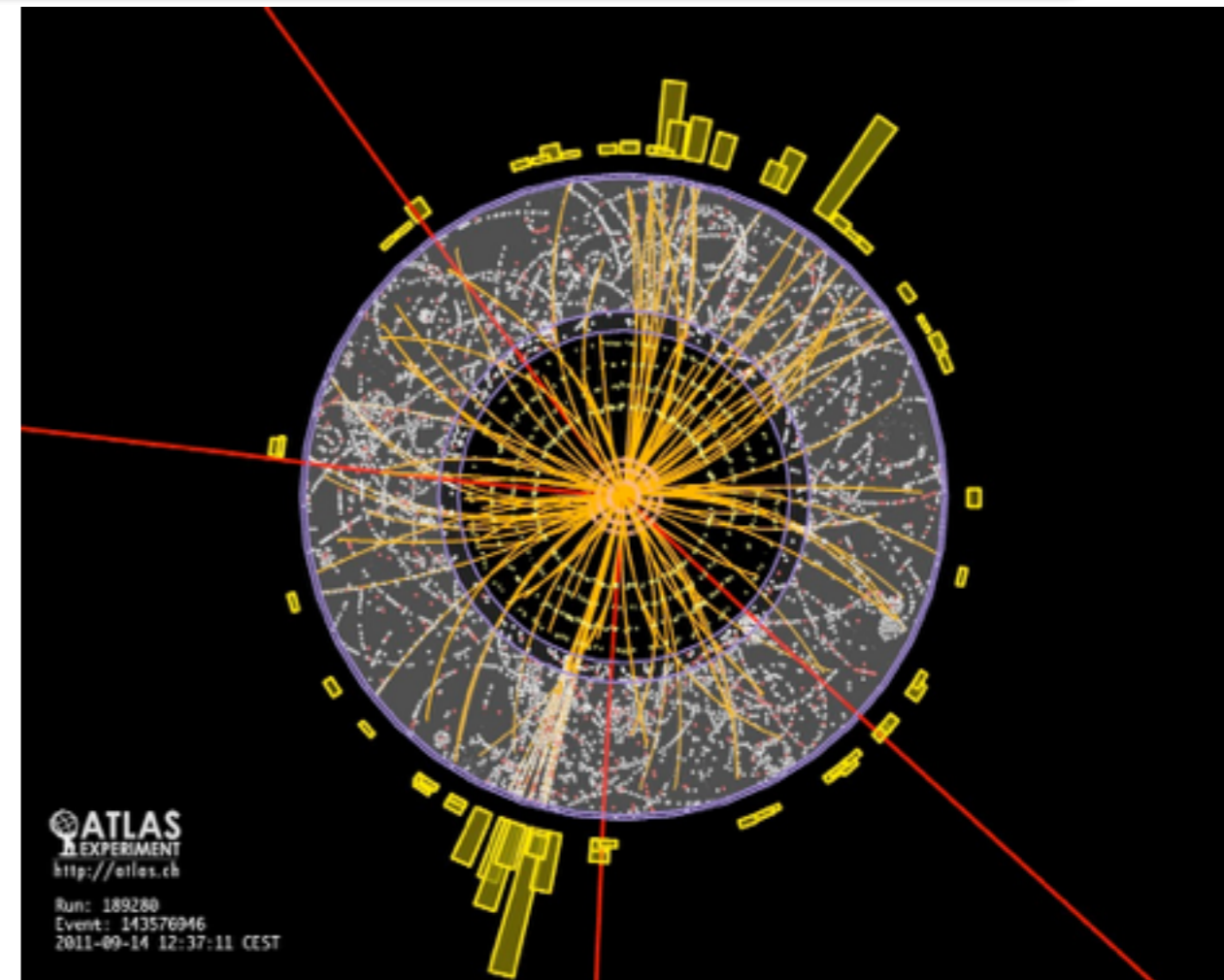
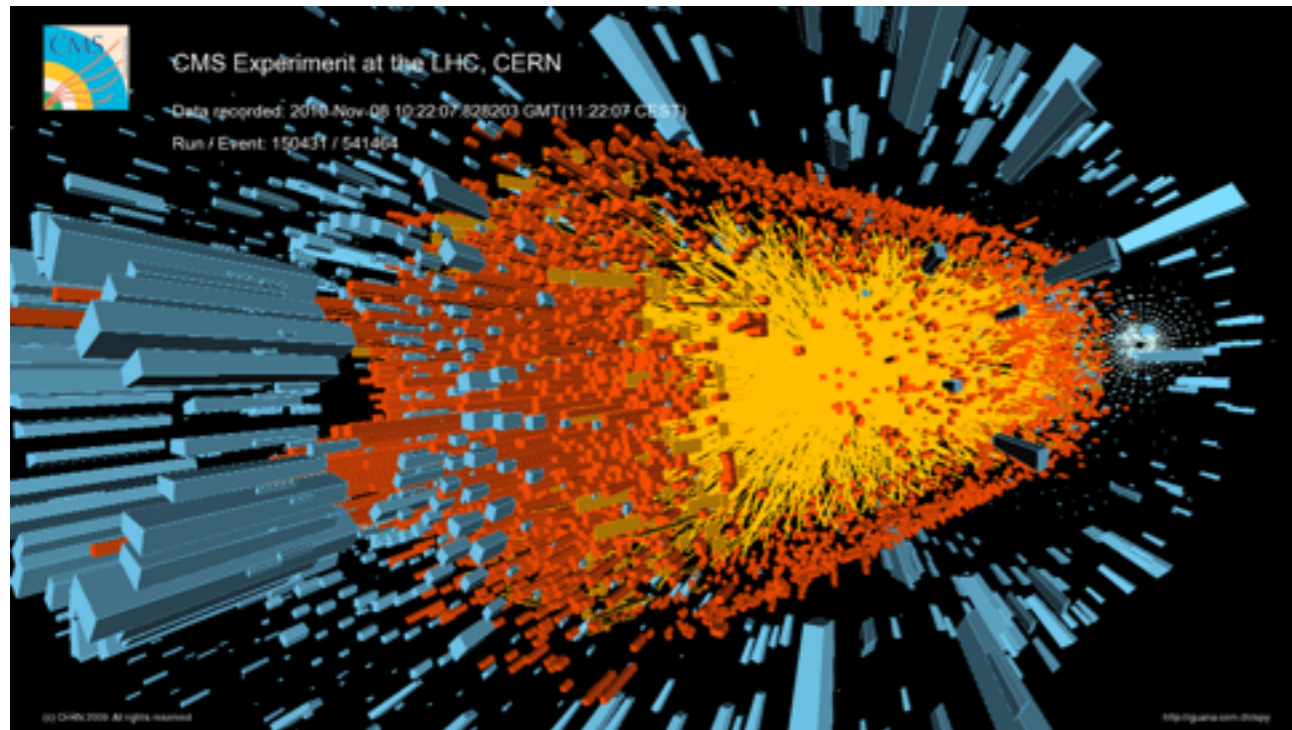
Topic

- Collider Physics
 - accelerating particle -> High Energy collision



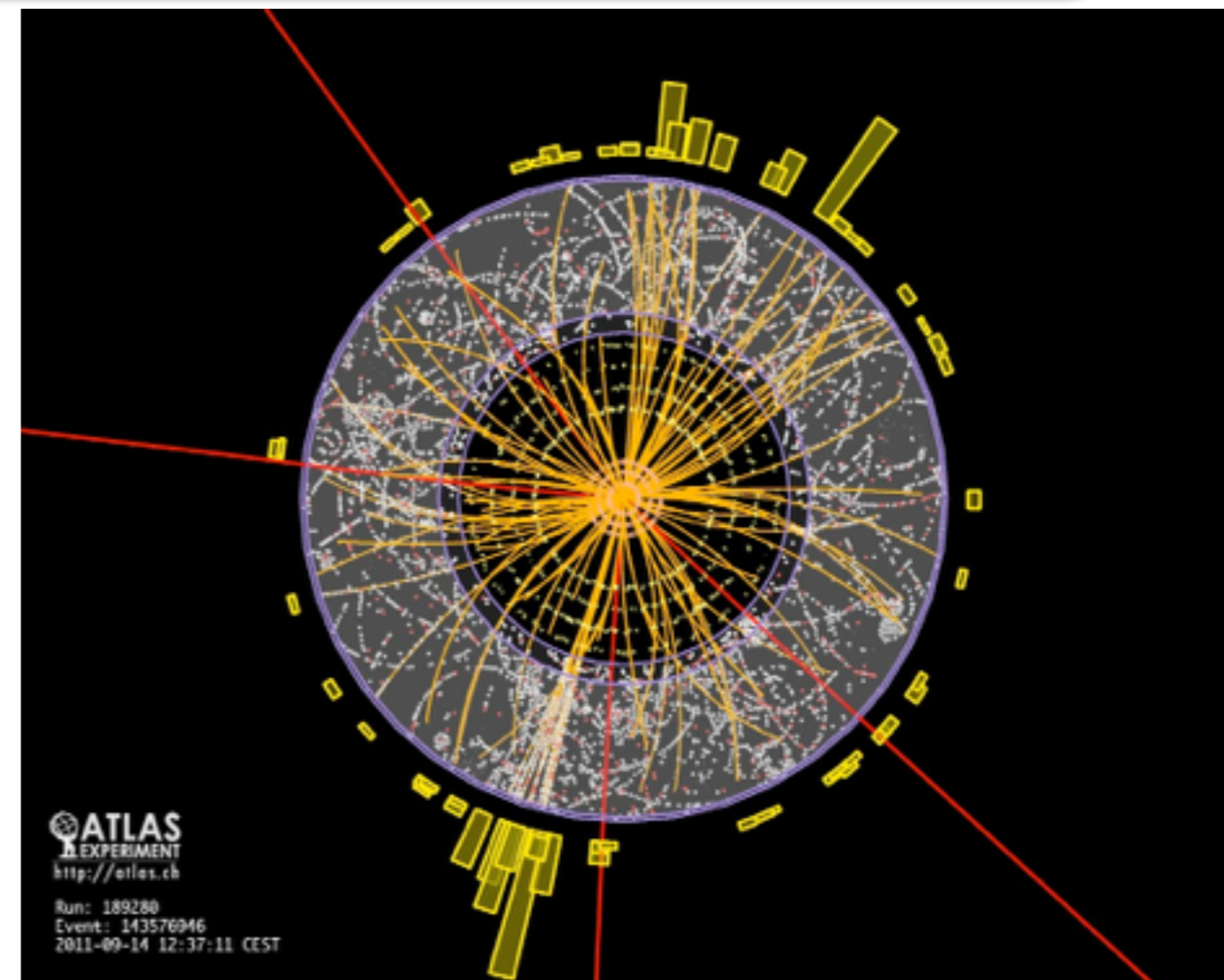
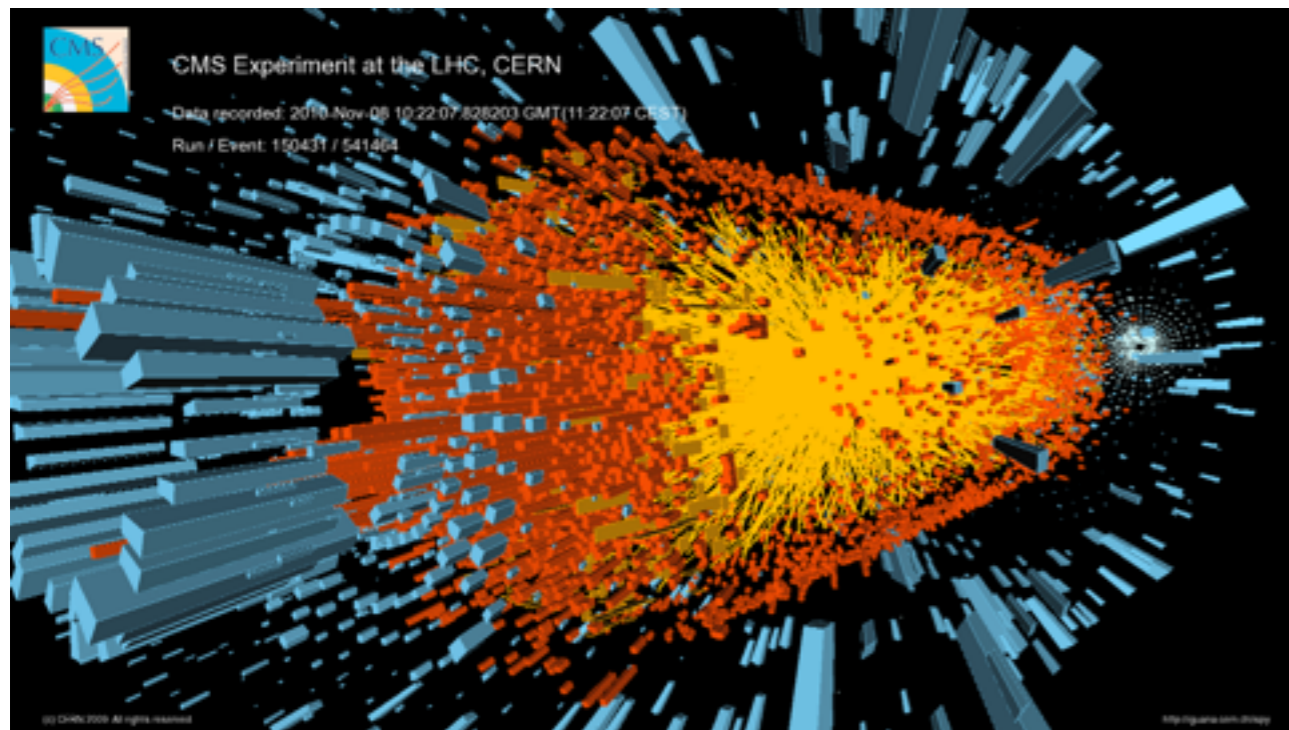
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- Collider Physics
 - accelerating particle -> High Energy collision

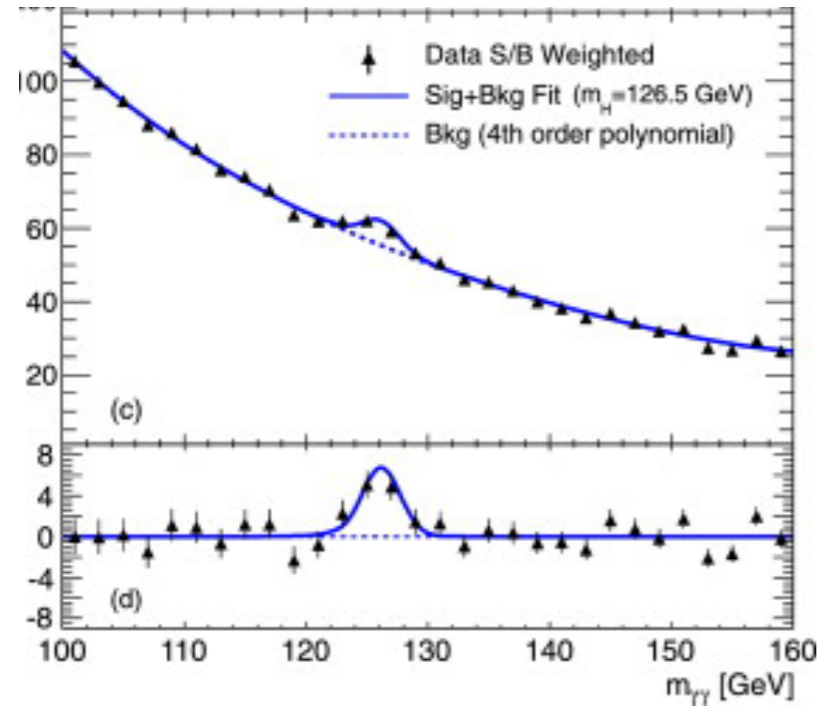


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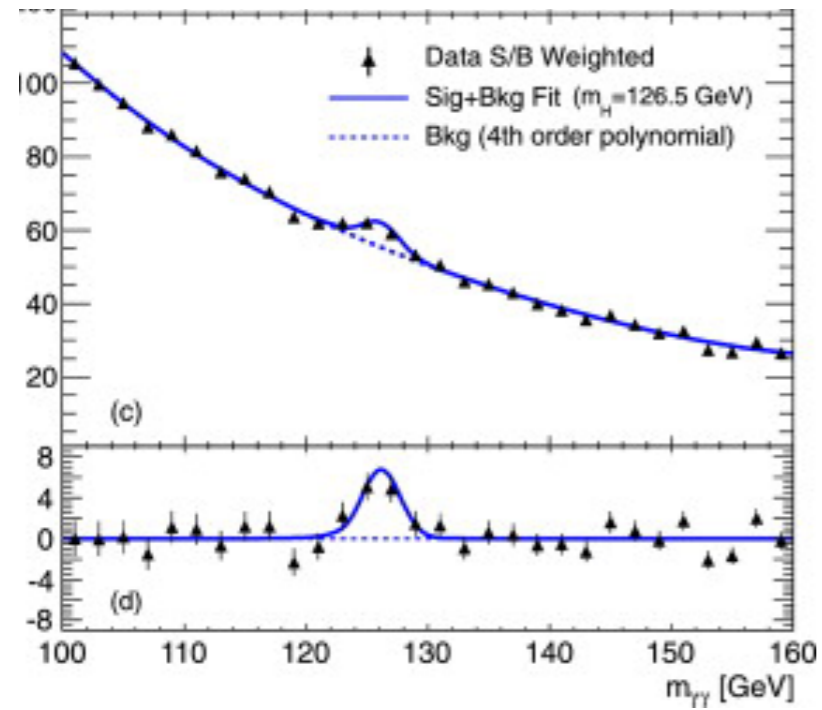
- Collider Physics
 - accelerating particle -> High Energy collision
- What do we need to predict/understand such collision?



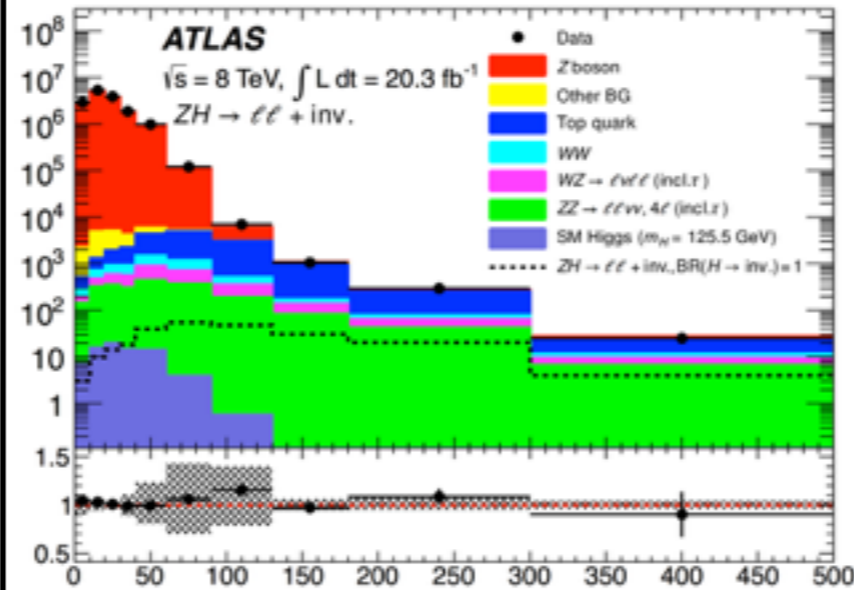
Peak



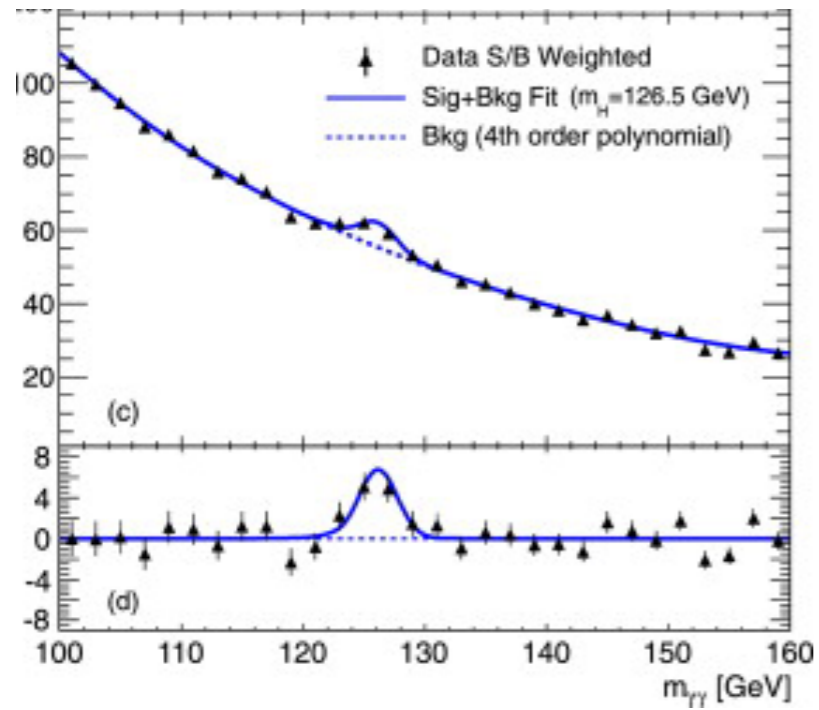
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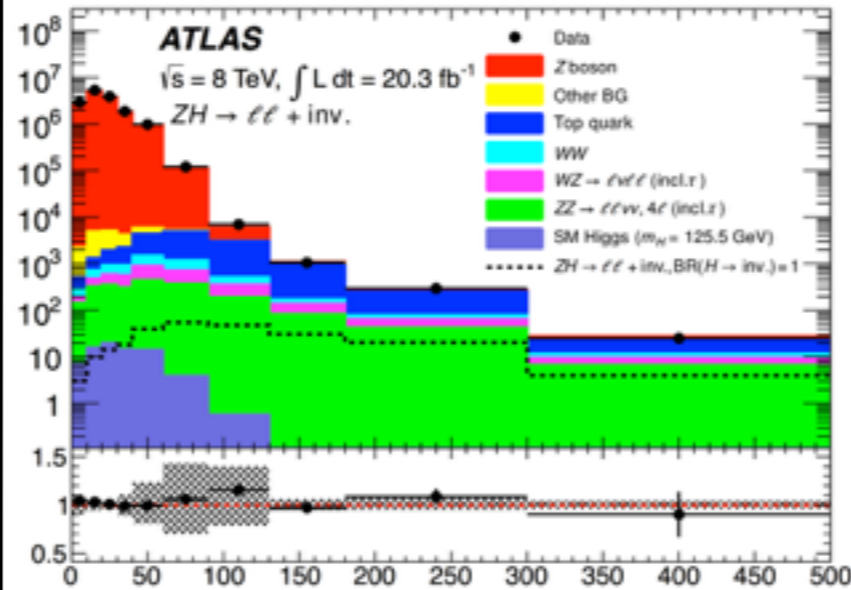
Shape



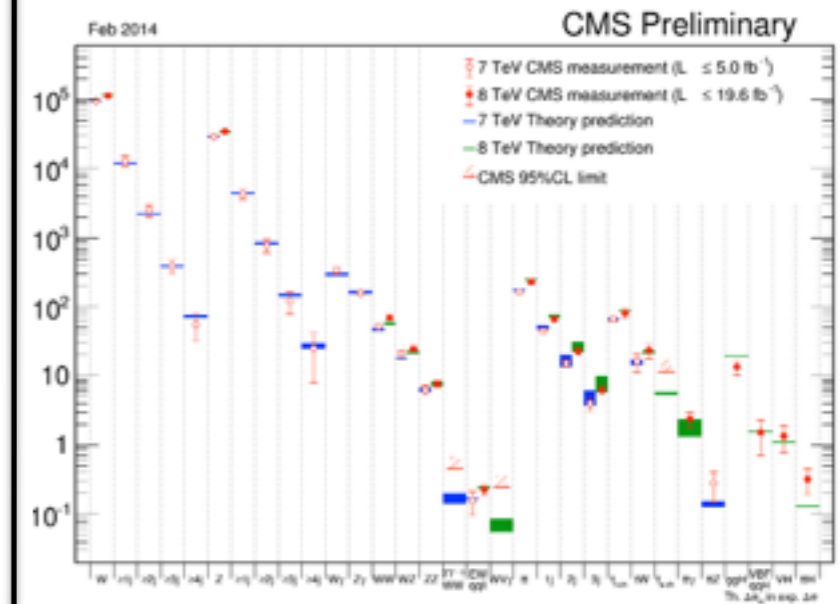
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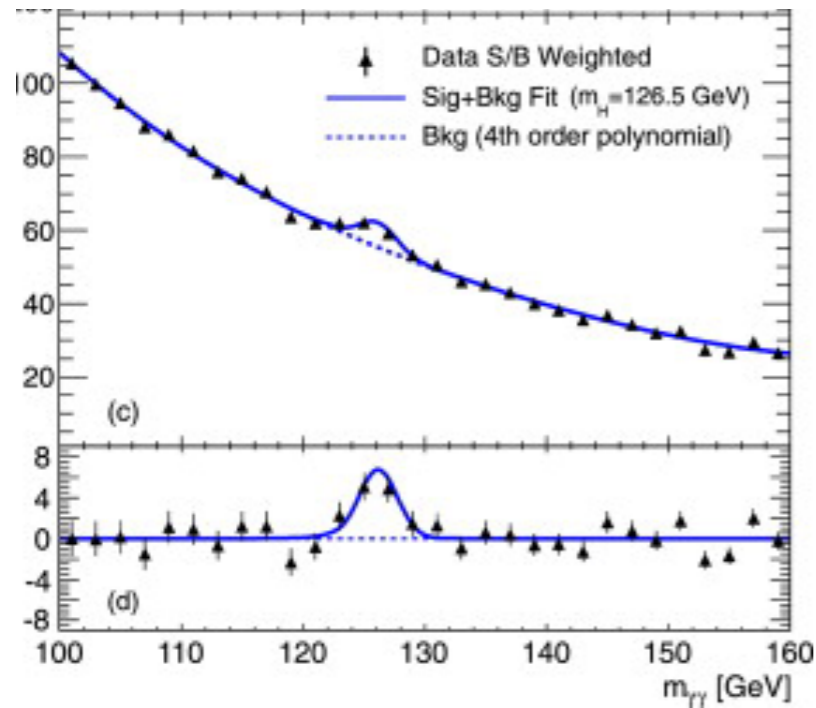
Shape



Rate

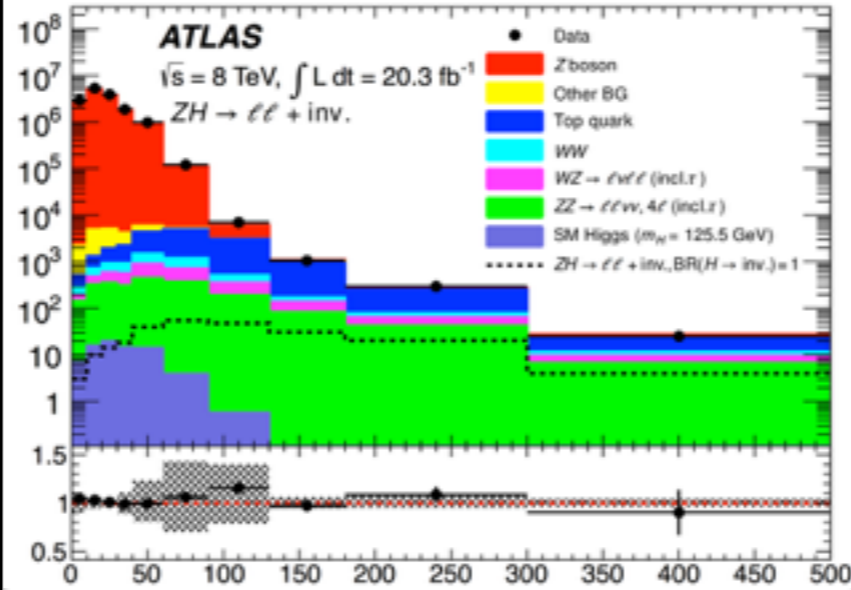


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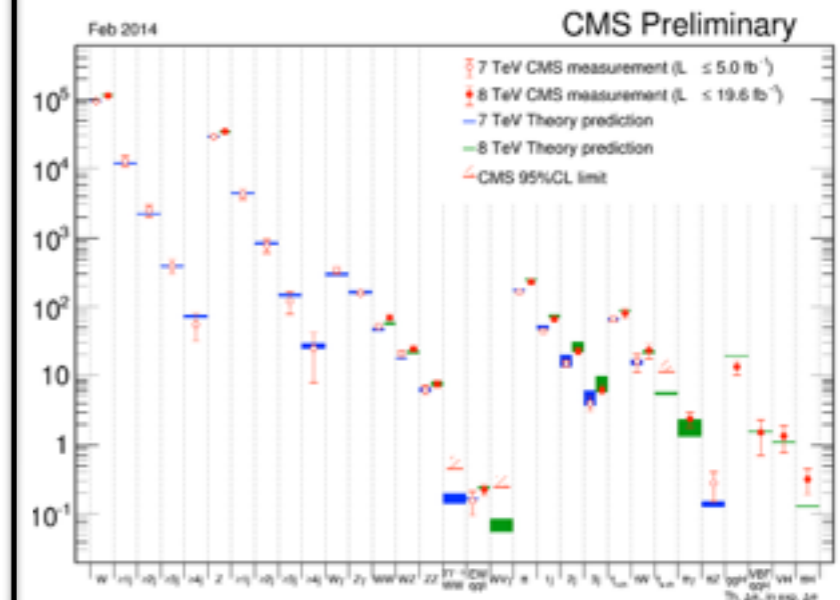
“EASY”

Shape



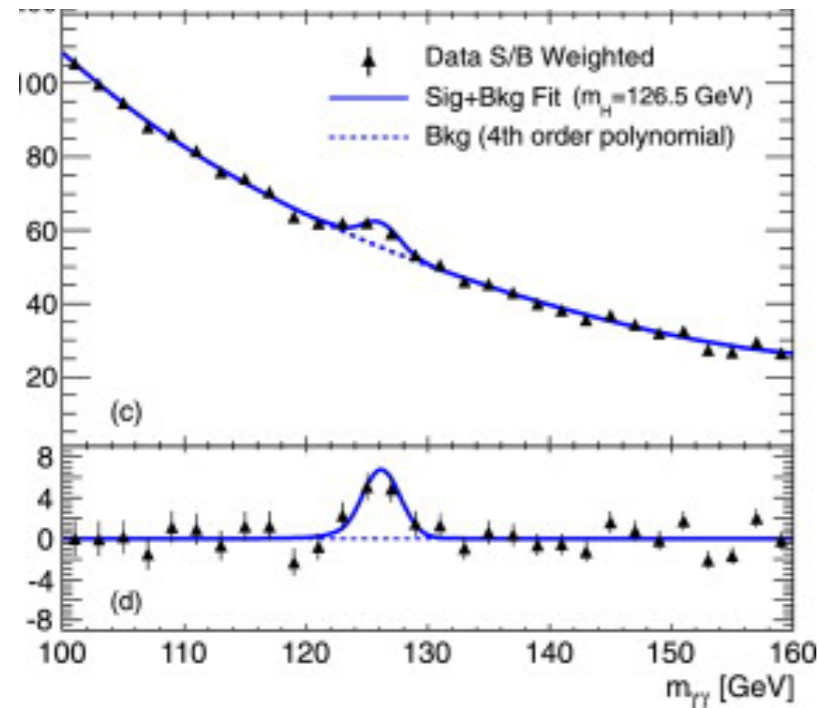
“HARD”

Rate



“VERY HARD”

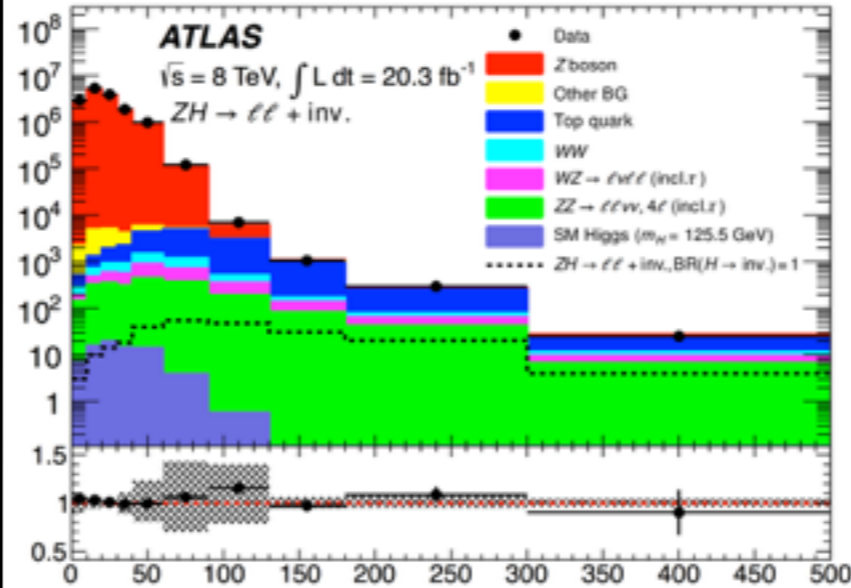
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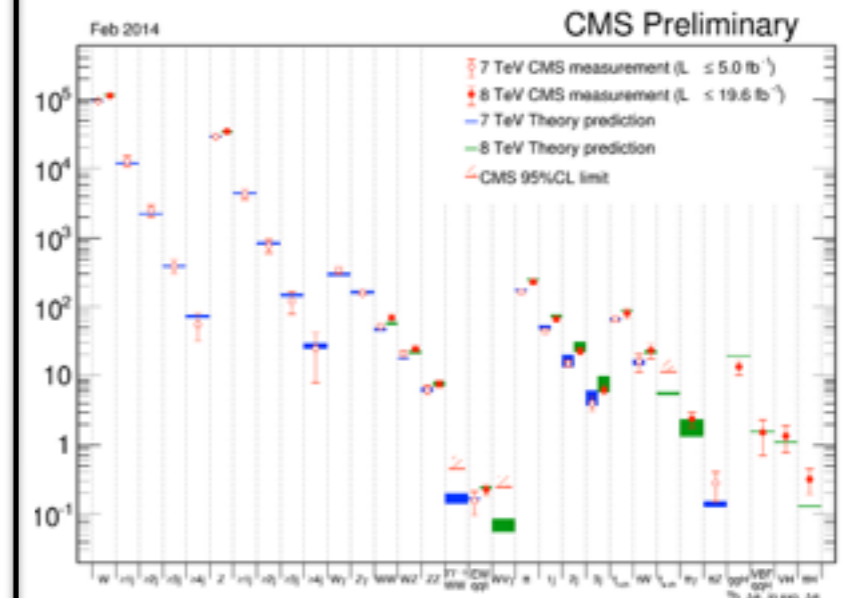
Background directly measured from **data**.
Theory needed only for parameter extraction

Shape



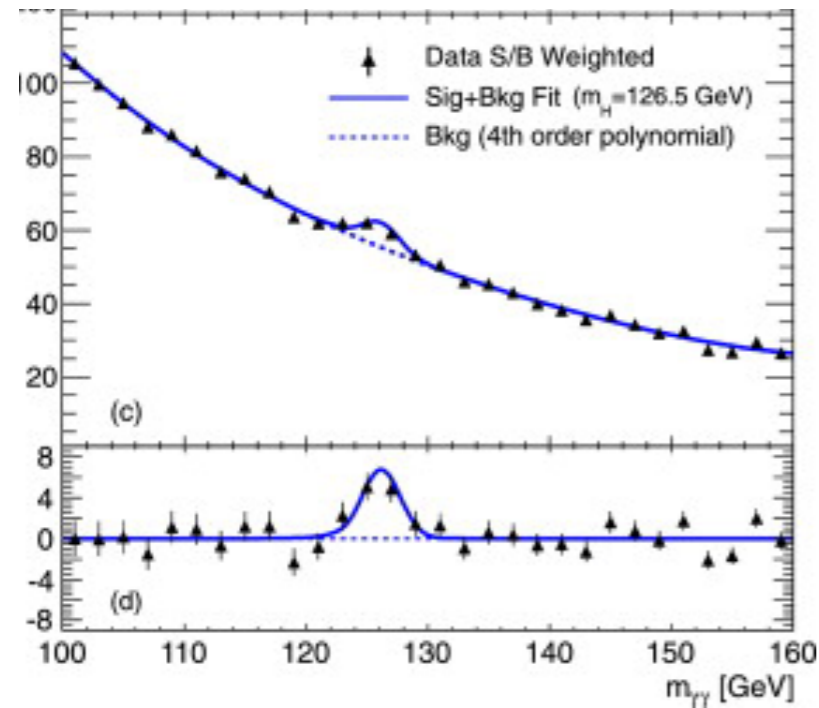
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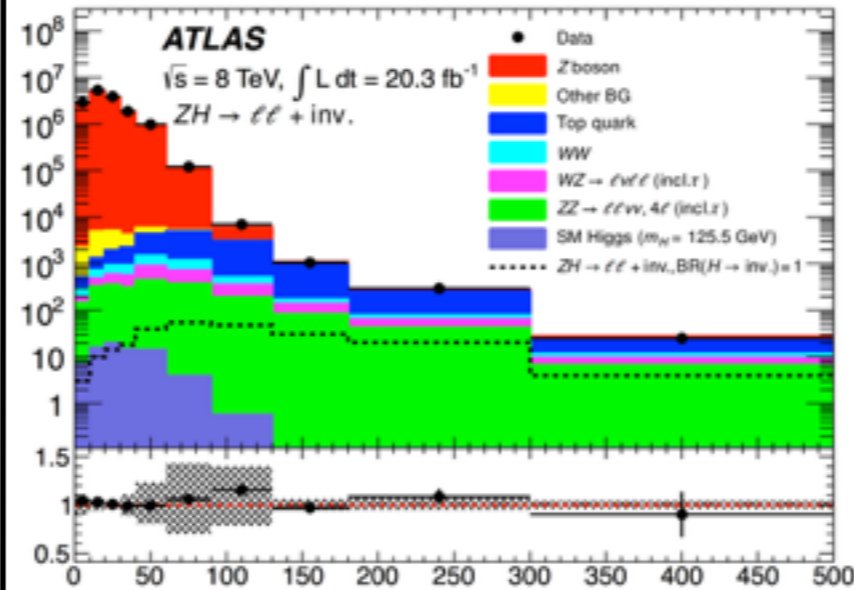
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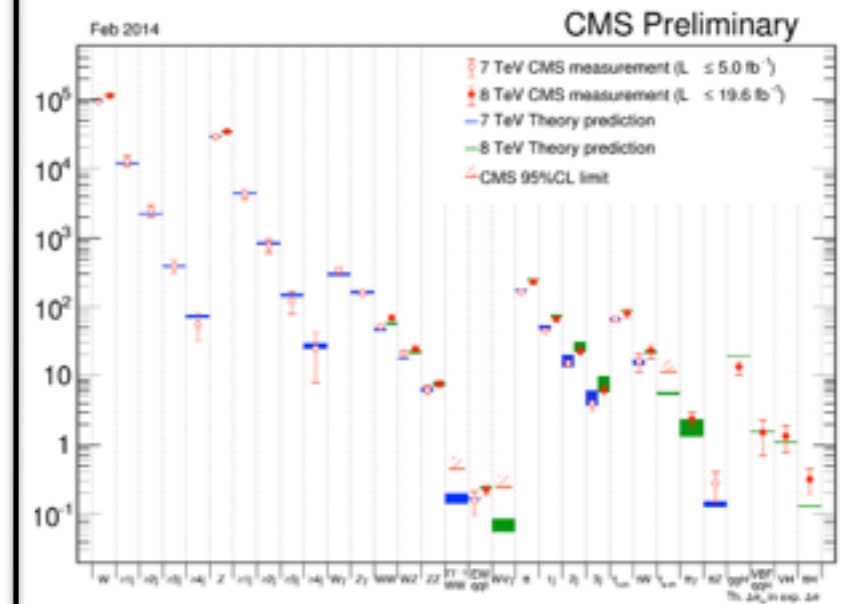
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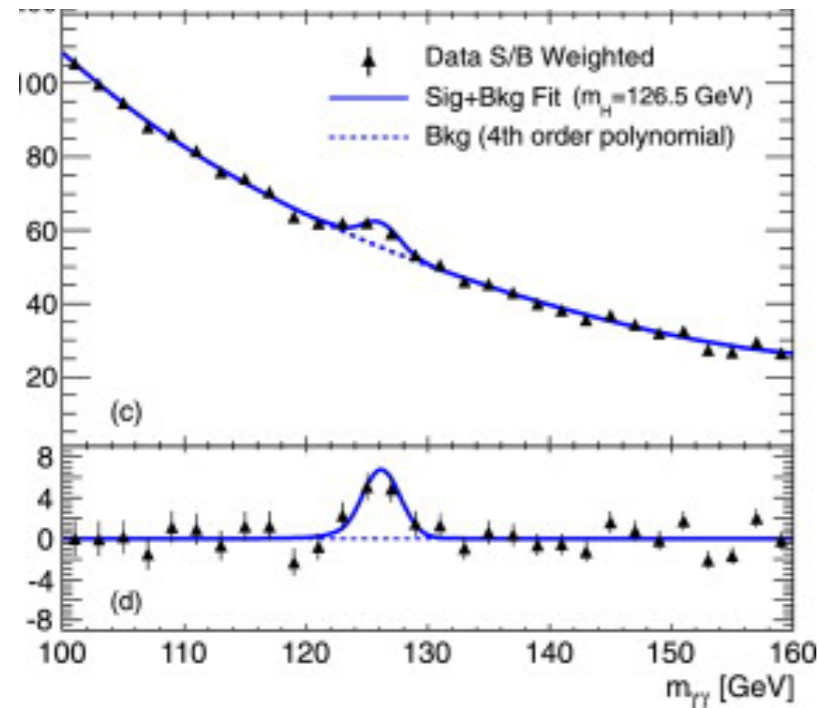
Background **SHAPE** needed.
Flexible MC for both signal and background validated and tuned to data

Rate



“VERY HARD”

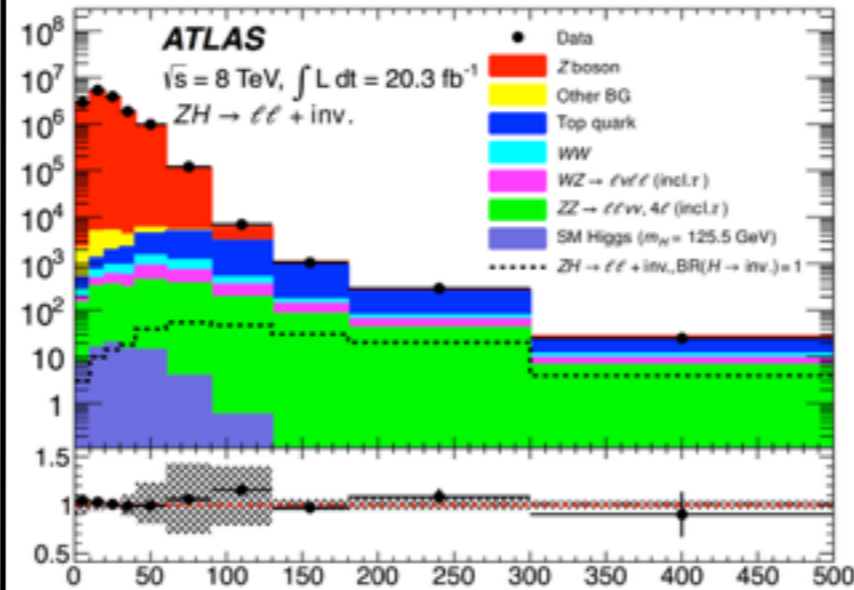
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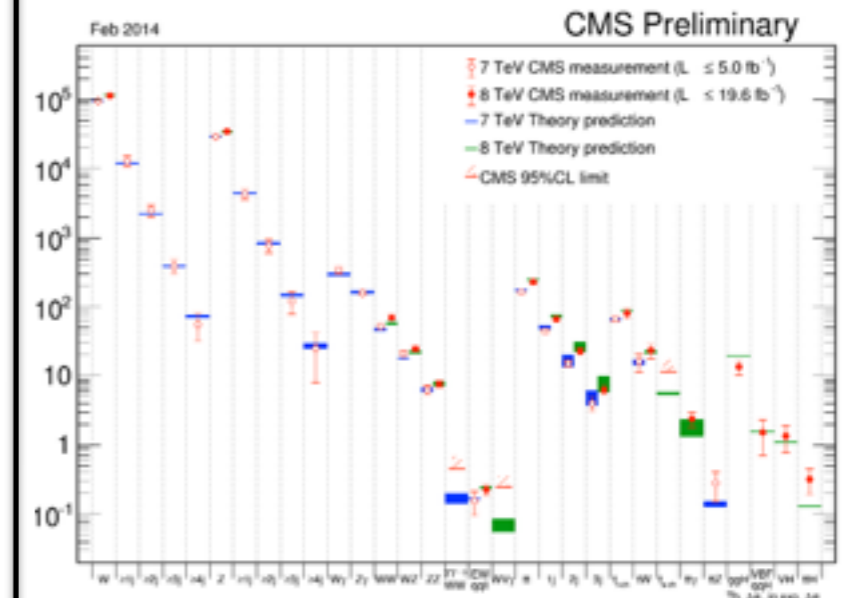
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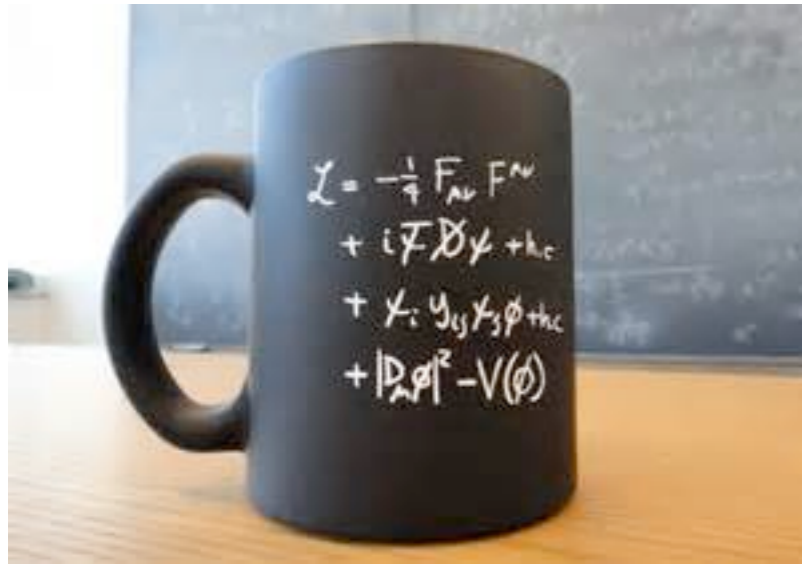
Rate



“VERY HARD”

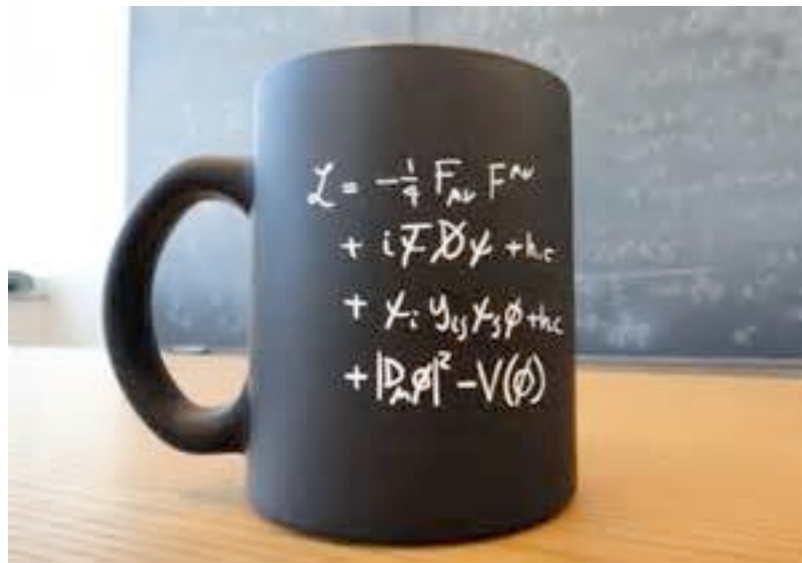
Relies on prediction for both **shape** and **normalization**.
 Complicated interplay of best simulations and data

Lagrangian



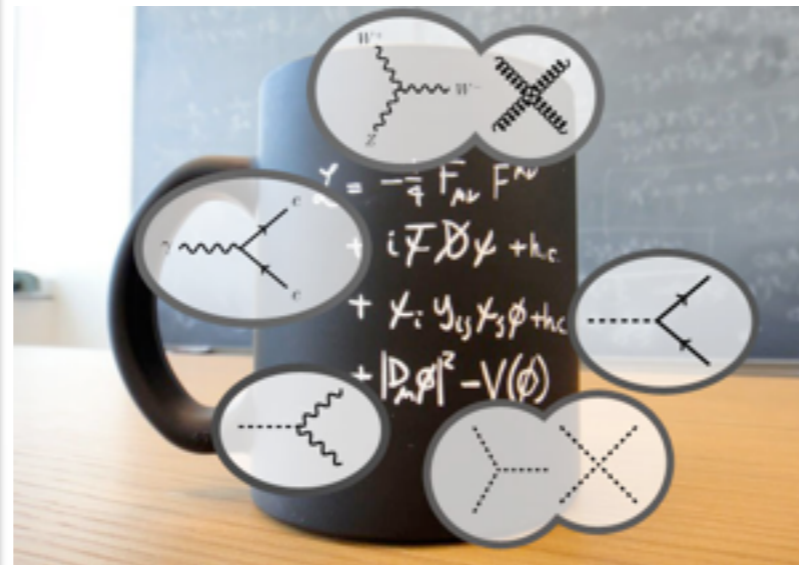
- This is Where the new idea are expressed

Lagrangian



- This is Where the new idea are expressed

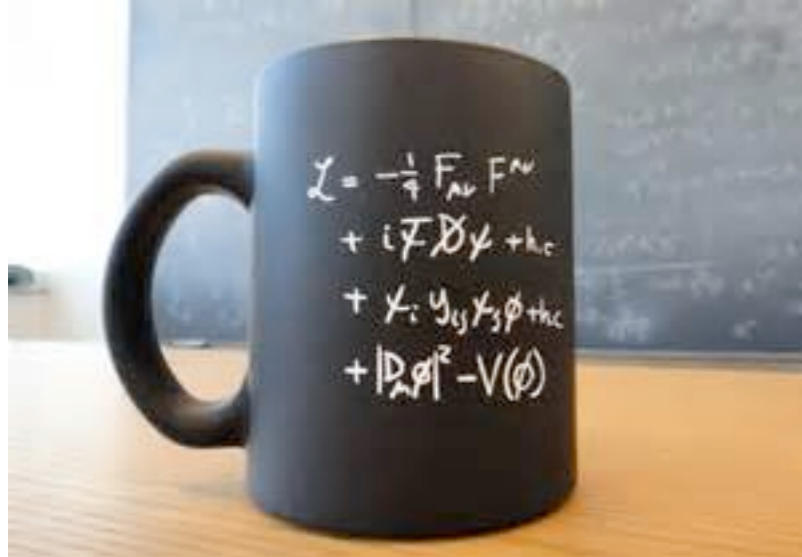
Feynman Rule



- Same information as the Lagrangian

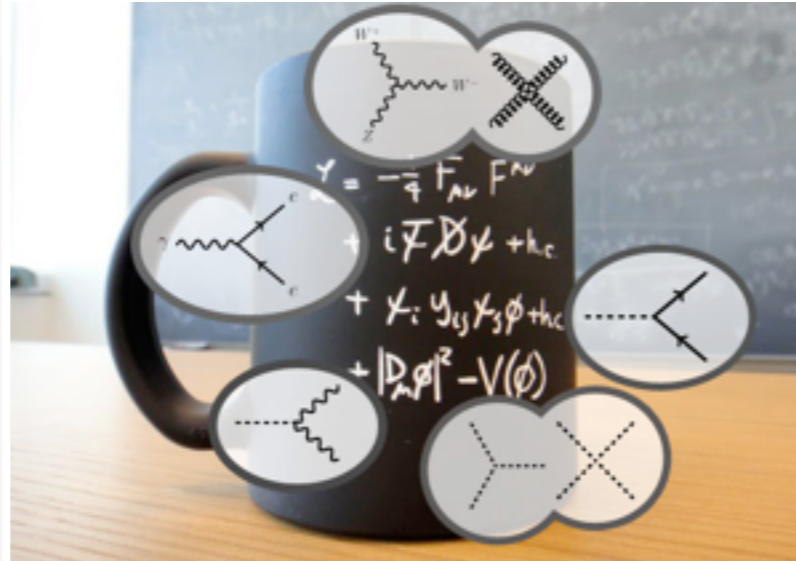
FeynRules

Lagrangian



- This is Where the new idea are expressed

Feynman Rule

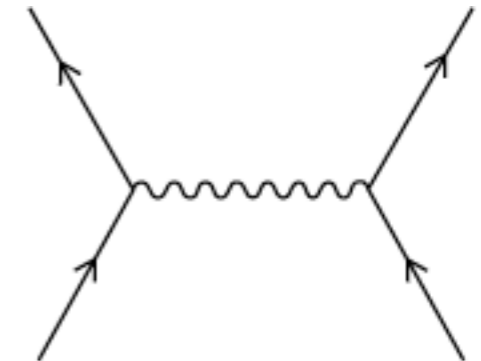


- Same information as the Lagrangian

FeynRules

Cross-section

$$\frac{d\sigma}{d\cos\theta} = \left(\frac{d\sigma}{d\cos\theta}\right)_R \left[1 + \frac{(1+\cos\theta)KE}{Mc^2}\right]$$



(a)



(b)



(c)



(d)

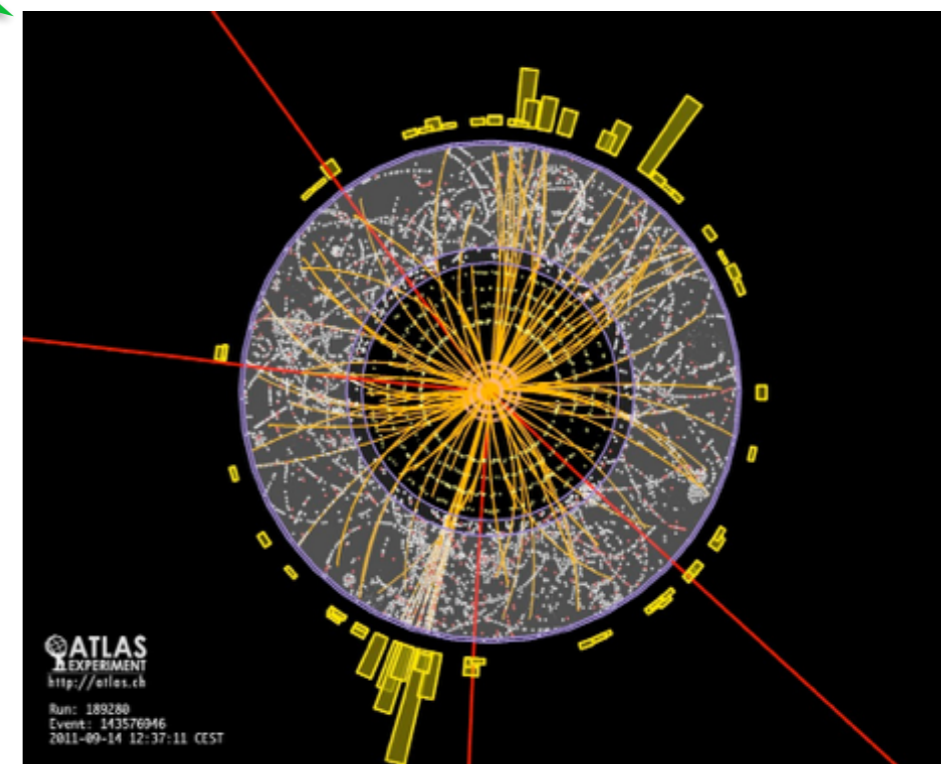
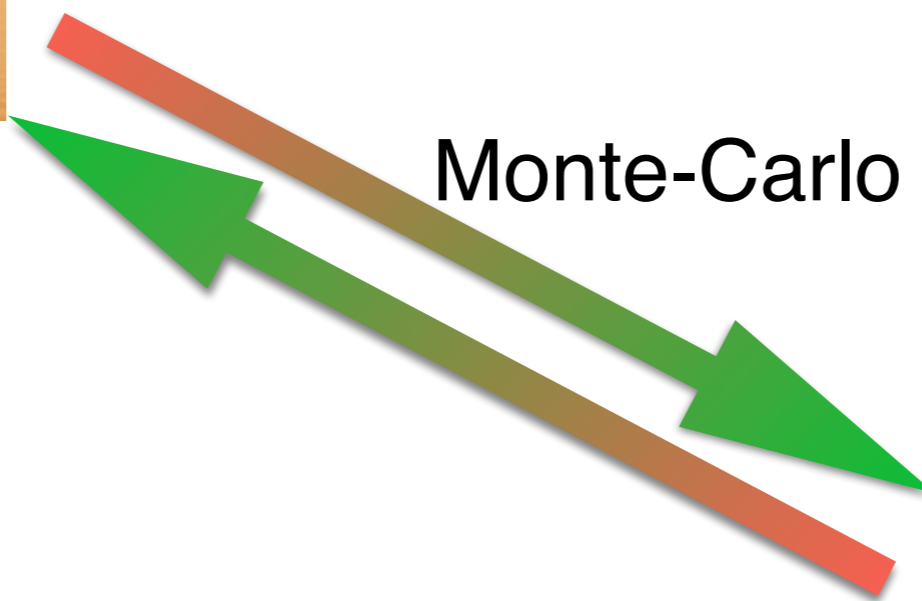
- What is the precision?



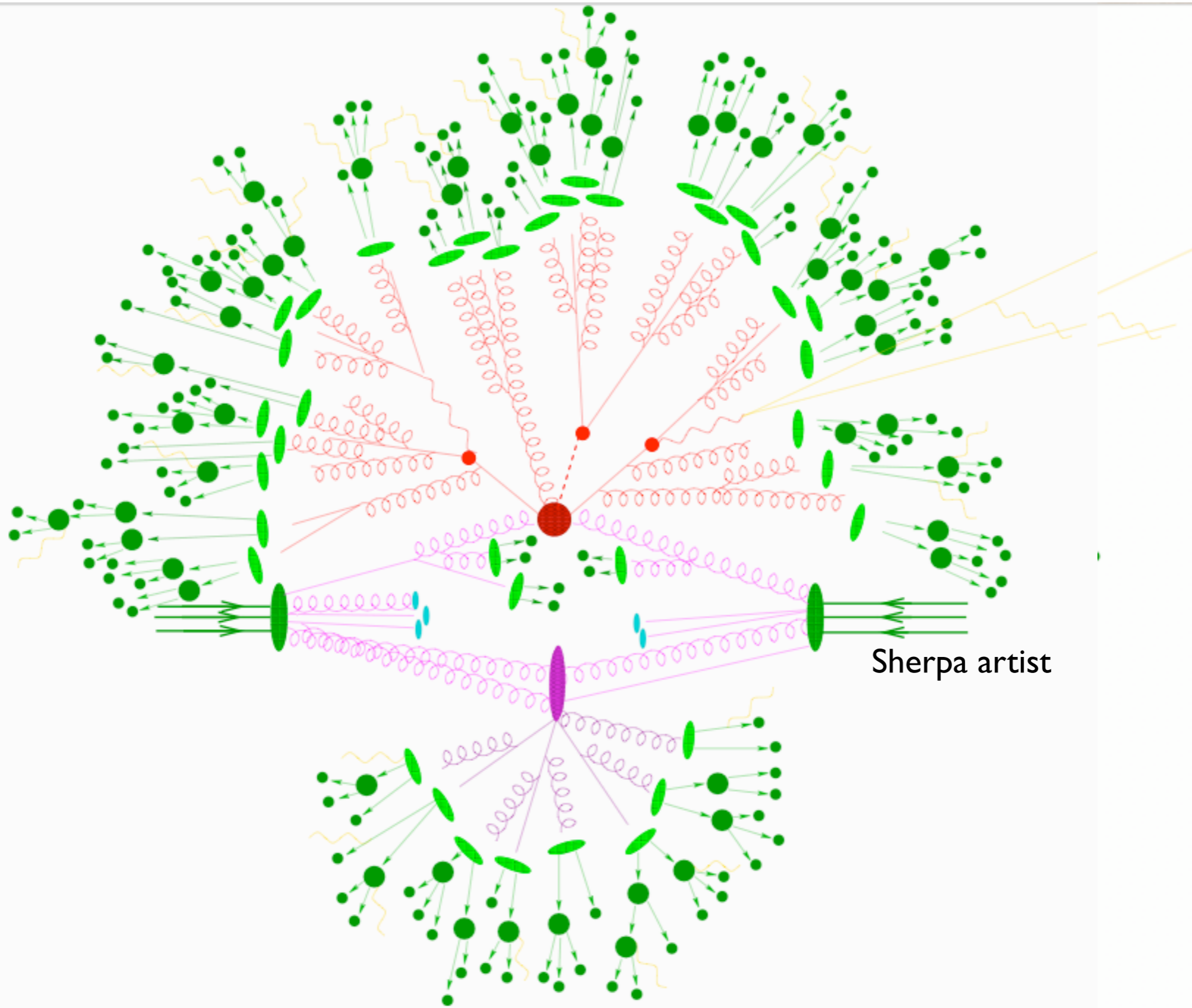
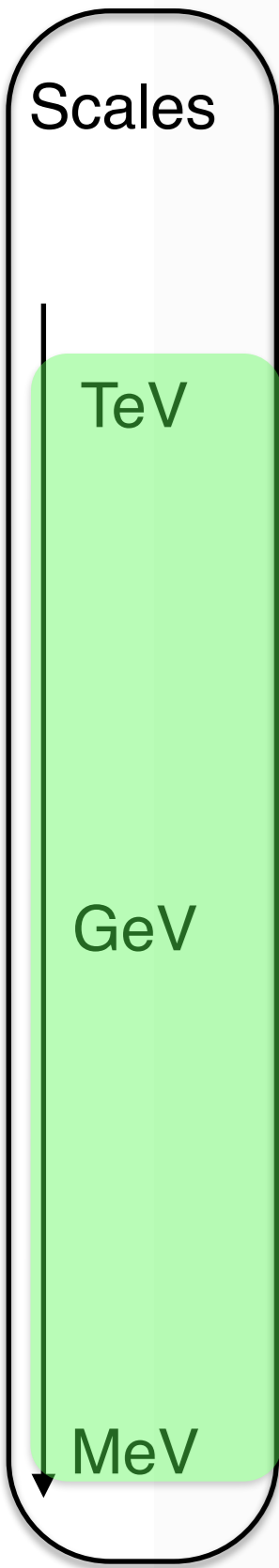
Monte-Carlo Physics



Monte-Carlo Physics



Simulation of collider events



Scales

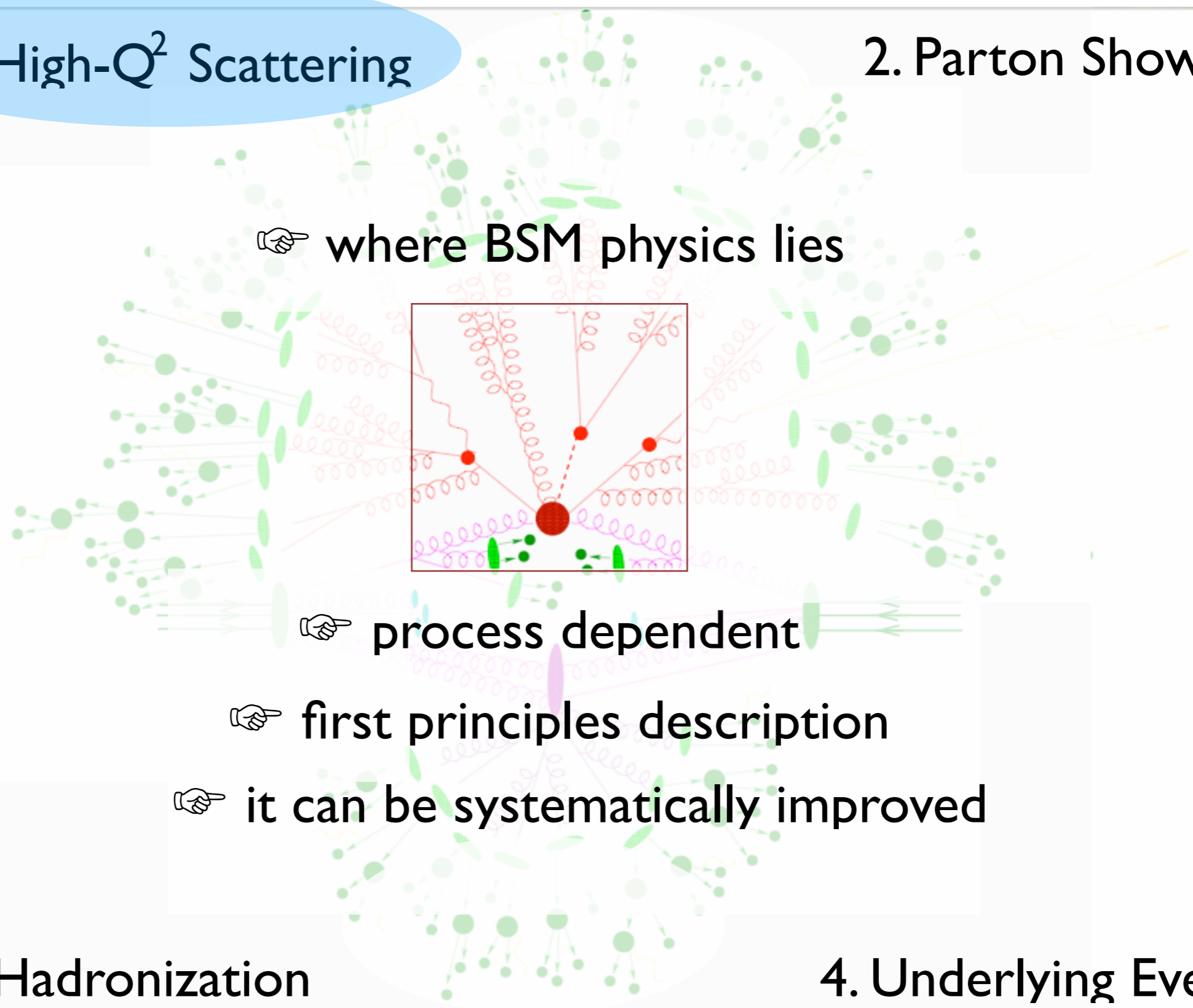
TeV

GeV

MeV

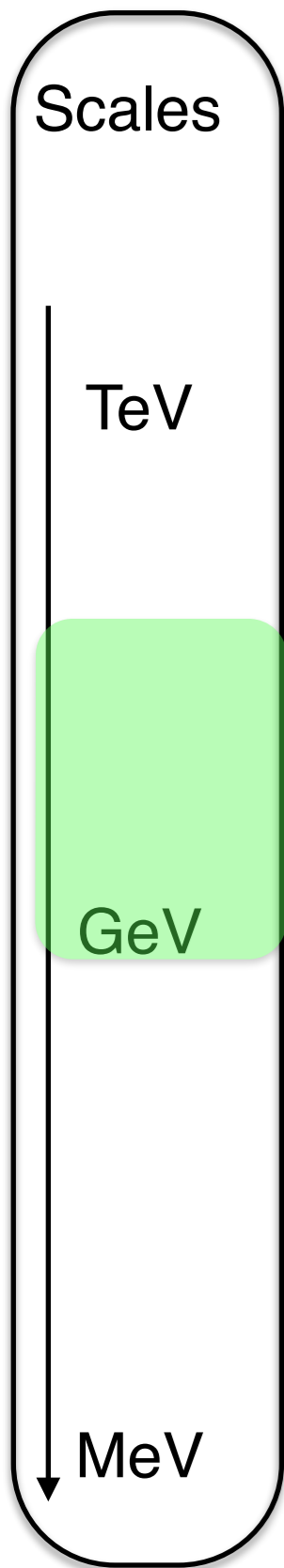
1. High- Q^2 Scattering

2. Parton Shower



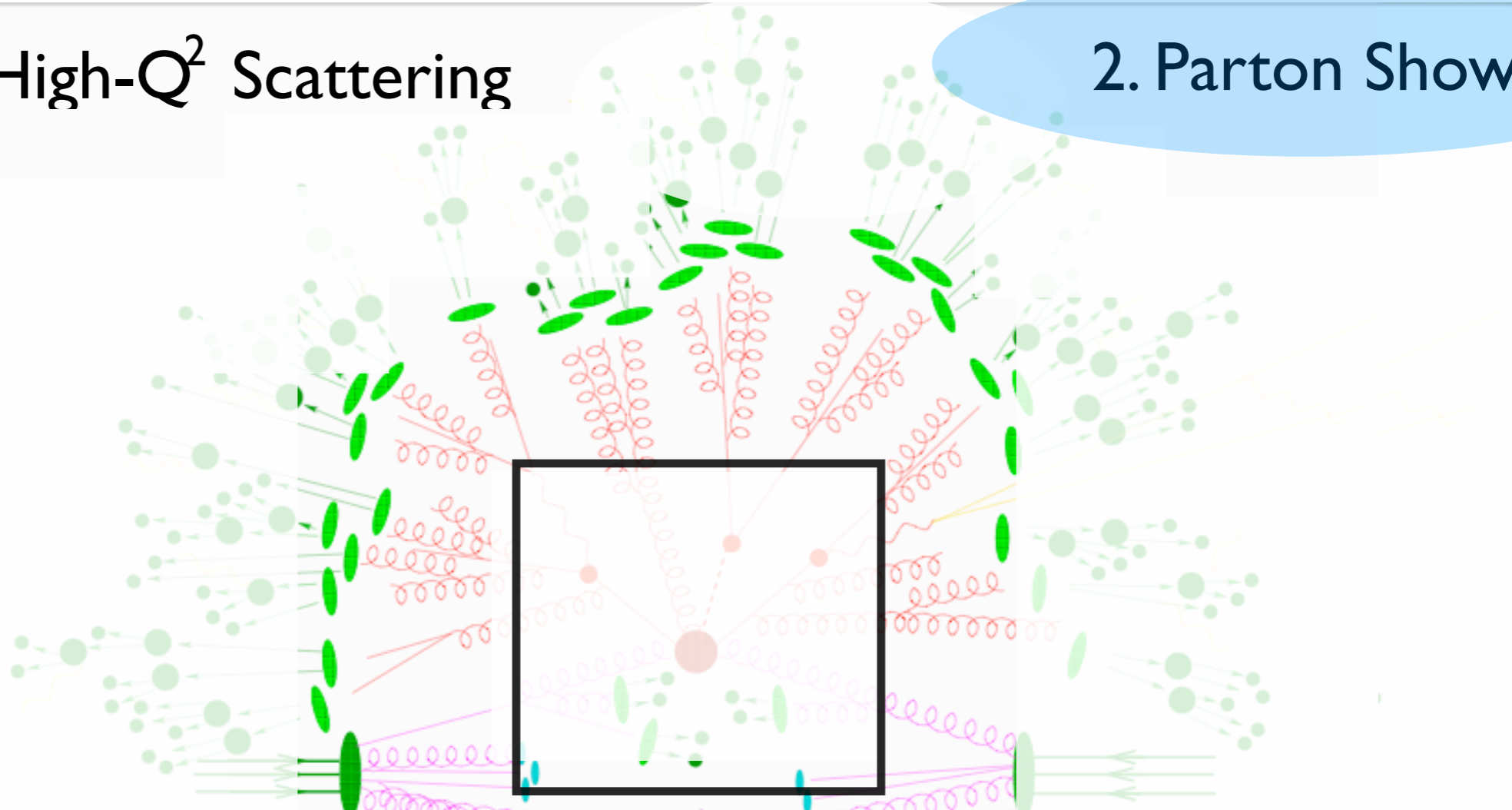
3. Hadronization

4. Underlying Event



1. High- Q^2 Scattering

2. Parton Shower



☞ QCD - "known physics"

☞ universal/ process independent

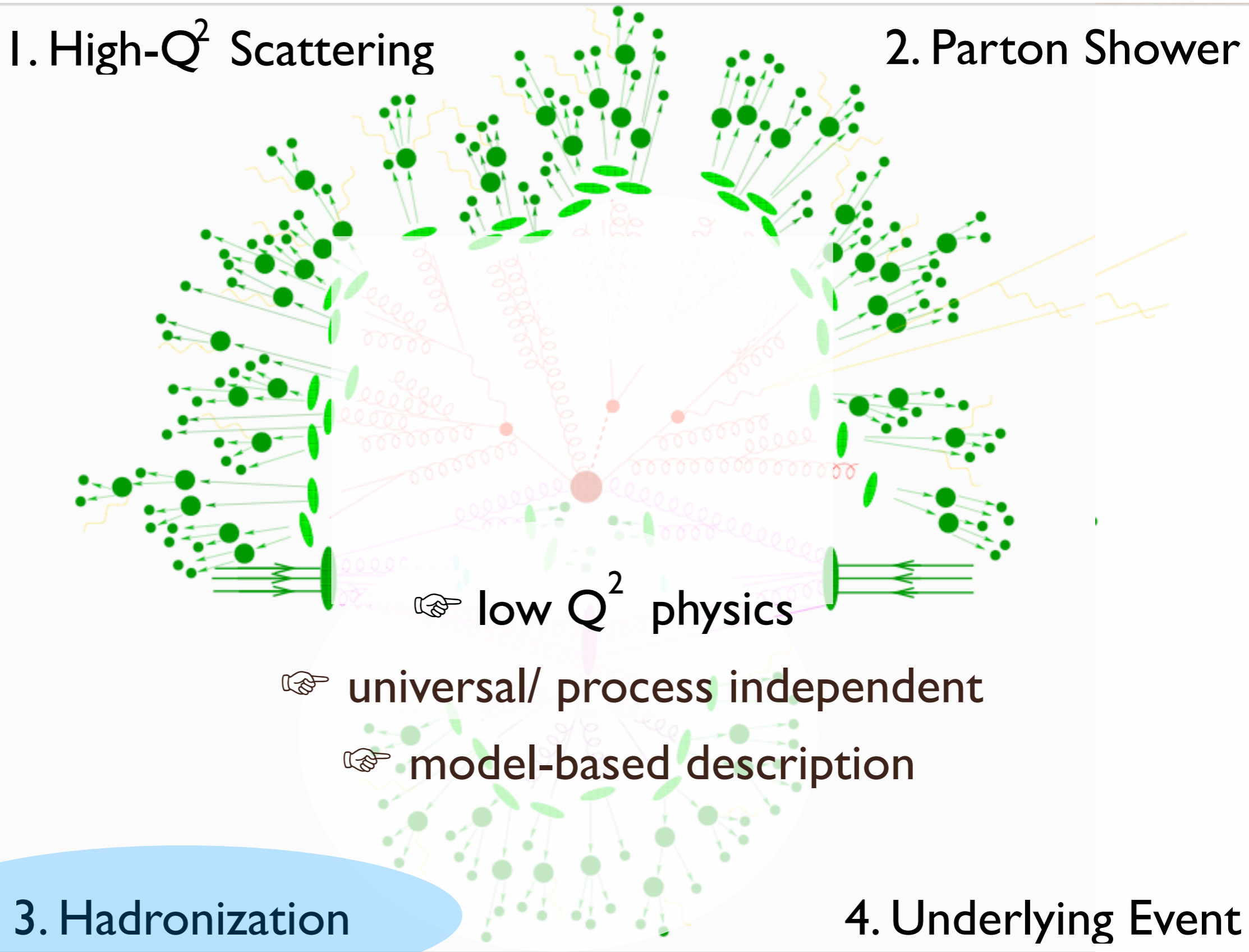
☞ first principles description

3. Hadronization

4. Underlying Event

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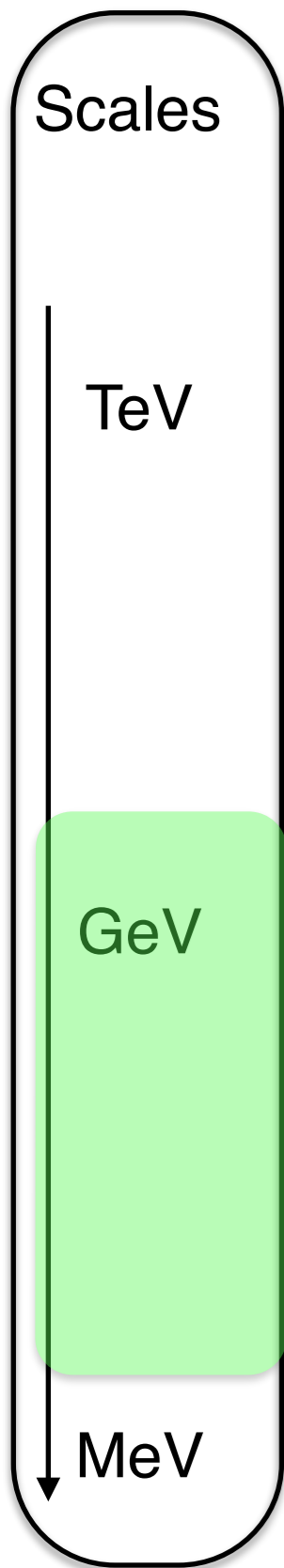


Scales

TeV

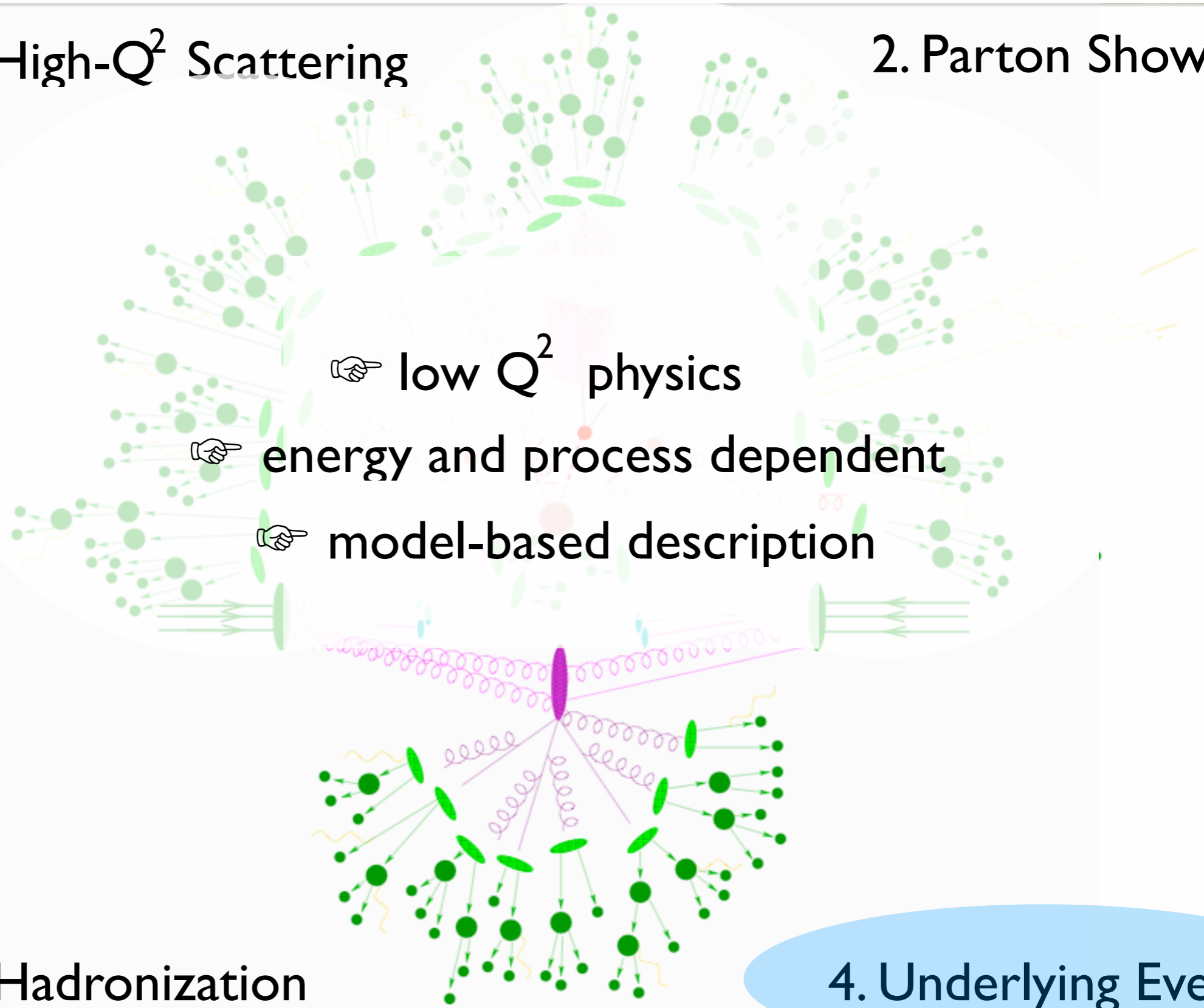
GeV

MeV



1. High- Q^2 Scattering

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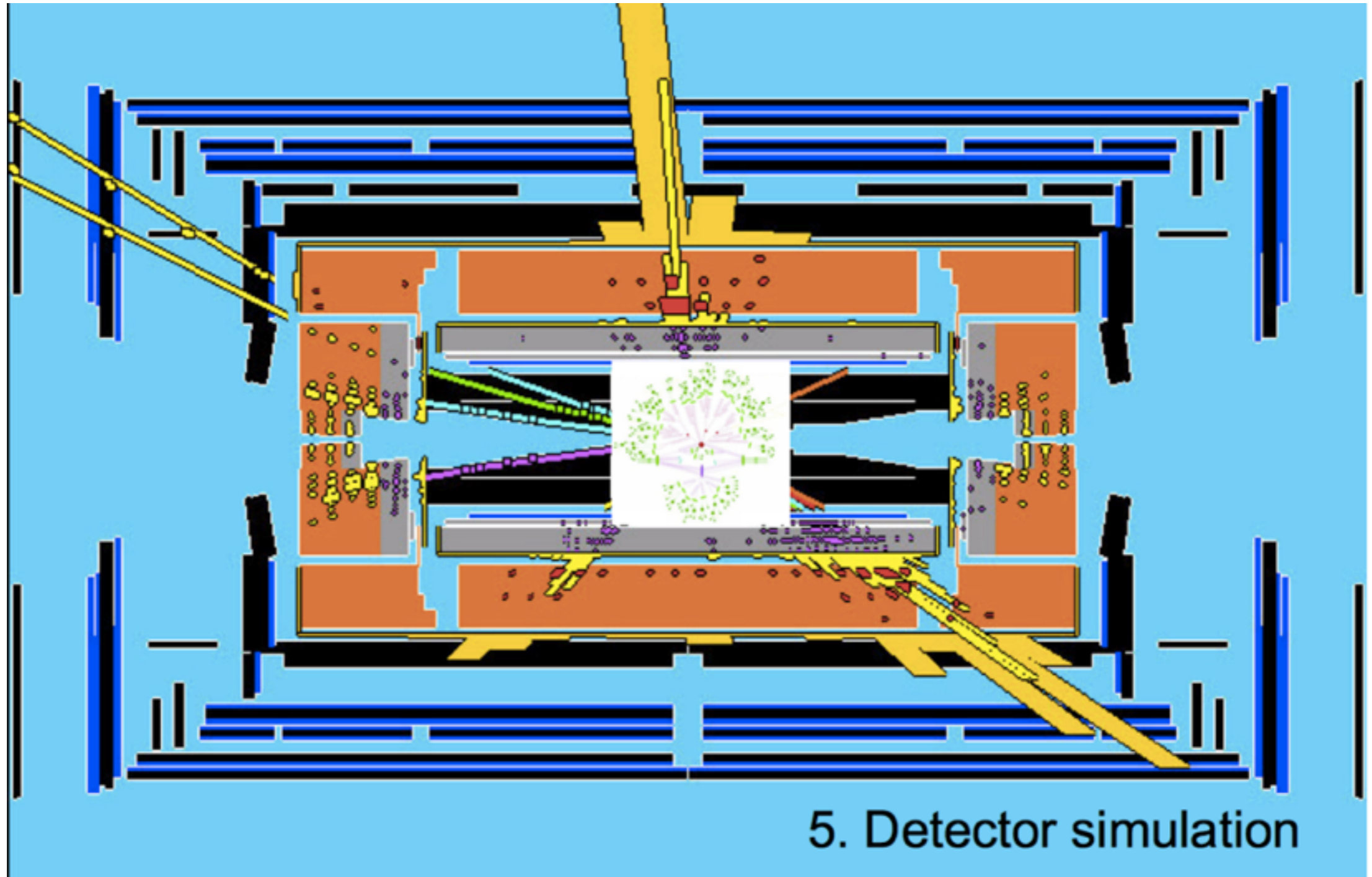
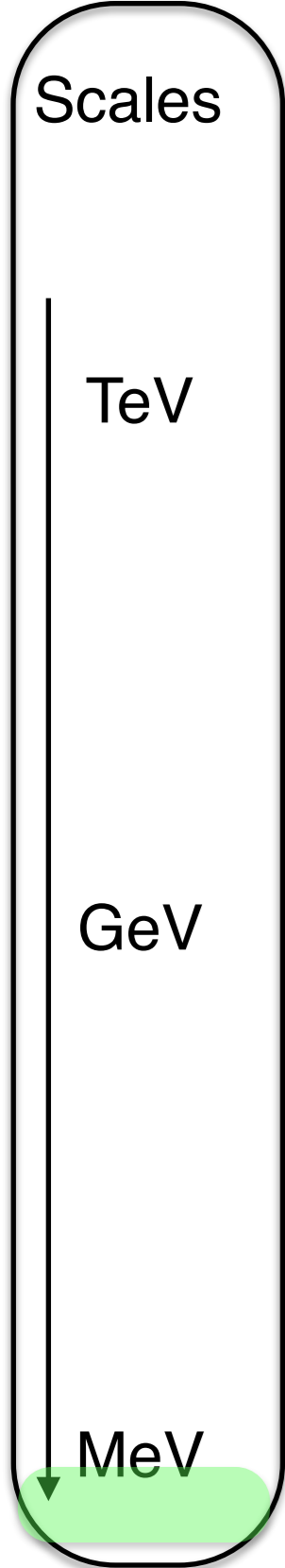
low Q^2 physics

energy and process dependent

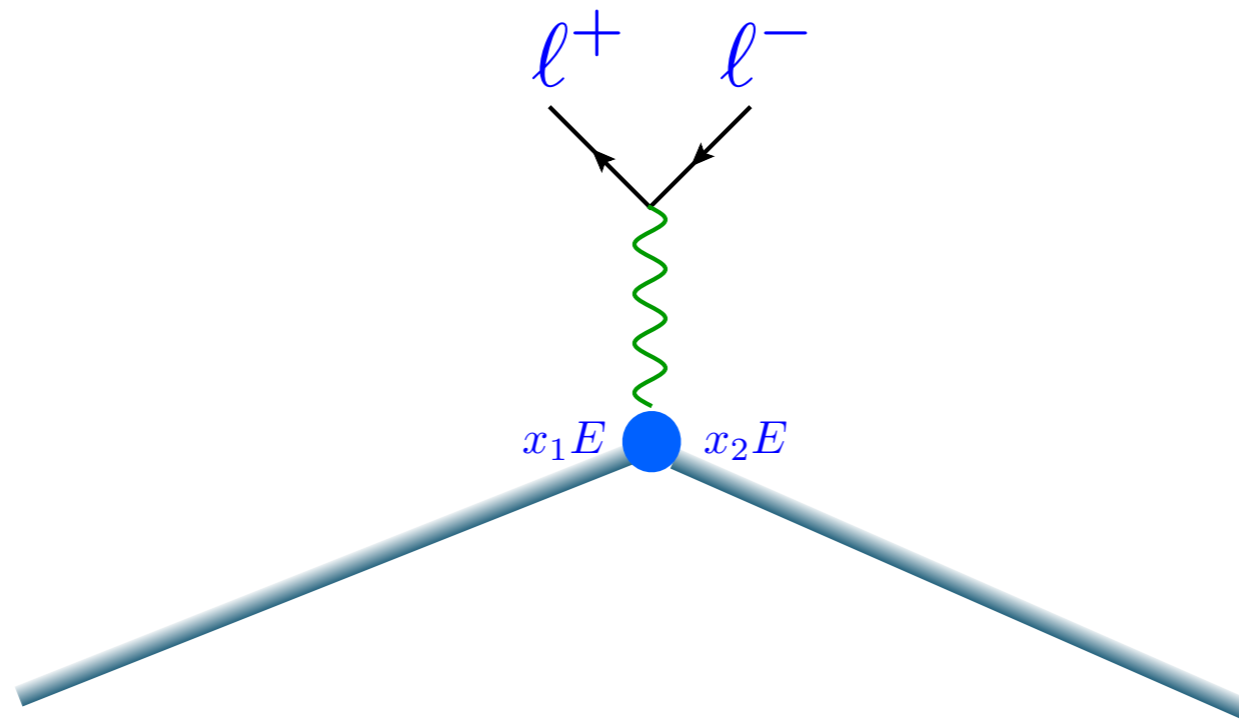
model-based description

3. Hadronization

4. Underlying Event

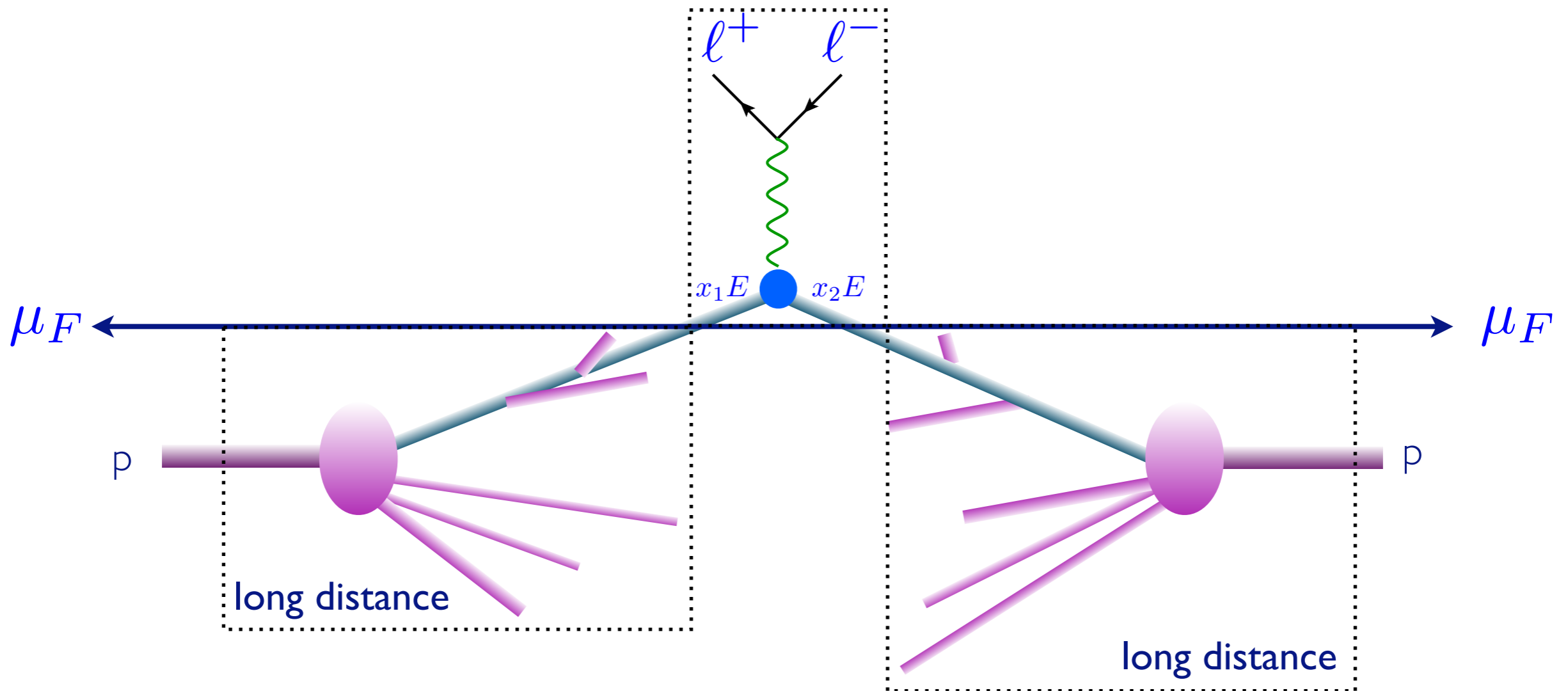


- Multi-scale problem
 - ➔ New physics visible only at High scale
 - ➔ Problem split in different scale



$$\hat{\sigma}_{ab \rightarrow X}(\hat{S}, \mu_F, \mu_R)$$

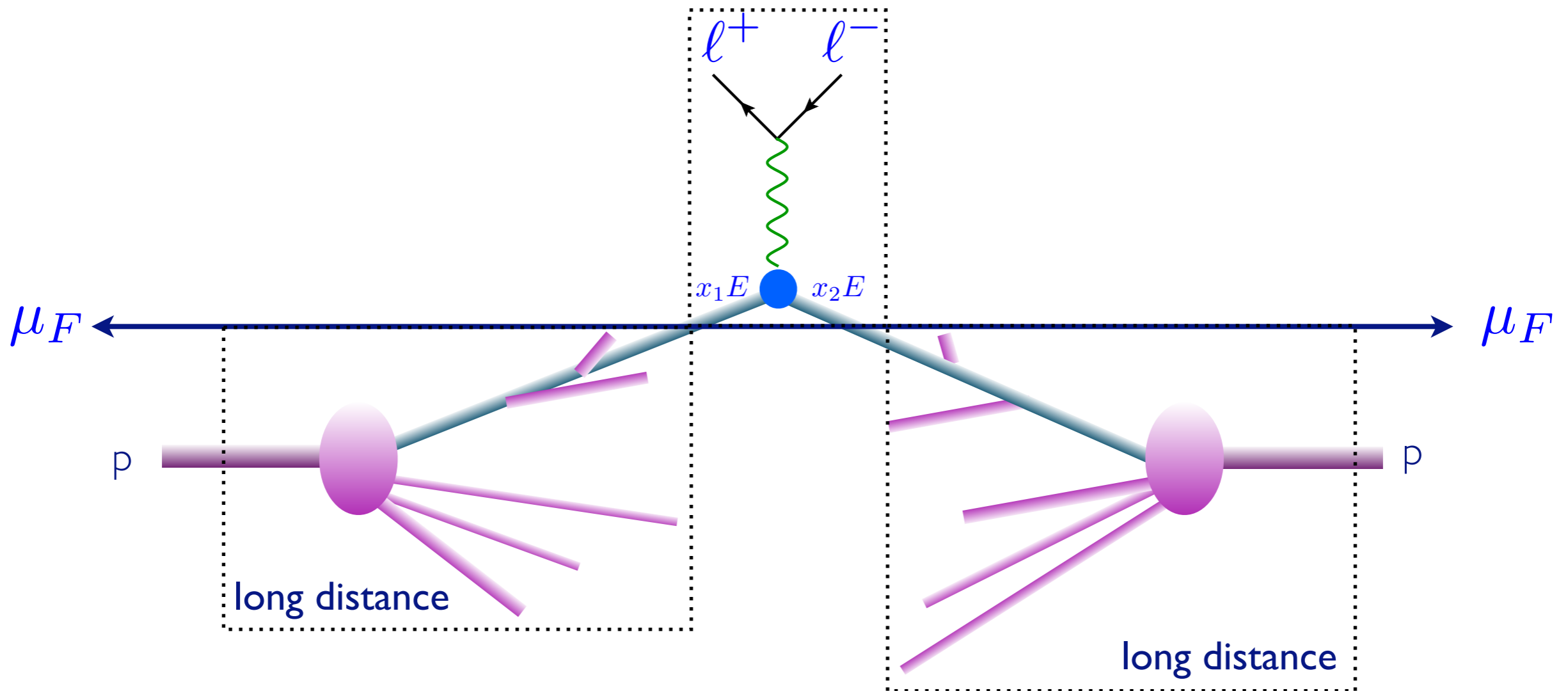
Parton-level cross
section



$$f_a(x_1, \mu_F) f_b(x_2, \mu_F) \hat{\sigma}_{ab \rightarrow X}(\hat{S}, \mu_F, \mu_R)$$

Parton density
functions

Parton-level cross
section



$$\sum_{a,b} \int dx_1 dx_2 d\Phi_{\text{FS}} f_a(x_1, \mu_F) f_b(x_2, \mu_F) \hat{\sigma}_{ab \rightarrow X}(\hat{S}, \mu_F, \mu_R)$$

Phase-space integral
Parton density functions
Parton-level cross section

$d\hat{\sigma}_{ab \rightarrow X}(\hat{s}, \mu_F, \mu_R)$ Parton-level cross section

- The parton-level cross section can be computed as a series in perturbation theory, using the coupling constant as an expansion parameter, schematically:

$$\hat{\sigma} = \sigma^{\text{Born}} \left(1 + \frac{\alpha_s}{2\pi} \sigma^{(1)} + \left(\frac{\alpha_s}{2\pi} \right)^2 \sigma^{(2)} + \left(\frac{\alpha_s}{2\pi} \right)^3 \sigma^{(3)} + \dots \right)$$

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LO
predictions

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LO
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NLO
corrections

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LO
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LO
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NLO
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NNLO
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N3LO or NNNLO
corrections

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LO
predictions

NLO
corrections

NNLO
corrections

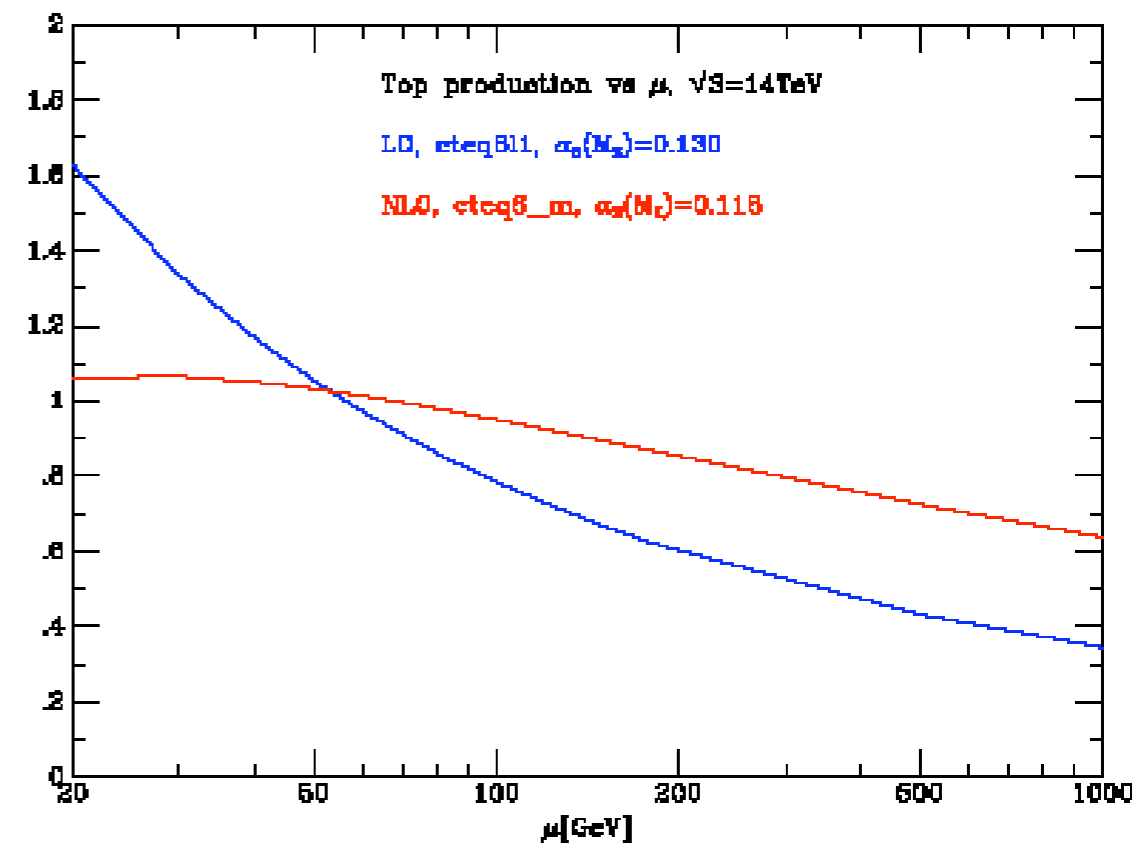
N3LO or NNNLO
corrections

- Including higher corrections improves predictions and reduces theoretical uncertainties

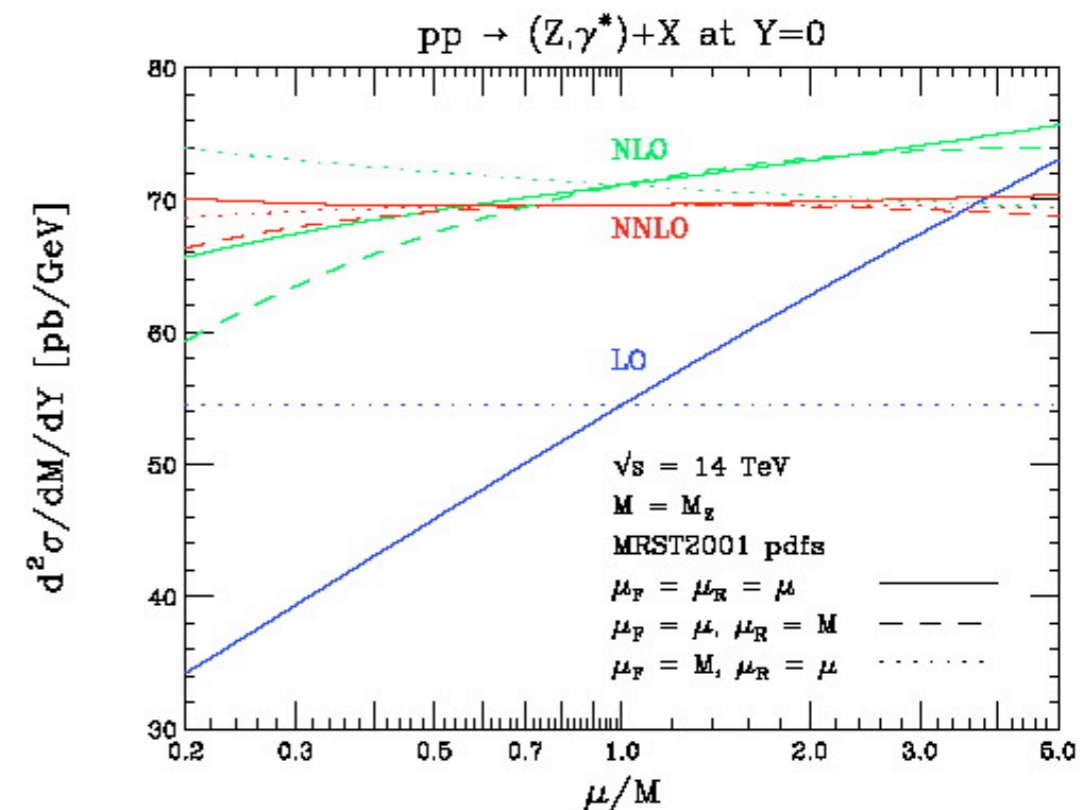
$$d\sigma = \sum_{a,b} \int dx_1 dx_2 f_a(x_1, \mu_F) f_b(x_2, \mu_F) d\hat{\sigma}_{ab \rightarrow X}(\hat{s}, \mu_F, \mu_R)$$

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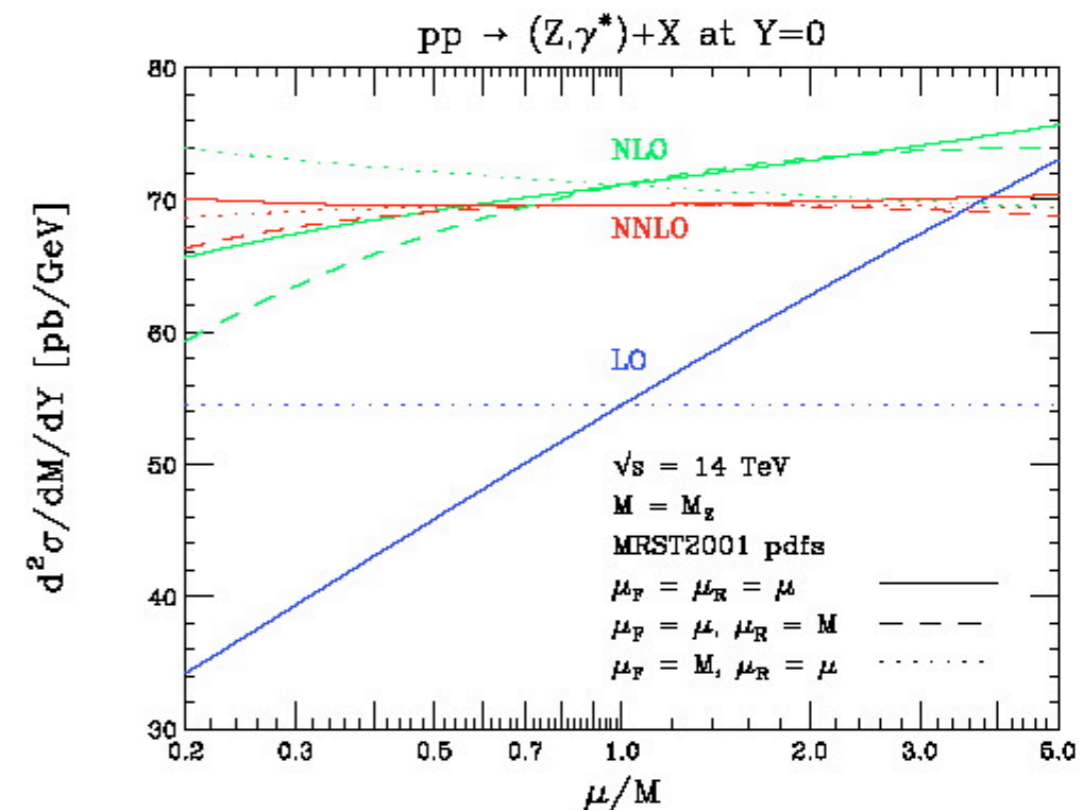
- Leading Order predictions can depend strongly on the renormalization and factorization scales
- Including higher order corrections reduces the dependence on these scales



- NNLO is the current state-of-the-art. There are only a few results available: Higgs, Drell-Yan, $t\bar{t}$
- Why do we need it?
 - control of the uncertainties in a calculation
 - It is “mandatory” if NLO corrections are very large to check the behavior of the perturbative series
 - It is needed for Standard Candles and very precise tests of perturbation theory, exploiting all the available information, e.g. for determining NNLO PDF sets

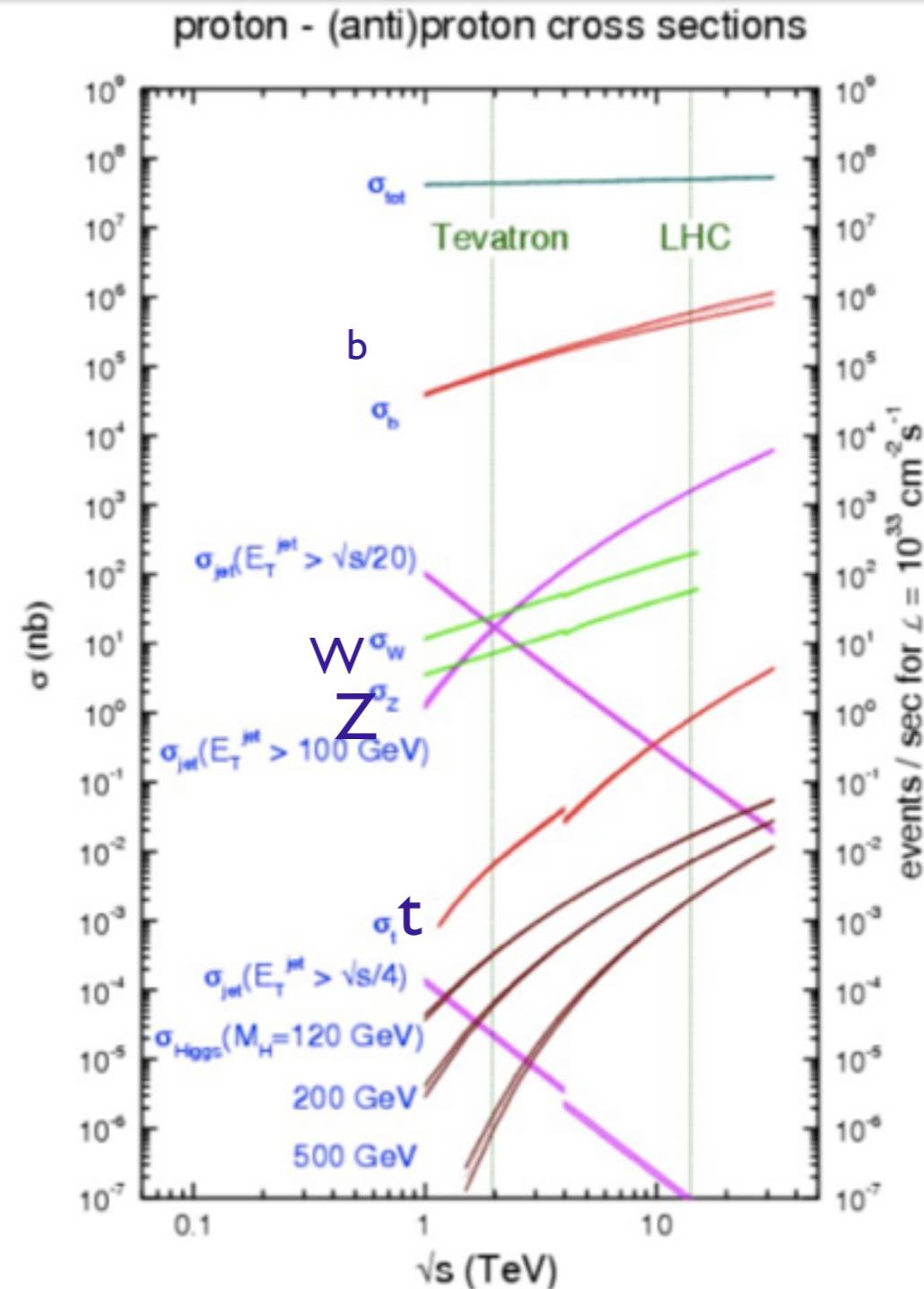


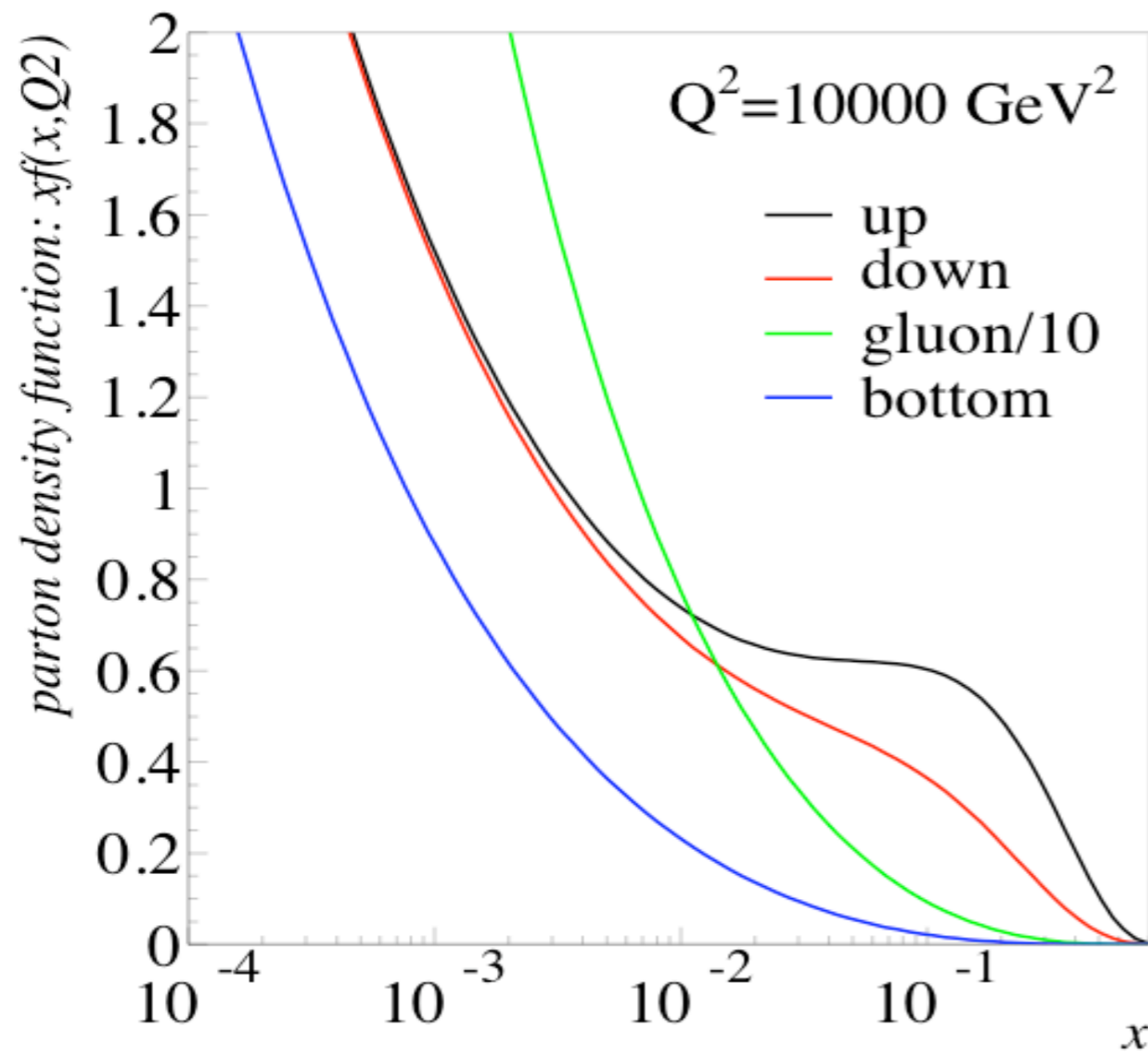
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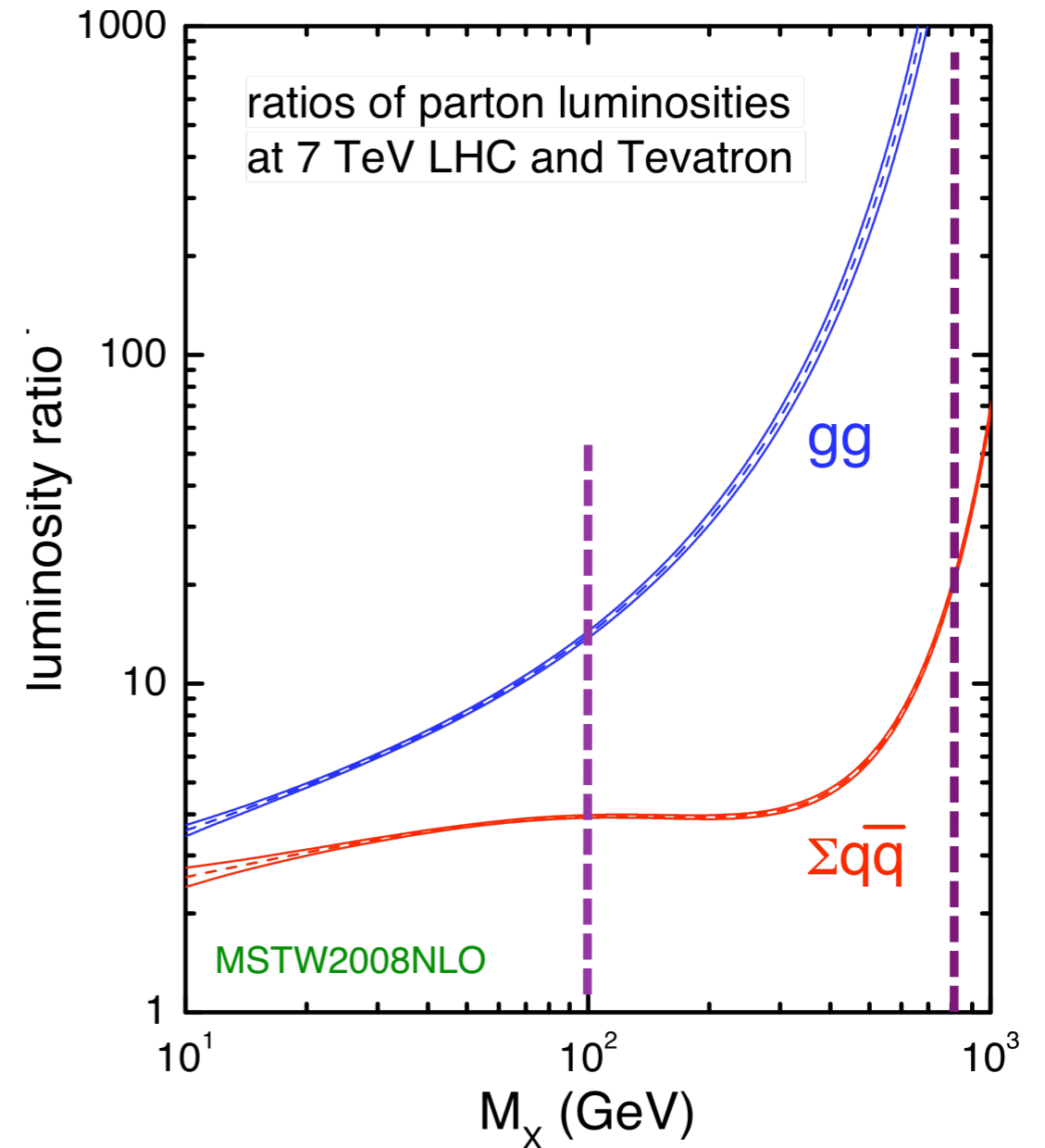
Let's focus on LO

$$\sum_{a,b} \int dx_1 dx_2 d\Phi_{\text{FS}} f_a(x_1, \mu_F) f_b(x_2, \mu_F) \hat{\sigma}_{ab \rightarrow X}(\hat{S}, \mu_F, \mu_R)$$

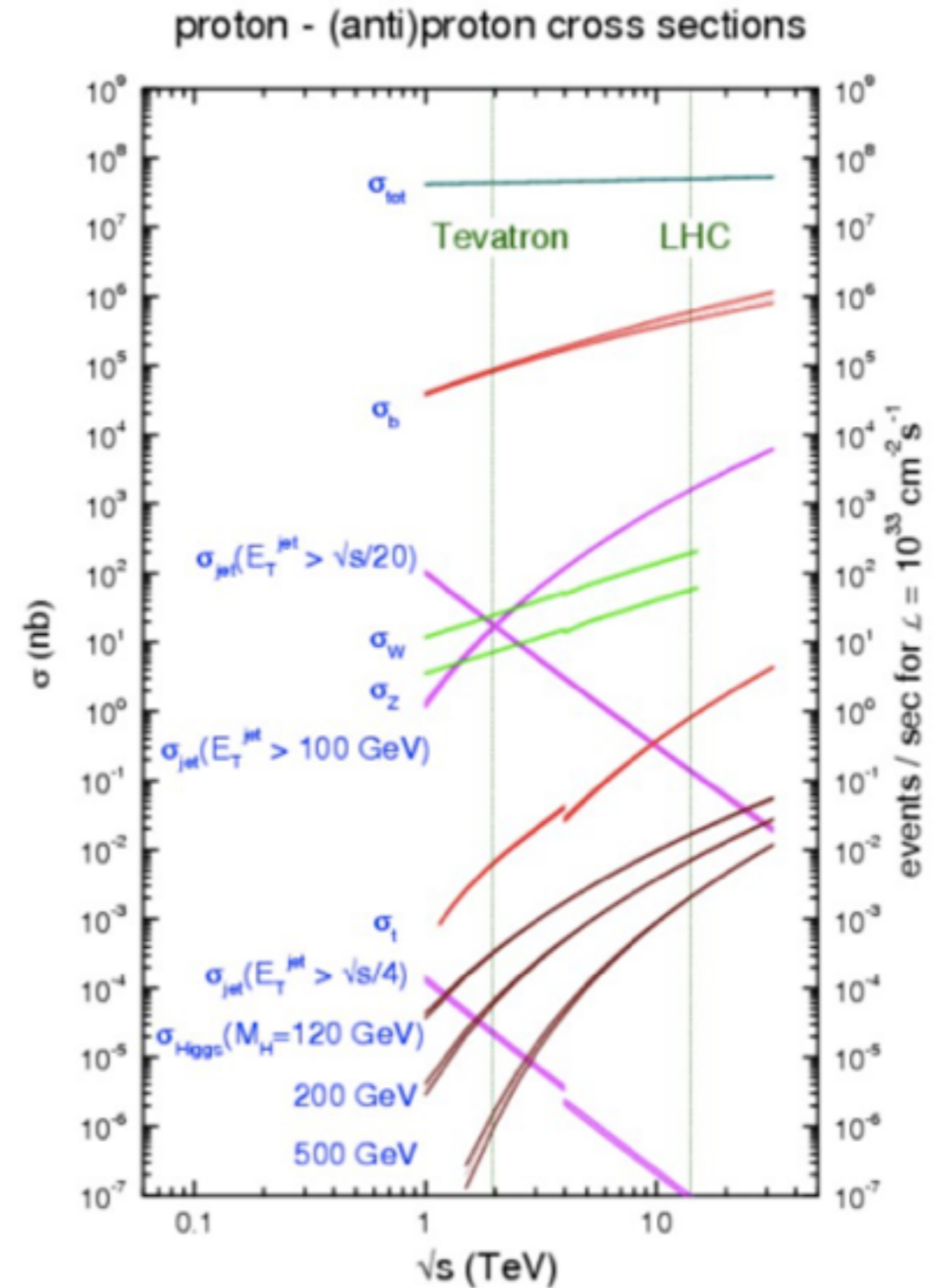
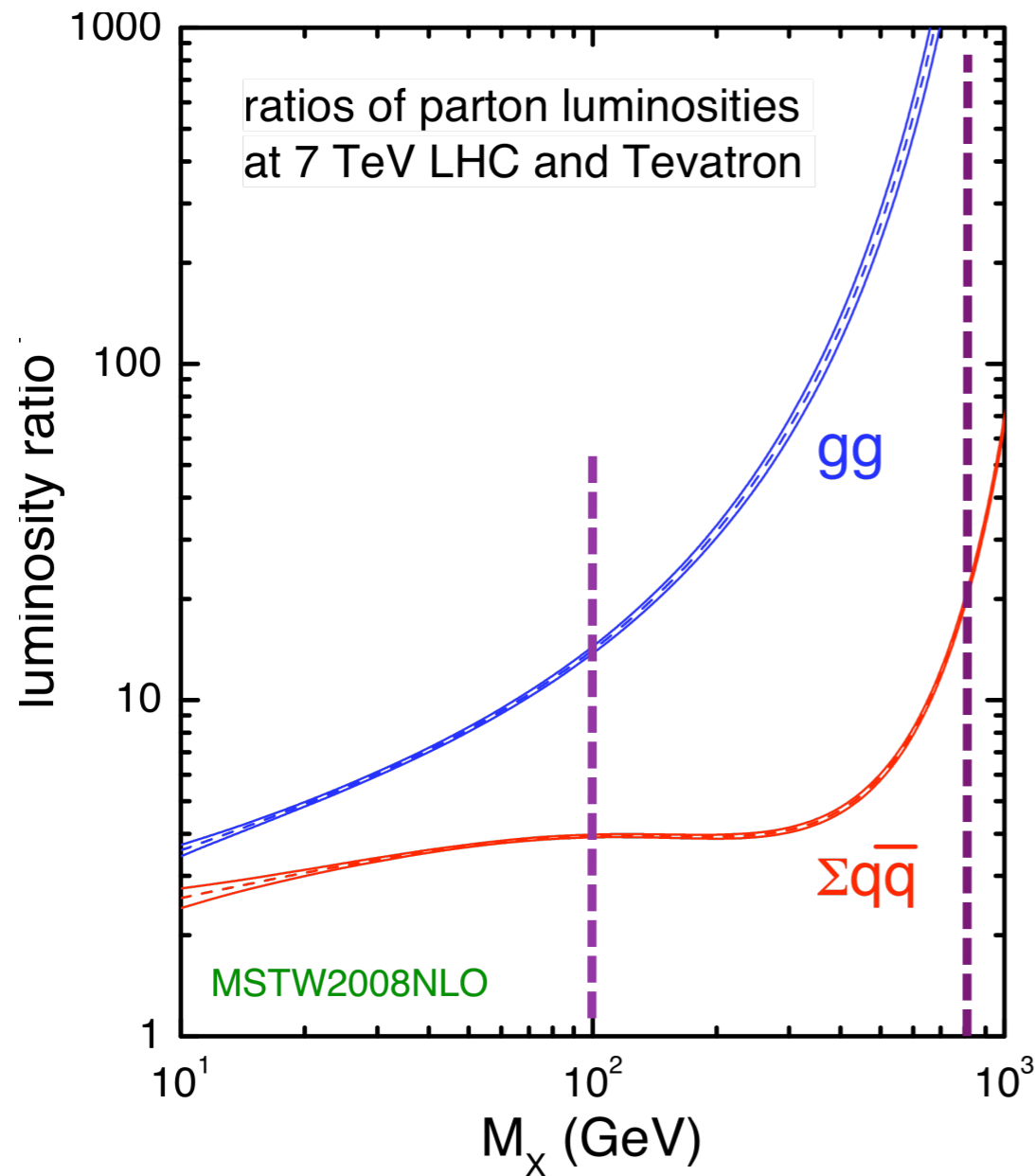


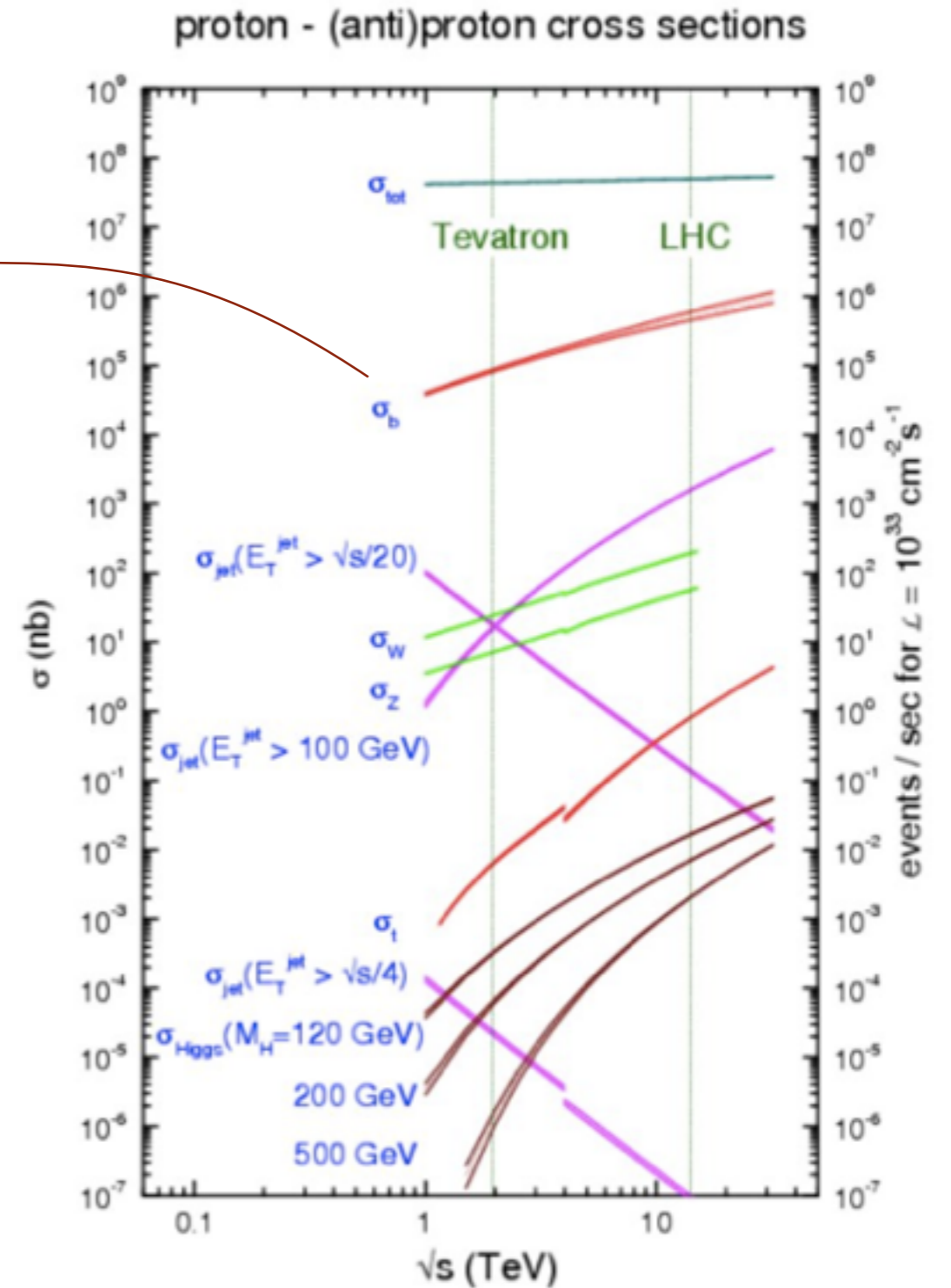
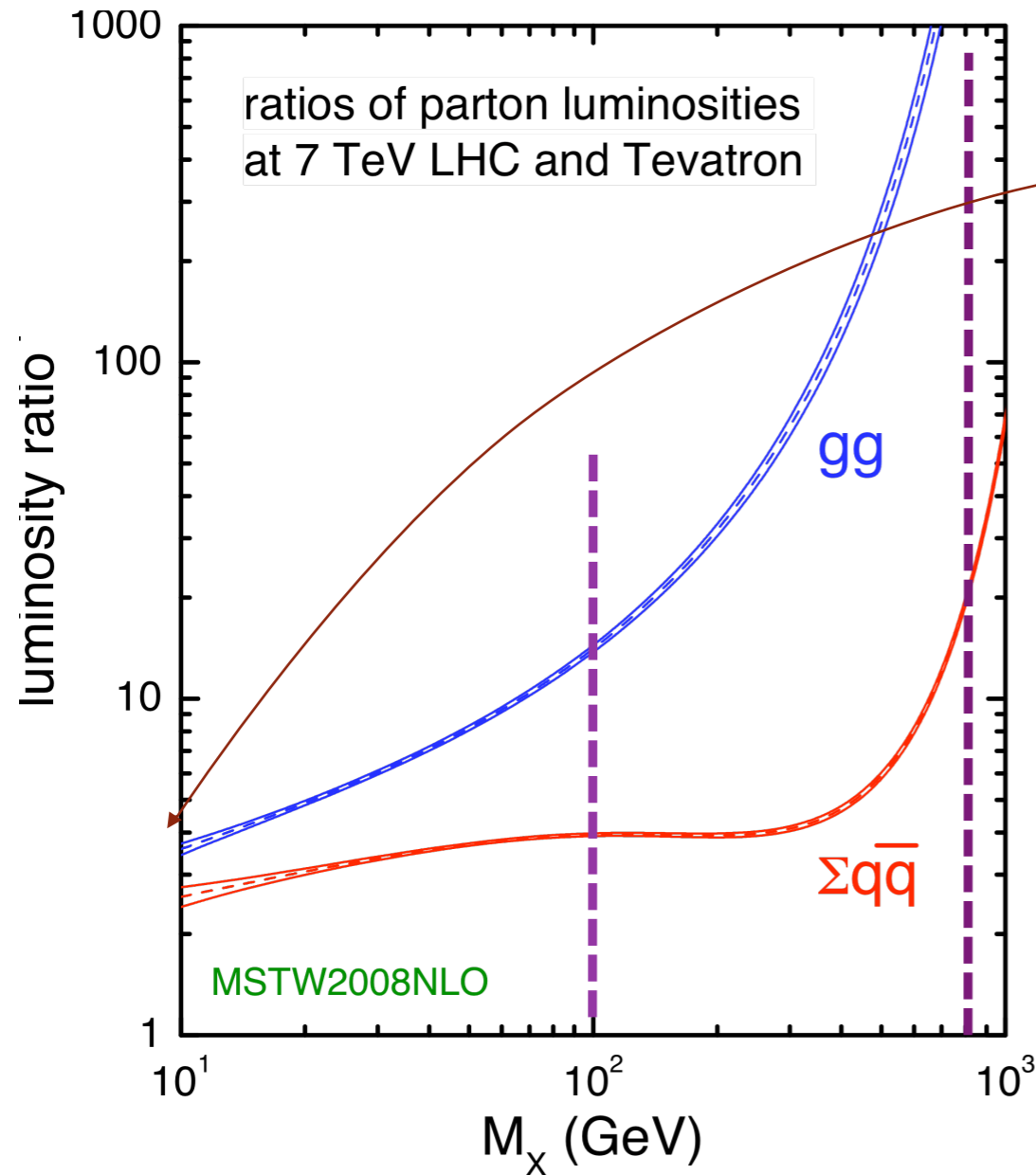


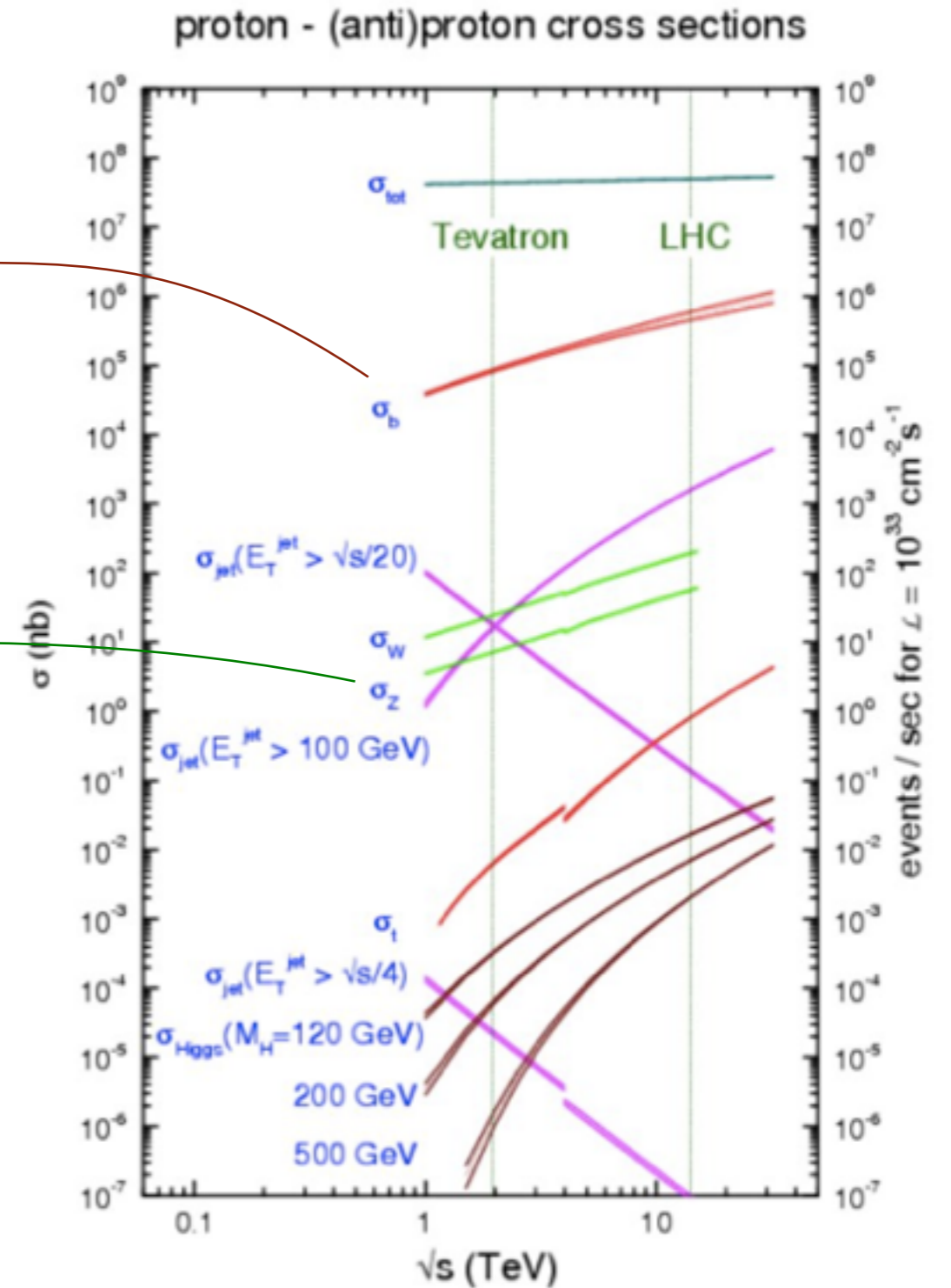
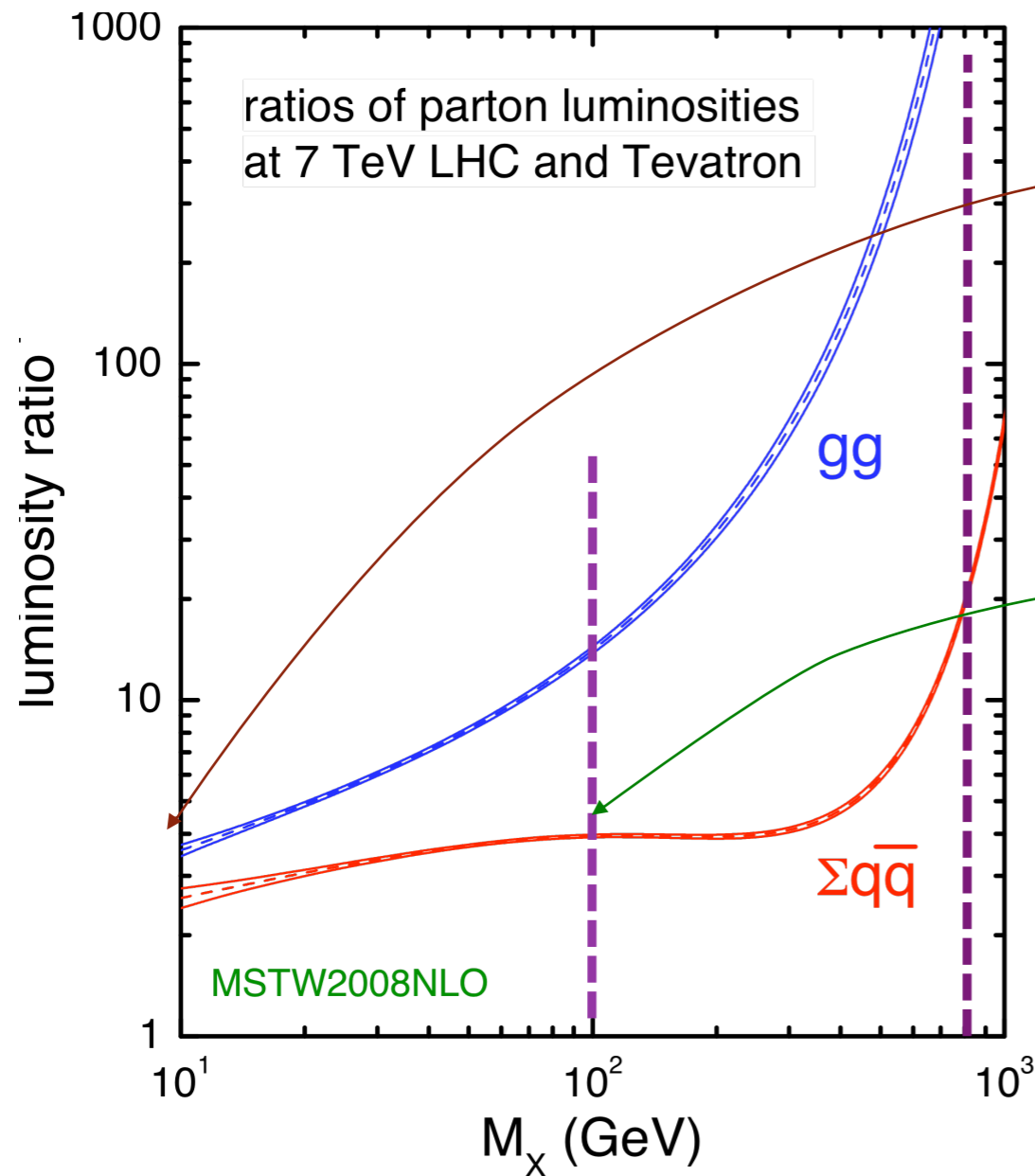
At small x (small \hat{s}), gluon domination.
At large x valence quarks

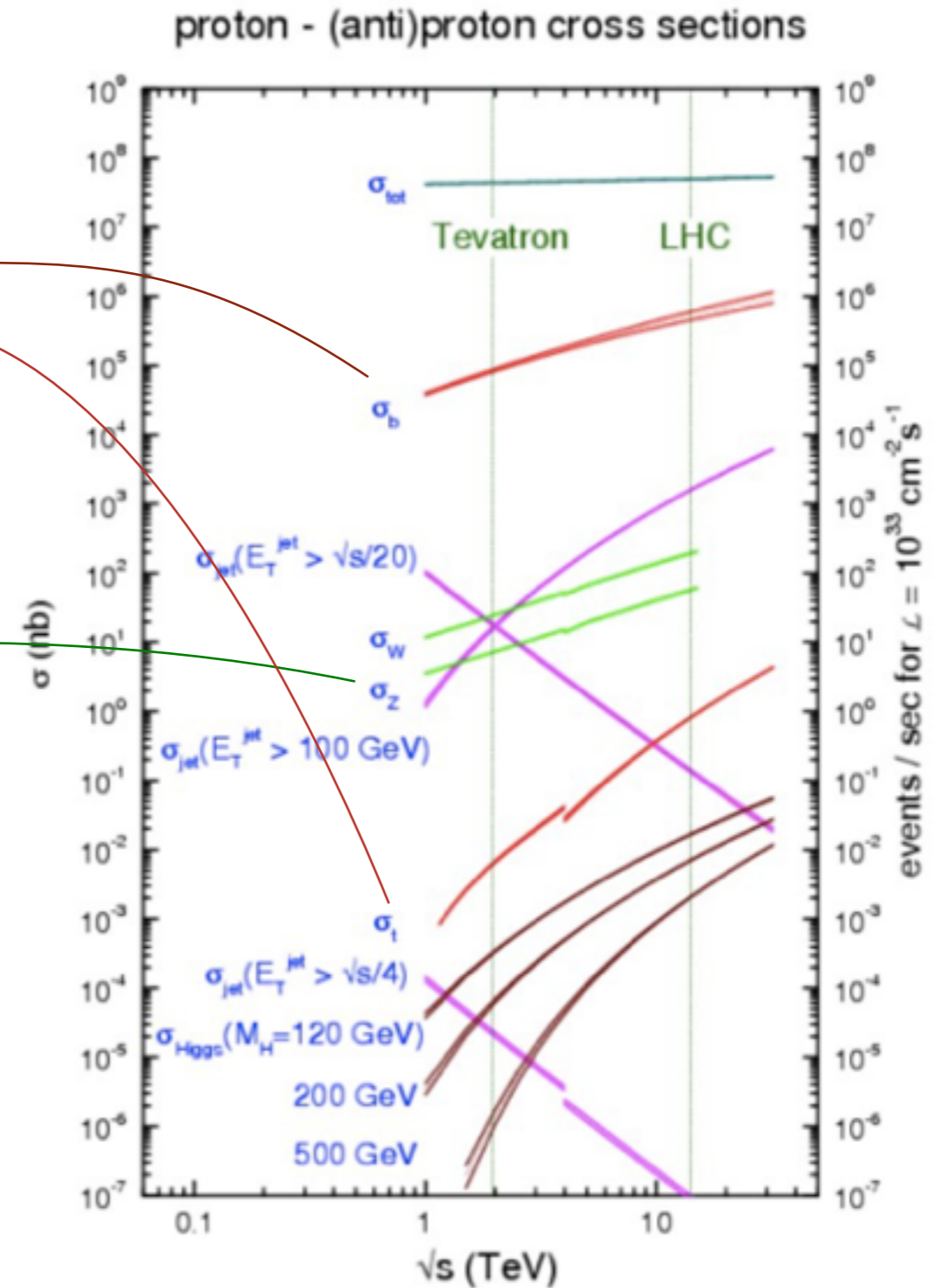
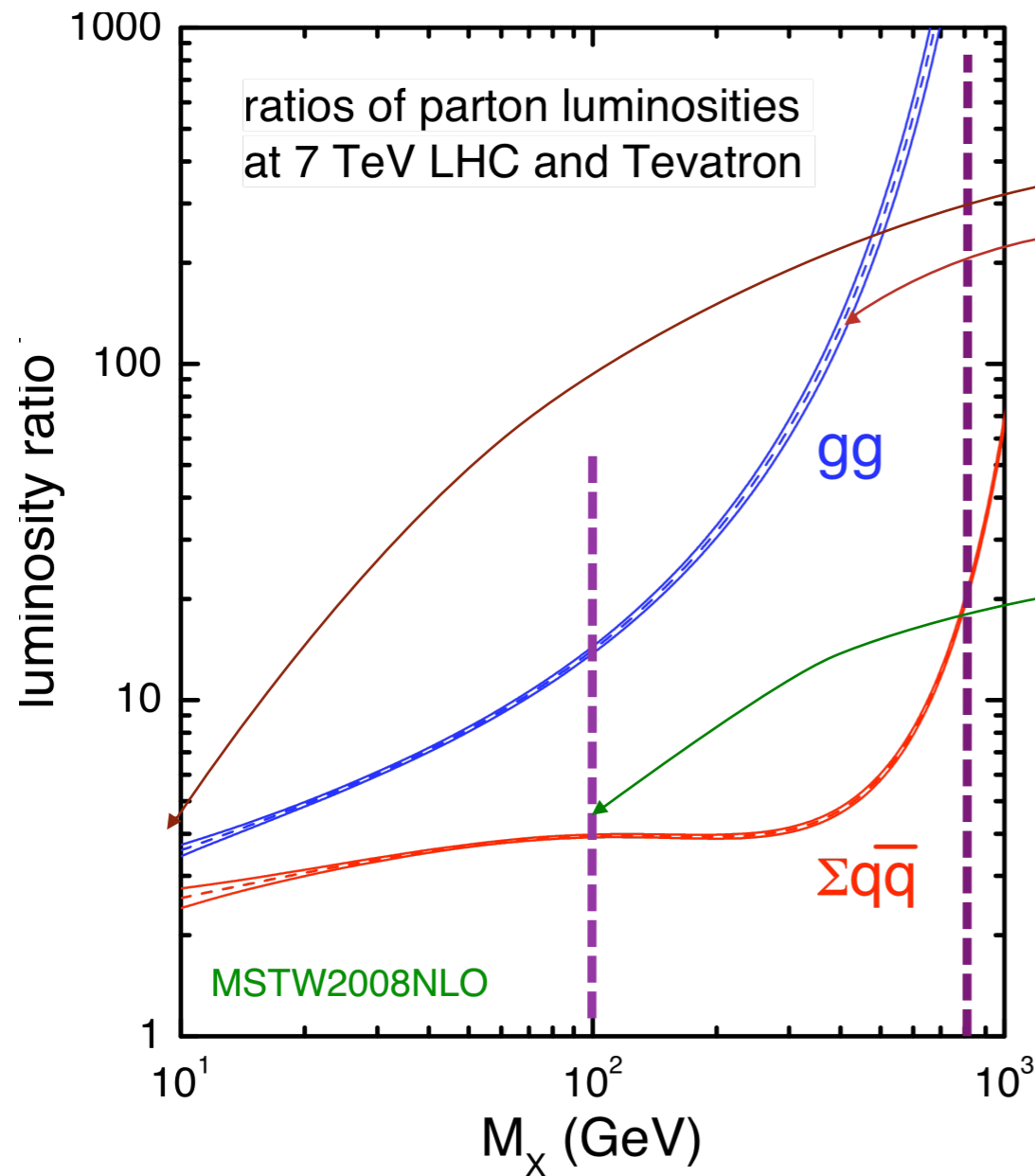


LHC formidable at large mass –
For low mass, Tevatron backgrounds smaller









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Phase-space integral
Parton density functions
Parton-level cross section

- PDF: content of the proton
 - ➔ Define the physics/processes that will dominate on your accelerator
- NLO/NNLO: Reduce scale uncertainty linked to your division of your multi-scale problem

Calculate a given process (e.g. gluino pair)

- Determine the production mechanism

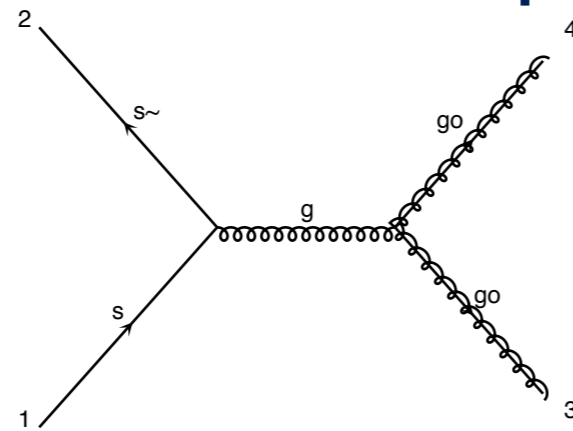


diagram 1 QCD=2, QED=0

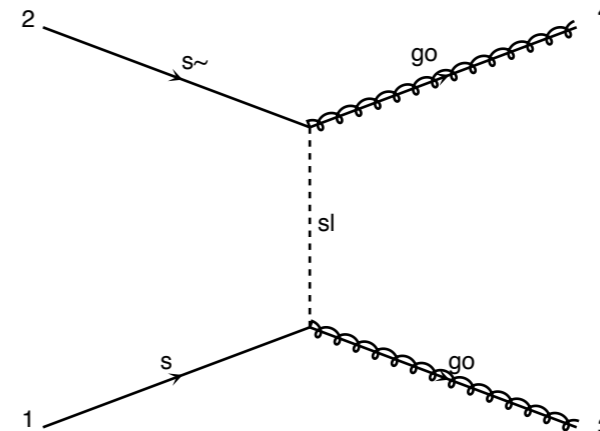


diagram 2 QCD=2, QED=0

- Evaluate the matrix-element

$$|\mathcal{M}|^2 \quad \Rightarrow \text{Need Feynman Rules!}$$

- Phase-Space Integration

$$\sigma = \frac{1}{2s} \int |\mathcal{M}|^2 d\Phi(n)$$

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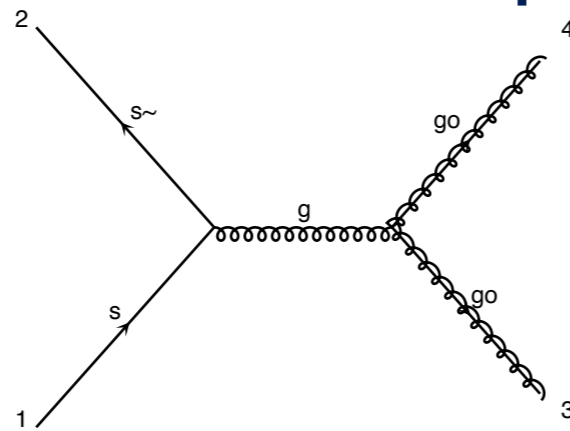


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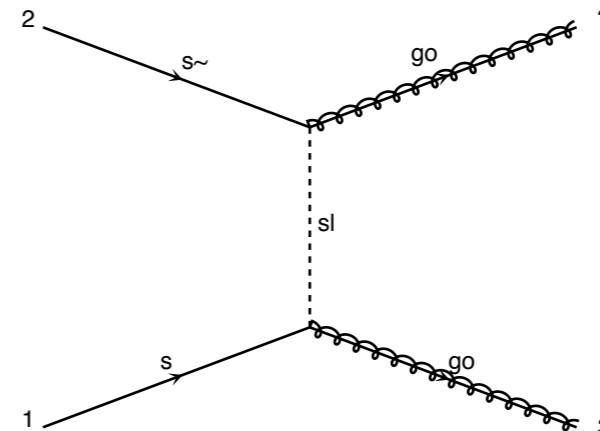


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Easy enough

Hard

Very Hard
(in general)

Calculate a given process (e.g. gluino pair)

- Determine the production mechanism

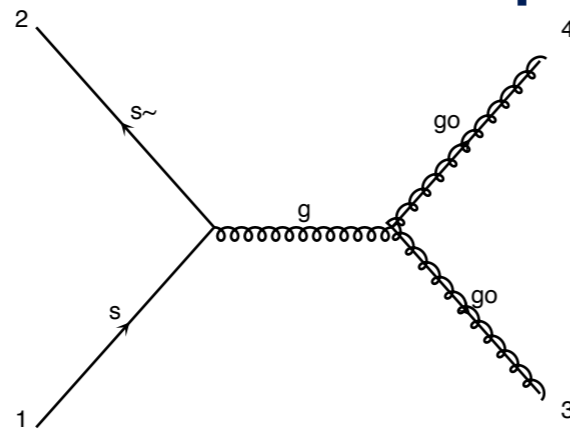


diagram 1 QCD=2, QED=0

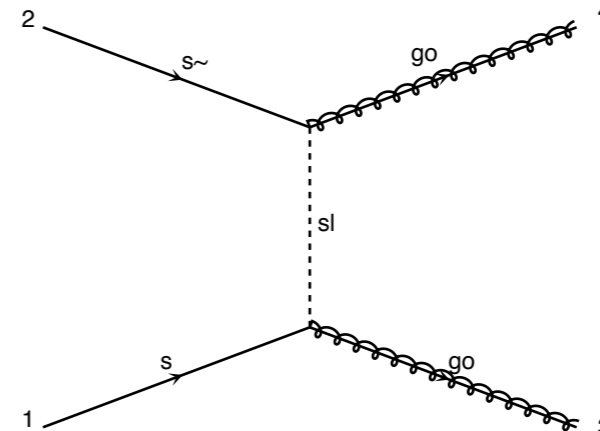


diagram 2 QCD=2, QED=0

Easy enough

- Evaluate the matrix-element

$$|\mathcal{M}|^2 \quad \rightarrow \text{Need Feynman Rules!}$$

Hard

Next

- Phase-Space Integration

$$\sigma = \frac{1}{2s} \int |\mathcal{M}|^2 d\Phi(n)$$

Very

Hard
(in general)

Calculate a given process (e.g. gluino pair)

- Determine the production mechanism

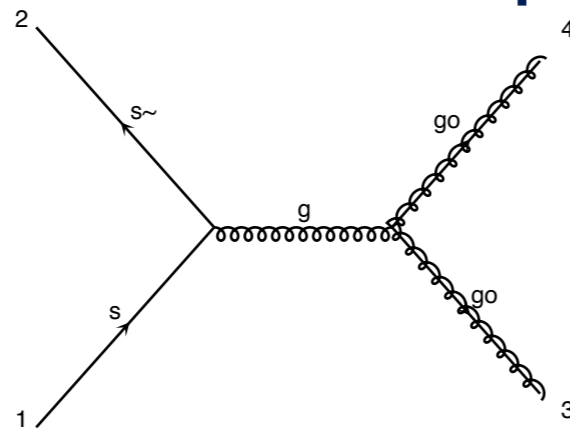


diagram 1 QCD=2, QED=0

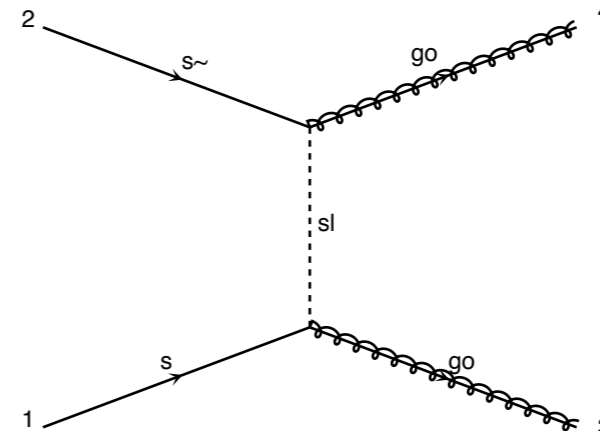


diagram 2 QCD=2, QED=0

Easy enough

- Evaluate the matrix-element

$$|\mathcal{M}|^2 \quad \rightarrow \text{Need Feynman Rules!}$$

Hard

Next

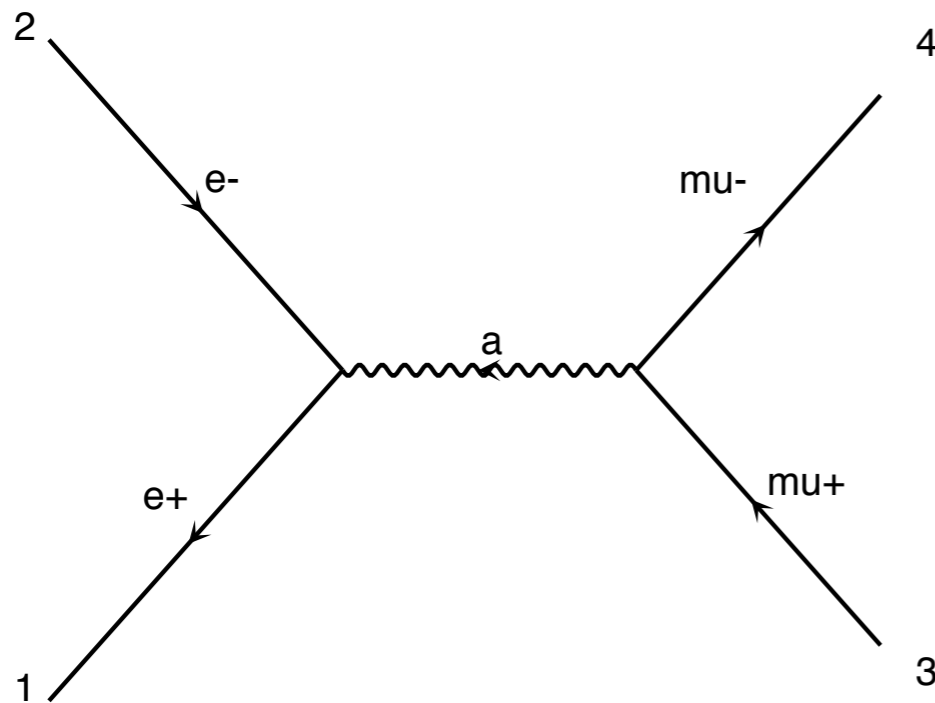
- Phase-Space Integration

$$\sigma = \frac{1}{2s} \int |\mathcal{M}|^2 d\Phi(n)$$

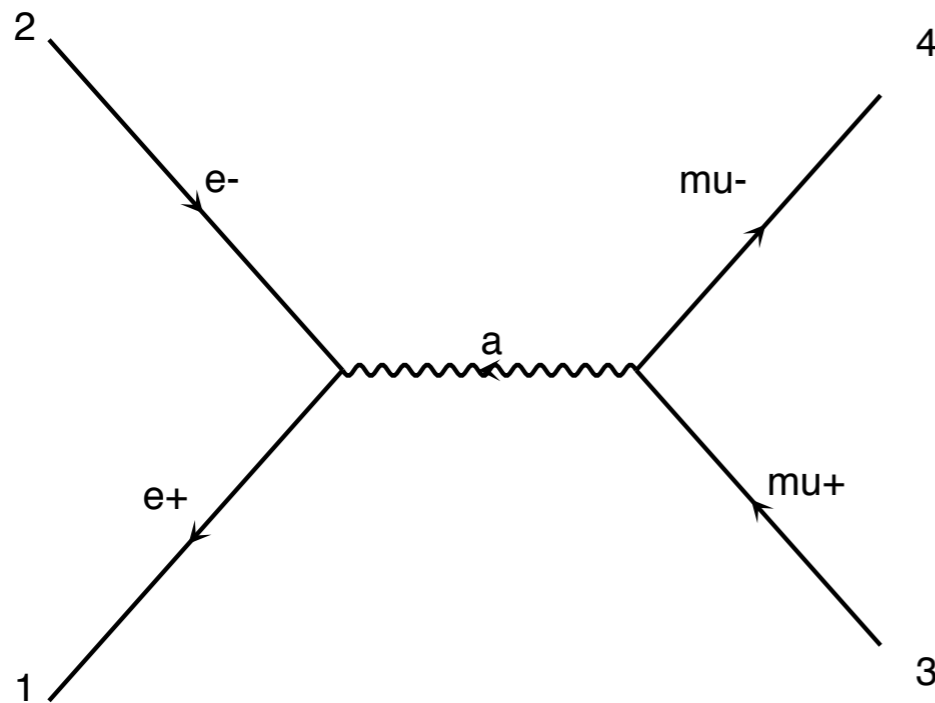
Very

Hard
(in general)

After

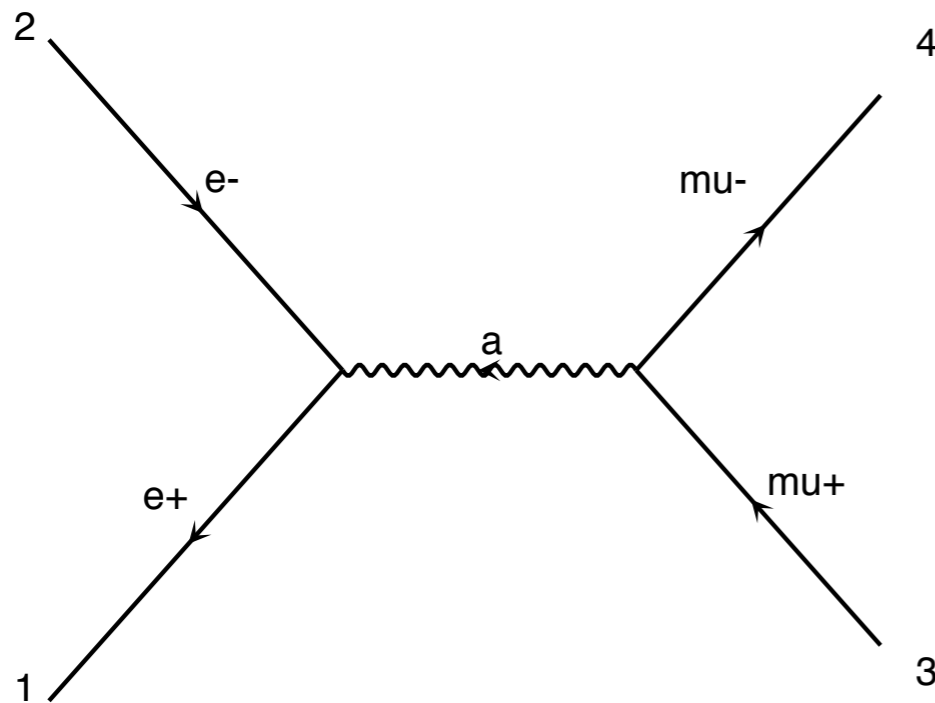


$$\mathcal{M} = e^2 (\bar{u} \gamma^\mu v) \frac{g_{\mu\nu}}{q^2} (\bar{v} \gamma^\nu u)$$



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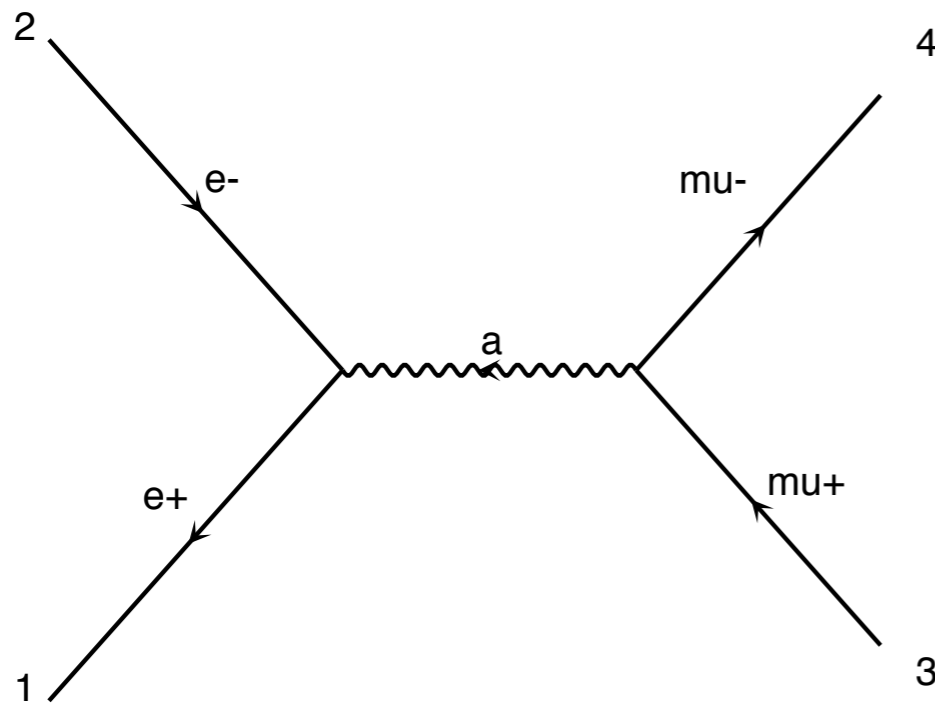
$$\frac{1}{4} \sum_{pol} |\mathcal{M}|^2 = \frac{1}{4} \sum_{pol} \mathcal{M}^* \mathcal{M}$$



$$\mathcal{M} = e^2 (\bar{u} \gamma^\mu v) \frac{g_{\mu\nu}}{q^2} (\bar{v} \gamma^\nu u)$$

$$\frac{1}{4} \sum_{pol} |\mathcal{M}|^2 = \frac{1}{4} \sum_{pol} \mathcal{M}^* \mathcal{M}$$

$$\sum_{pol} \bar{u} u = \not{p} + m$$

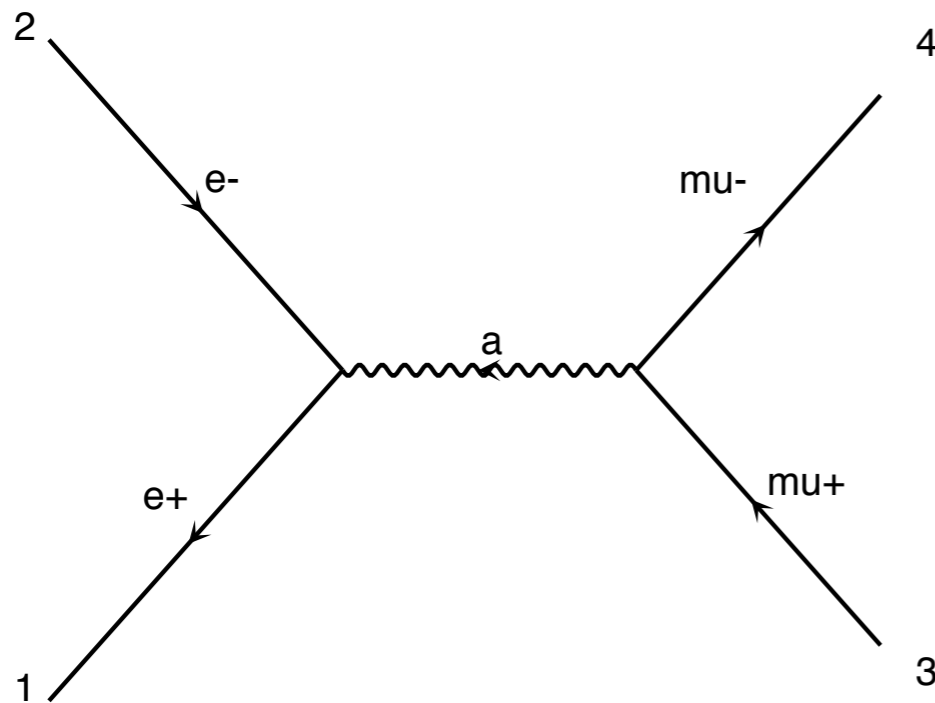


$$\mathcal{M} = e^2 (\bar{u} \gamma^\mu v) \frac{g_{\mu\nu}}{q^2} (\bar{v} \gamma^\nu u)$$

$$\frac{1}{4} \sum_{pol} |\mathcal{M}|^2 = \frac{1}{4} \sum_{pol} \mathcal{M}^* \mathcal{M}$$

$$\sum_{pol} \bar{u} u = \not{p} + m$$

$$\rightarrow \frac{e^4}{4q^4} \text{Tr}[\not{p}_1 \gamma^\mu \not{p}_2 \gamma^\nu] \text{Tr}[\not{p}_3 \gamma_\mu \not{p}_4 \gamma_\nu]$$



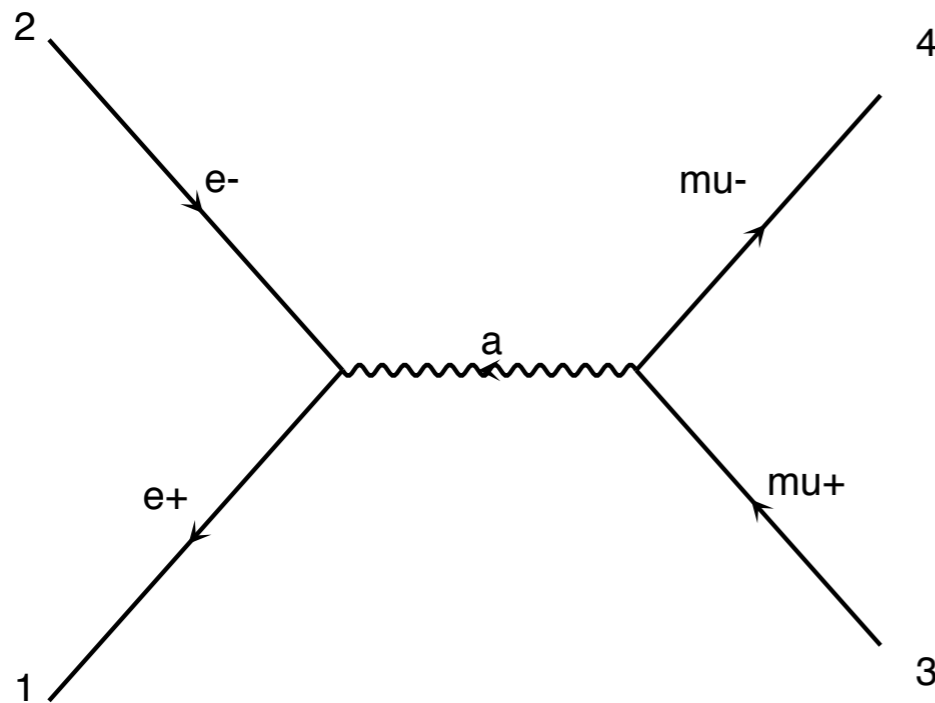
$$\mathcal{M} = e^2 (\bar{u} \gamma^\mu v) \frac{g_{\mu\nu}}{q^2} (\bar{v} \gamma^\nu u)$$

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$$\rightarrow \frac{e^4}{4q^4} Tr[\not{p}_1 \gamma^\mu \not{p}_2 \gamma^\nu] Tr[\not{p}_3 \gamma_\mu \not{p}_4 \gamma_\nu]$$

$$\rightarrow \frac{8e^4}{q^4} [(p_1 \cdot p_3)(p_2 \cdot p_4) + (p_1 \cdot p_4)(p_2 \cdot p_3)]$$



$$\mathcal{M} = e^2 (\bar{u} \gamma^\mu v) \frac{g_{\mu\nu}}{q^2} (\bar{v} \gamma^\nu u)$$

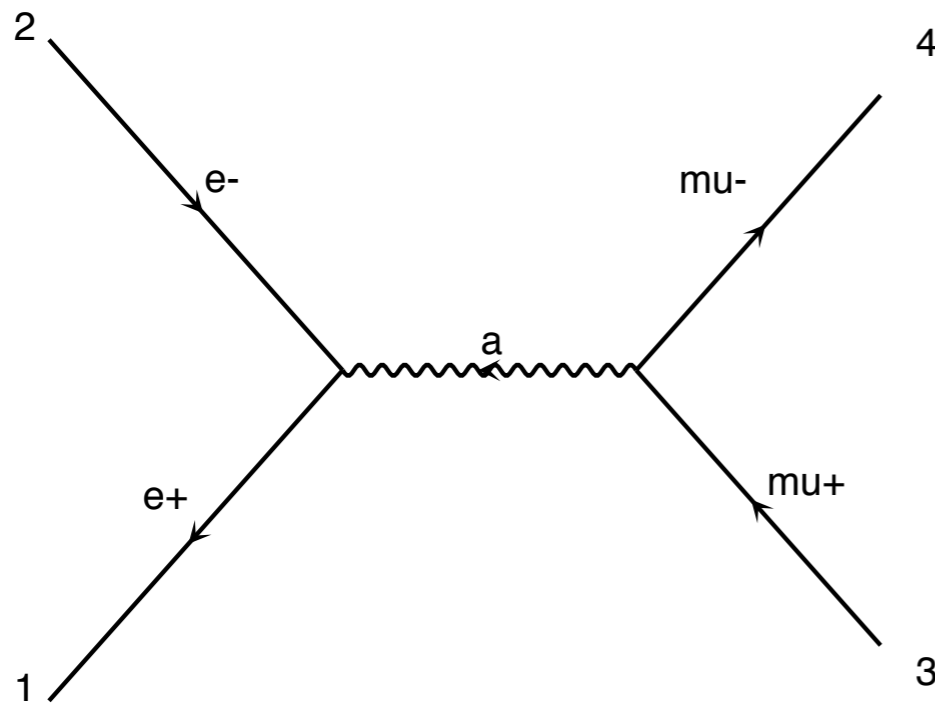
$$\frac{1}{4} \sum_{pol} |\mathcal{M}|^2 = \frac{1}{4} \sum_{pol} \mathcal{M}^* \mathcal{M}$$

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$$\rightarrow \frac{8e^4}{q^4} [(p_1 \cdot p_3)(p_2 \cdot p_4) + (p_1 \cdot p_4)(p_2 \cdot p_3)]$$

Very Efficient !!!



$$\mathcal{M} = e^2 (\bar{u} \gamma^\mu v) \frac{g_{\mu\nu}}{q^2} (\bar{v} \gamma^\nu u)$$

$$\frac{1}{4} \sum_{pol} |\mathcal{M}|^2 = \frac{1}{4} \sum_{pol} \mathcal{M}^* \mathcal{M}$$

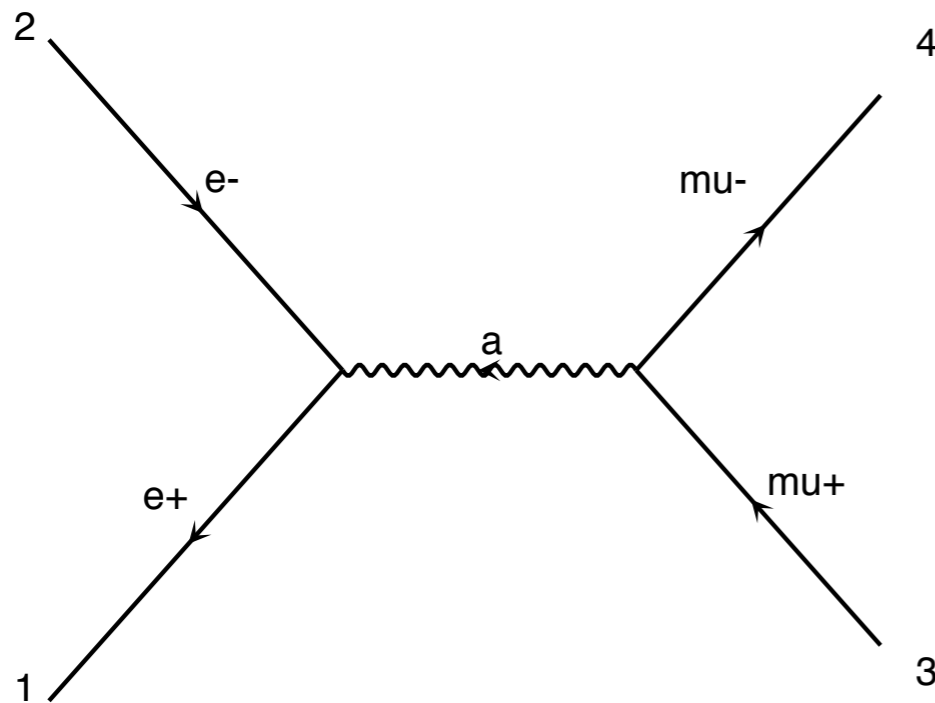
$$\sum_{pol} \bar{u} u = \not{p} + m$$

$$\rightarrow \frac{e^4}{4q^4} Tr[\not{p}_1 \gamma^\mu \not{p}_2 \gamma^\nu] Tr[\not{p}_3 \gamma_\mu \not{p}_4 \gamma_\nu]$$

$$\rightarrow \frac{8e^4}{q^4} [(p_1 \cdot p_3)(p_2 \cdot p_4) + (p_1 \cdot p_4)(p_2 \cdot p_3)]$$

Very Efficient !!!

Only for $2 \rightarrow 2$ and $2 \rightarrow 3$



$$\mathcal{M} = e^2 (\bar{u} \gamma^\mu v) \frac{g_{\mu\nu}}{q^2} (\bar{v} \gamma^\nu u)$$

$$\frac{1}{4} \sum_{pol} |\mathcal{M}|^2 = \frac{1}{4} \sum_{pol} \mathcal{M}^* \mathcal{M}$$

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$$\rightarrow \frac{8e^4}{q^4} [(p_1 \cdot p_3)(p_2 \cdot p_4) + (p_1 \cdot p_4)(p_2 \cdot p_3)]$$

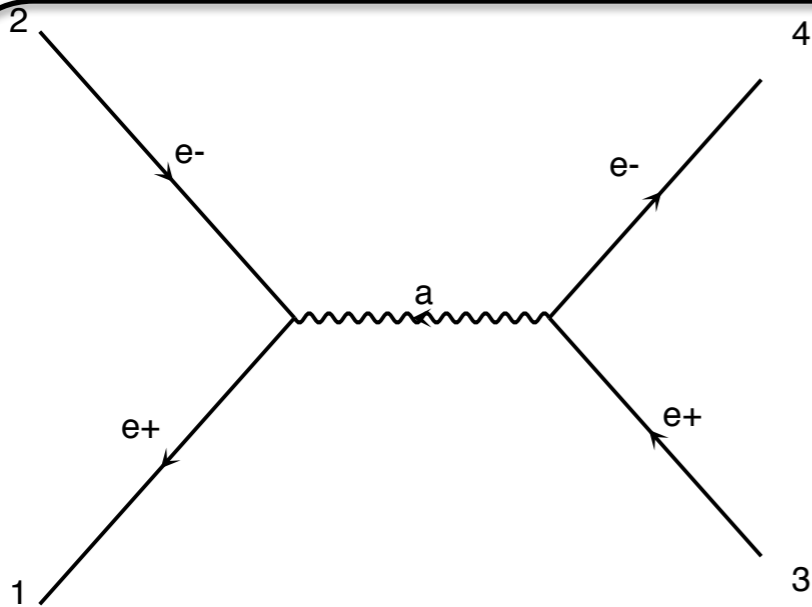
Very Efficient !!!

Only for $2 \rightarrow 2$ and $2 \rightarrow 3$

Because the number of terms rises as N^2

Idea

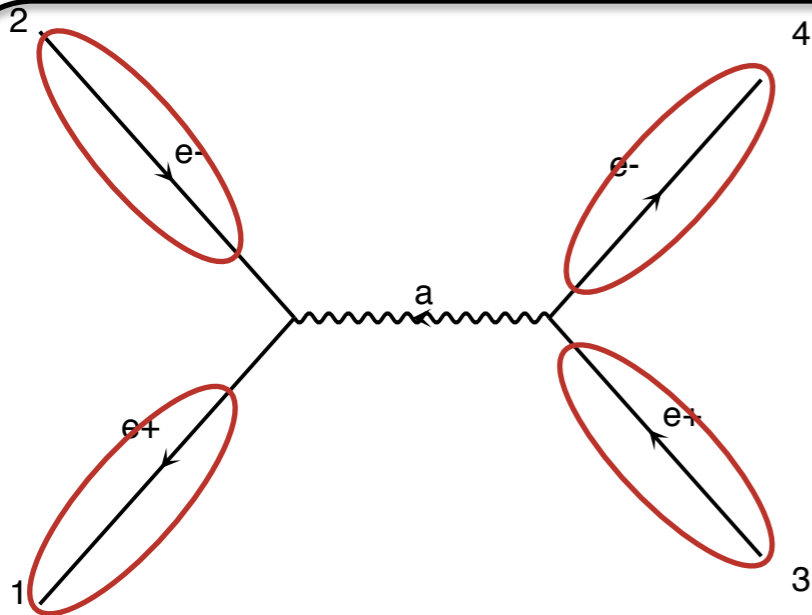
- Evaluate \mathcal{M} for fixed helicity of external particles
 - ➔ Multiply \mathcal{M} with \mathcal{M}^* $\rightarrow |\mathcal{M}|^2$
 - ➔ Loop on Helicity and average the results



$$\mathcal{M} = (\bar{u} e \gamma^\mu v) \frac{g_{\mu\nu}}{q^2} (\bar{v} e \gamma^\nu u)$$

Idea

- Evaluate \mathcal{M} for fixed helicity of external particles
 - ➔ Multiply \mathcal{M} with \mathcal{M}^* $\rightarrow |\mathcal{M}|^2$
 - ➔ Loop on Helicity and average the results

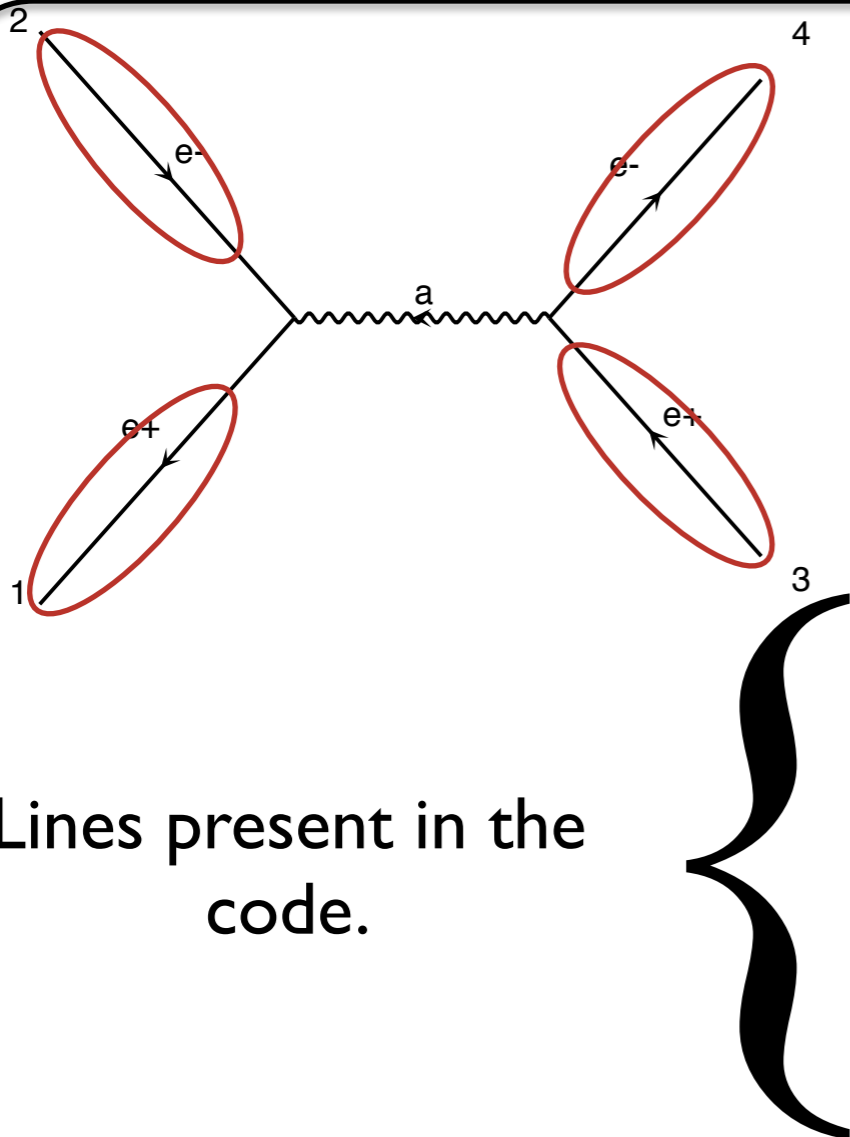


$$\mathcal{M} = (\bar{u} e \gamma^\mu v) \frac{g_{\mu\nu}}{q^2} (\bar{v} e \gamma^\nu u)$$

Numbers for given helicity and momenta

Idea

- Evaluate \mathcal{M} for fixed helicity of external particles
 - ➔ Multiply \mathcal{M} with \mathcal{M}^* -> $|\mathcal{M}|^2$
 - ➔ Loop on Helicity and average the results



Lines present in the code.

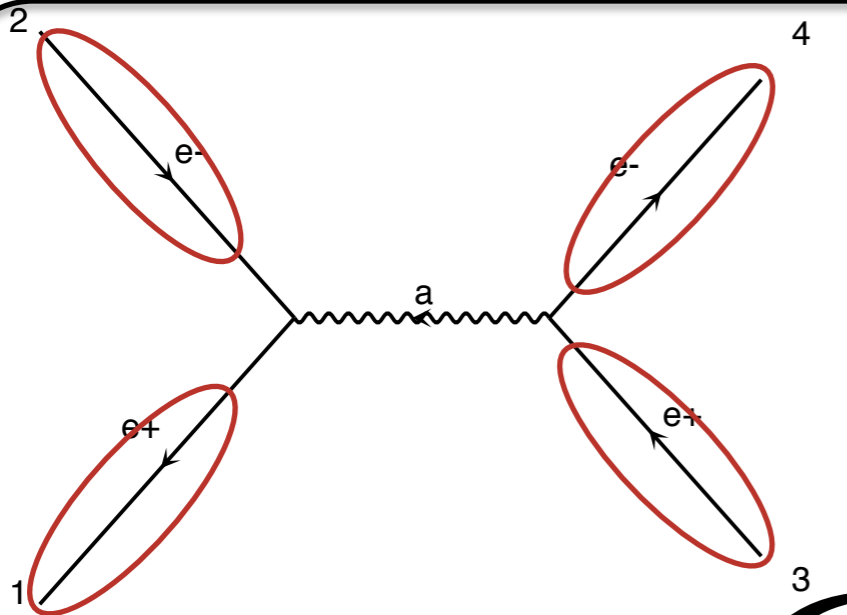
$$\left. \begin{aligned} \bar{v}_1 &= fct(\vec{p}_1, m_1) \\ u_2 &= fct(\vec{p}_2, m_2) \\ v_3 &= fct(\vec{p}_3, m_3) \\ \bar{u}_4 &= fct(\vec{p}_4, m_4) \end{aligned} \right\}$$

$$\mathcal{M} = (\bar{u} e \gamma^\mu v) \frac{g_{\mu\nu}}{q^2} (\bar{v} e \gamma^\nu u)$$

Numbers for given helicity and momenta

Idea

- Evaluate \mathcal{M} for fixed helicity of external particles
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Numbers for given helicity and momenta

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$$\left. \begin{aligned} \bar{v}_1 &= fct(\vec{p}_1, m_1) \\ u_2 &= fct(\vec{p}_2, m_2) \\ v_3 &= fct(\vec{p}_3, m_3) \\ \bar{u}_4 &= fct(\vec{p}_4, m_4) \end{aligned} \right\}$$

$$u(p) = \begin{pmatrix} \omega_{-\lambda}(p) \chi_\lambda(\vec{p}) \\ \omega_\lambda(p) \chi_\lambda(\vec{p}) \end{pmatrix}$$

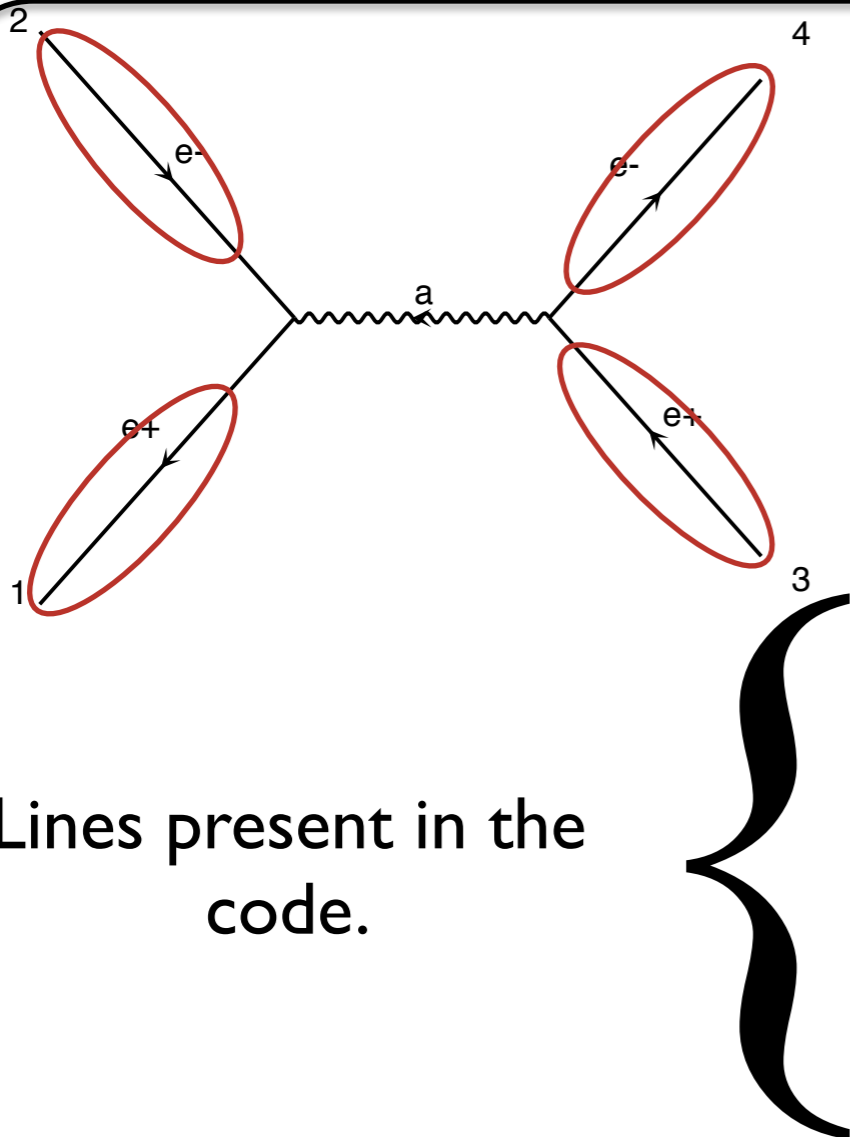
$$\omega_\pm(p) \equiv \sqrt{E \pm |\vec{p}|}$$

$$\chi_+(\vec{p}) = \frac{1}{\sqrt{2|\vec{p}|(|\vec{p}| + p_z)}} \begin{pmatrix} |\vec{p}| + p_z \\ p_x + ip_y \end{pmatrix},$$

$$\chi_-(\vec{p}) = \frac{1}{\sqrt{2|\vec{p}|(|\vec{p}| + p_z)}} \begin{pmatrix} -p_x + ip_y \\ |\vec{p}| + p_z \end{pmatrix}.$$

Idea

- Evaluate \mathcal{M} for fixed helicity of external particles
 - ➔ Multiply \mathcal{M} with \mathcal{M}^* -> $|\mathcal{M}|^2$
 - ➔ Loop on Helicity and average the results



Lines present in the code.

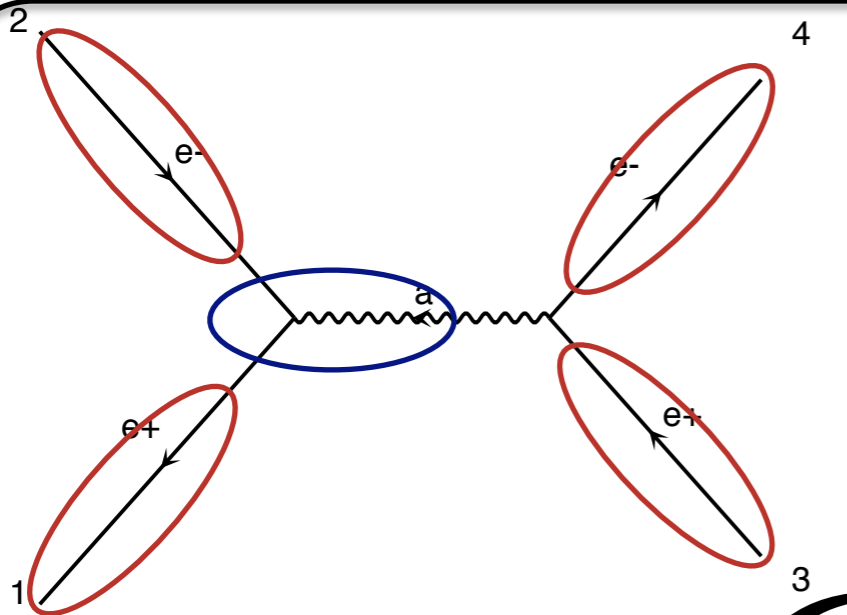
$$\left. \begin{aligned} \bar{v}_1 &= fct(\vec{p}_1, m_1) \\ u_2 &= fct(\vec{p}_2, m_2) \\ v_3 &= fct(\vec{p}_3, m_3) \\ \bar{u}_4 &= fct(\vec{p}_4, m_4) \end{aligned} \right\}$$

$$\mathcal{M} = (\bar{u} e \gamma^\mu v) \frac{g_{\mu\nu}}{q^2} (\bar{v} e \gamma^\nu u)$$

Numbers for given helicity and momenta

Idea

- Evaluate \mathcal{M} for fixed helicity of external particles
 - ➔ Multiply \mathcal{M} with \mathcal{M}^* -> $|\mathcal{M}|^2$
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$$\mathcal{M} = (\bar{u} e \gamma^\mu v) \frac{g_{\mu\nu}}{q^2} (\bar{v} e \gamma^\nu u)$$

Numbers for given helicity and momenta

Calculate propagator wavefunctions

Lines present in the code.

$$\bar{v}_1 = fct(\vec{p}_1, m_1)$$

$$u_2 = fct(\vec{p}_2, m_2)$$

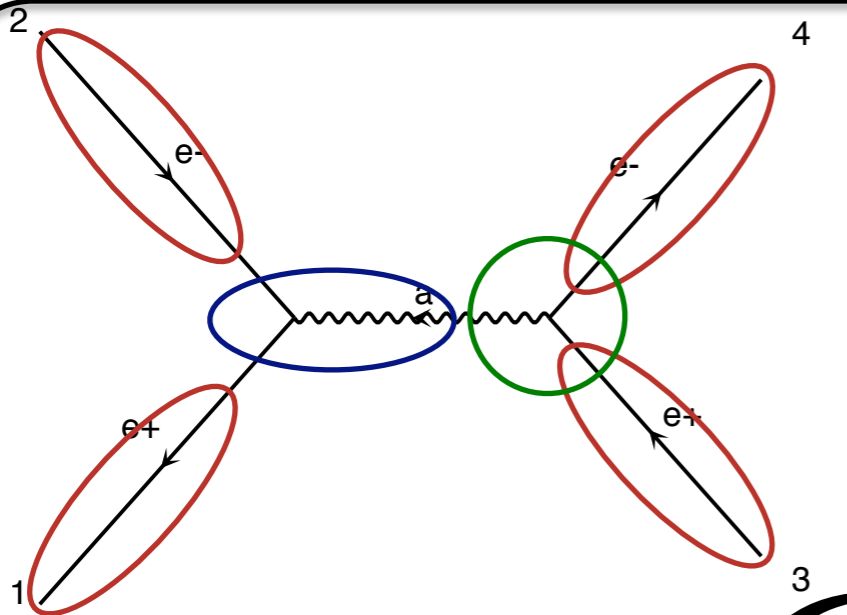
$$v_3 = fct(\vec{p}_3, m_3)$$

$$\bar{u}_4 = fct(\vec{p}_4, m_4)$$

$$W_a = fct(\bar{v}_1, u_2, m_a, \Gamma_a)$$

Idea

- Evaluate \mathcal{M} for fixed helicity of external particles
 - ➔ Multiply \mathcal{M} with \mathcal{M}^* $\rightarrow |\mathcal{M}|^2$
 - ➔ Loop on Helicity and average the results



$$\mathcal{M} = \bar{u} e \gamma^\mu v \frac{g_{\mu\nu}}{q^2} \bar{v} e \gamma^\nu u$$

Numbers for given helicity and momenta

Calculate propagator wavefunctions

Finally evaluate amplitude (c-number)

Lines present in the code.

$$\bar{v}_1 = fct(\vec{p}_1, m_1)$$

$$u_2 = fct(\vec{p}_2, m_2)$$

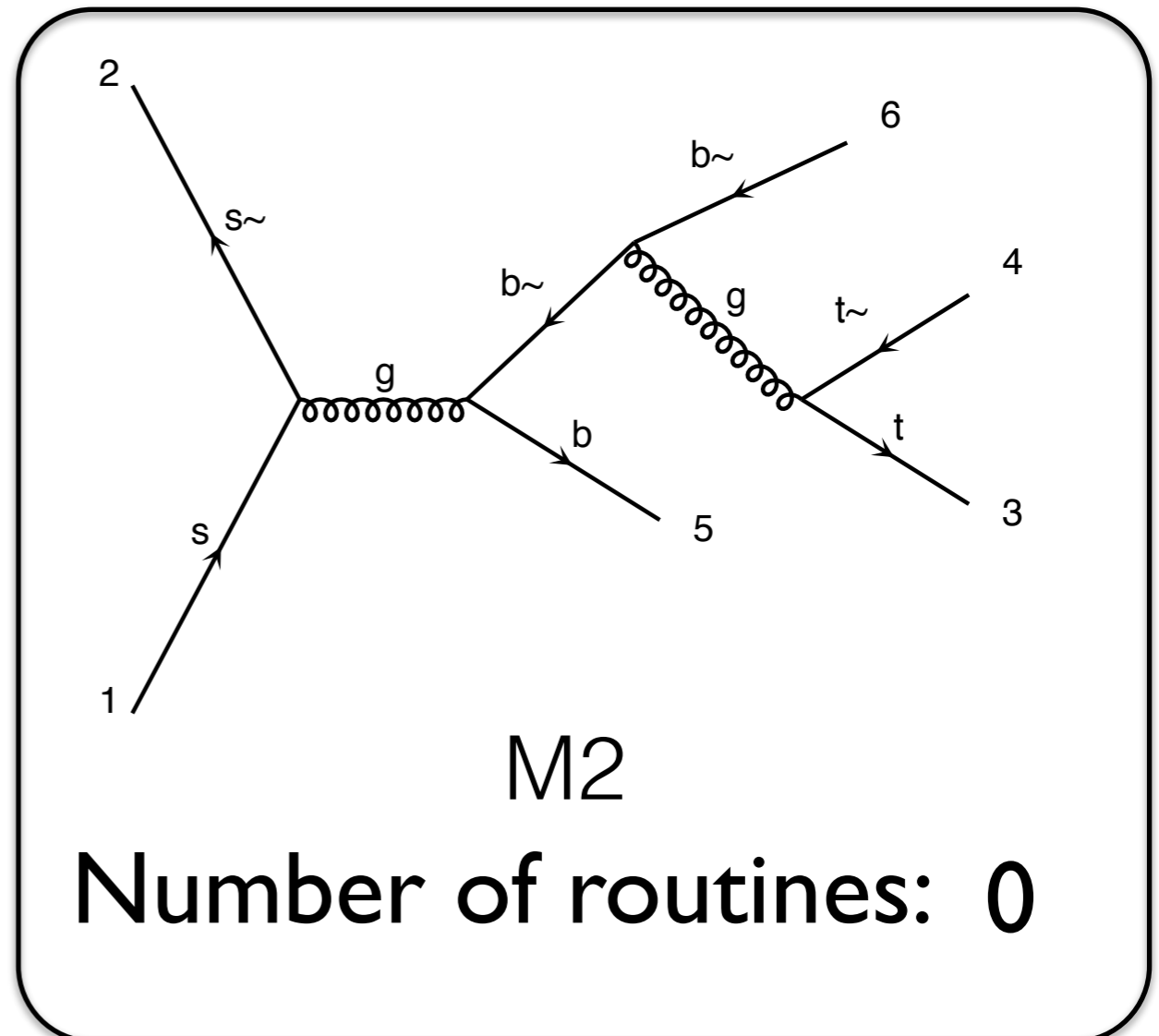
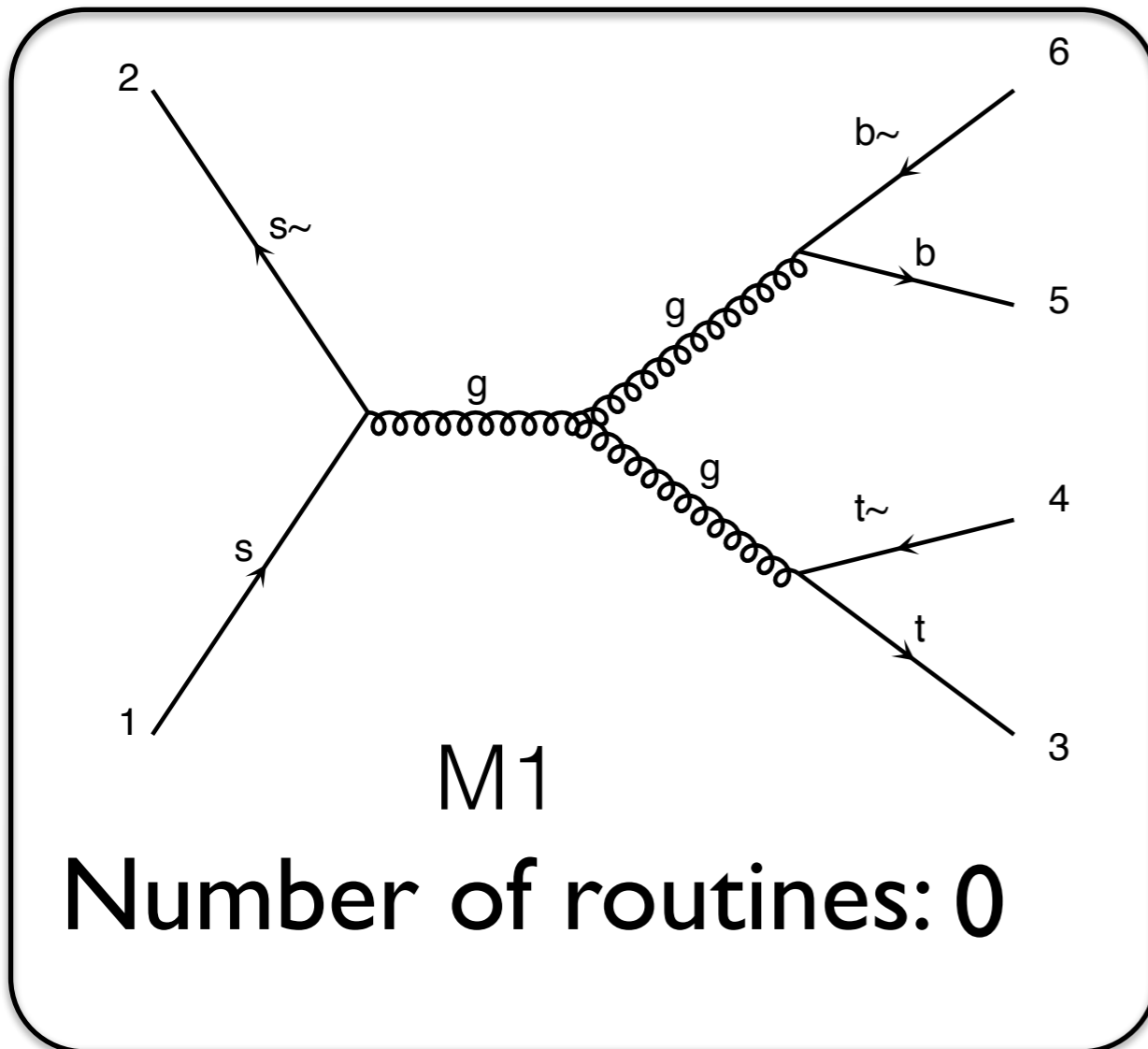
$$v_3 = fct(\vec{p}_3, m_3)$$

$$\bar{u}_4 = fct(\vec{p}_4, m_4)$$

$$W_a = fct(\bar{v}_1, u_2, m_a, \Gamma_a)$$

$$\mathcal{M} = fct(v_3, \bar{u}_4, W_a)$$

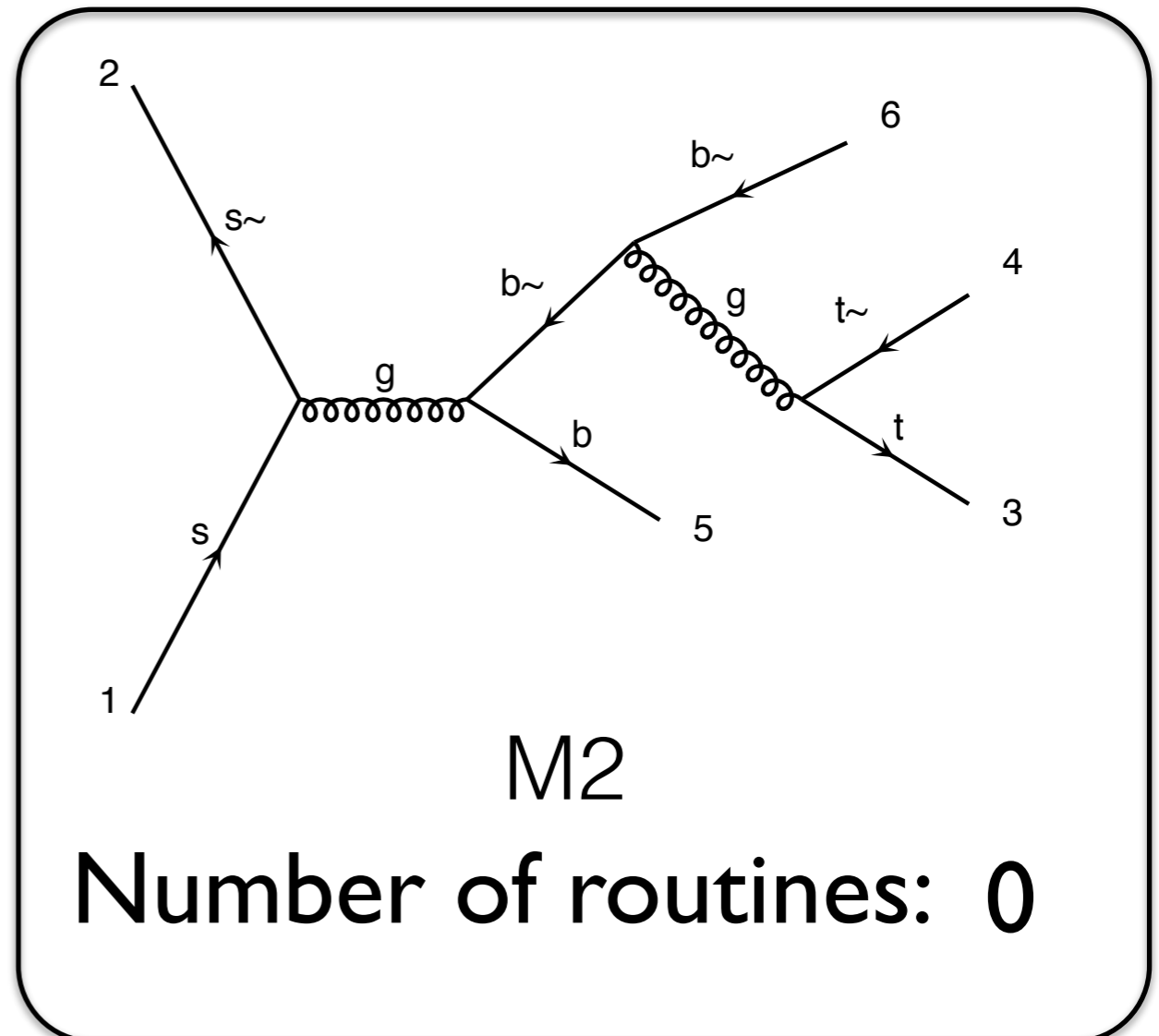
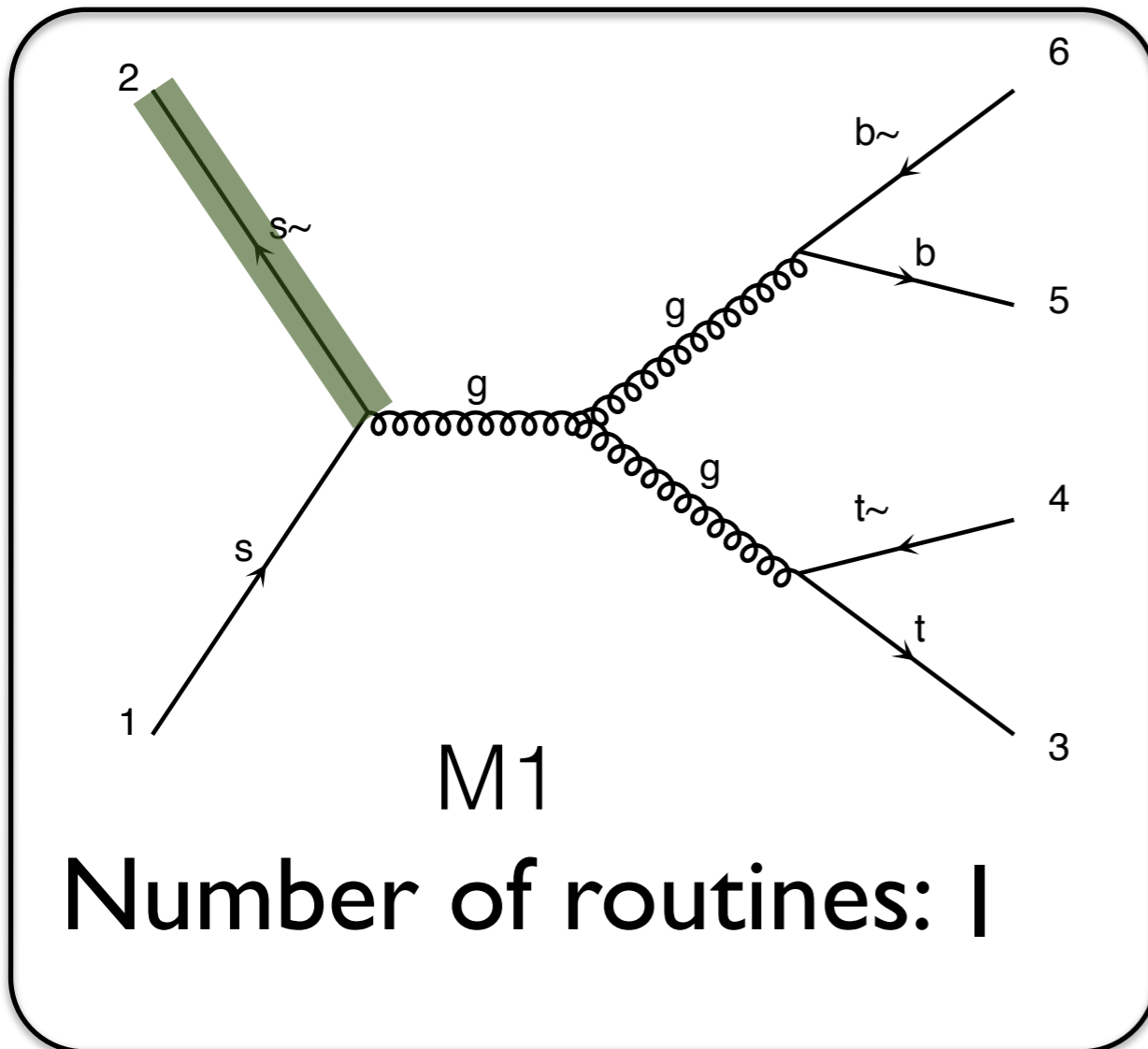
Known



Number of routines for both: 0

$$|M|^2 = |M_1 + M_2|^2$$

Known

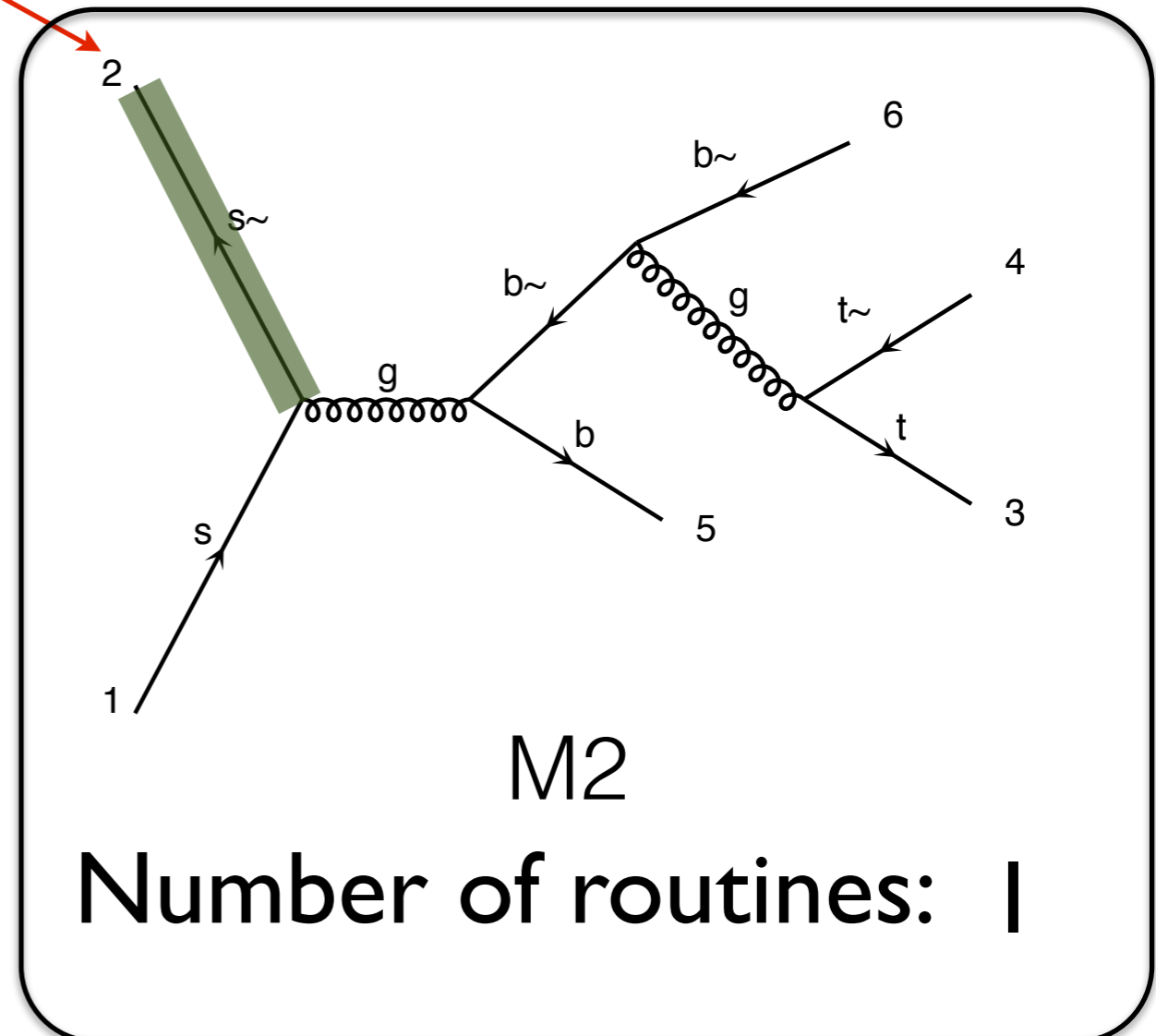
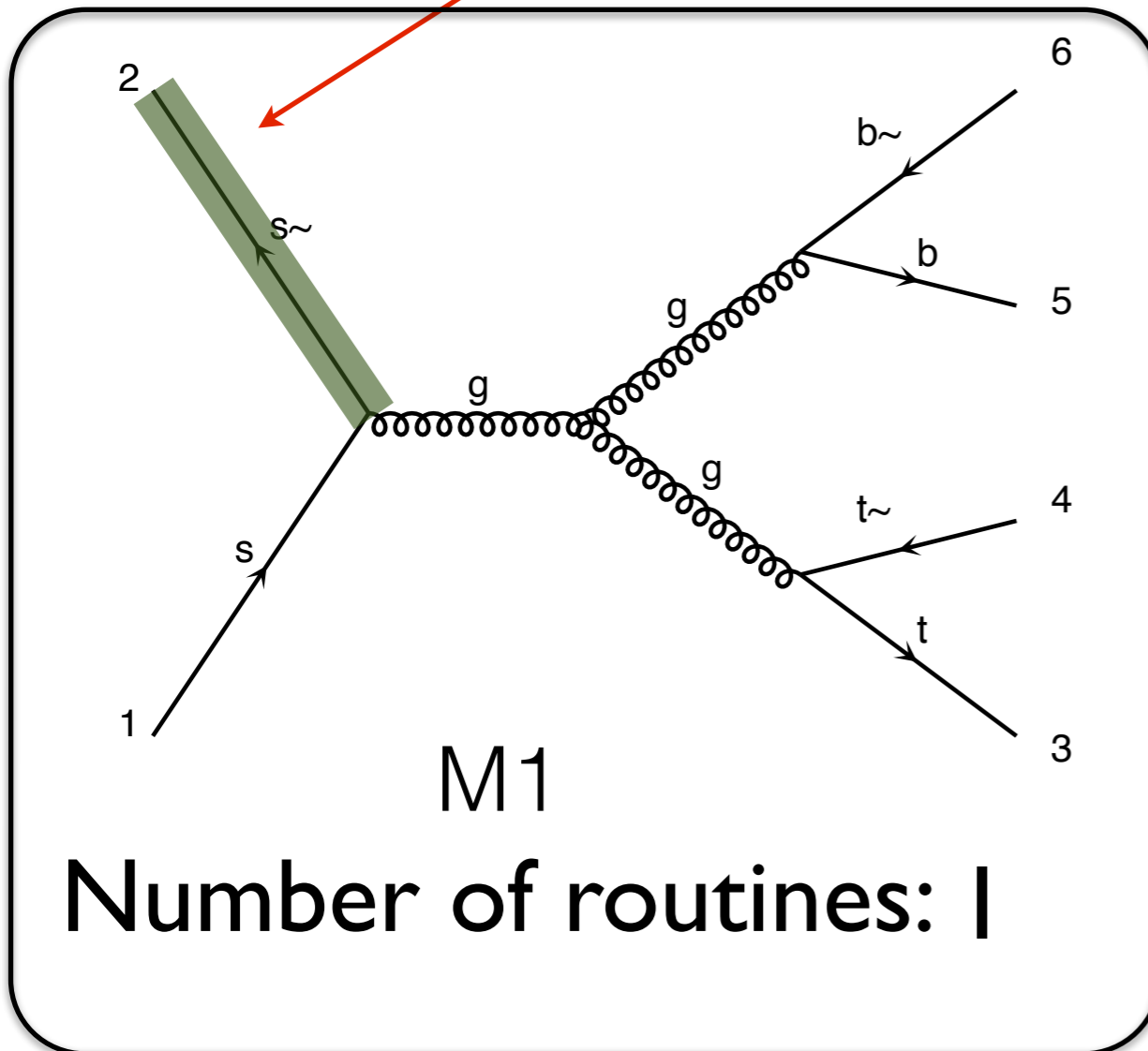


Number of routines for both: 1

$$|M|^2 = |M_1 + M_2|^2$$

Identical

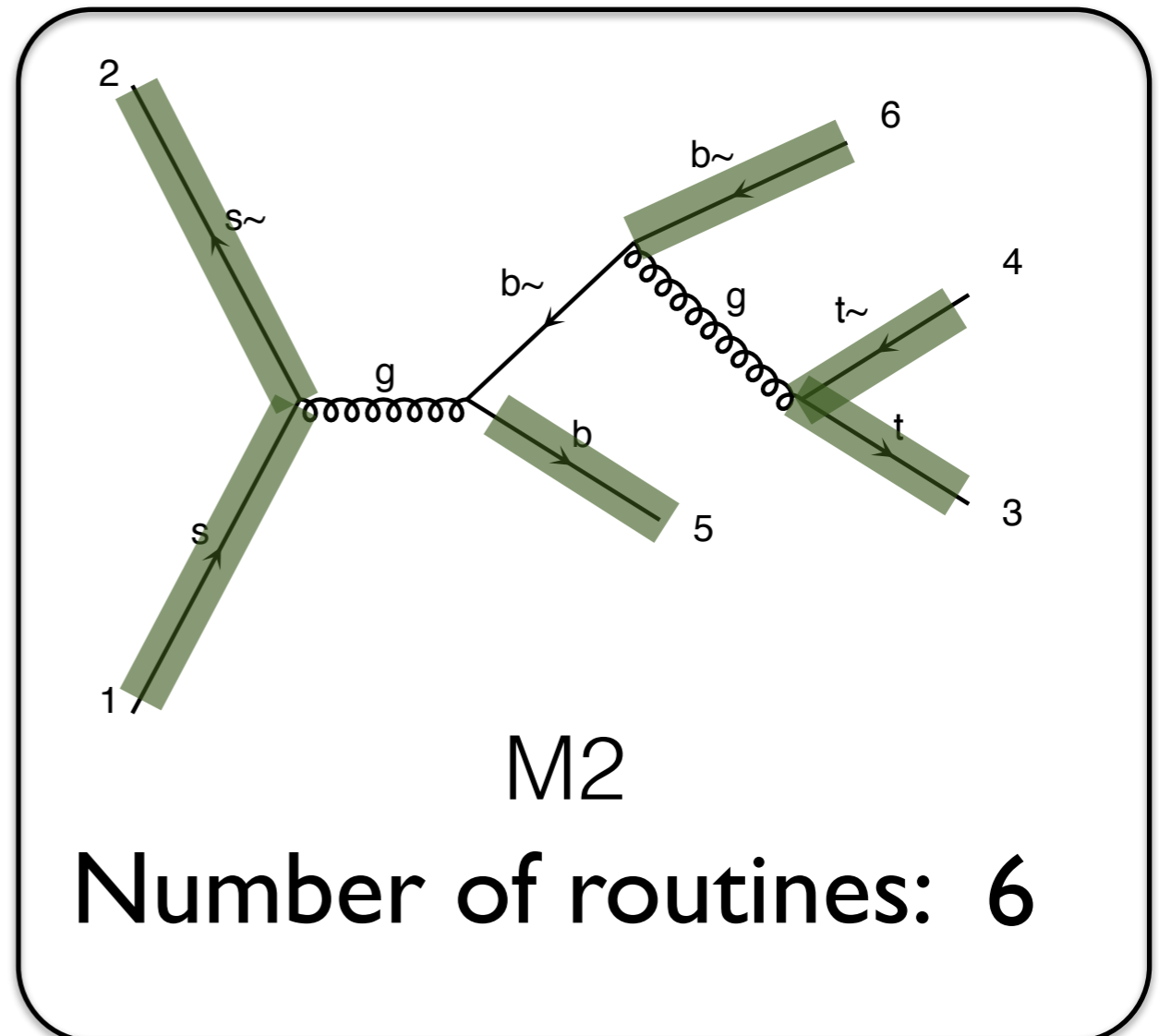
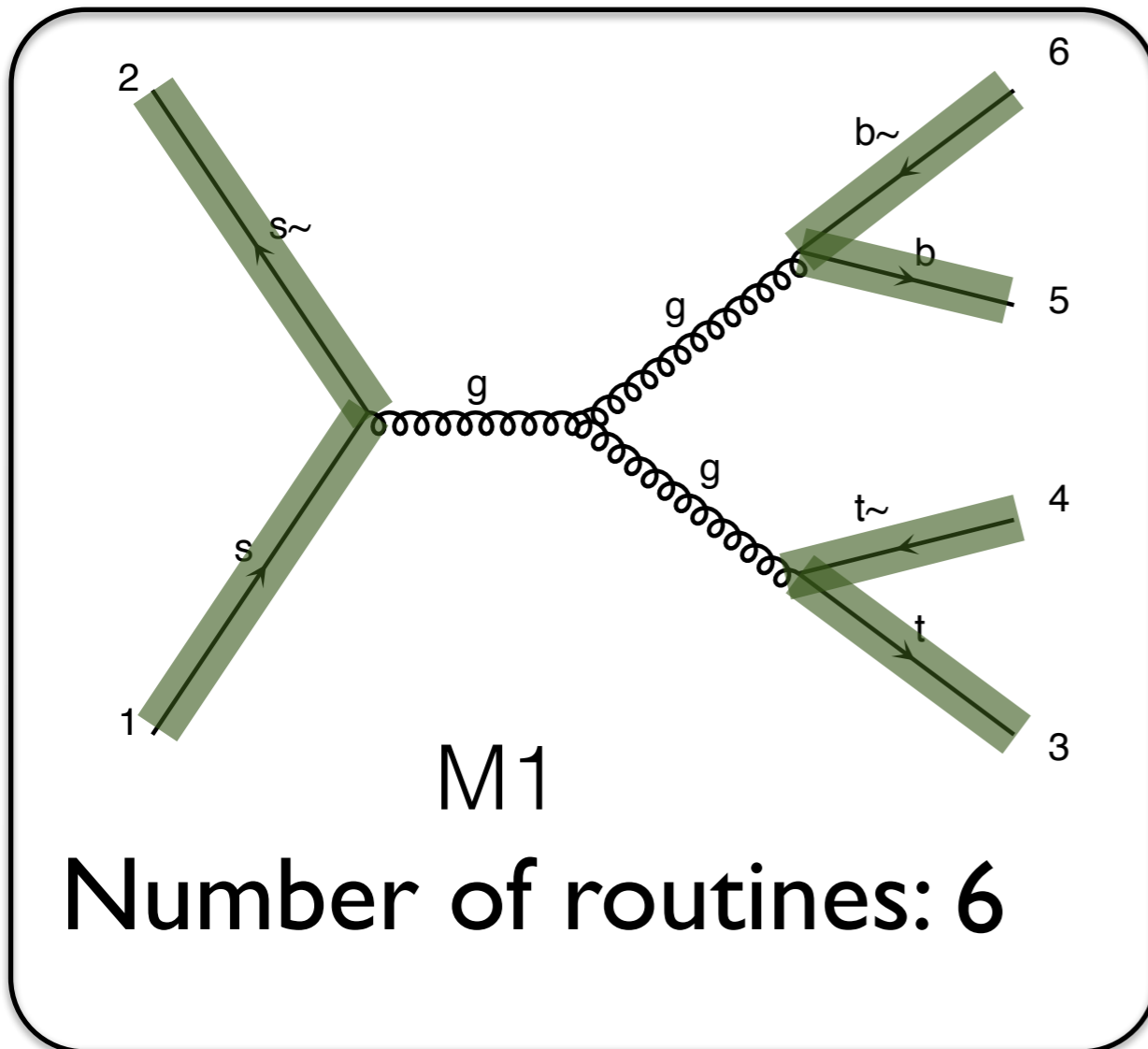
Known



Number of routines for both: 1

$$|M|^2 = |M_1 + M_2|^2$$

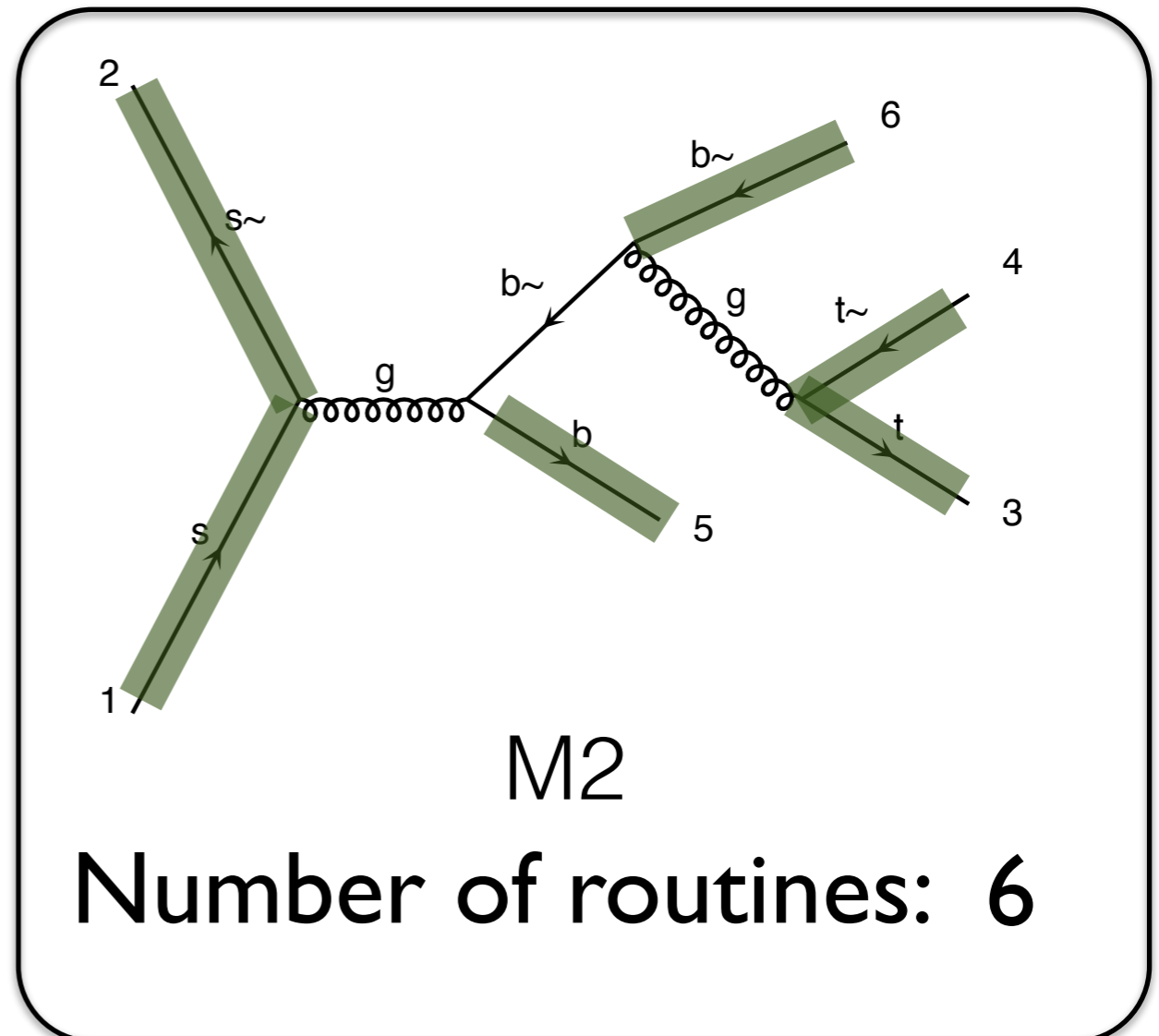
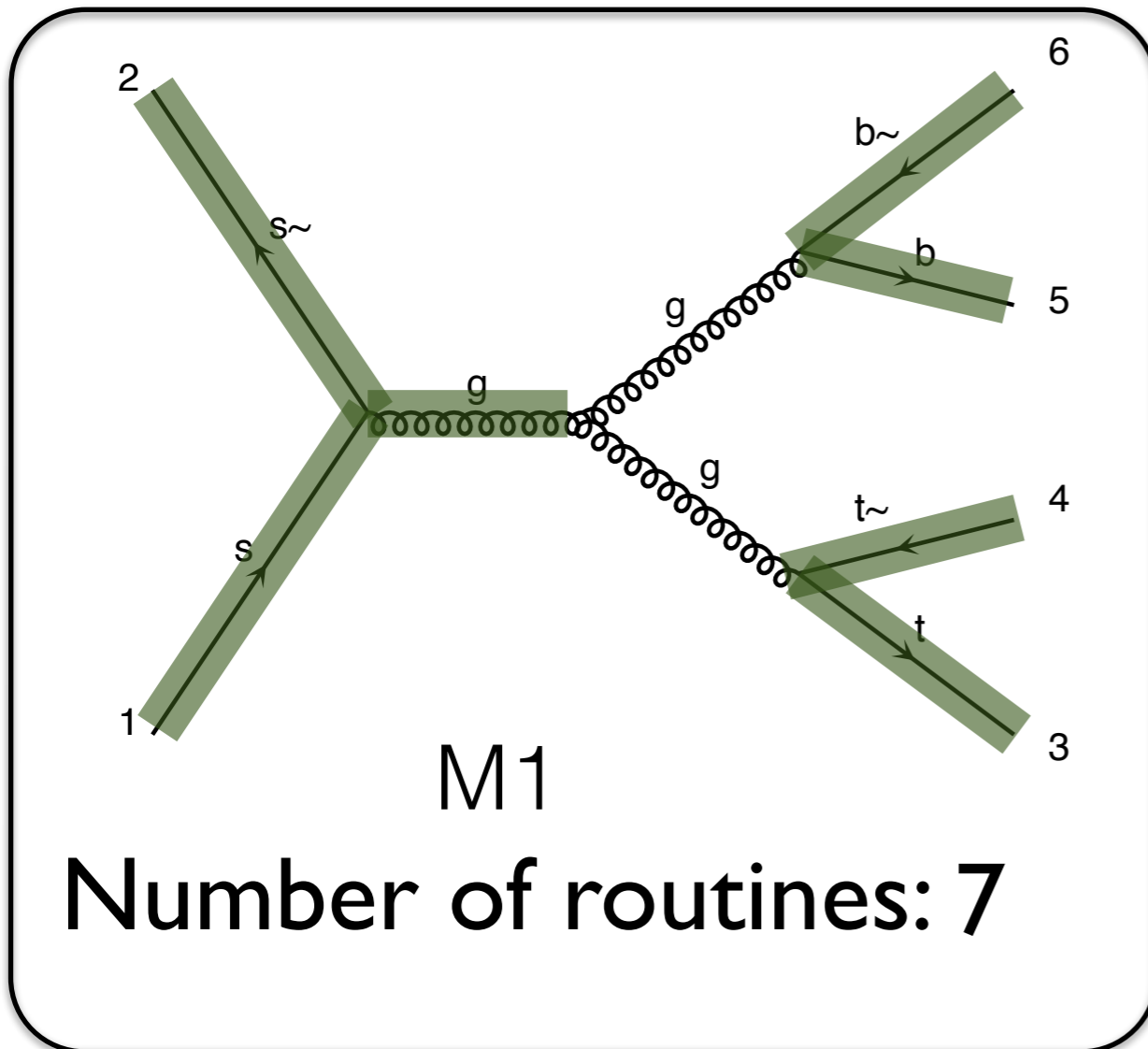
Known



Number of routines for both: 6

$$|M|^2 = |M_1 + M_2|^2$$

Known

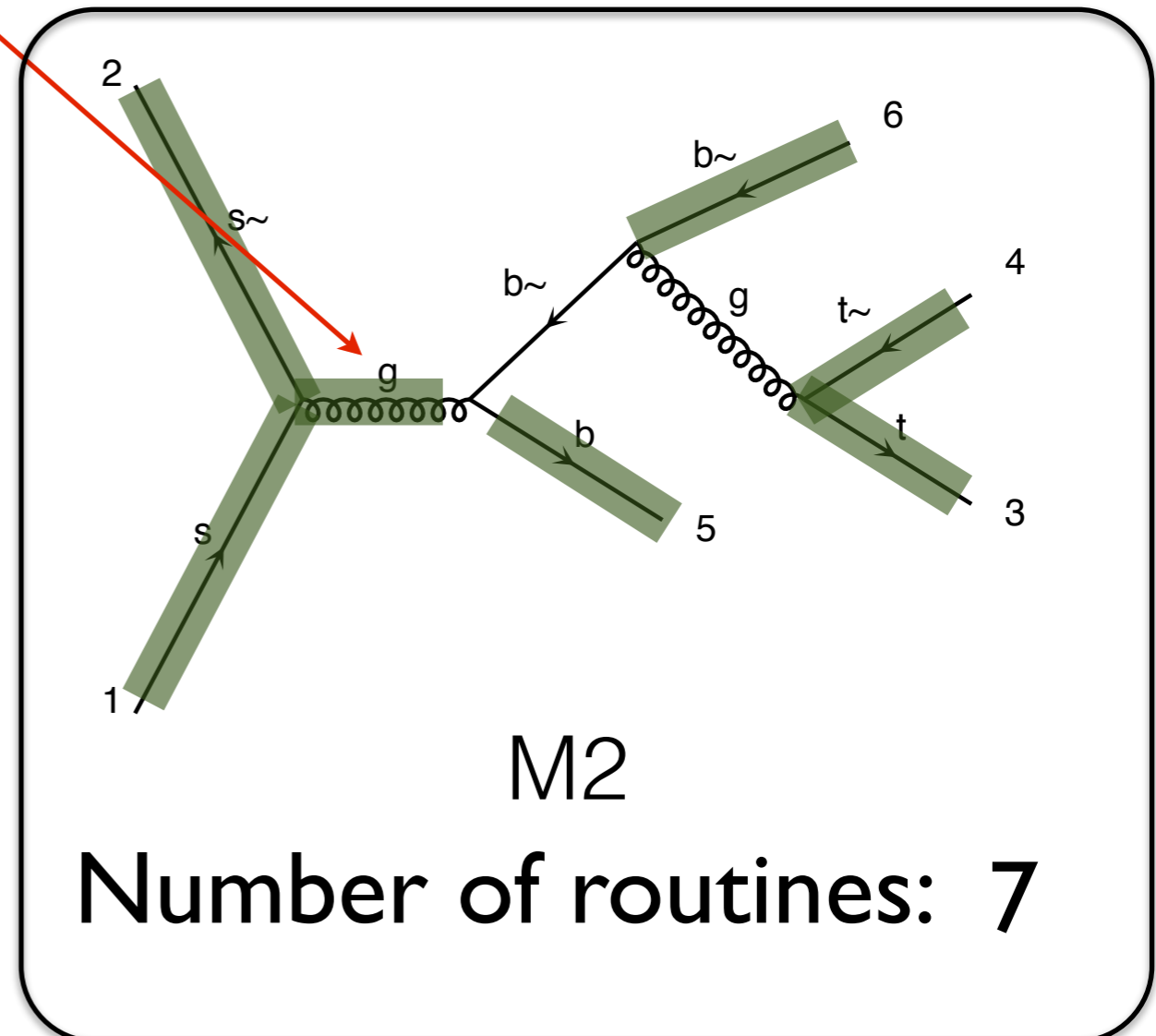
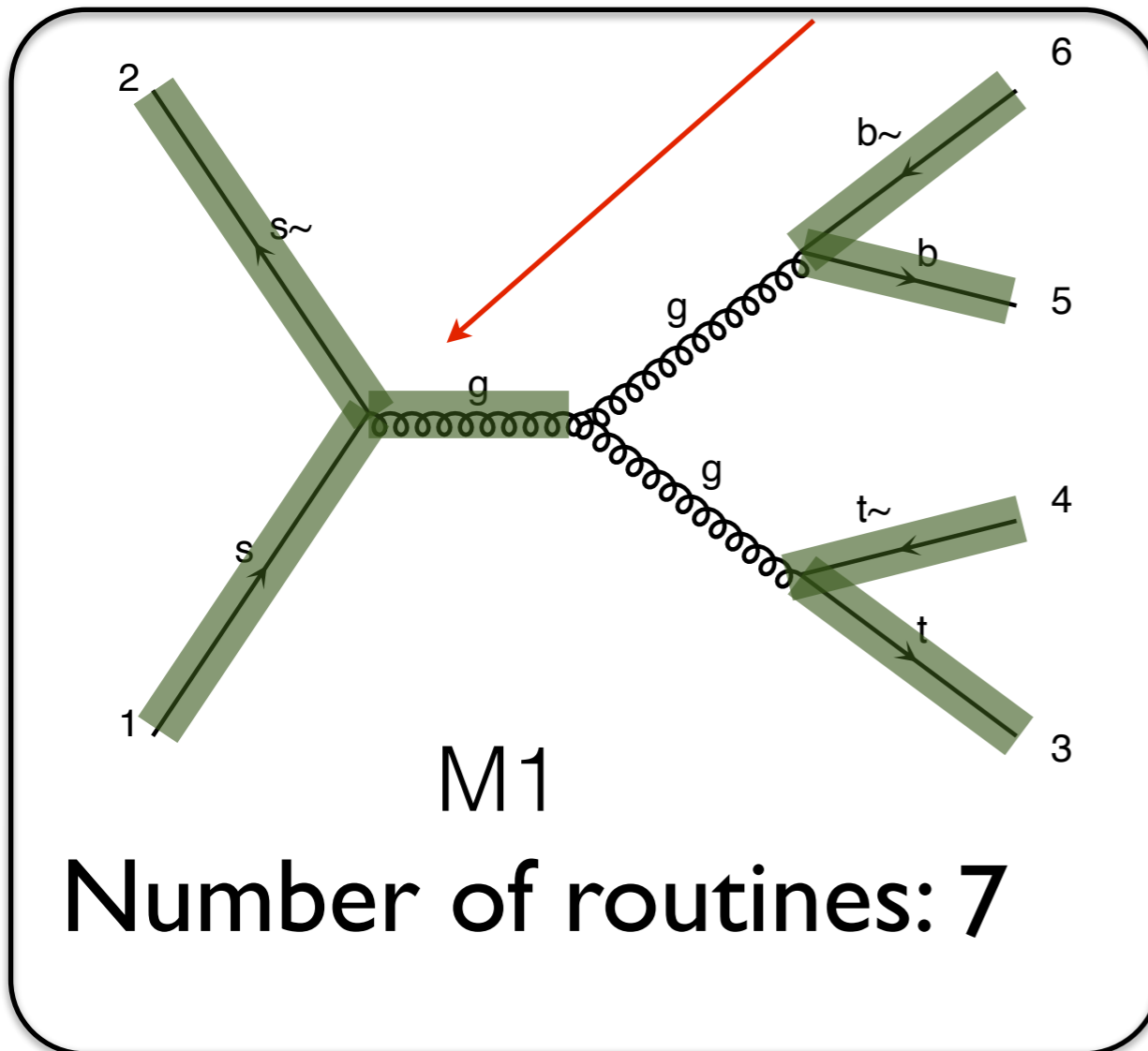


Number of routines for both: 7

$$|M|^2 = |M_1 + M_2|^2$$

Known

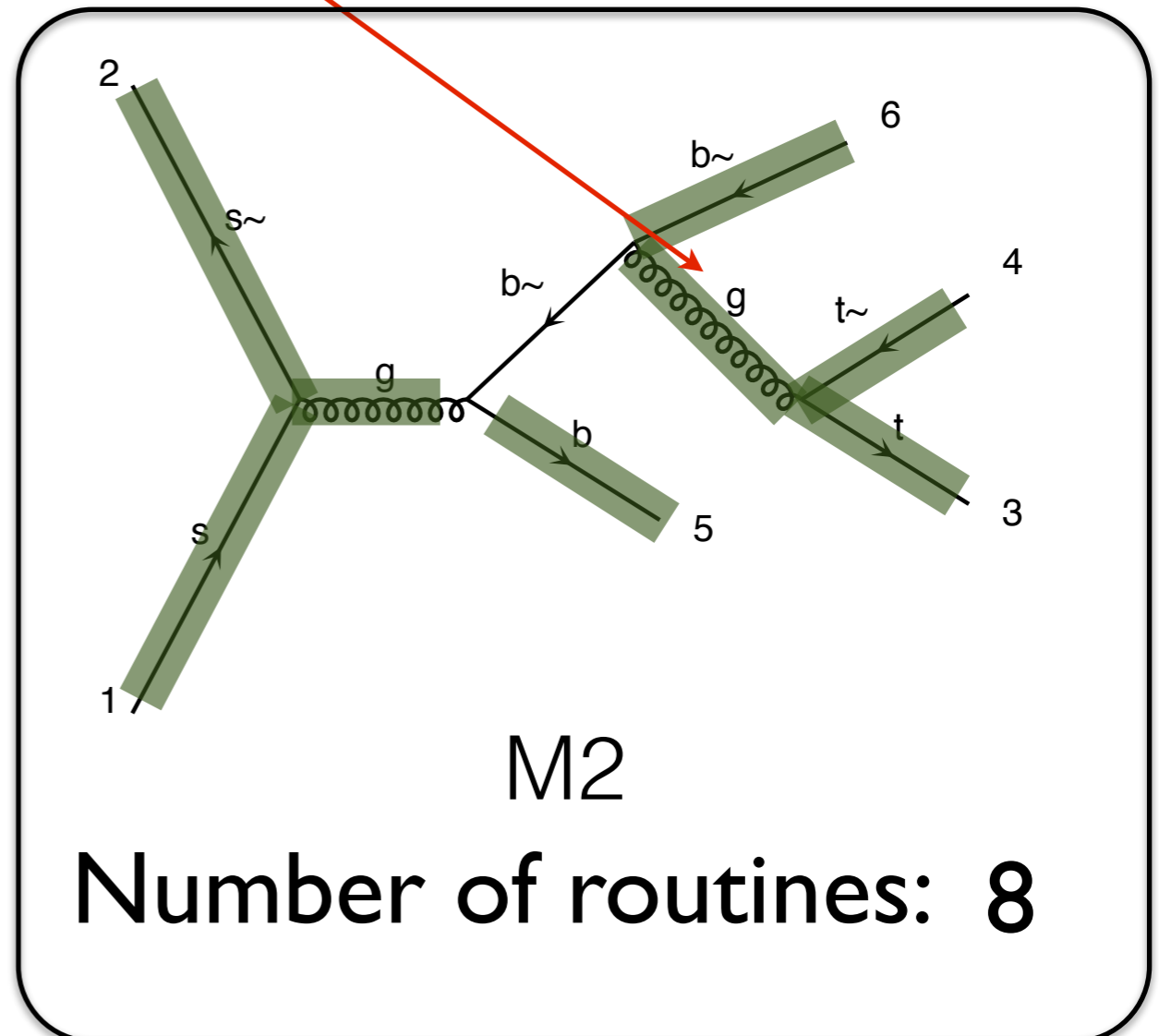
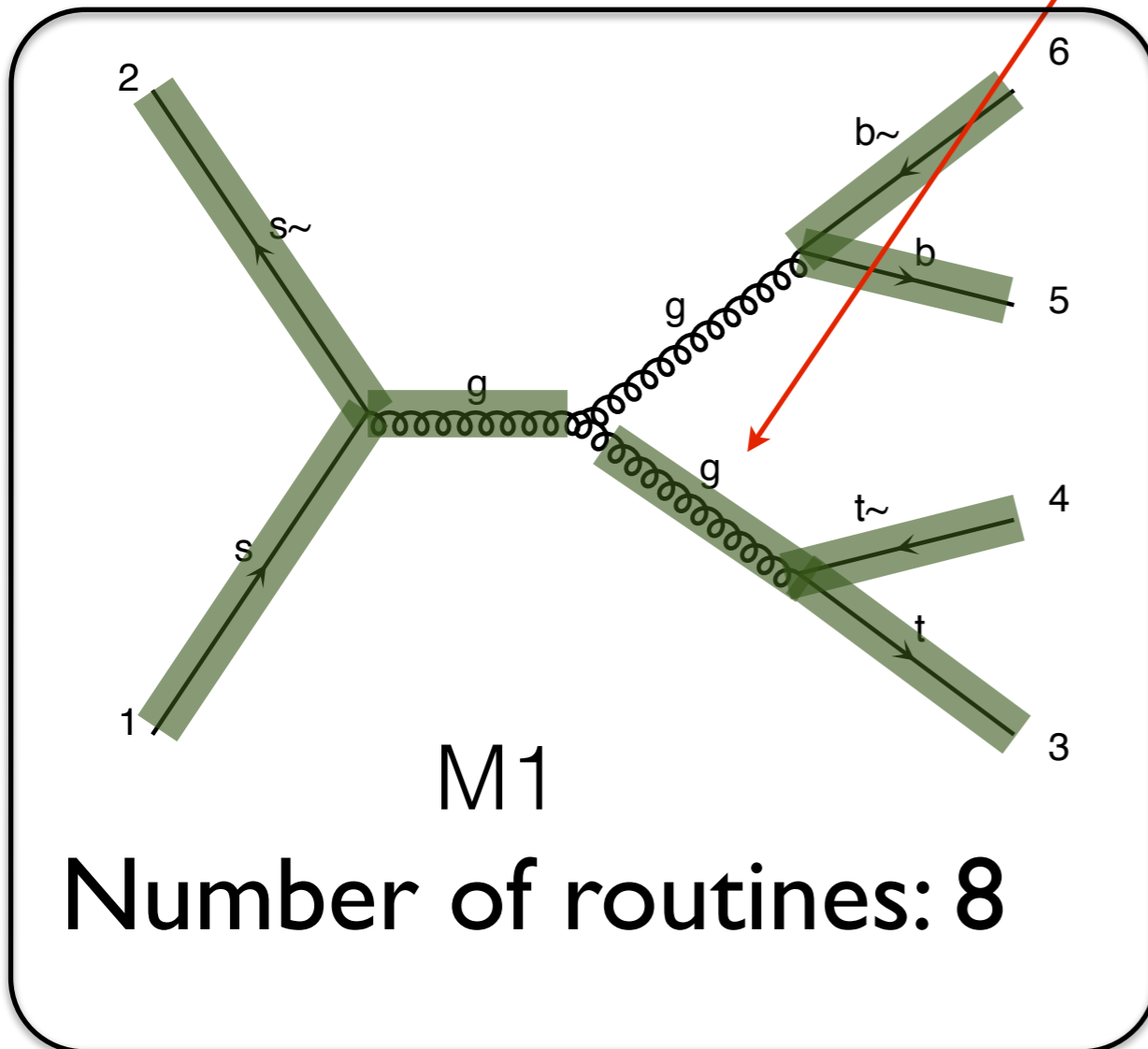
Identical



Number of routines for both: 7

$$|M|^2 = |M_1 + M_2|^2$$

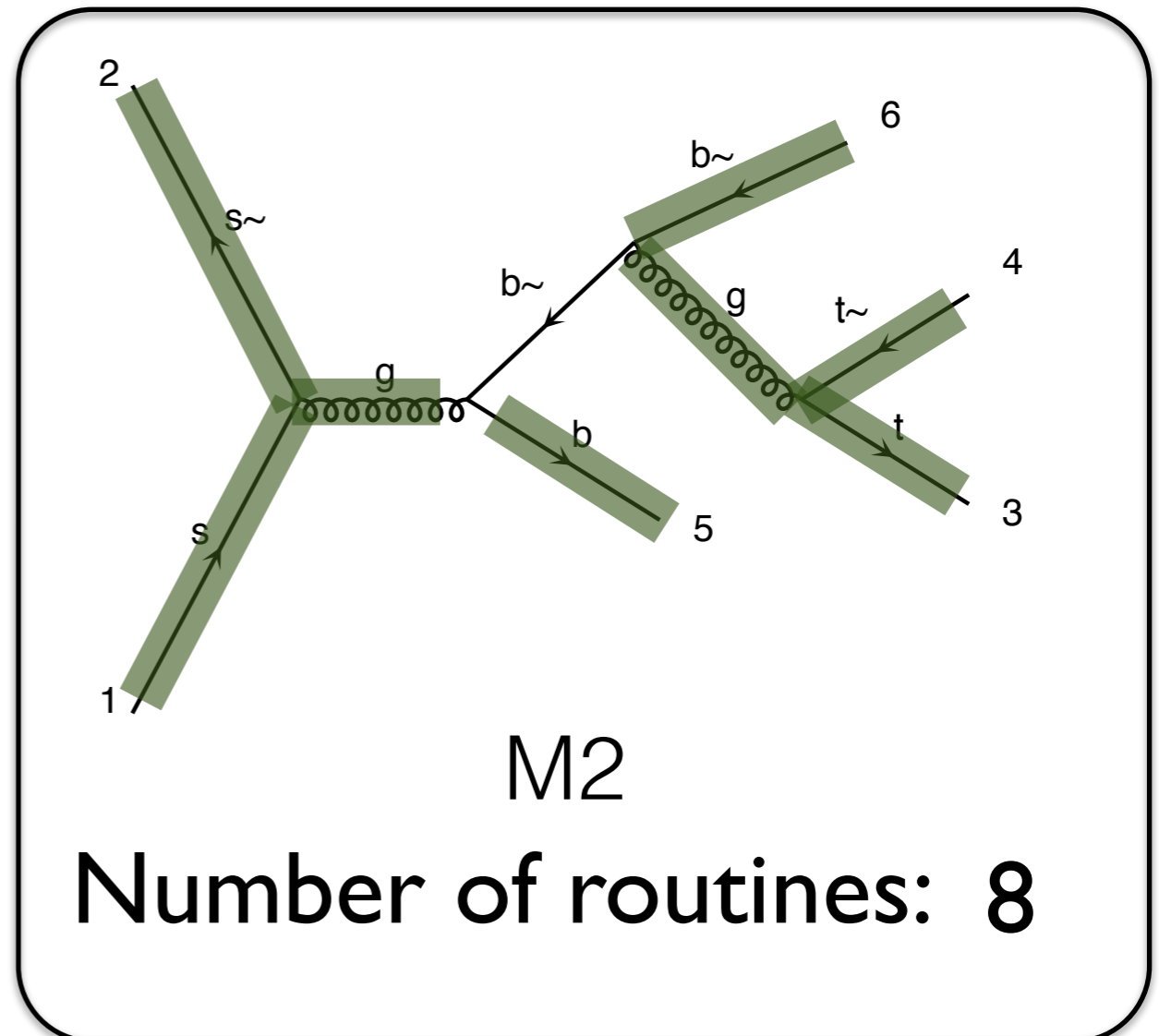
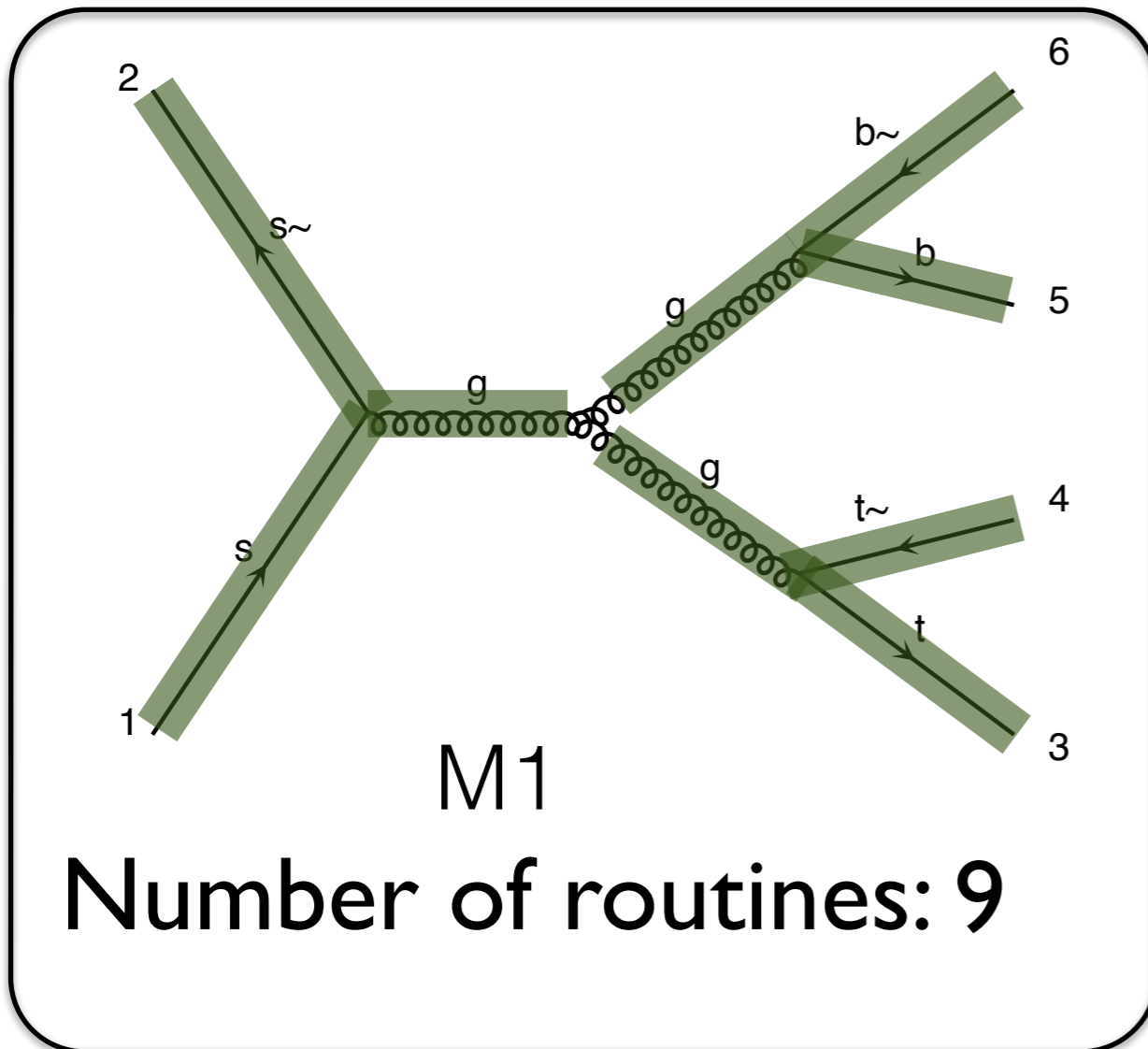
Identical Known



Number of routines for both: 8

$$|M|^2 = |M_1 + M_2|^2$$

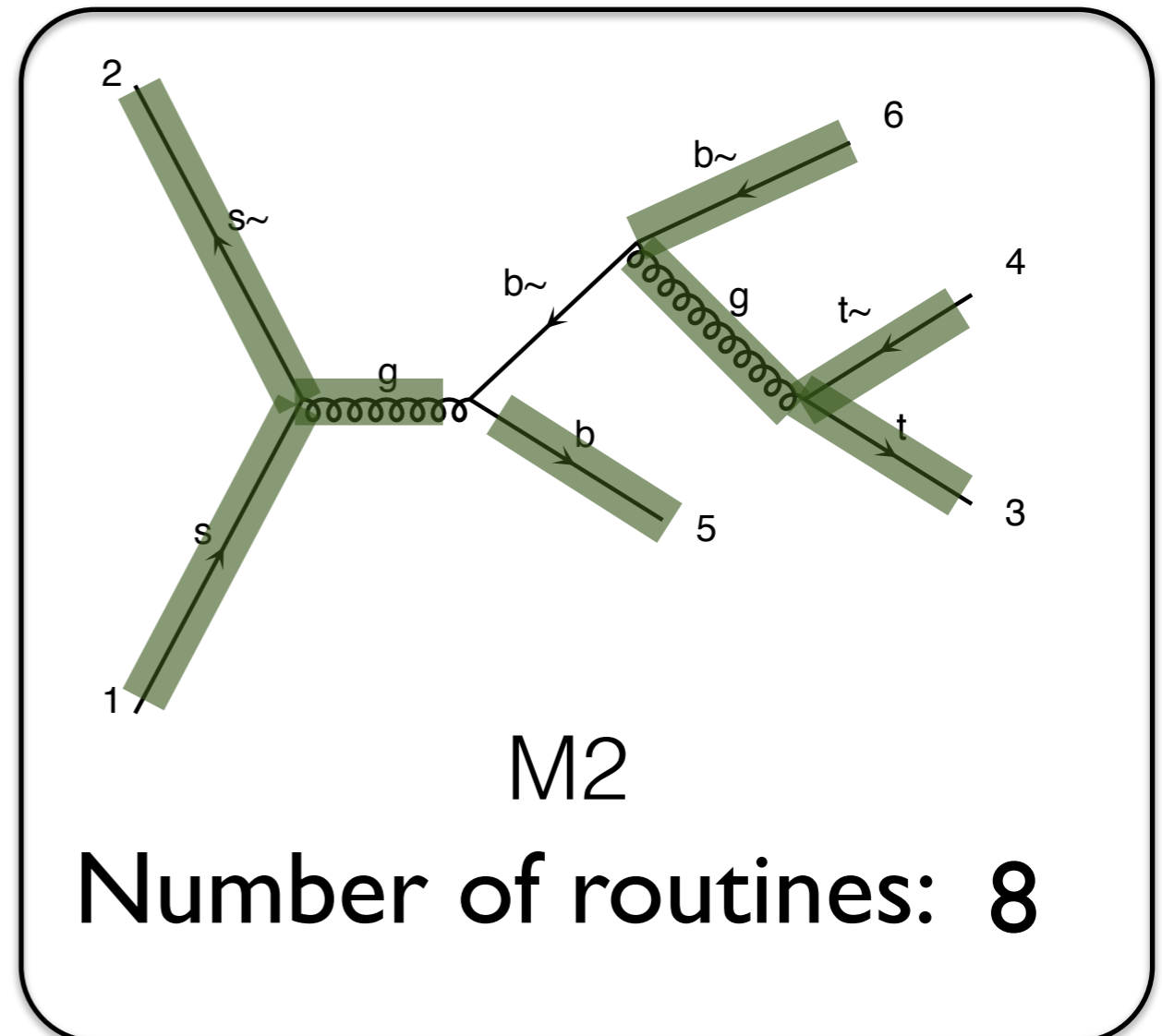
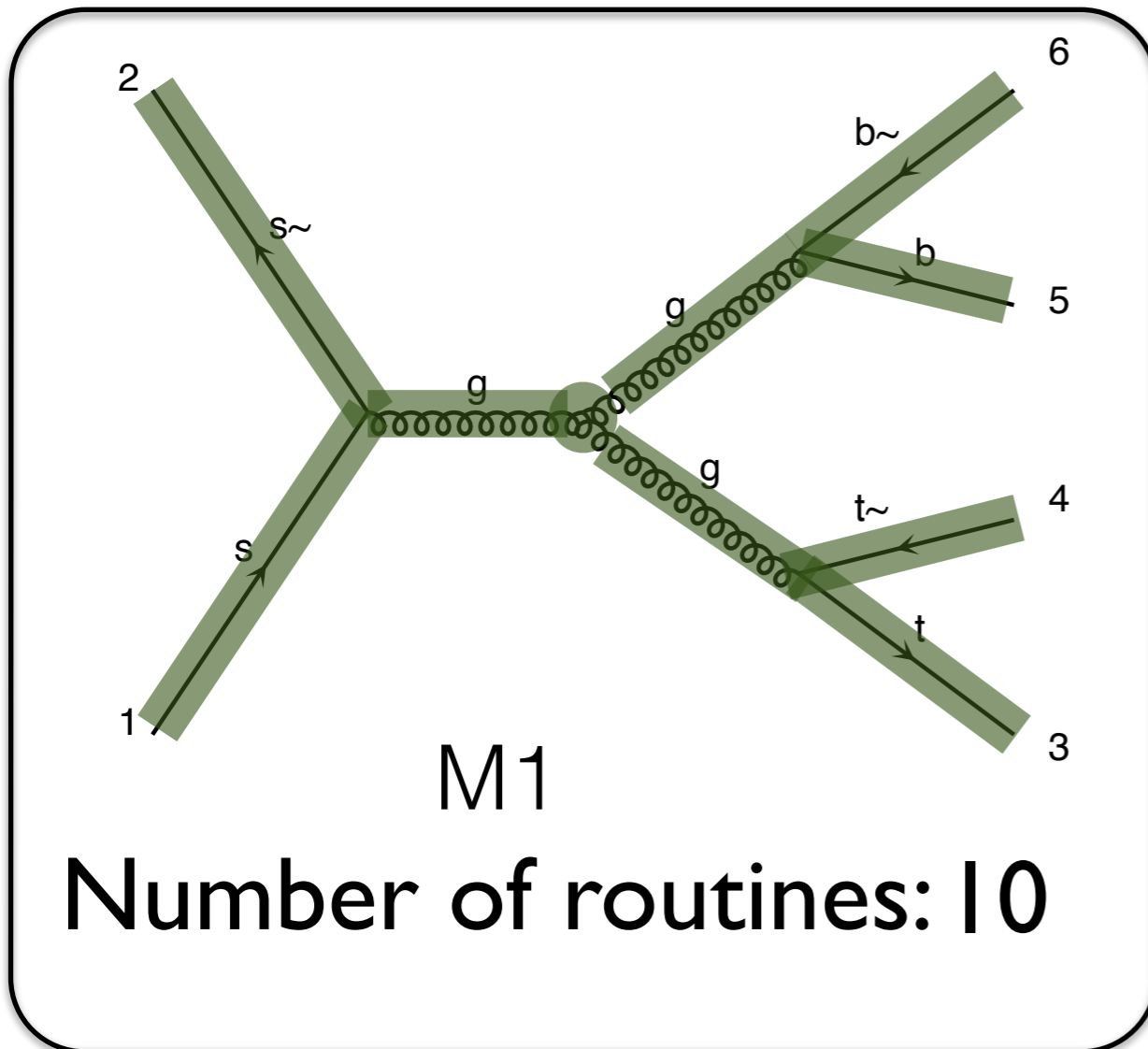
Known



Number of routines for both: 9

$$|M|^2 = |M_1 + M_2|^2$$

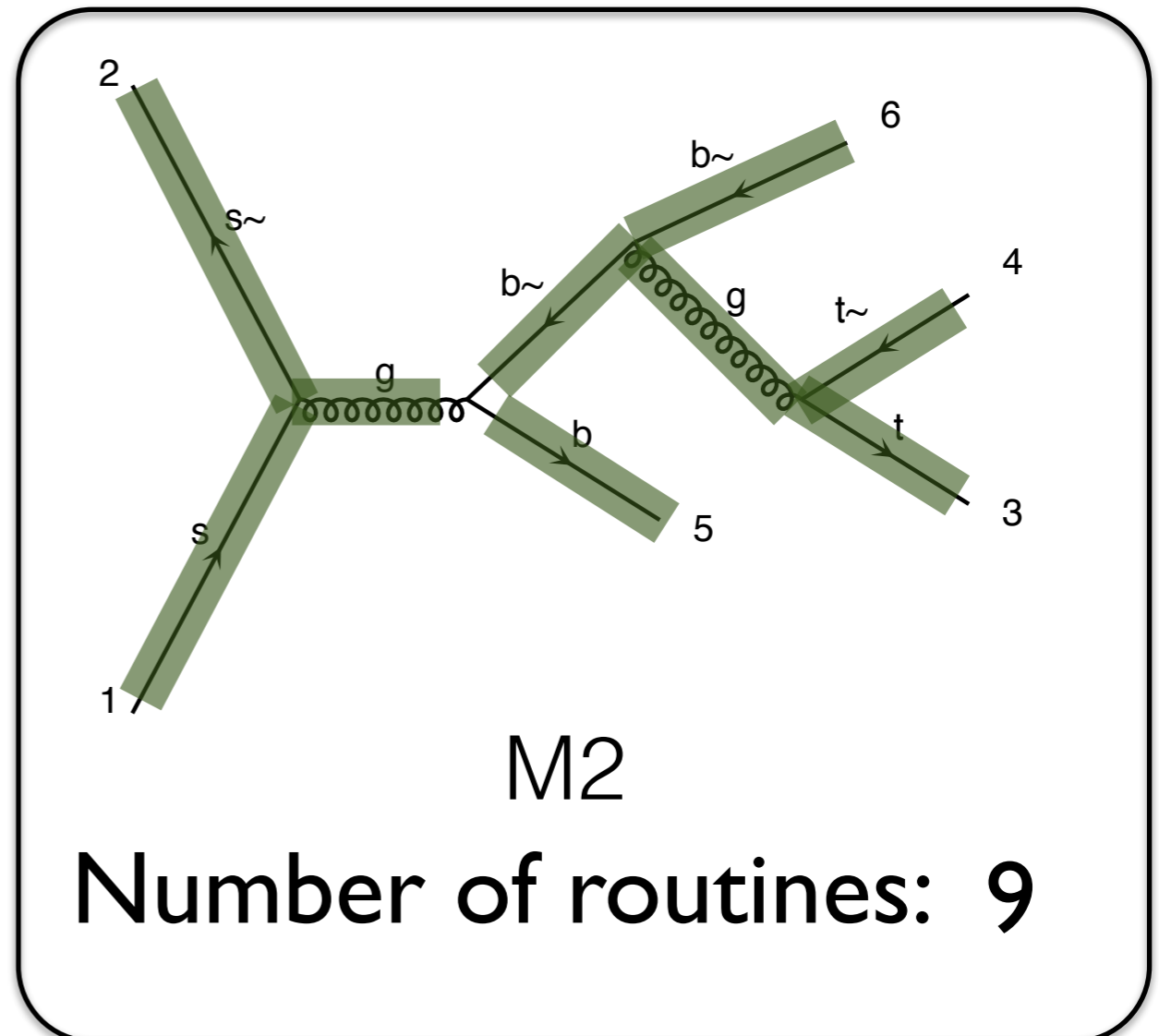
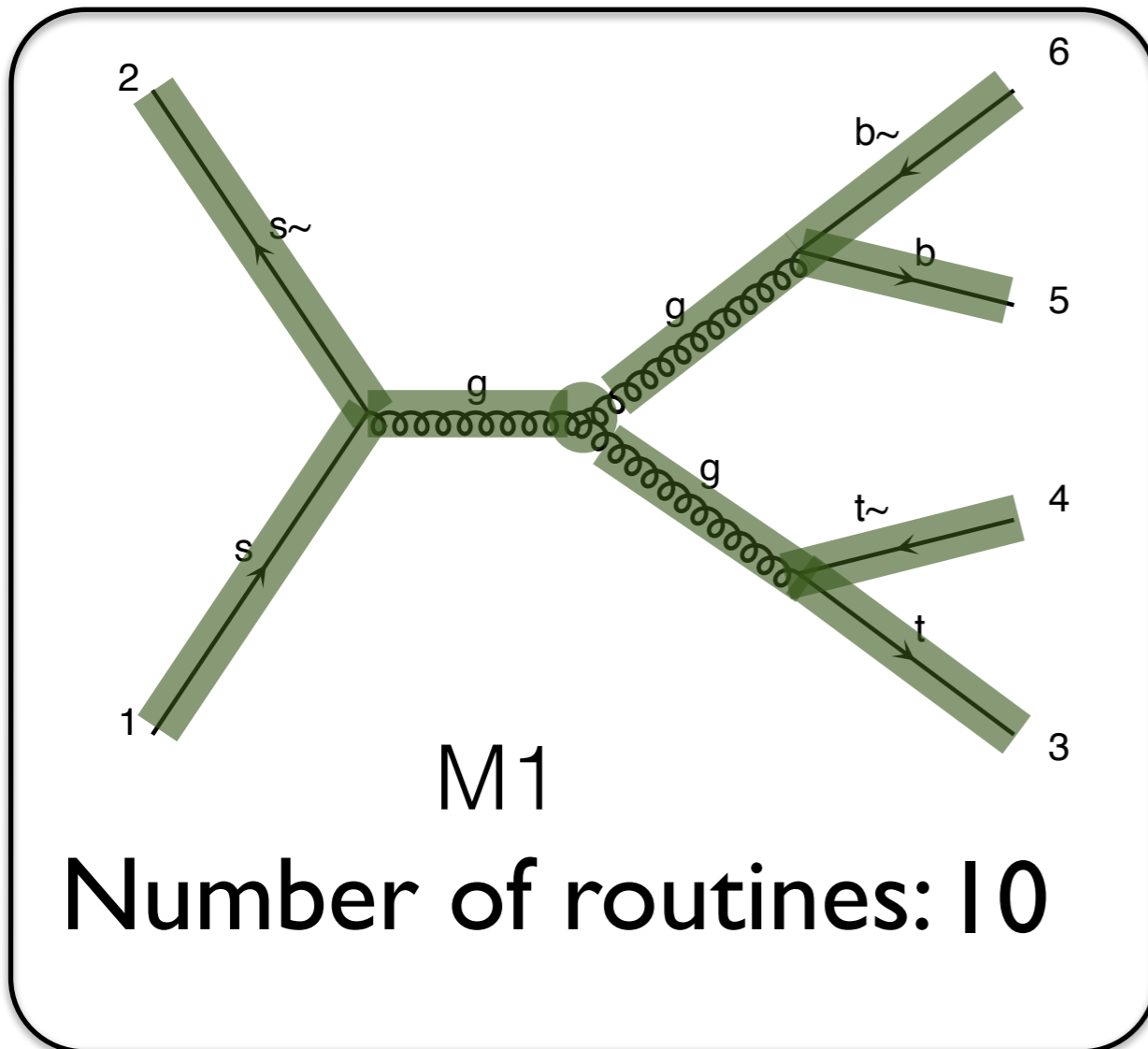
Known



Number of routines for both: 10

$$|M|^2 = |M_1 + M_2|^2$$

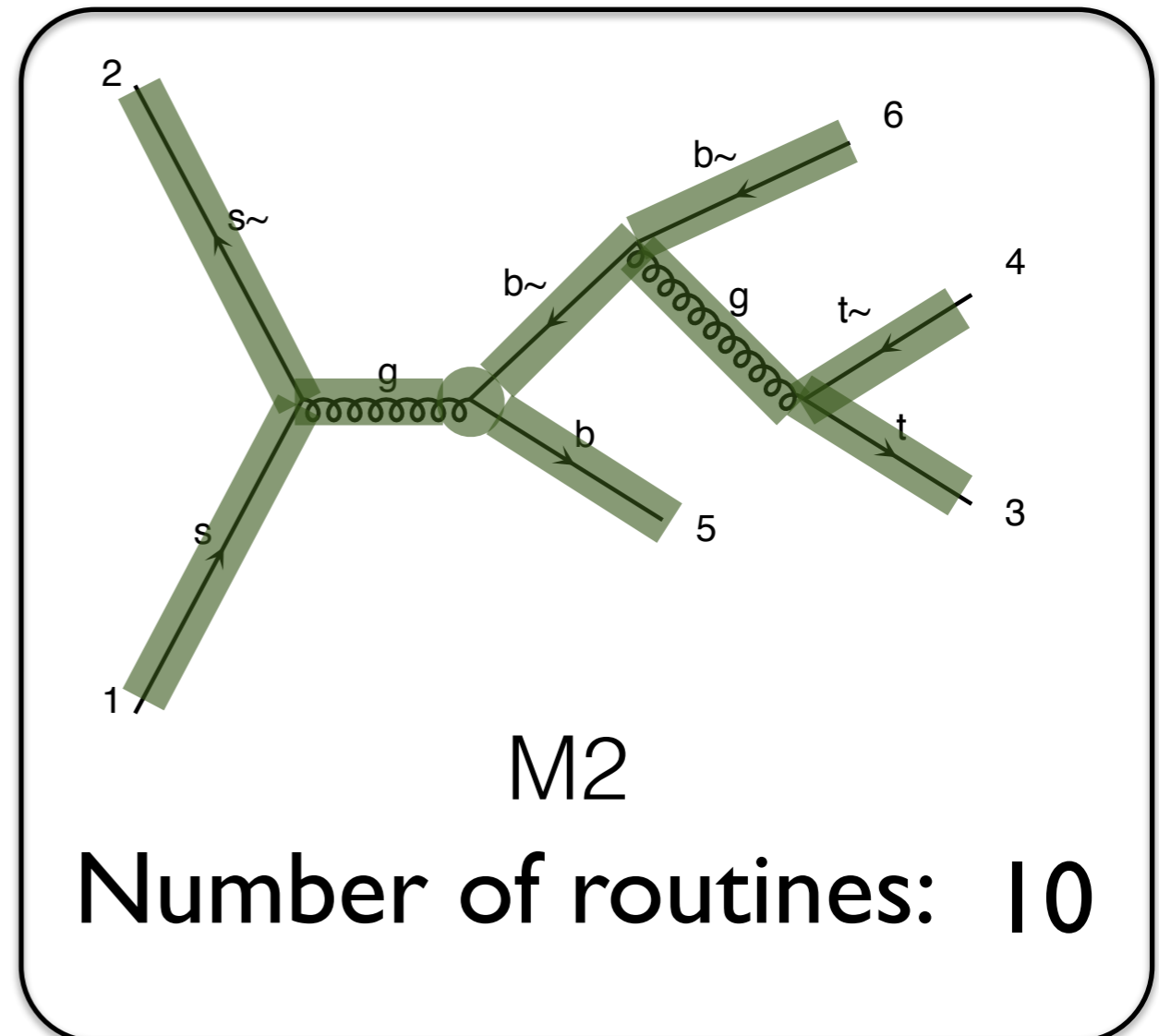
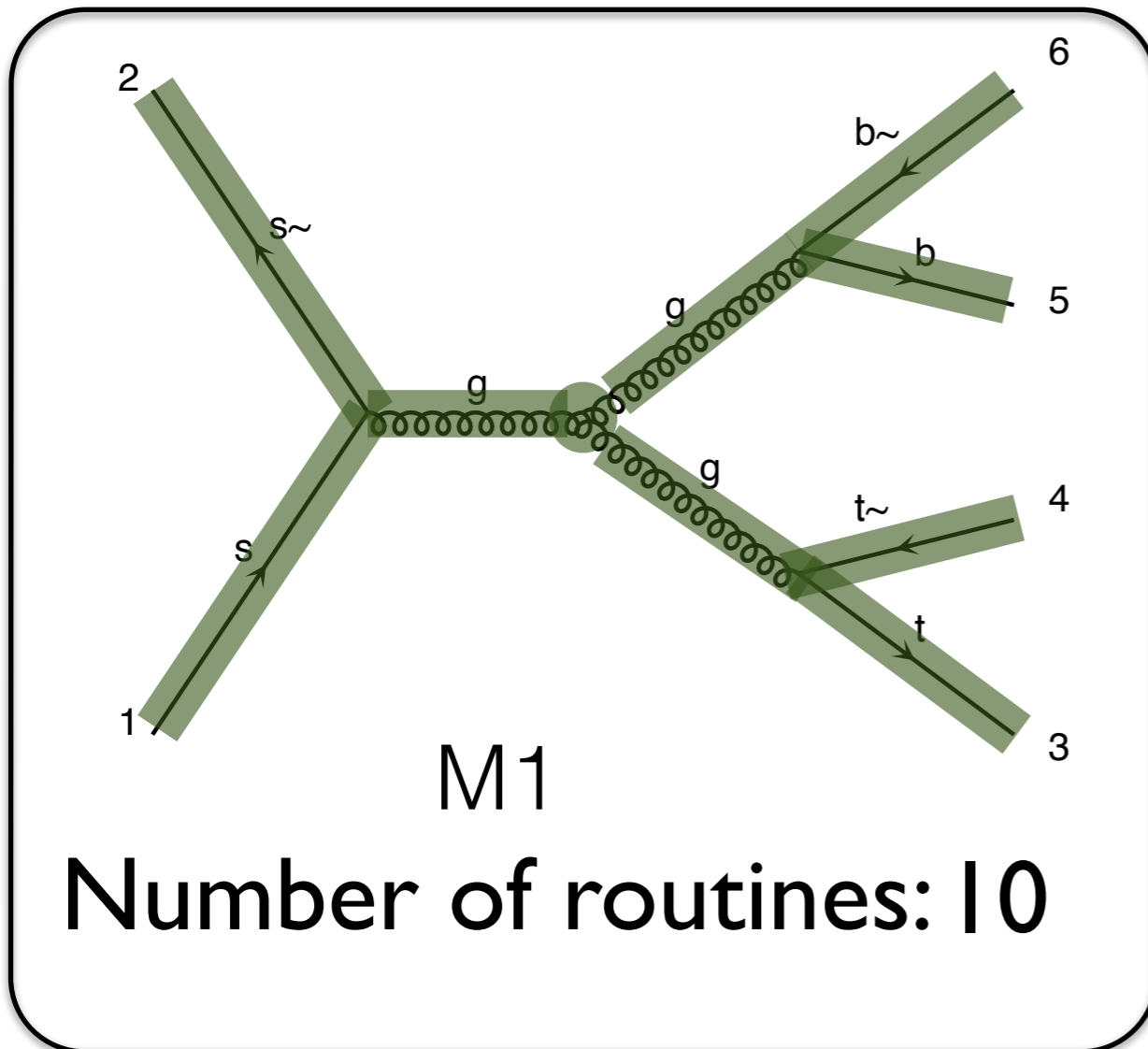
Known



Number of routines for both: 11

$$|M|^2 = |M_1 + M_2|^2$$

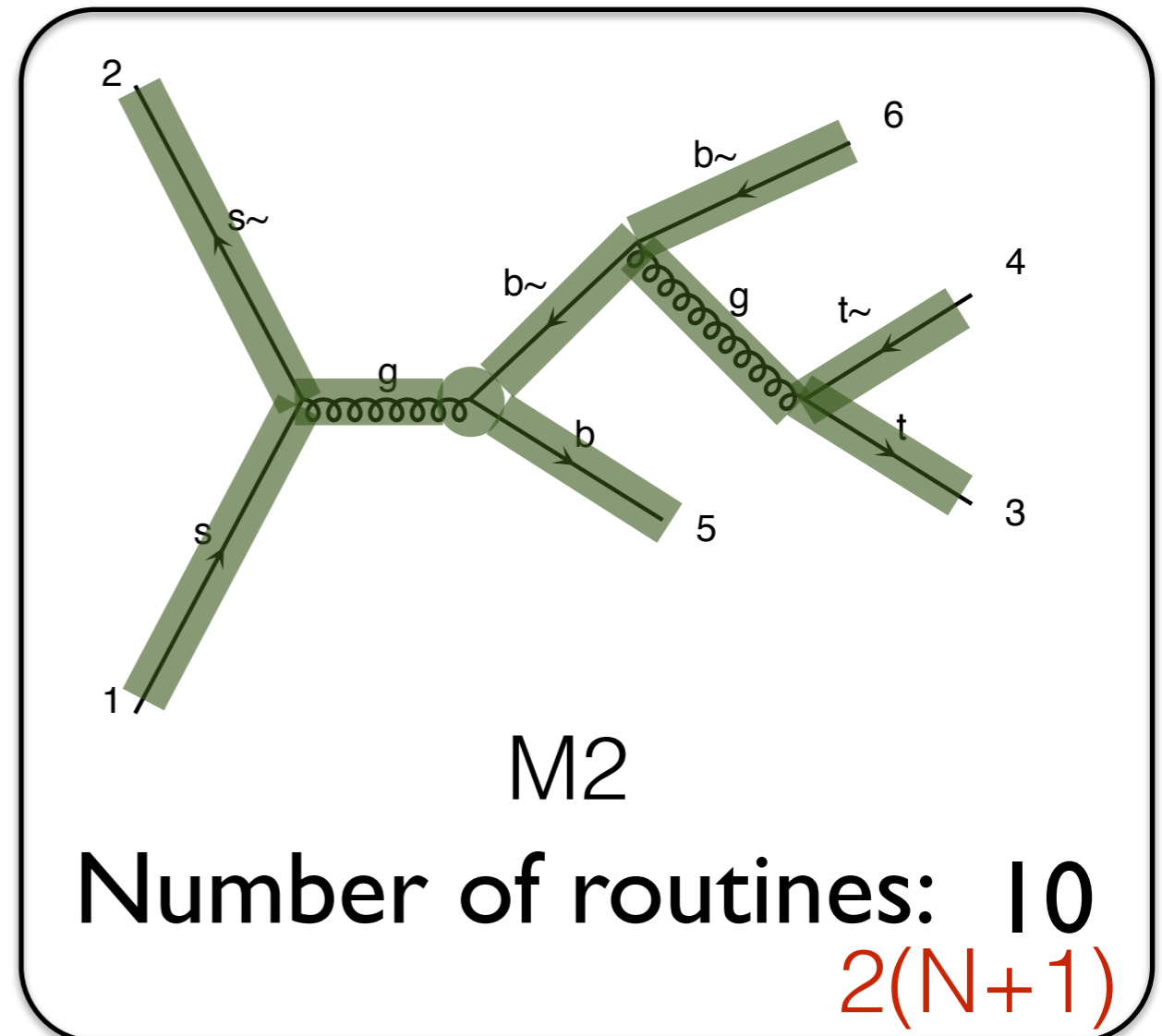
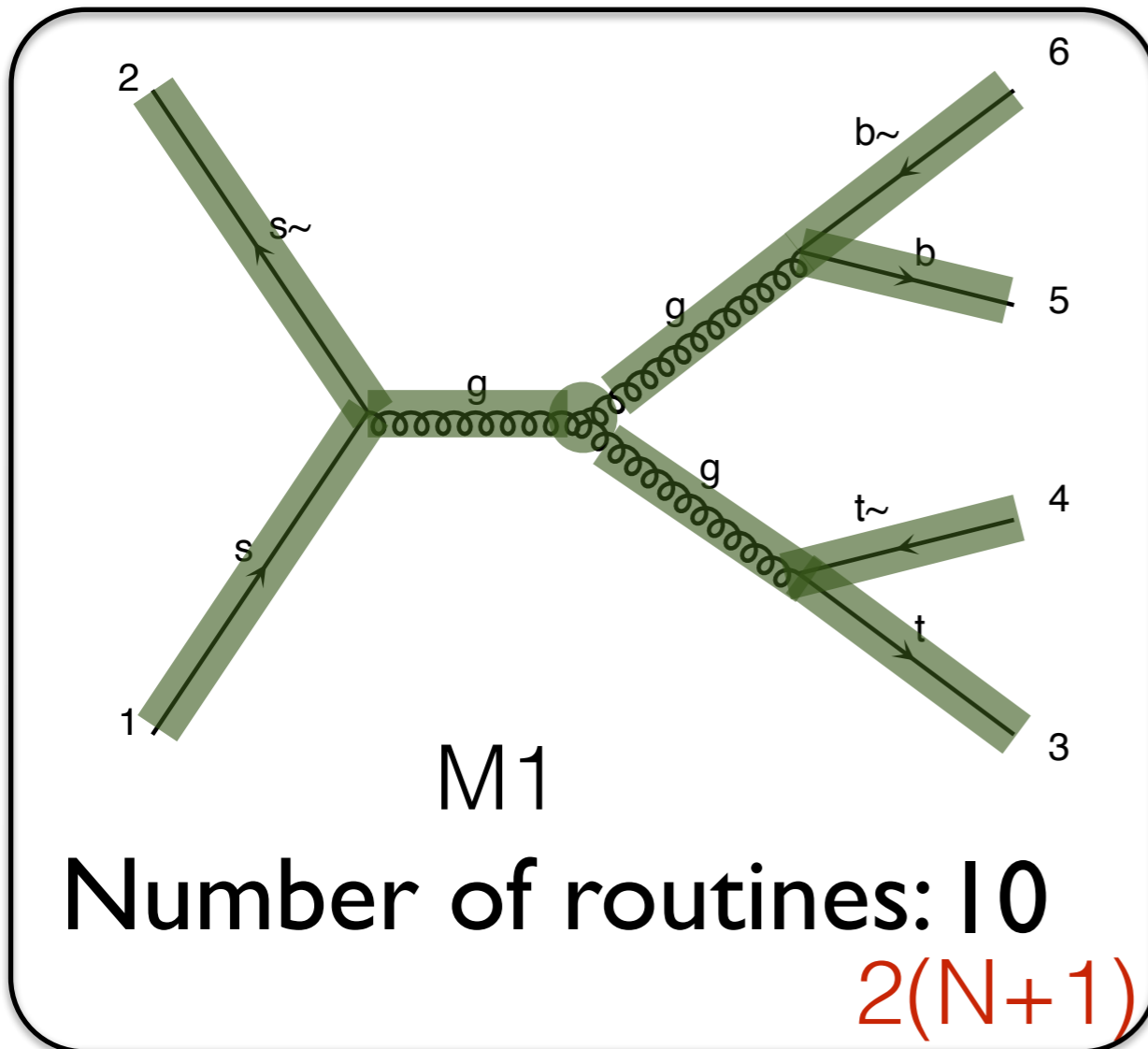
Known



Number of routines for both: 12

$$|M|^2 = |M_1 + M_2|^2$$

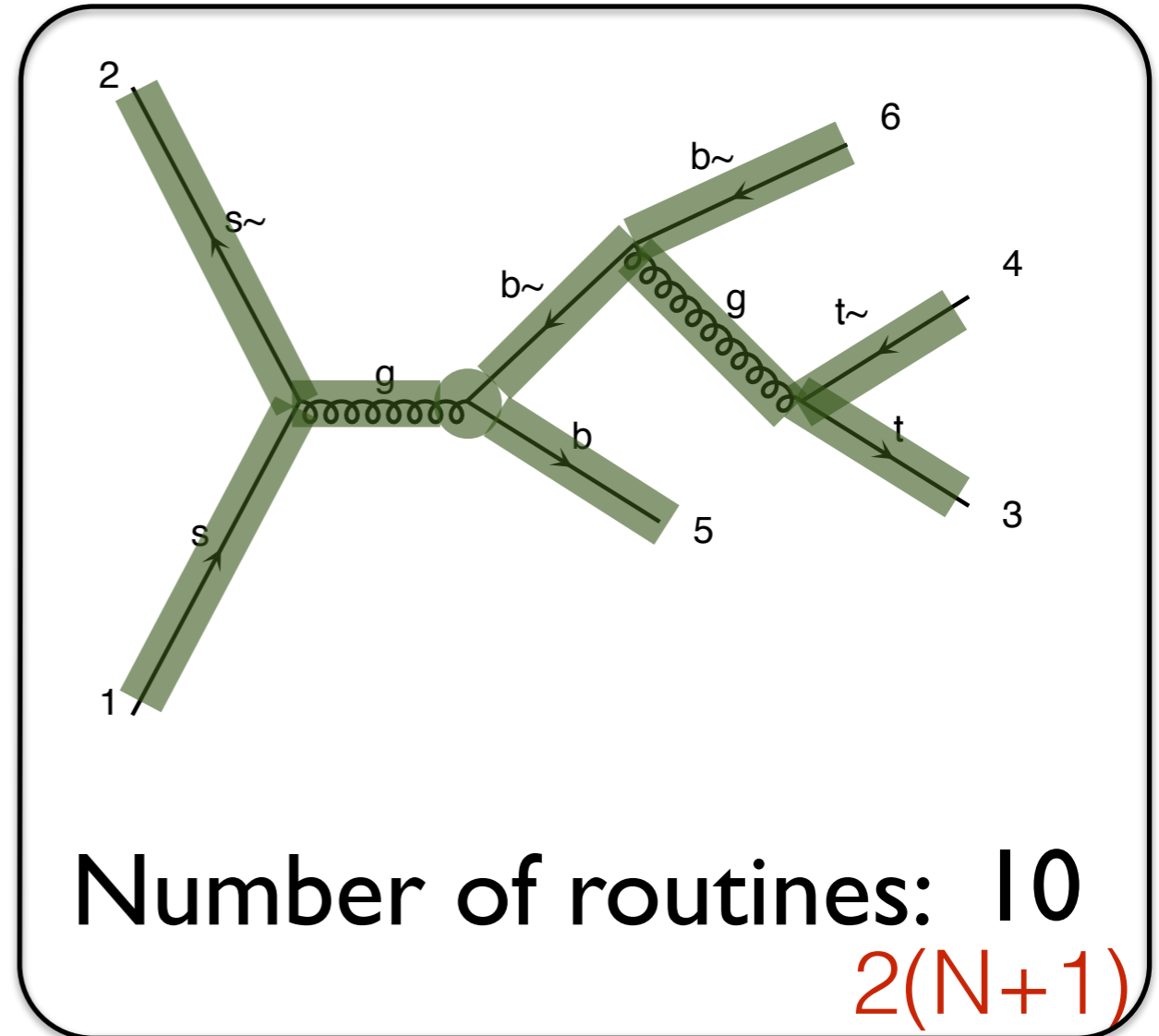
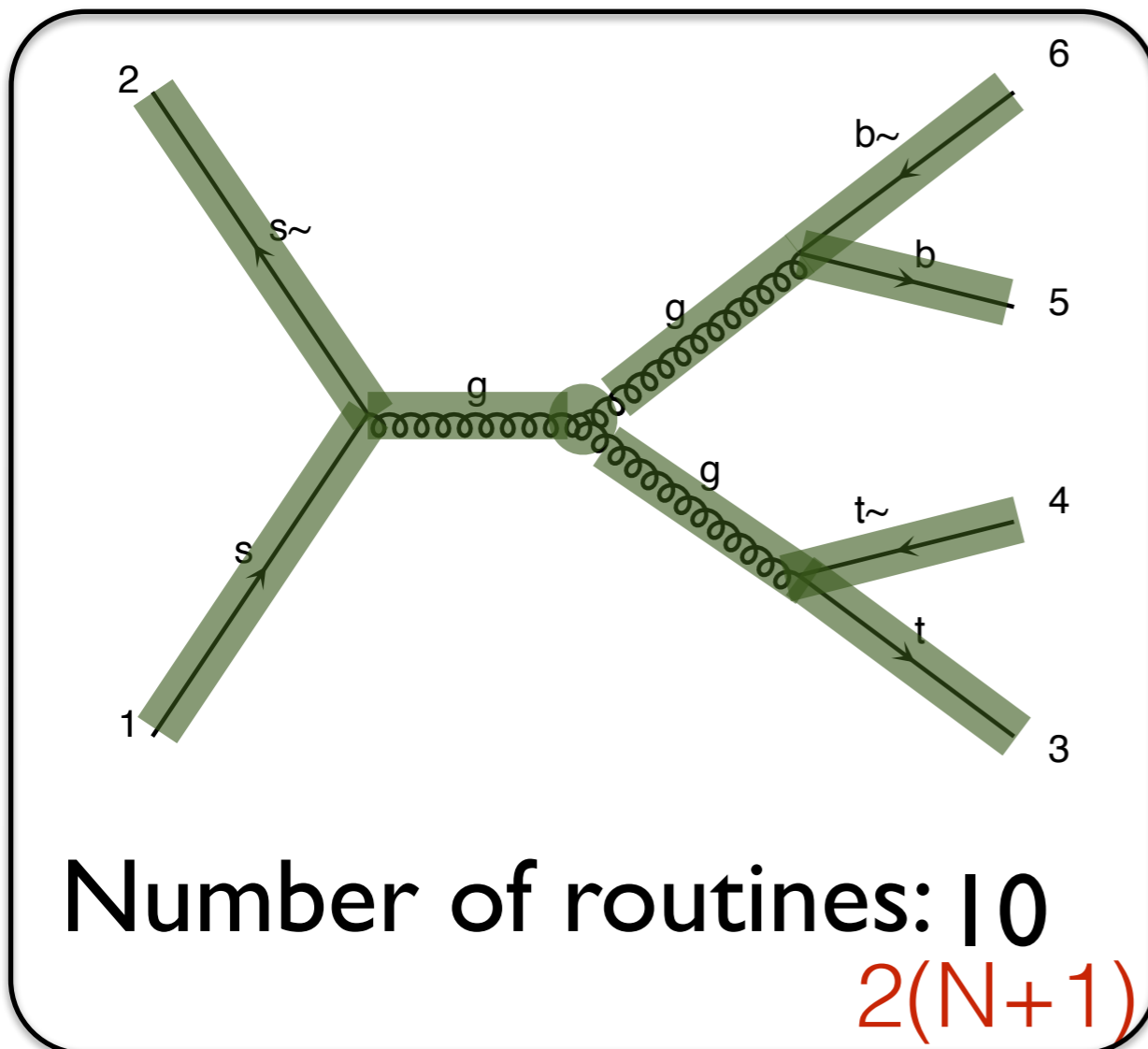
Known



Number of routines for both: 12

$$|M|^2 = |M_1 + M_2|^2$$

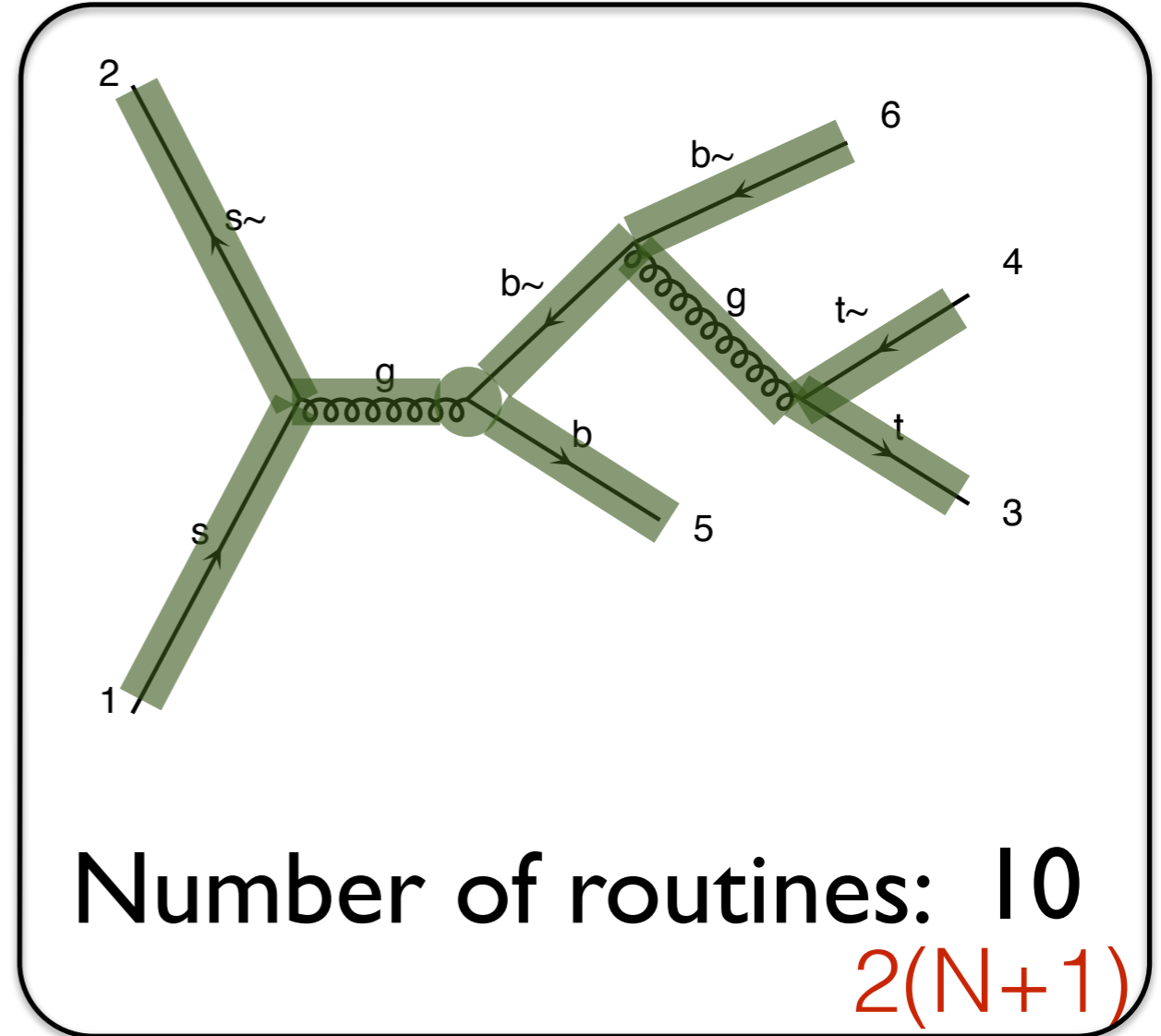
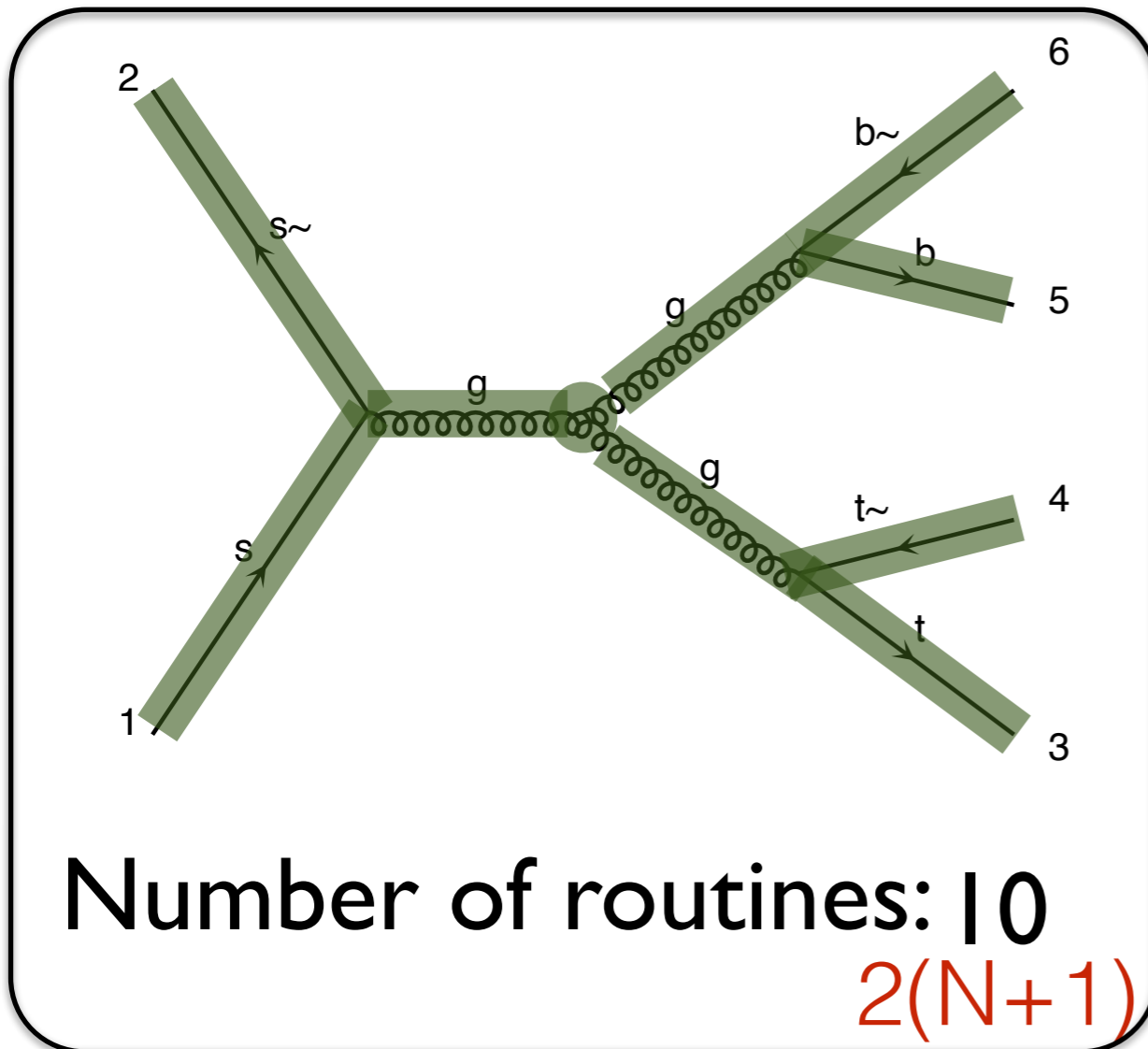
Known



Number of routines for both: 12

$N! * 2(N+1) \longrightarrow N!$

Known



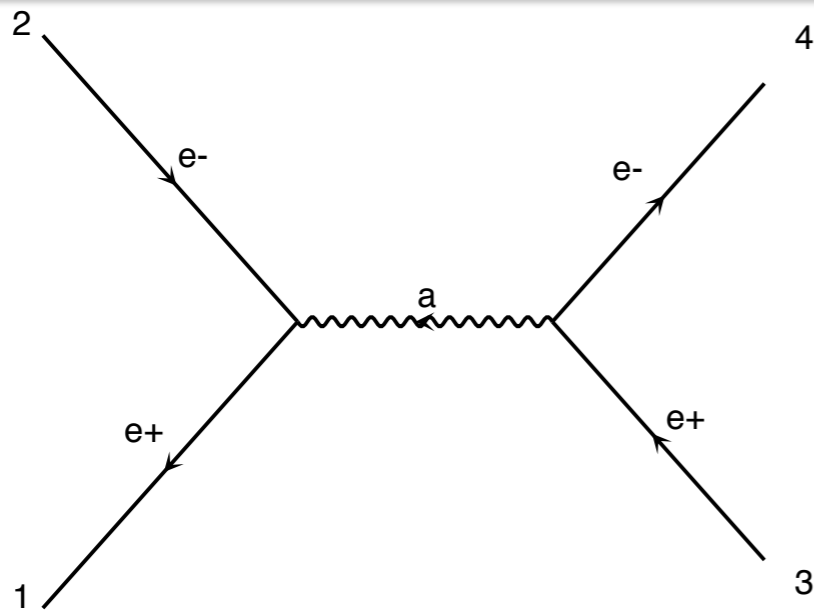
Number of routines for both: 12

$N! * 2(N+1) \longrightarrow N! \xrightarrow{\text{recursion}} 2^N$

- Thanks to new diagram generation algorithm, wf recycling much more efficient in MG5 than MG4

Process	Amplitudes	Wavefunctions		Run time		
		MG 4	MG 5	MG 4	MG 5	
$u\bar{u} \rightarrow e^+e^-$	2	6	6	$< 6\mu\text{s}$	$< 6\mu\text{s}$	
$u\bar{u} \rightarrow e^+e^-e^+e^-$	48	62	32	0.22 ms	0.14 ms	
$u\bar{u} \rightarrow e^+e^-e^+e^-e^+e^-$	3474	3194	301	46.5 ms	19.0 ms	300,000
$u\bar{u} \rightarrow d\bar{d}$	1	5	5	$< 4\mu\text{s}$	$< 4\mu\text{s}$	
$u\bar{u} \rightarrow d\bar{d}g$	5	11	11	27 μs	27 μs	
$u\bar{u} \rightarrow d\bar{d}gg$	38	47	29	0.42 ms	0.31 ms	
$u\bar{u} \rightarrow d\bar{d}ggg$	393	355	122	10.8 ms	6.75 ms	
$u\bar{u} \rightarrow u\bar{u}gg$	76	84	40	1.24 ms	0.80 ms	
$u\bar{u} \rightarrow u\bar{u}ggg$	786	682	174	35.7 ms	17.2 ms	
$u\bar{u} \rightarrow d\bar{d}d\bar{d}$	14	28	19	84 μs	83 μs	
$u\bar{u} \rightarrow d\bar{d}d\bar{d}g$	132	178	65	1.88 ms	1.15 ms	
$u\bar{u} \rightarrow d\bar{d}d\bar{d}gg$	1590	1782	286	141 ms	34.4 ms	
$u\bar{u} \rightarrow d\bar{d}d\bar{d}d\bar{d}$	612	758	141	42.5 ms	6.6 ms	5500

Time for matrix element evaluation on a Sony Vaio TZ laptop



$$\mathcal{M} = \bar{u}\gamma^\mu v P_{\mu\nu} \bar{u}\gamma^\nu v$$

$$\bar{u}_1 = fct(\vec{p}_1, m)$$

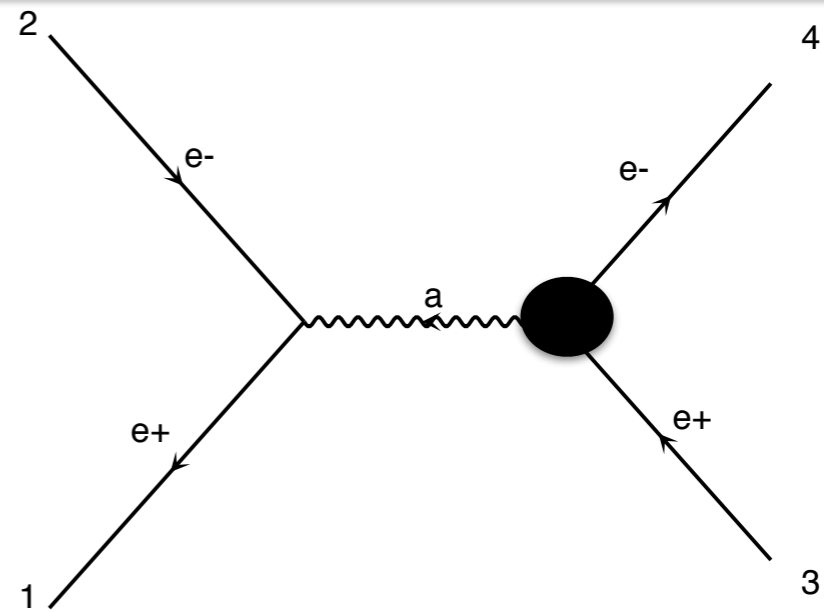
$$v_2 = fct(\vec{p}_2, m)$$

$$\bar{u}_3 = fct(\vec{p}_3, m)$$

$$v_4 = fct(\vec{p}_4, m)$$

$$W_a = fct(\bar{u}_1, v_2, M_a, \Gamma_a)$$

$$\mathcal{M} = fct(\bar{u}_3, v_4, W_a)$$



$$\mathcal{M} = \bar{u}\gamma^\mu v P_{\mu\nu} \bar{u}\Gamma^\mu v$$

$$\bar{u}_1 = fct(\vec{p}_1, m)$$

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[Murayama, Watanabe, Hagiwara]

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	Chiral Perturbation	BNV Model
SLIH	Effective Field Theory	NMSSM
Full HEFT	Chromo-magnetic operator	Black Holes

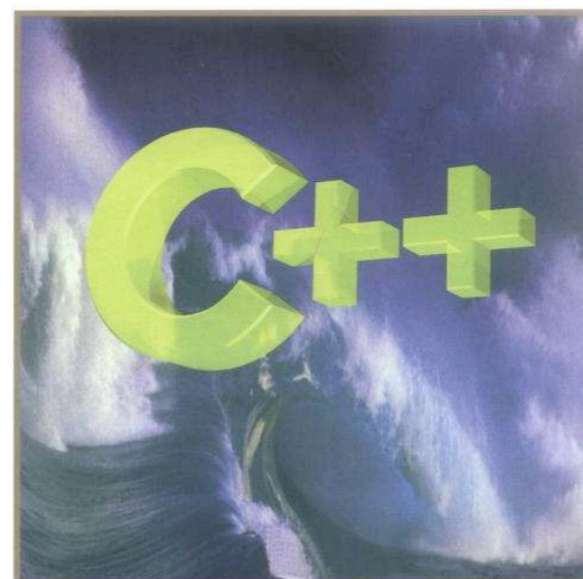
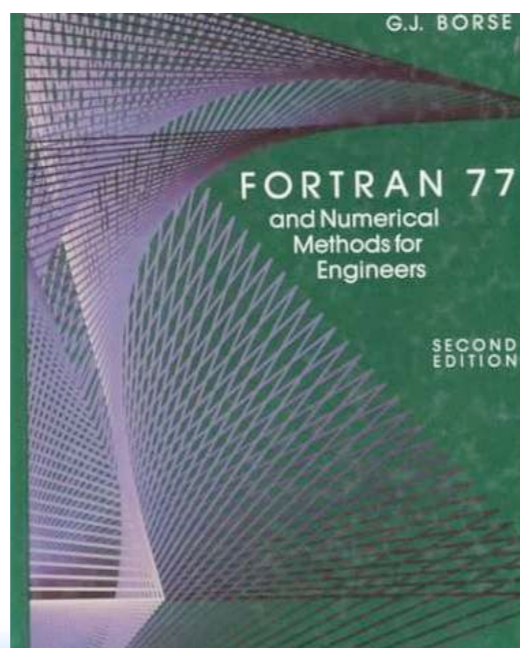


ALOHA

~~ALOHA~~
~~Google~~ translate

From: [UFO] To: Helicity [Translate]

Type text or a website address or [translate a document](#).





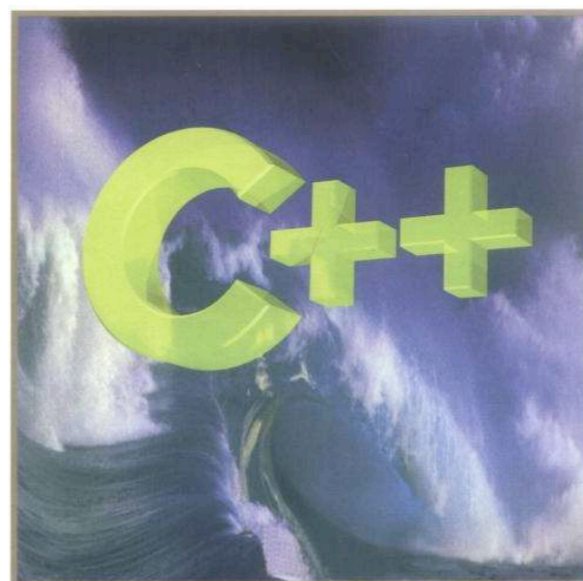
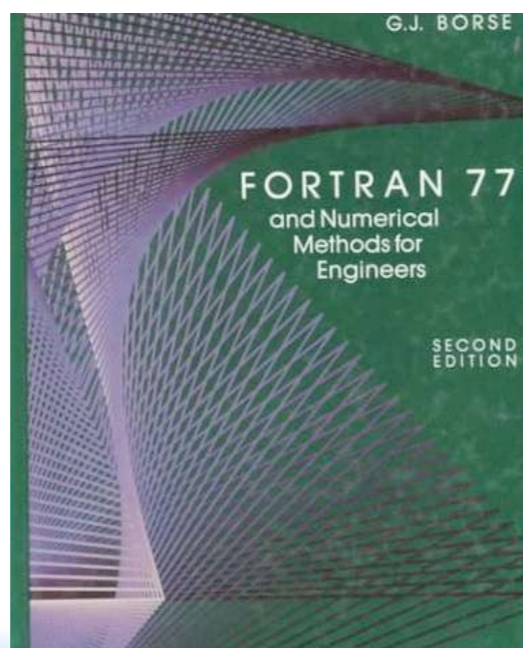
ALOHA

~~ALOHA~~
~~Google~~ translate

From: [UFO] To: Helicity [Translate]

Basically, any new operator can be handle by MG5/Pythia8 out of the box!

Type text or a website address or [translate a document](#).



Input

```
FFV1 = Lorentz(name = 'FFV1',
               spins = [ 2, 2, 3 ],
               structure = 'Gamma(3,2,1)')
```

Output

```
C      This File is Automatically generated by ALOHA
C      The process calculated in this file is:
C      Gamma(3,2,1)
C
SUBROUTINE FFV1_0(F1,F2,V3,C,VERTEX)
IMPLICIT NONE
DOUBLE COMPLEX F1(6)
DOUBLE COMPLEX F2(6)
DOUBLE COMPLEX V3(6)
DOUBLE COMPLEX C
DOUBLE COMPLEX VERTEX

      VERTEX = C*( (F2(1)*( (F1(3)*( (0, -1)*V3(1)+(0, 1)*V3(4)))
$ +(F1(4)*( (0, 1)*V3(2)+V3(3)))))+( (F2(2)*( (F1(3)*( (0, 1)
$ *V3(2)-V3(3))))+(F1(4)*( (0, -1)*V3(1)+(0, -1)*V3(4))))))
$ +( (F2(3)*( (F1(1)*( (0, -1)*V3(1)+(0, -1)*V3(4)))+(F1(2)
$ *( (0, -1)*V3(2)-V3(3)))))+(F2(4)*( (F1(1)*( (0, -1)*V3(2)
$ +V3(3)))+(F1(2)*( (0, -1)*V3(1)+(0, 1)*V3(4)))))))))

      END
```


- Compute those Function Analytically
- Code in Python
- Can handle
 - all spin up to 2
 - custom propagator
 - majorana (but in 4 fermion operator)
 - Any dimensional operator
- Only use in MadGraph5_aMC@NLO
- Plan to have similar tools for the other generator

- Numerical computation faster than analytical computation
- We are able to compute matrix-element
 - ➔ for large number of final state
 - ➔ for any BSM theory
 - ➔ actually also for loop

Monte Carlo Integration and Generation

Calculations of cross section or decay widths involve integrations over high-dimension phase space of very peaked functions:

$$\sigma = \frac{1}{2s} \int |\mathcal{M}|^2 d\Phi(n)$$

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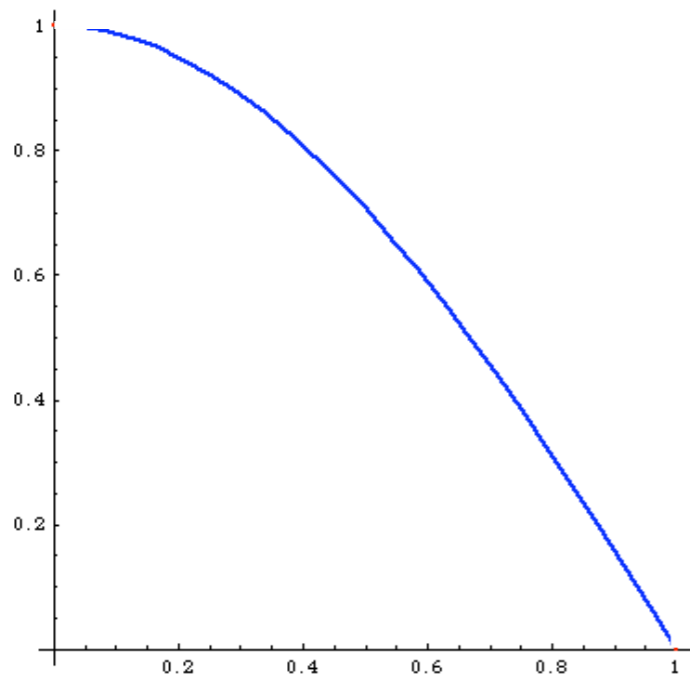
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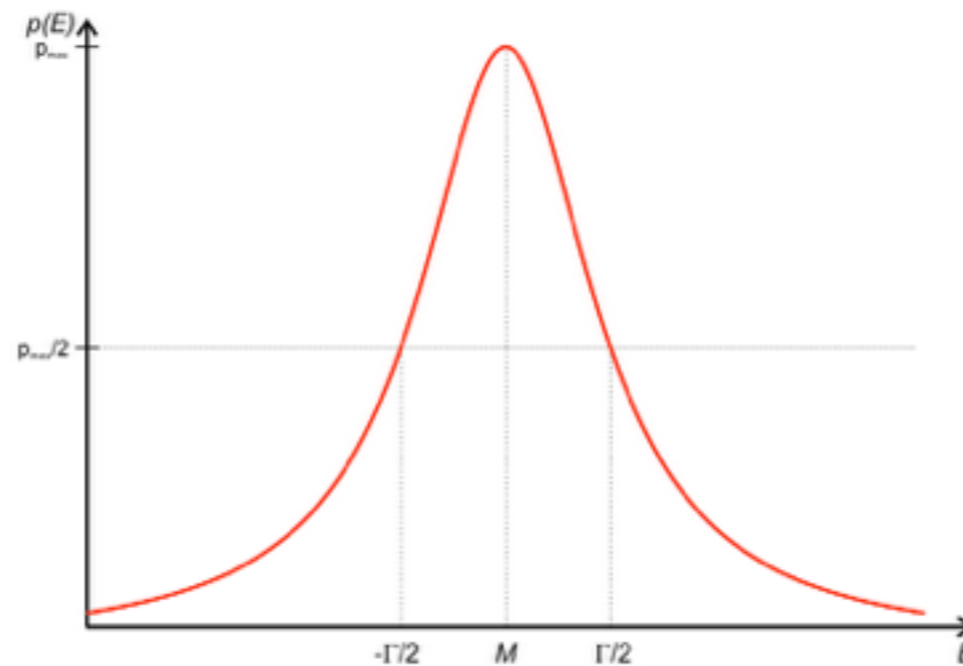
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General and flexible method is needed

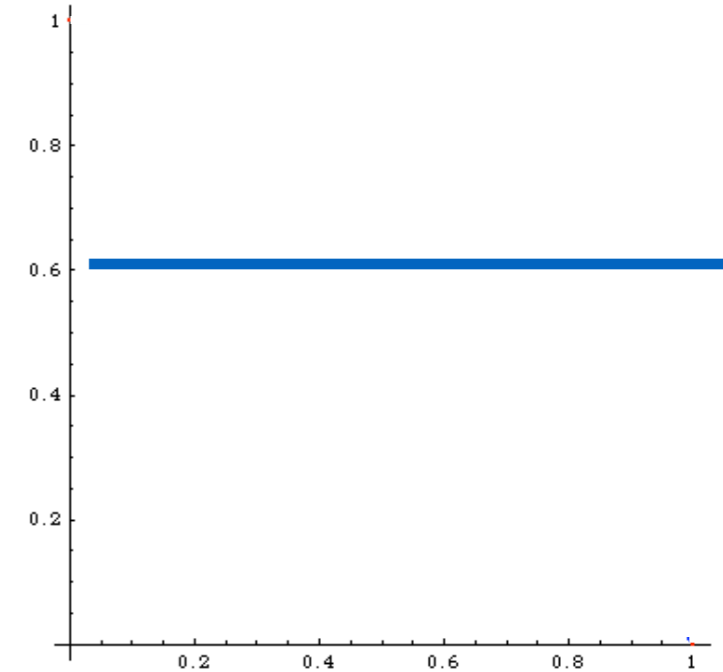
$$I = \int_0^1 dx \cos \frac{\pi}{2} x$$



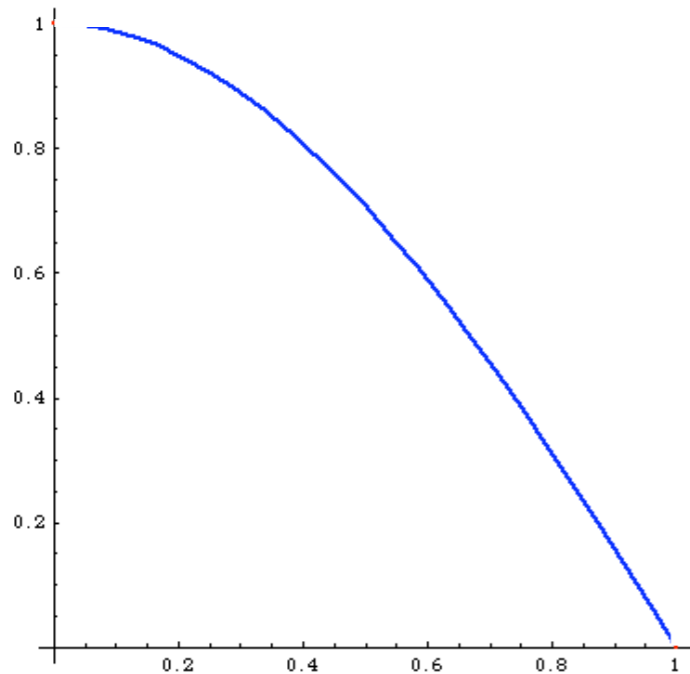
$$\int \frac{dq^2}{(q^2 - M^2 + iM\Gamma)^2}$$



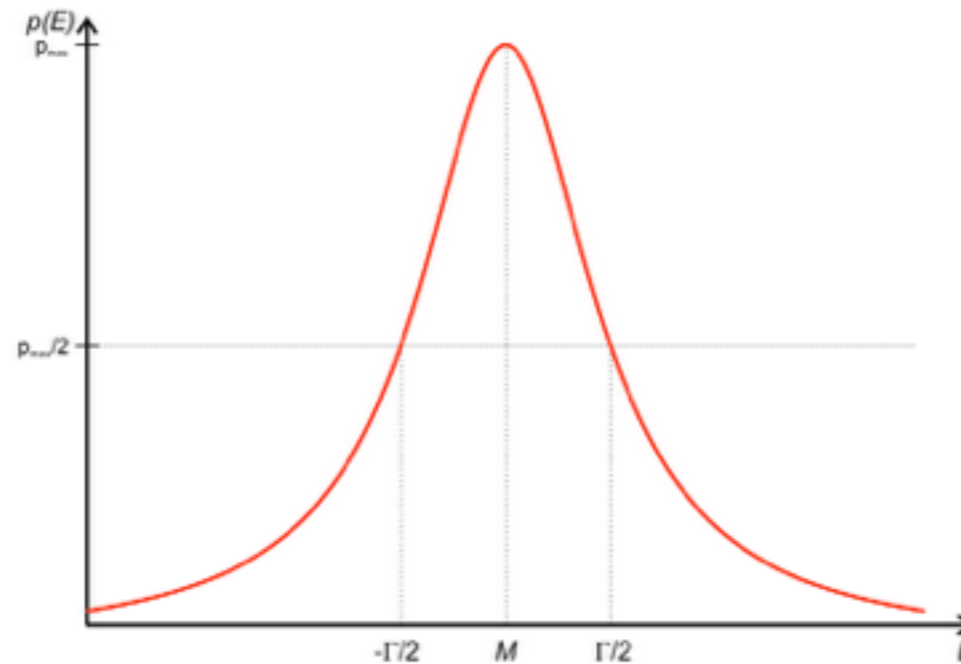
$$\int dx C$$



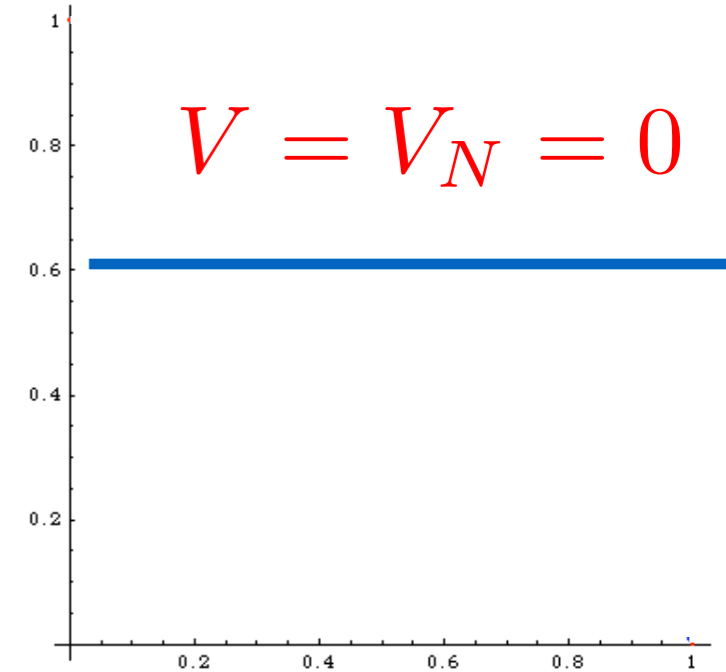
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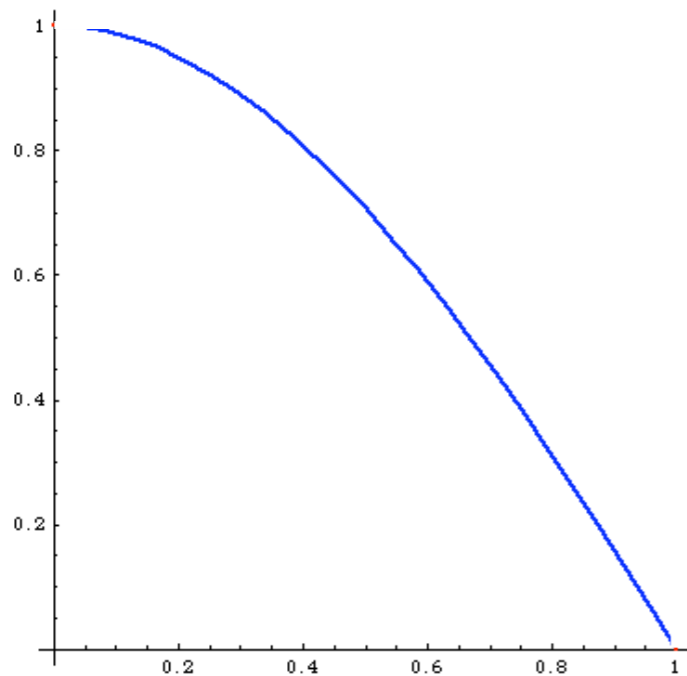
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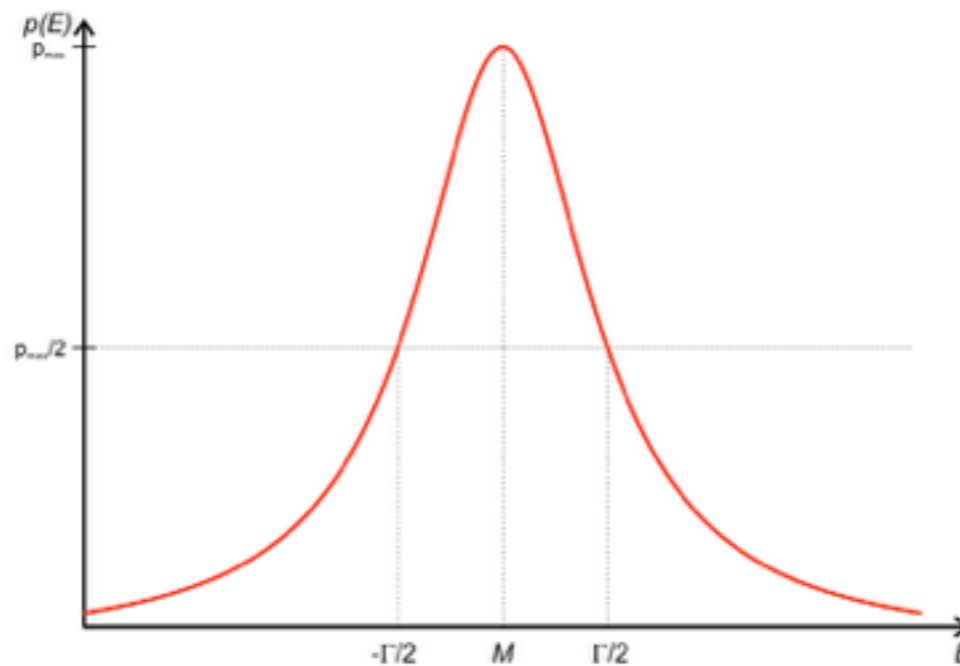
Method of evaluation

- MonteCarlo $1/\sqrt{N}$
- Trapezium $1/N^2$
- Simpson $1/N^4$

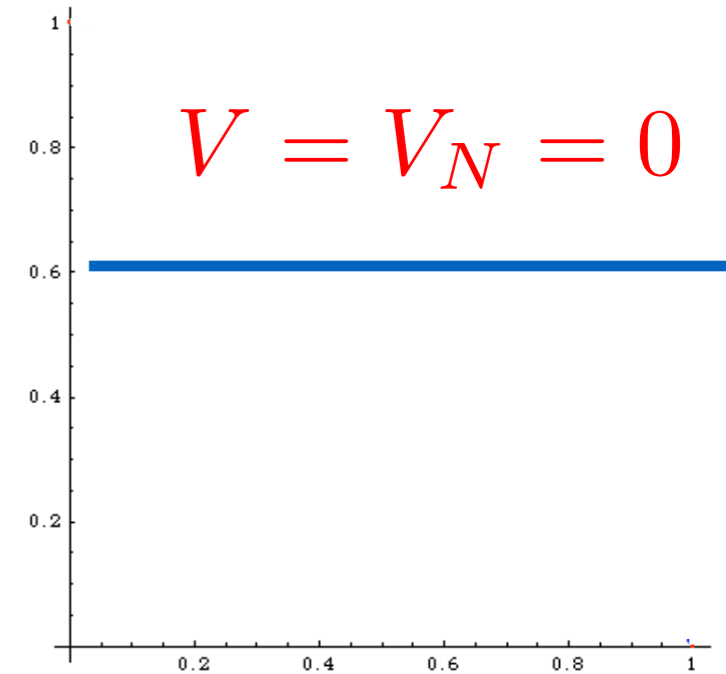
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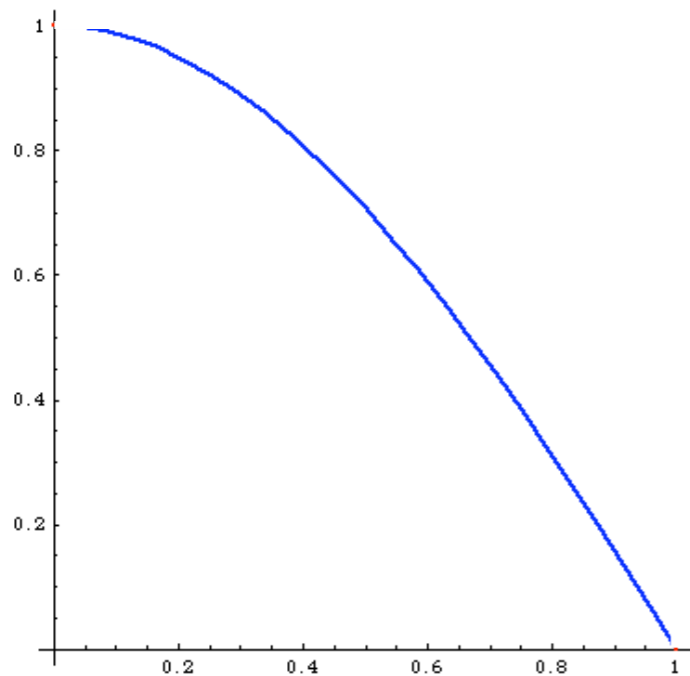


	simpson	MC
3	0,638	0,3
5	0,6367	0,8
20	0,63662	0,6
100	0,636619	0,65
1000	0,636619	0,636

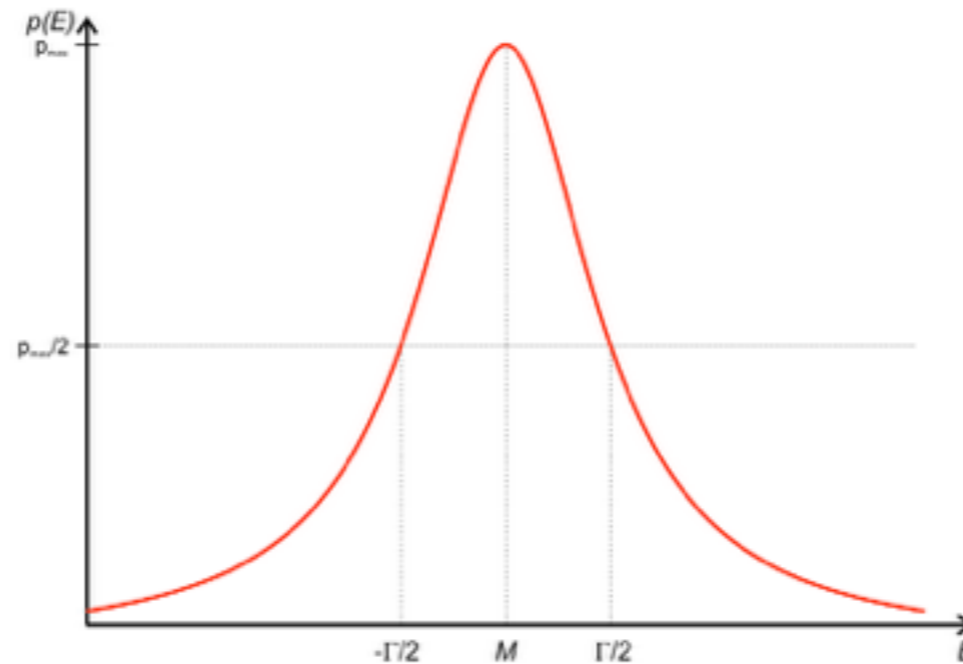
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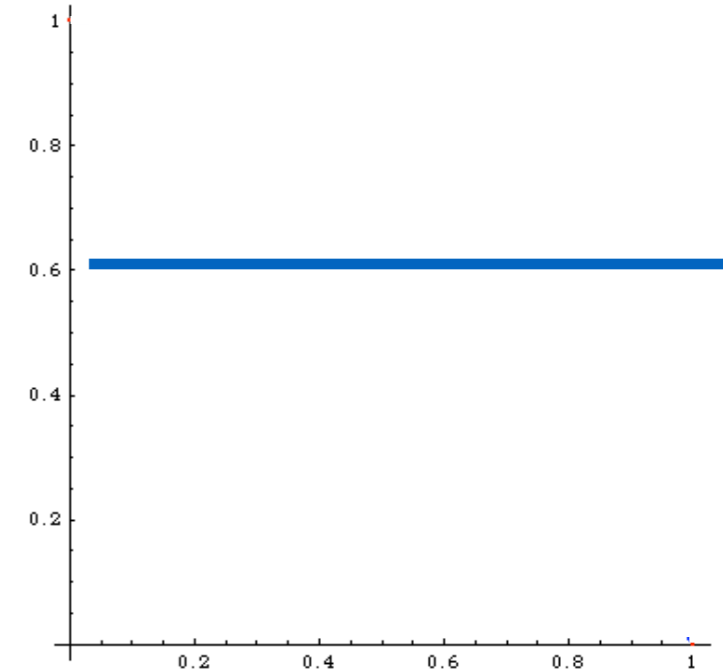
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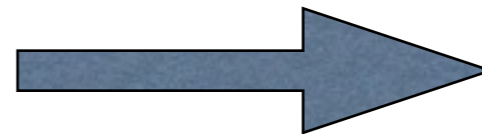
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More Dimension

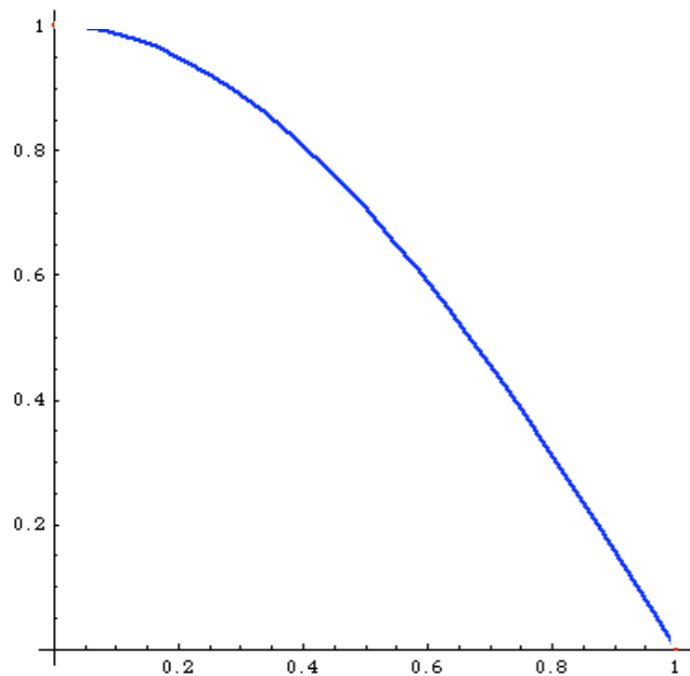


$$1/\sqrt{N}$$

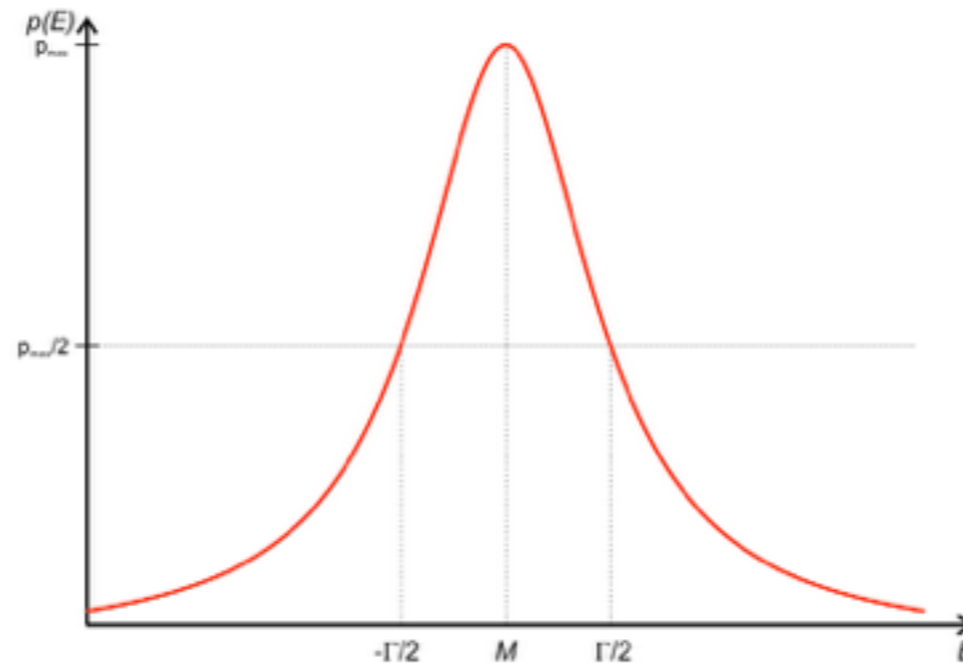
$$1/N^{2/d}$$

$$1/N^{4/d}$$

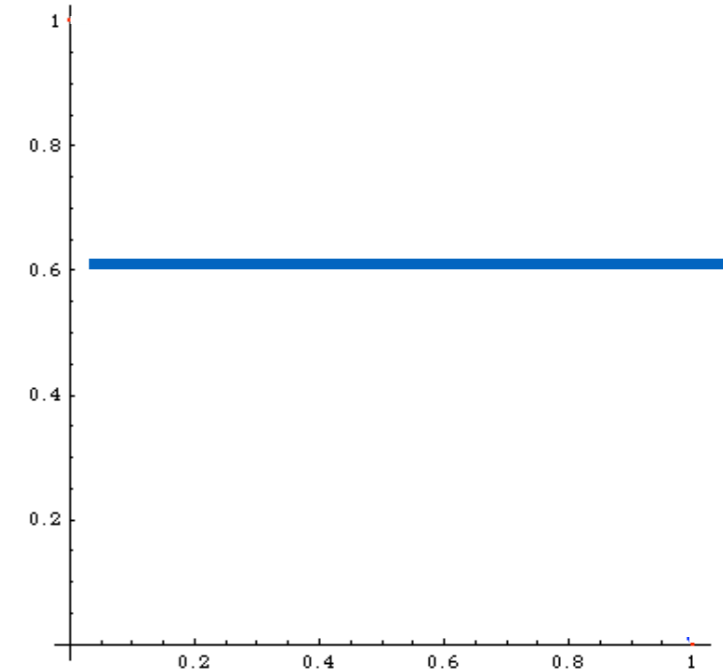
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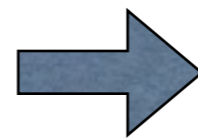
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$$\int dx C$$

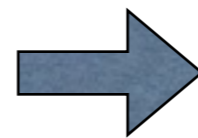


$$I = \int_{x_1}^{x_2} f(x) dx$$



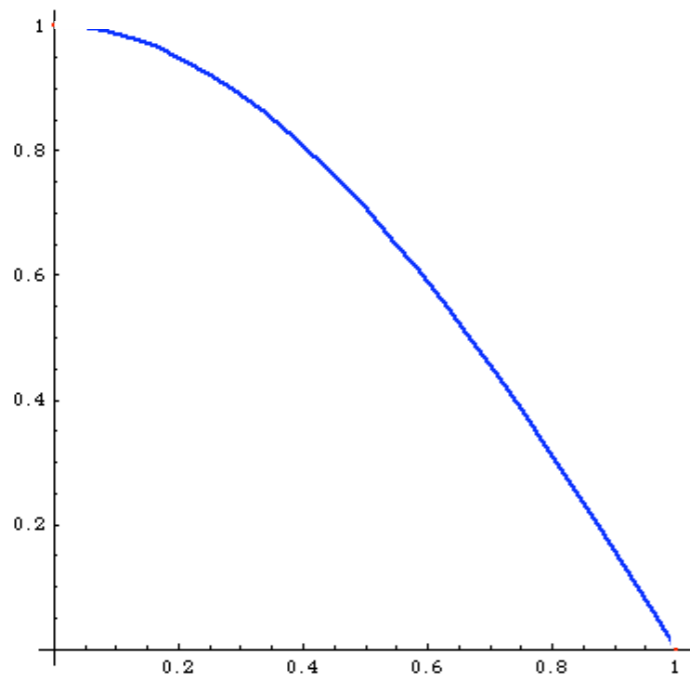
$$I_N = (x_2 - x_1) \frac{1}{N} \sum_{i=1}^N f(x)$$

$$V = (x_2 - x_1) \int_{x_1}^{x_2} [f(x)]^2 dx - I^2$$



$$V_N = (x_2 - x_1)^2 \frac{1}{N} \sum_{i=1}^N [f(x)]^2 - I_N^2$$

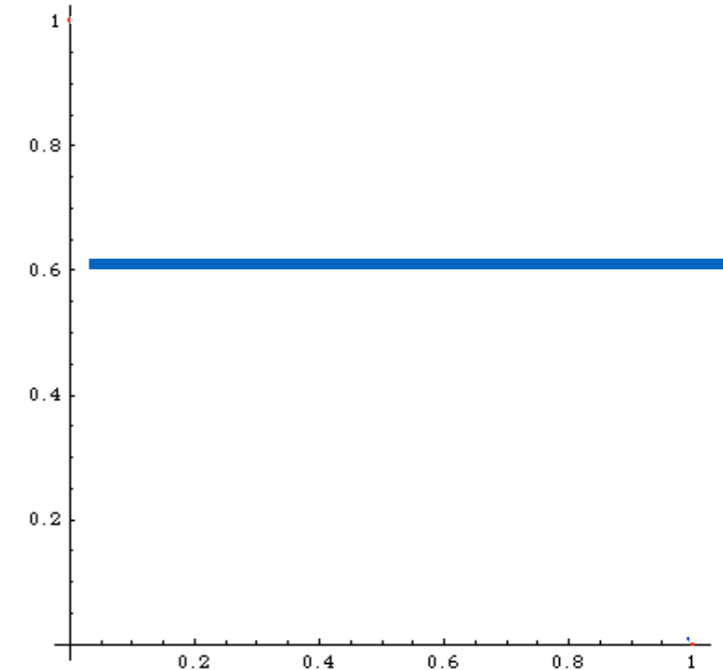
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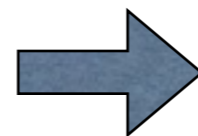
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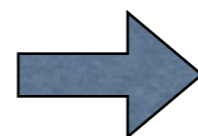


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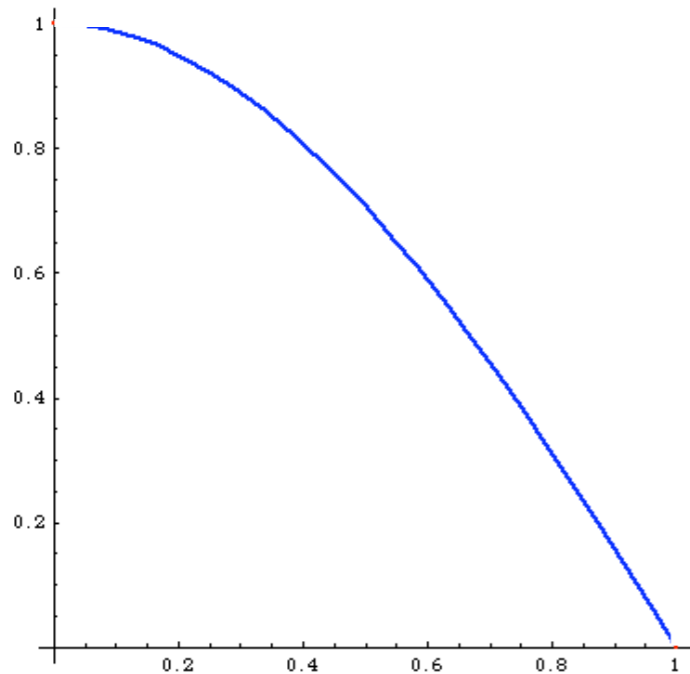
$$V = (x_2 - x_1) \int_{x_1}^{x_2} [f(x)]^2 dx - I^2$$



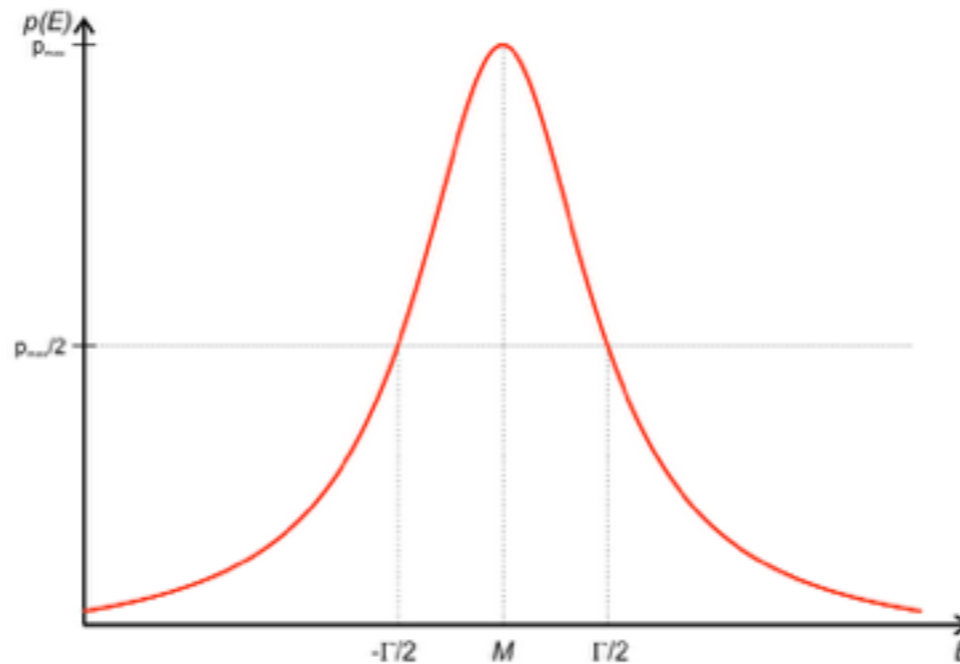
$$V_N = (x_2 - x_1)^2 \frac{1}{N} \sum_{i=1}^N [f(x)]^2 - I_N^2$$

$$I = I_N \pm \sqrt{V_N/N}$$

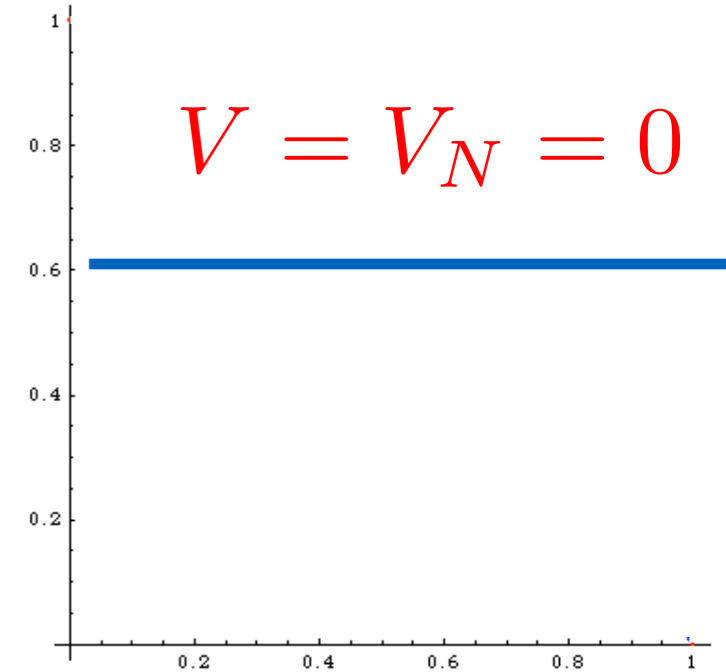
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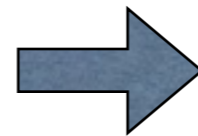
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$$\int dx C$$

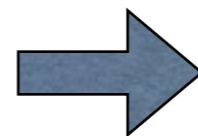


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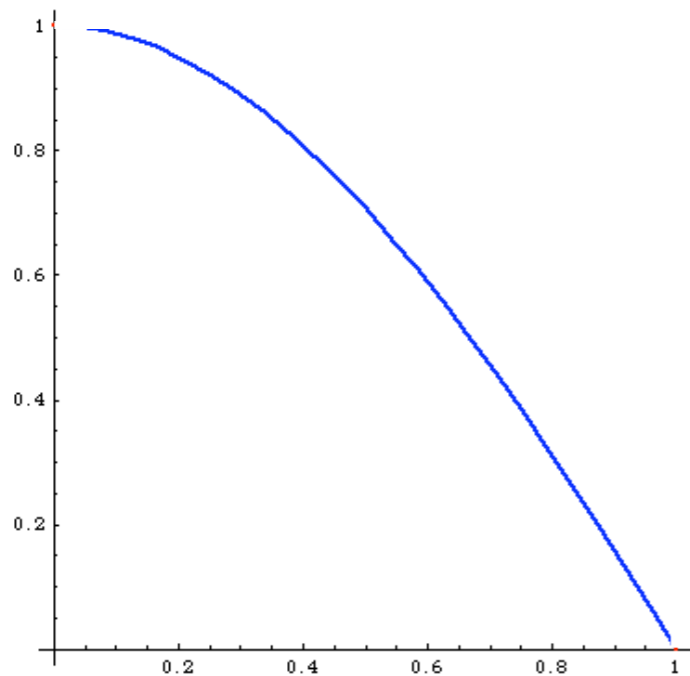
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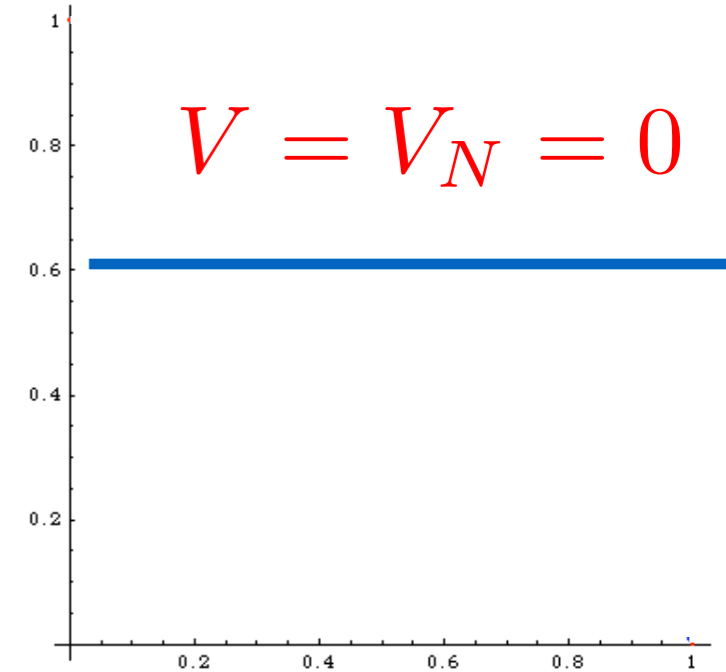
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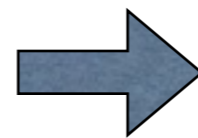
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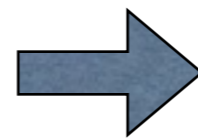


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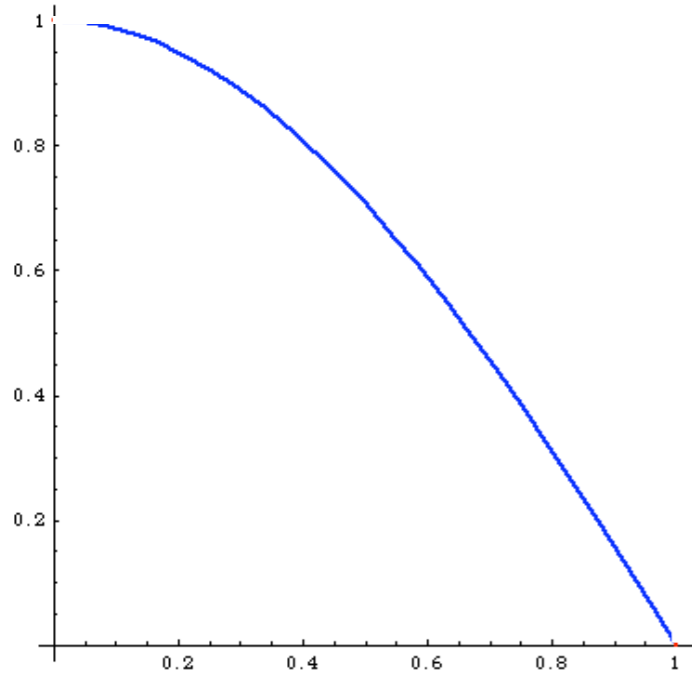
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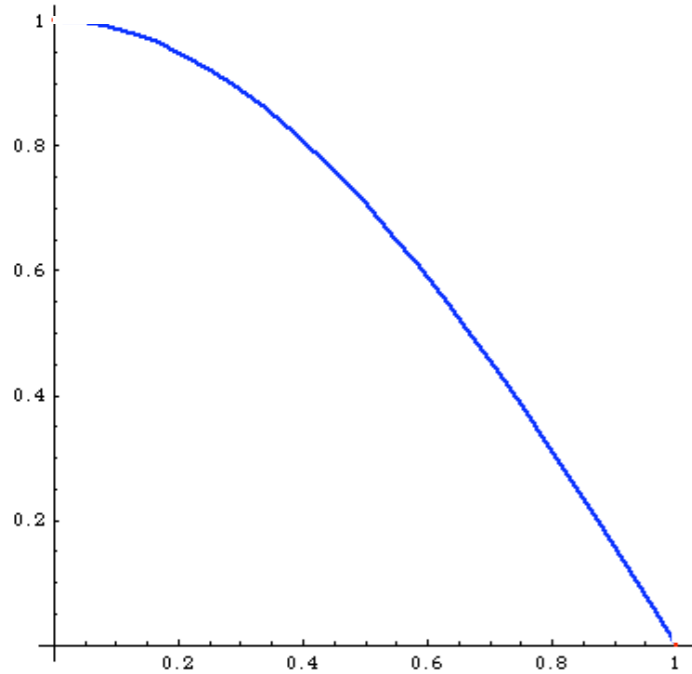
$$I = I_N \pm \sqrt{V_N/N}$$

Can be minimized!



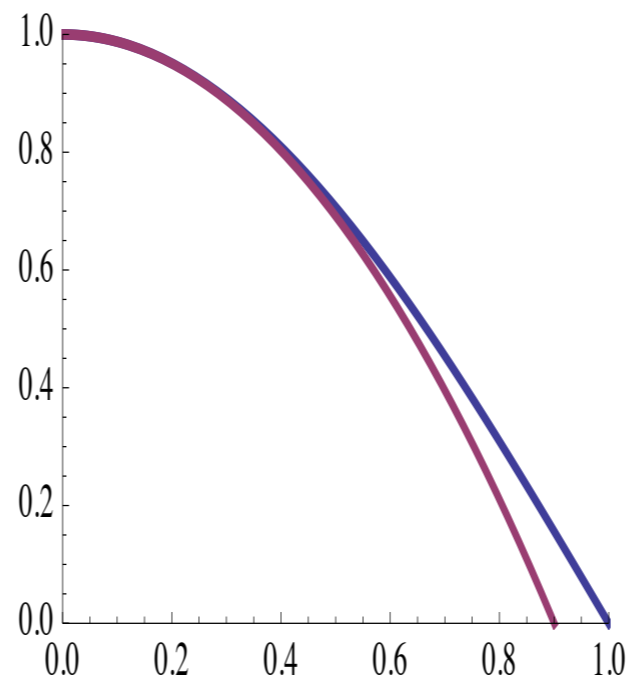
$$I = \int_0^1 dx \cos \frac{\pi}{2} x$$

$$I_N = 0.637 \pm 0.307/\sqrt{N}$$

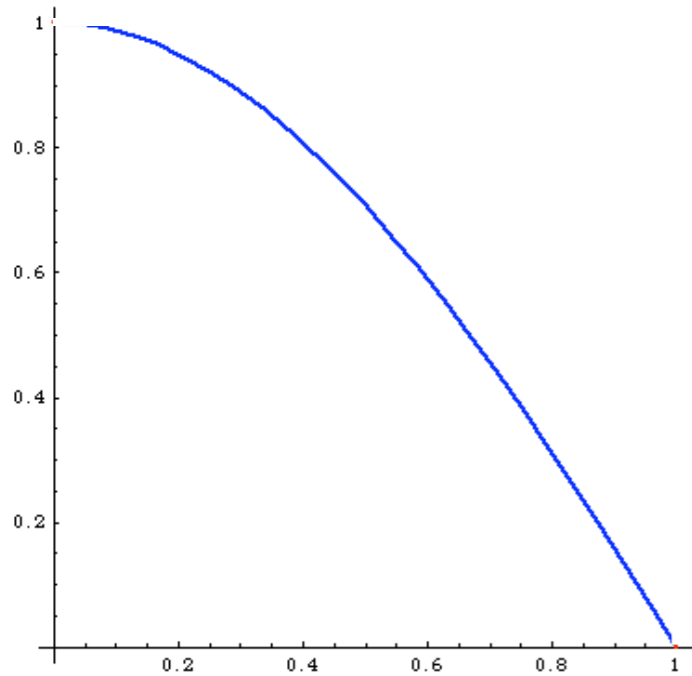


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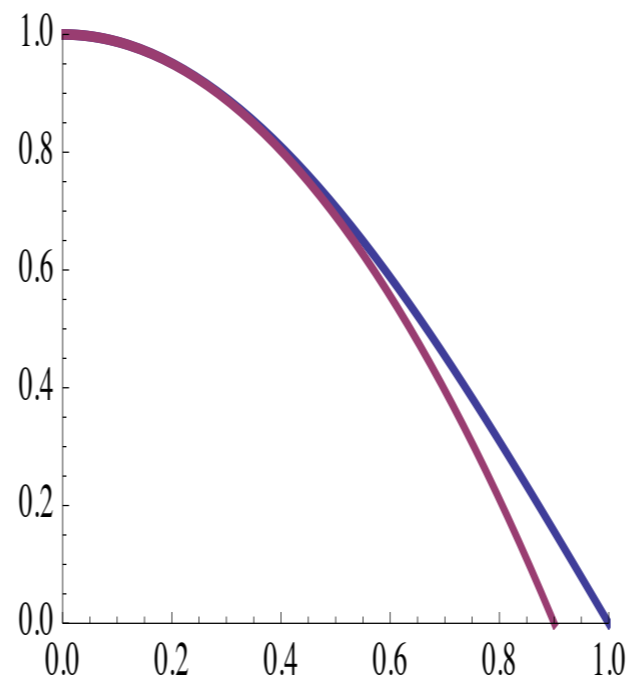


$$I = \int_0^1 dx (1 - cx^2) \frac{\cos \left(\frac{\pi}{2} x \right)}{(1 - cx^2)}$$

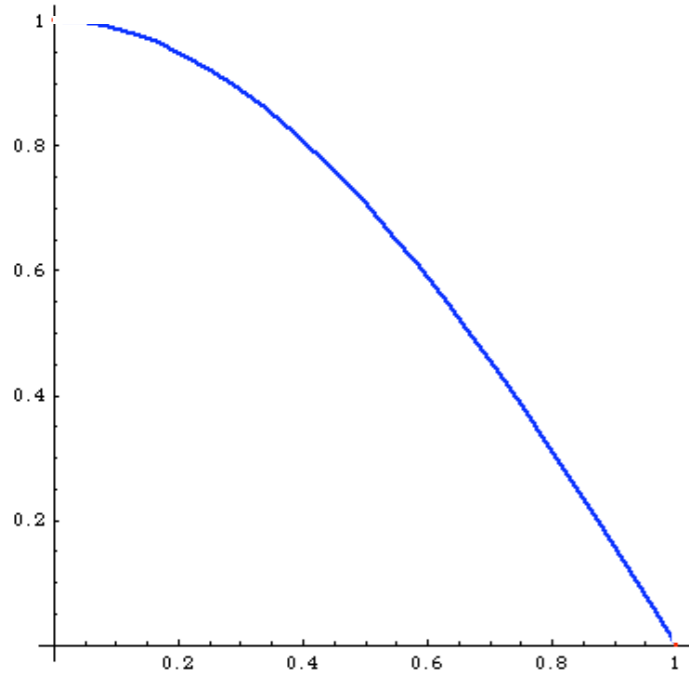


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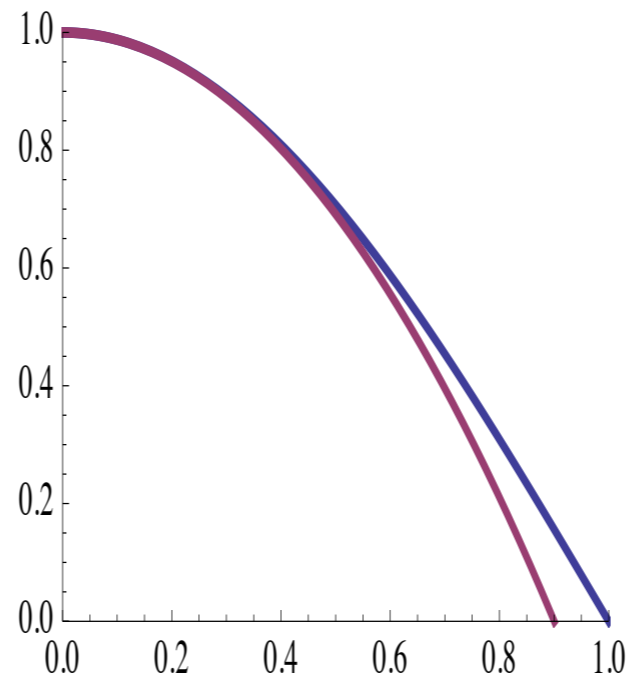


$$I = \int_0^1 dx (1 - cx^2) \frac{\cos\left(\frac{\pi}{2}x\right)}{(1 - cx^2)} = \int_{\xi_1}^{\xi_2} d\xi \frac{\cos \frac{\pi}{2} x[\xi]}{1 - x[\xi]^2 c}$$

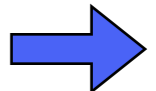


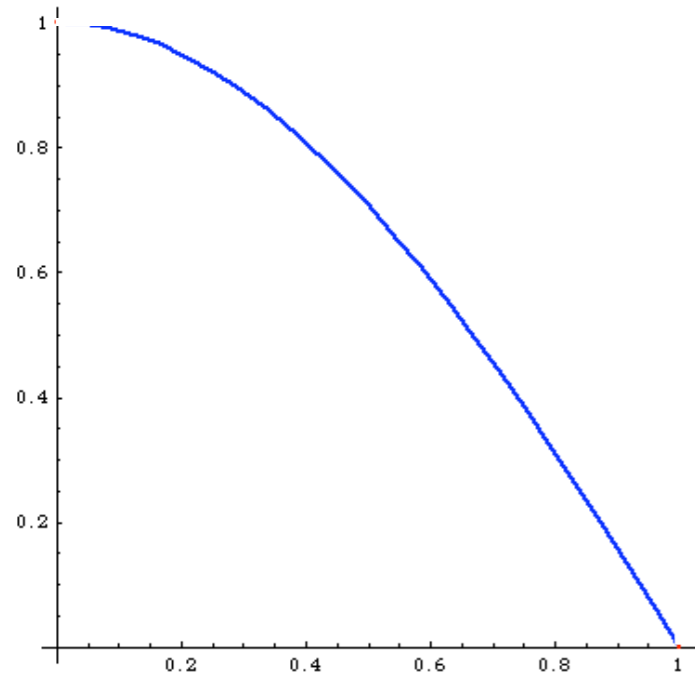
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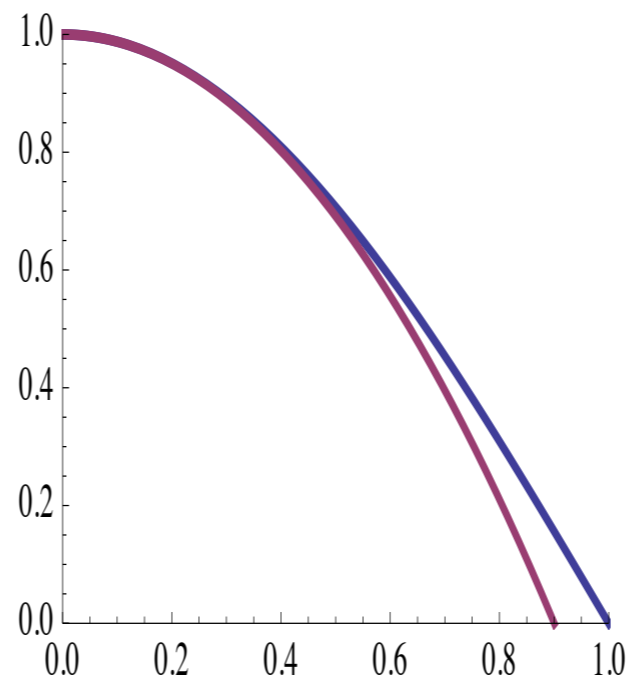
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 $\simeq 1$



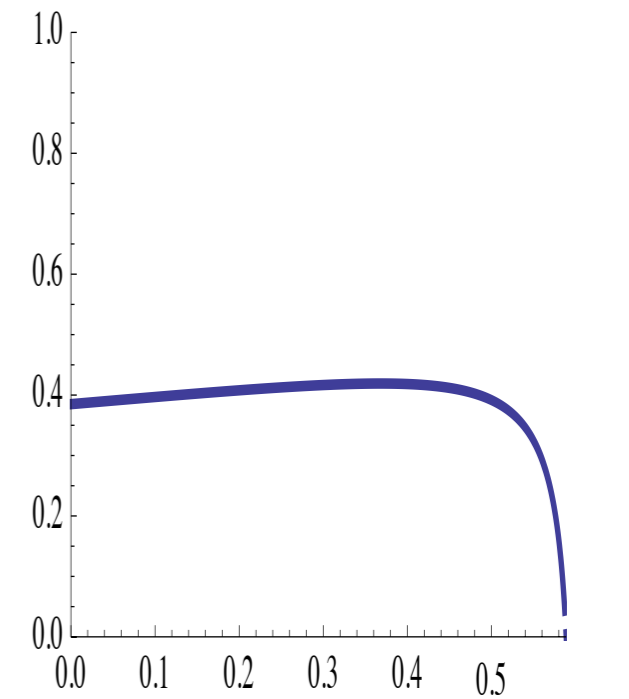
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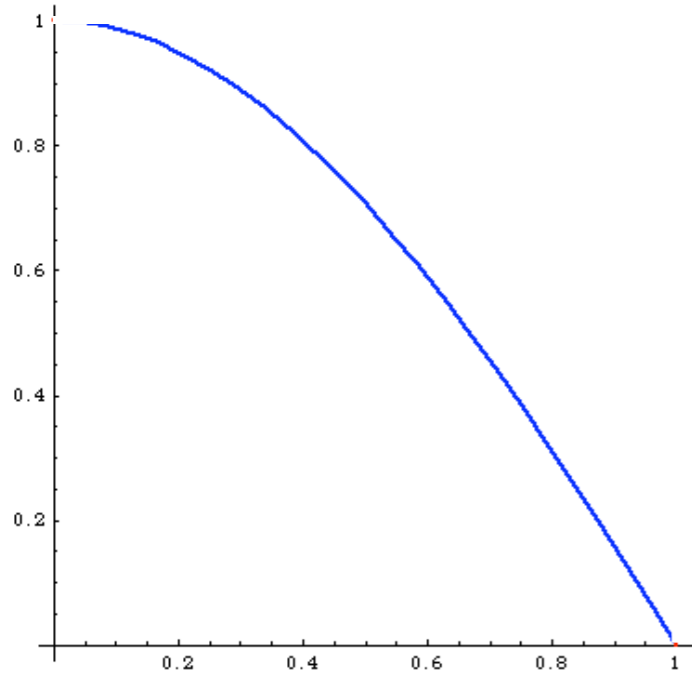
$$I_N = 0.637 \pm 0.307/\sqrt{N}$$



$$I = \int_0^1 dx (1 - cx^2) \frac{\cos(\frac{\pi}{2} x)}{(1 - cx^2)} = \int_{\xi_1}^{\xi_2} d\xi \frac{\cos \frac{\pi}{2} x[\xi]}{1 - x[\xi]^2 c}$$

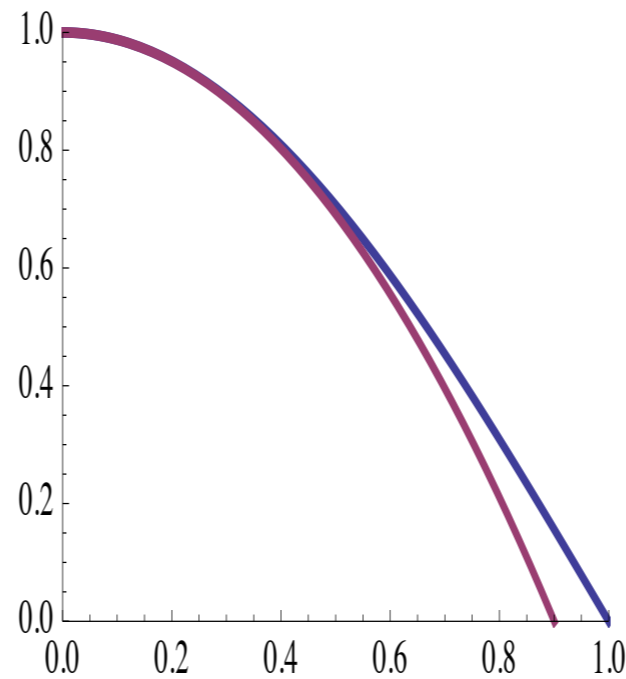
$$\rightarrow \simeq 1$$





$$I = \int_0^1 dx \cos \frac{\pi}{2} x$$

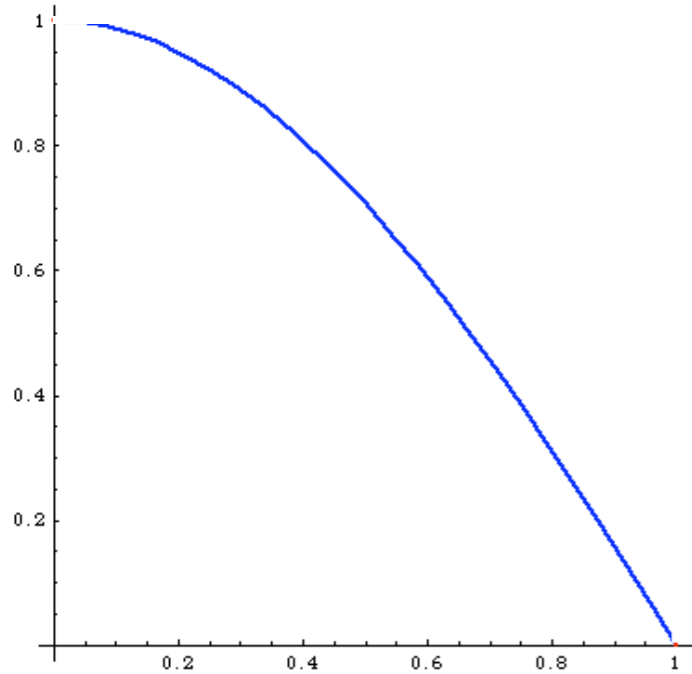
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$$I = \int_0^1 dx (1 - cx^2) \frac{\cos(\frac{\pi}{2} x)}{(1 - cx^2)} = \int_{\xi_1}^{\xi_2} d\xi \frac{\cos \frac{\pi}{2} x[\xi]}{1 - x[\xi]^2 c}$$

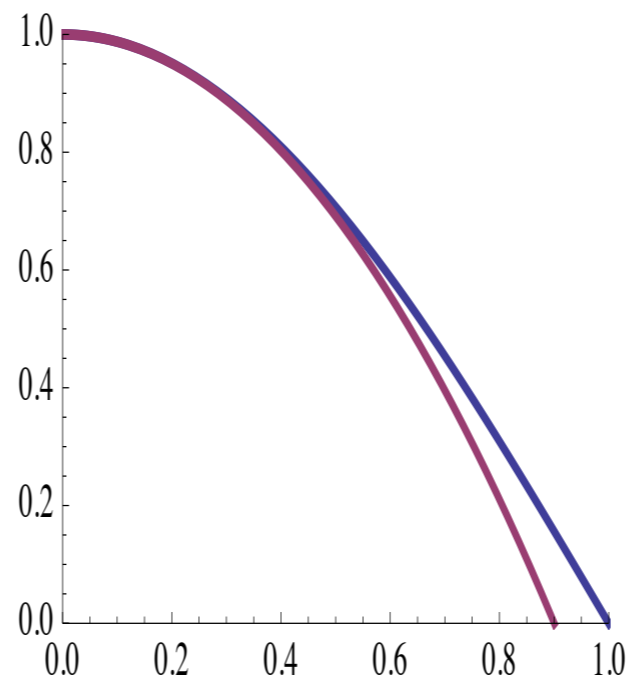
$$\rightarrow \simeq 1$$

$$I_N = 0.637 \pm 0.031/\sqrt{N}$$



$$I = \int_0^1 dx \cos \frac{\pi}{2} x$$

$$I_N = 0.637 \pm 0.307/\sqrt{N}$$



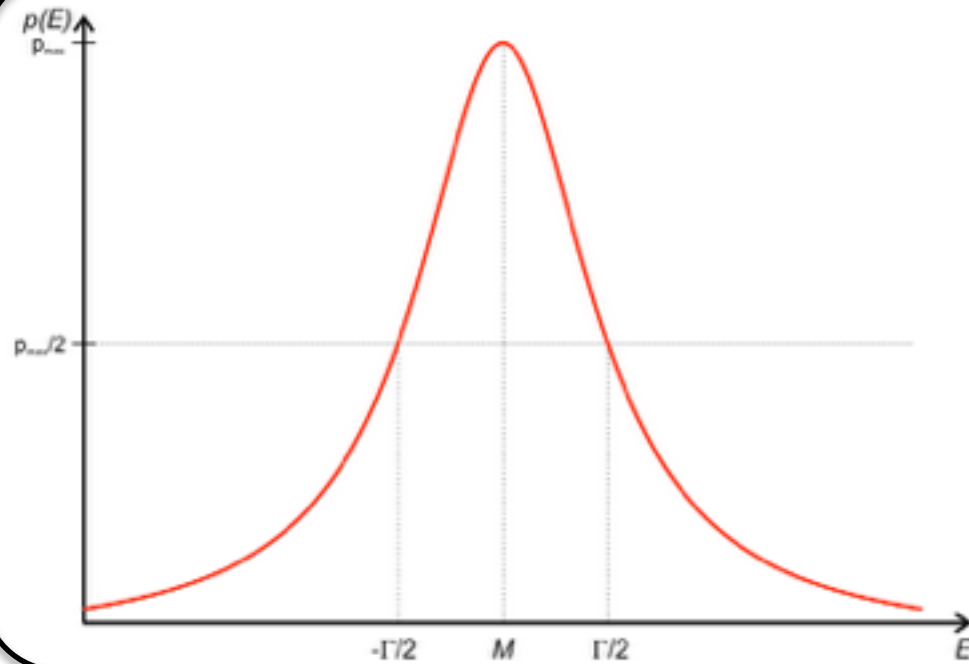
$$I = \int_0^1 dx (1 - cx^2) \frac{\cos(\frac{\pi}{2} x)}{(1 - cx^2)} = \int_{\xi_1}^{\xi_2} d\xi \frac{\cos \frac{\pi}{2} x[\xi]}{1 - x[\xi]^2 c}$$

$$\rightarrow \simeq 1$$

$$I_N = 0.637 \pm 0.031/\sqrt{N}$$

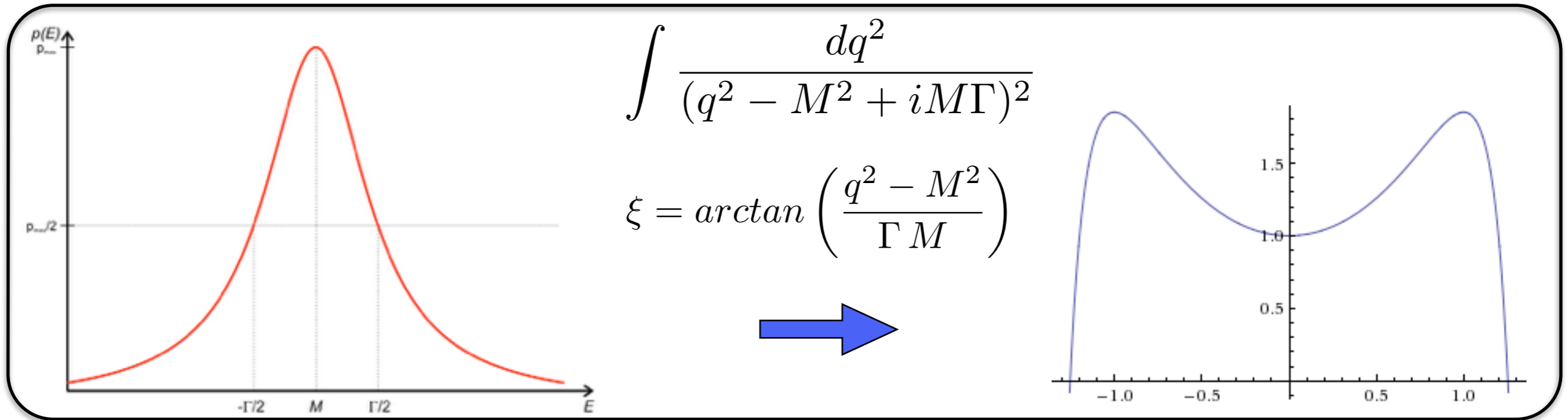
The Phase-Space parametrization is important to have an efficient computation!

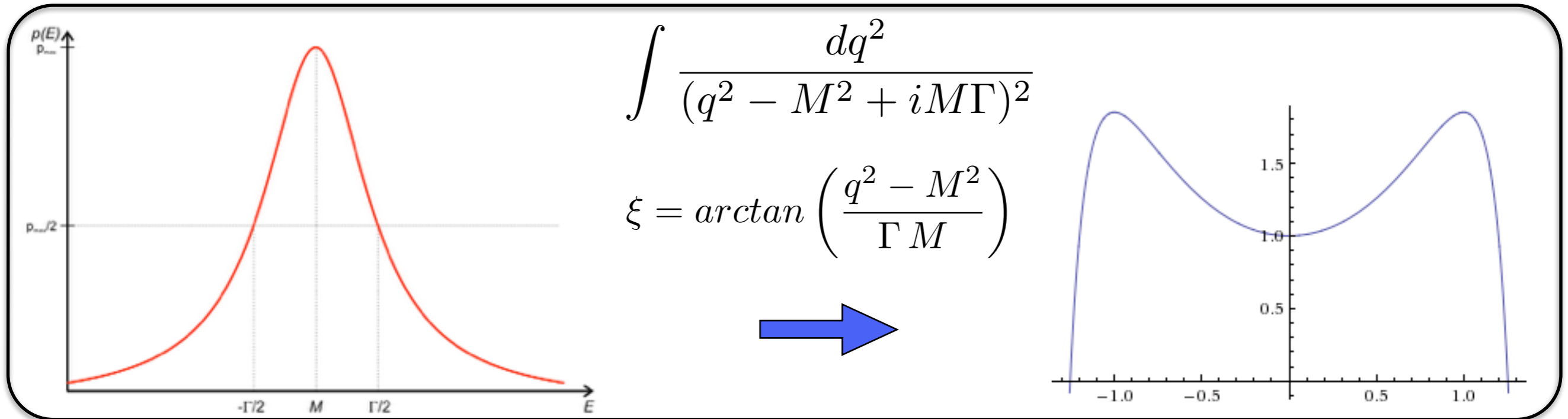




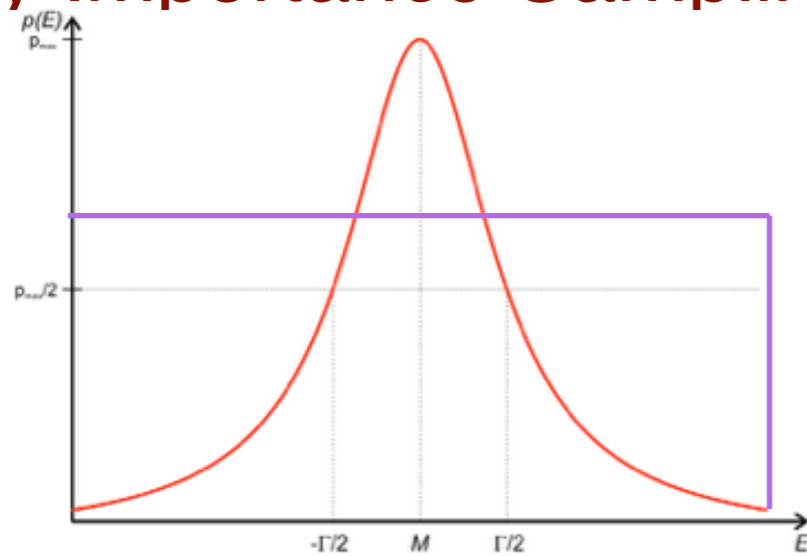
$$\int \frac{dq^2}{(q^2 - M^2 + iM\Gamma)^2}$$

$$\xi = \arctan\left(\frac{q^2 - M^2}{\Gamma M}\right)$$

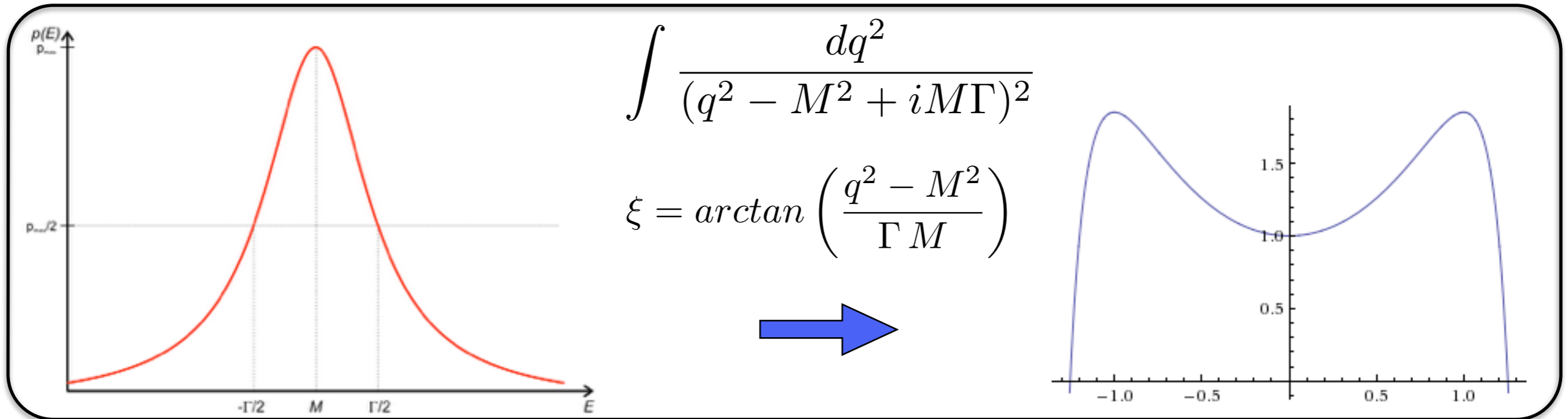




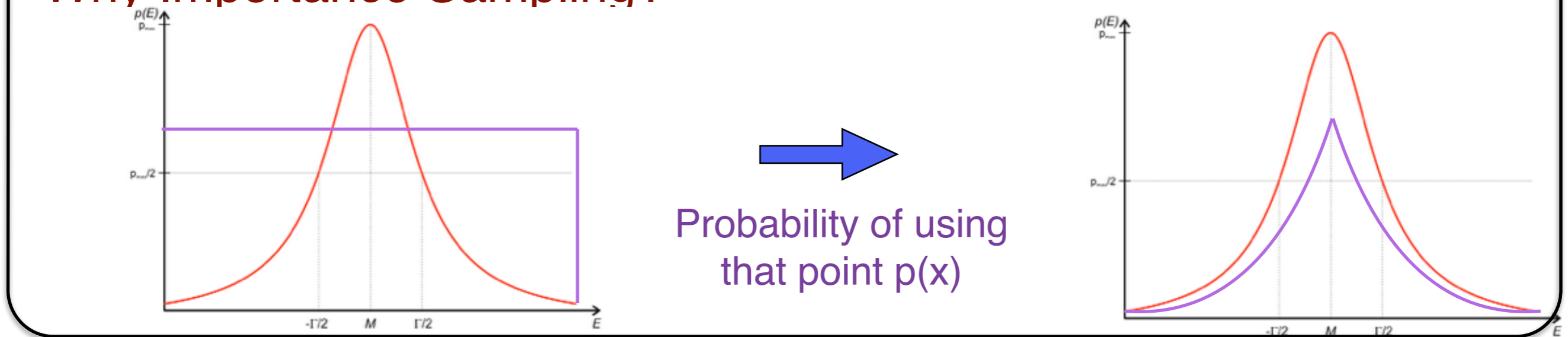
Why Importance Sampling?

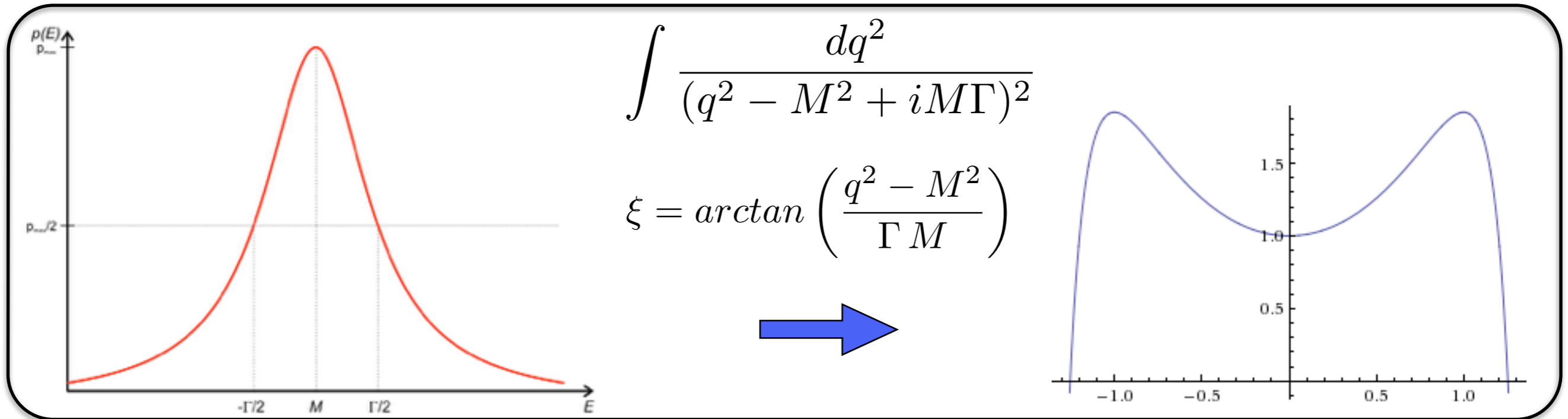


Probability of using that point $p(x)$

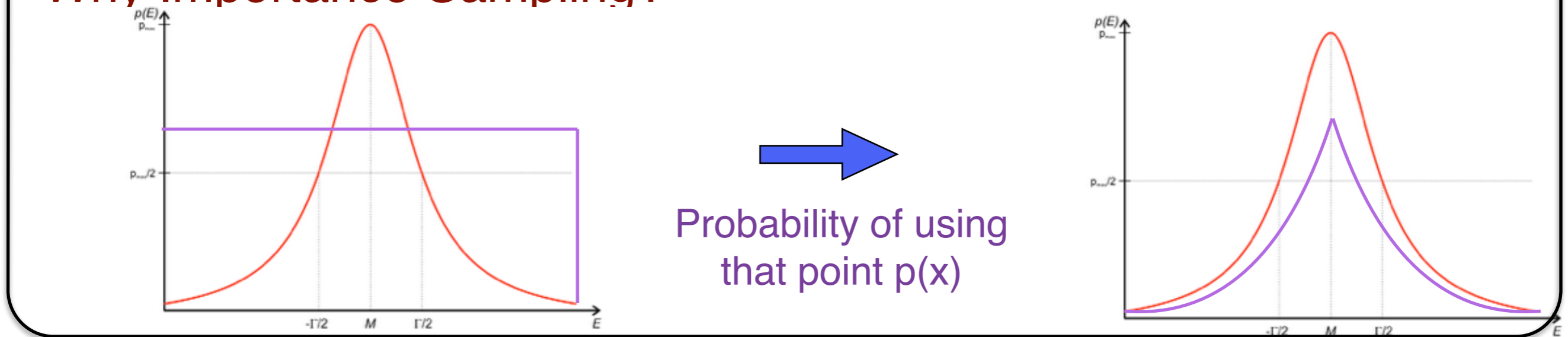


Why Importance Sampling?





Why Importance Sampling?



The change of variable ensure that the evaluation of the function is done where the function is the largest!

Key Point

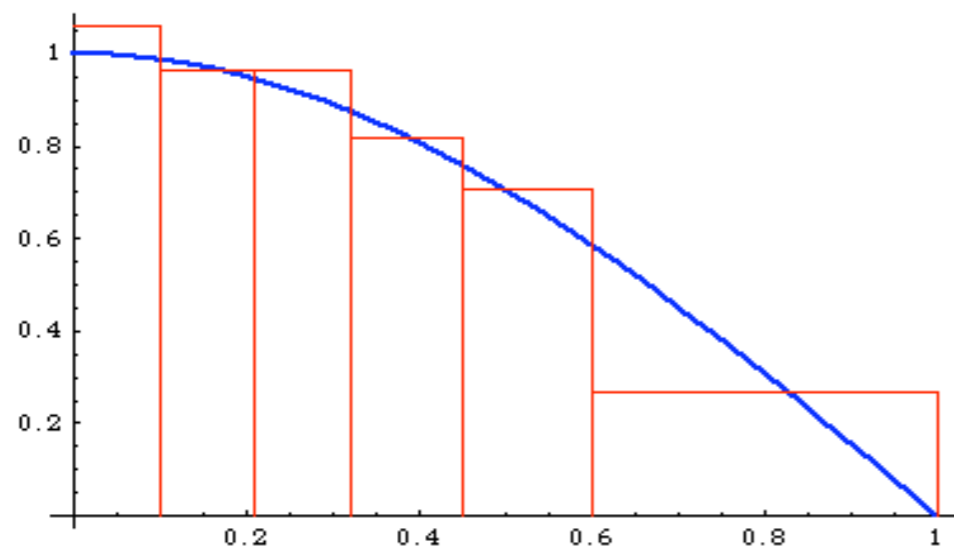
- Generate the **random point** in a distribution which is close to the function to integrate.
- This is a change of variable, such that the function is **flatter** in this new variable.
- Needs to know an **approximate function**.

Adaptative Monte-Carlo

- Create an approximation of the function on the flight!

Adaptative Monte-Carlo

- Create an approximation of the function on the flight!



Algorithm

1. Creates bin such that each of them have the same contribution.
 - ➔ Many bins where the function is large
2. Use the approximate for the importance sampling method.

More than one Dimension

- VEGAS works only with 1 (few) dimension
 - ➔ memory problem

More than one Dimension

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Solution

- Use projection on the axis

$$\vec{p}(x) = p(x) \cdot p(y) \cdot p(z) \dots$$

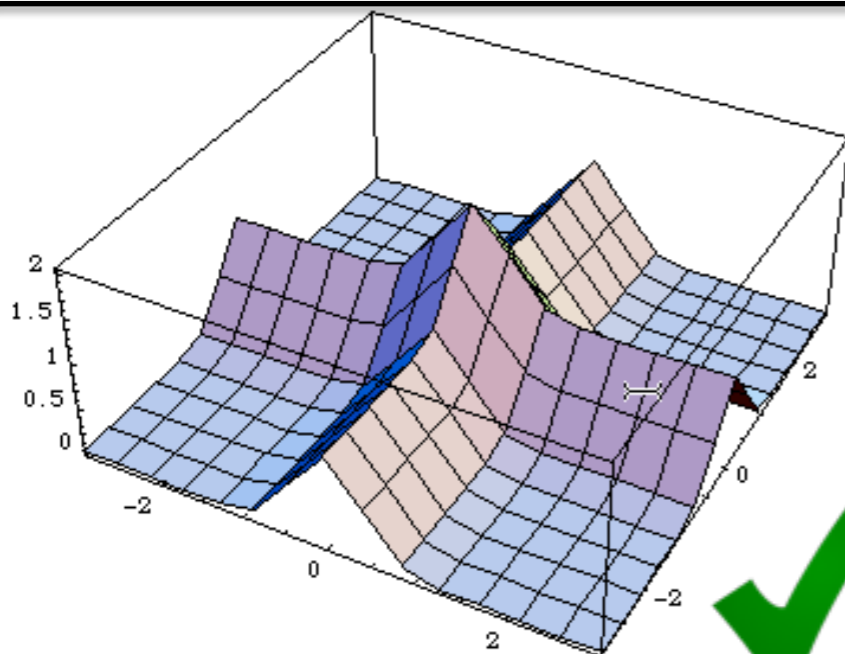
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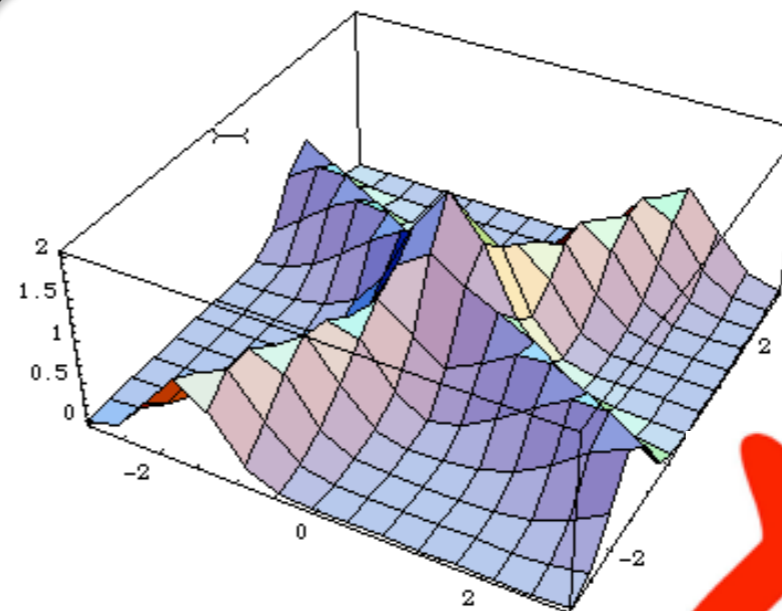
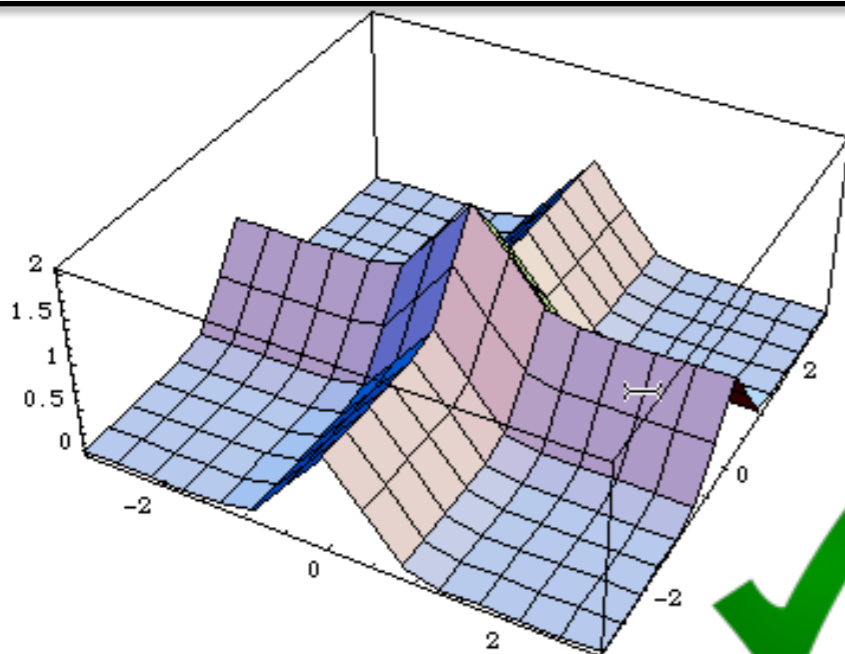
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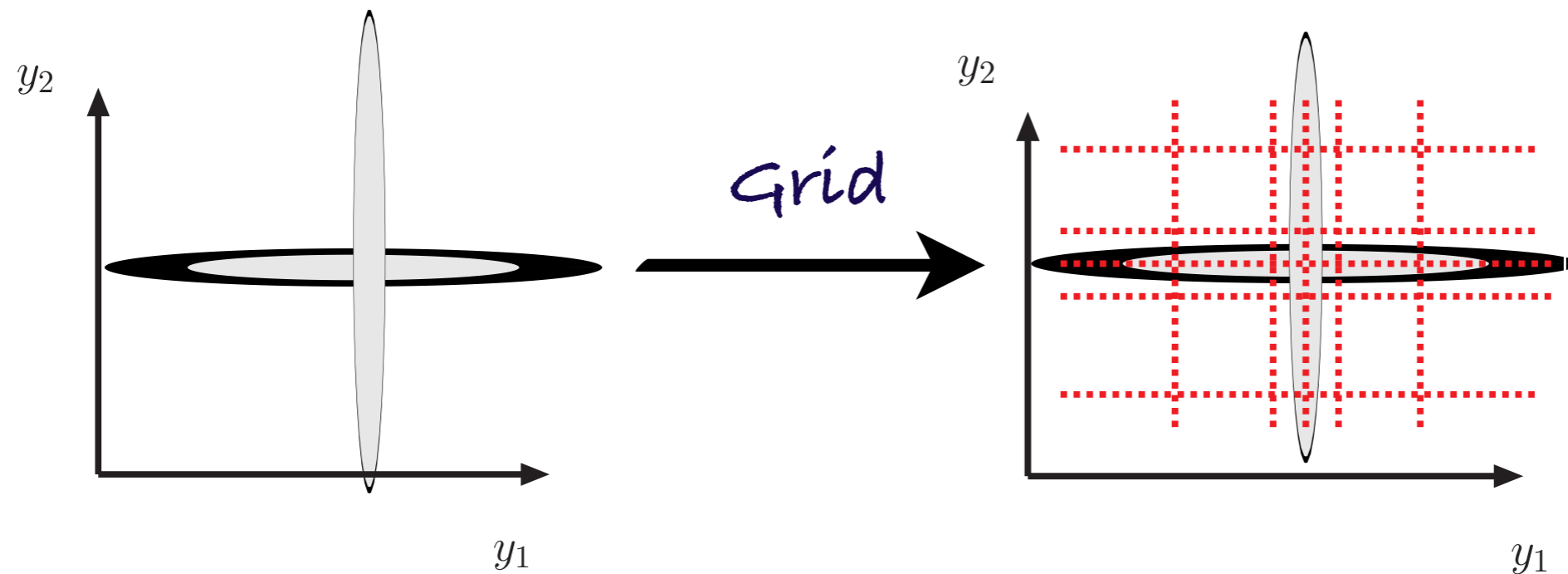
$$\vec{p}(x) = p(x) \cdot p(y) \cdot p(z) \dots$$



- We need to ensure the factorization !

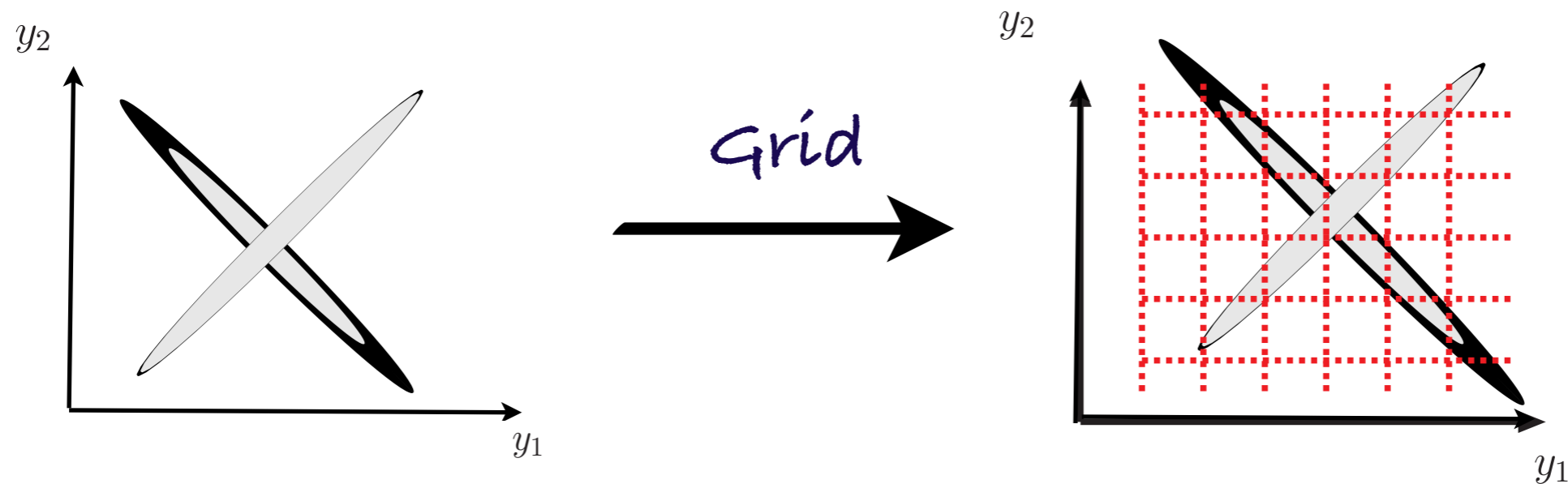
➔ Additional change of variable

- The choice of the parameterisation has a strong **impact** on the efficiency



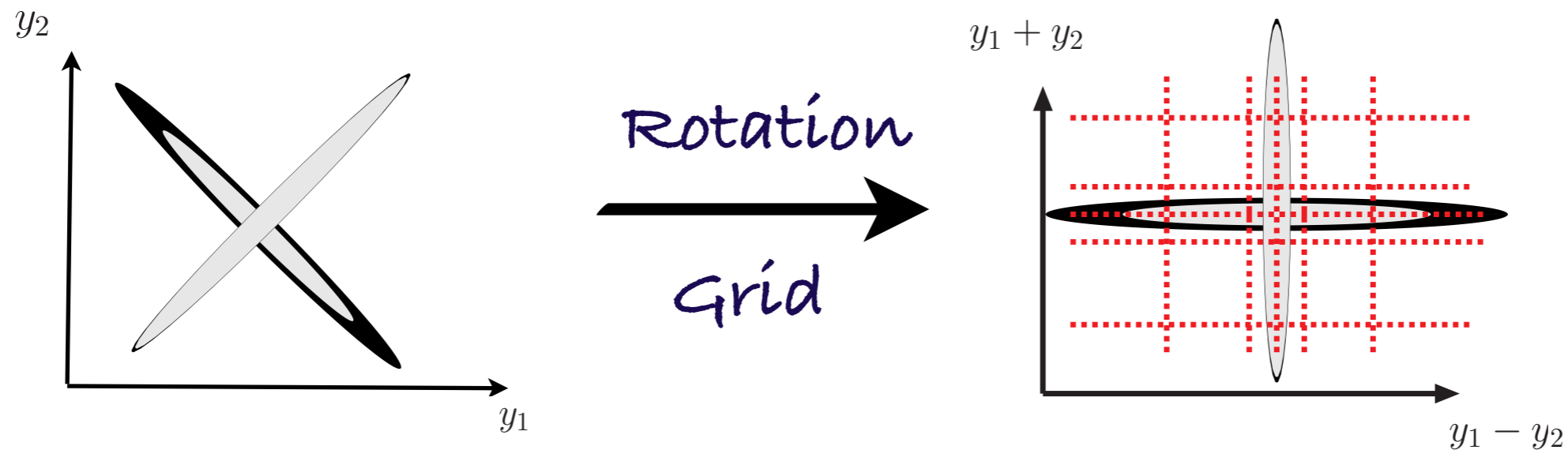
- The **adaptive** Monte-Carlo Technique picks point in interesting areas
→ The technique is **efficient**

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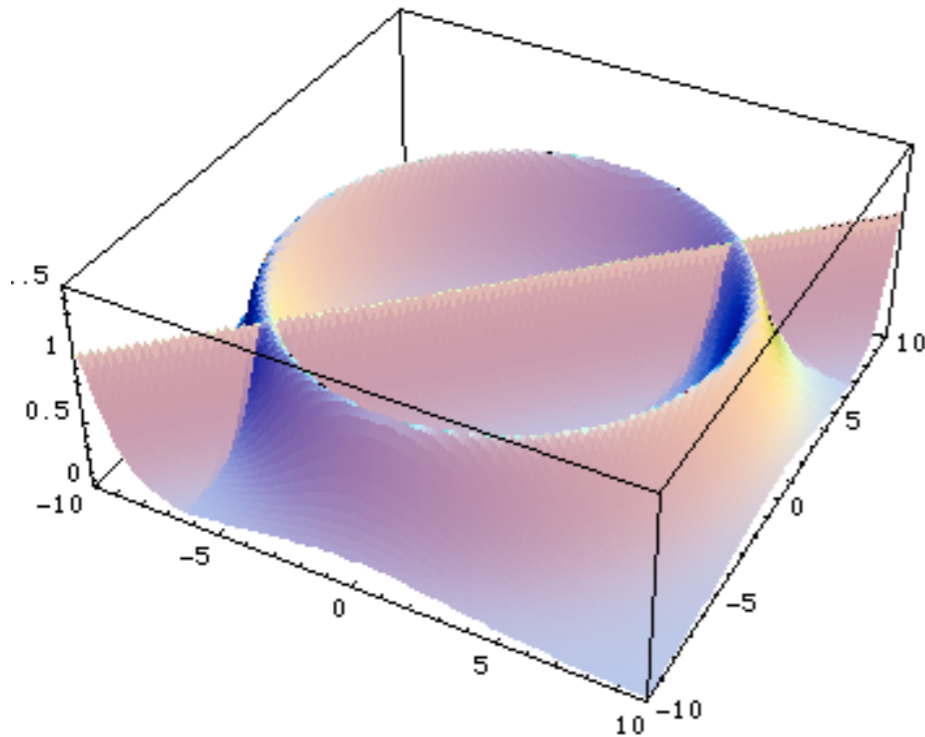


- The **adaptive** Monte-Carlo Techniques picks points everywhere
→ The integral converges **slowly**

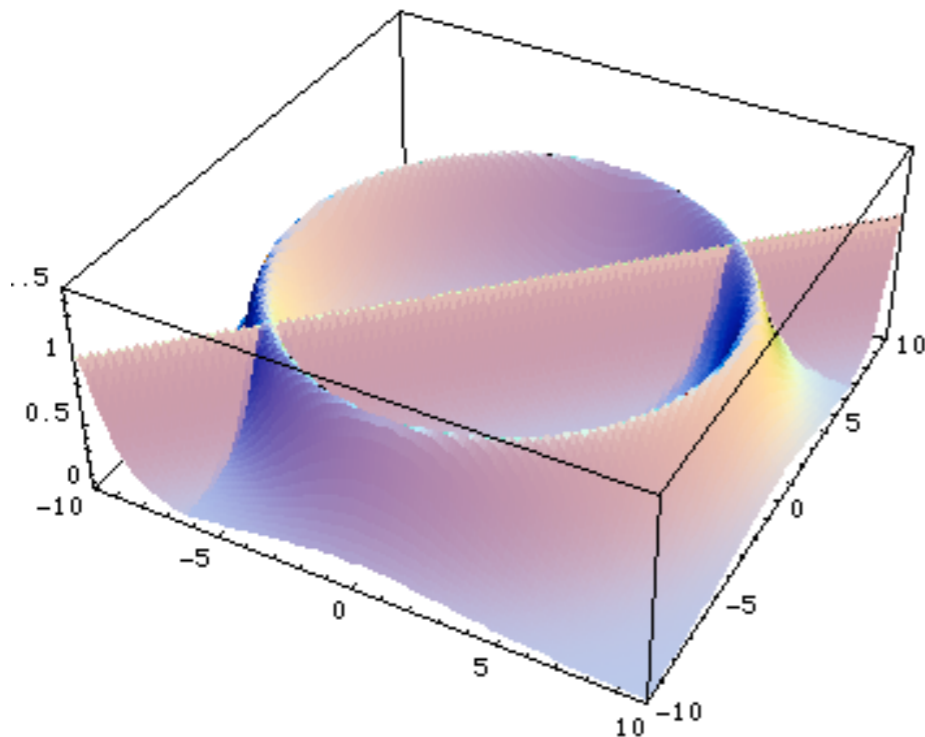
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What do we do if there is no transformation that aligns all integrand peaks to the chosen axes?
Vegas is bound to fail!

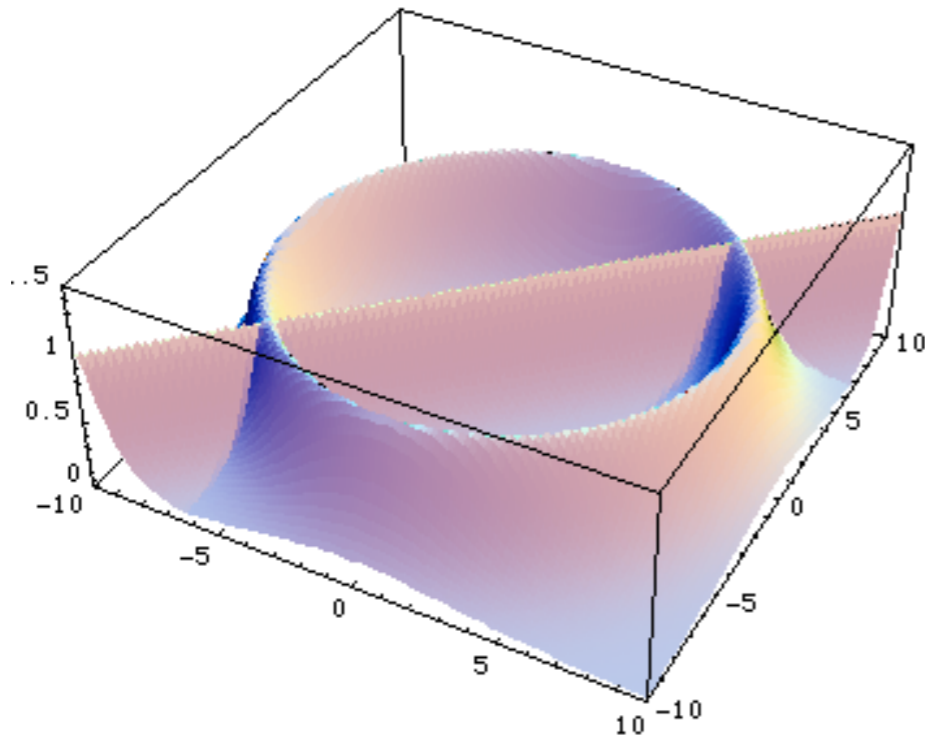


What do we do if there is no transformation that aligns all integrand peaks to the chosen axes?
Vegas is bound to fail!

Solution: use different transformations = channels

$$p(x) = \sum_{i=1}^n \alpha_i p_i(x) \quad \text{with} \quad \sum_{i=1}^n \alpha_i = 1$$

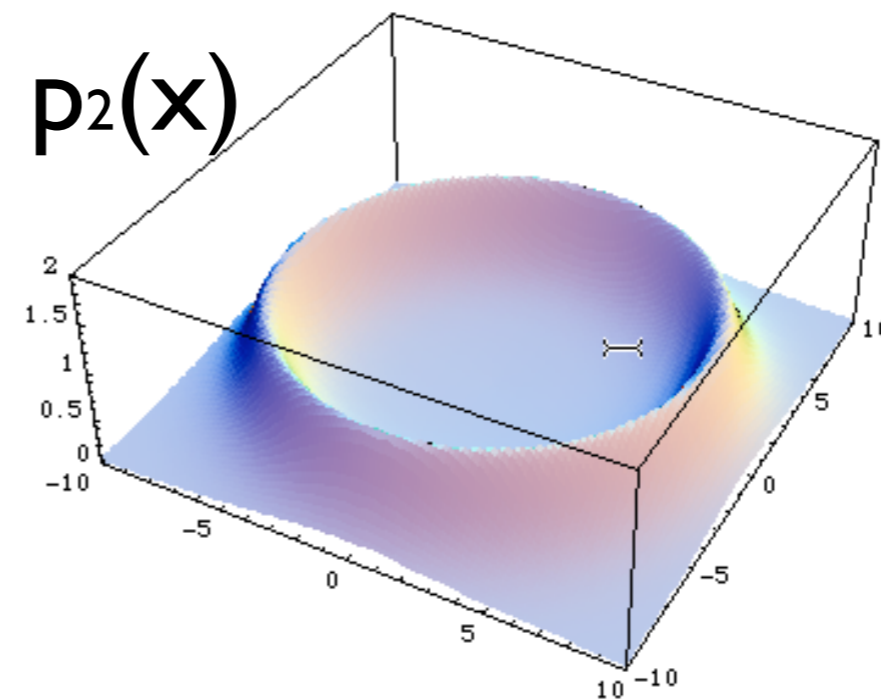
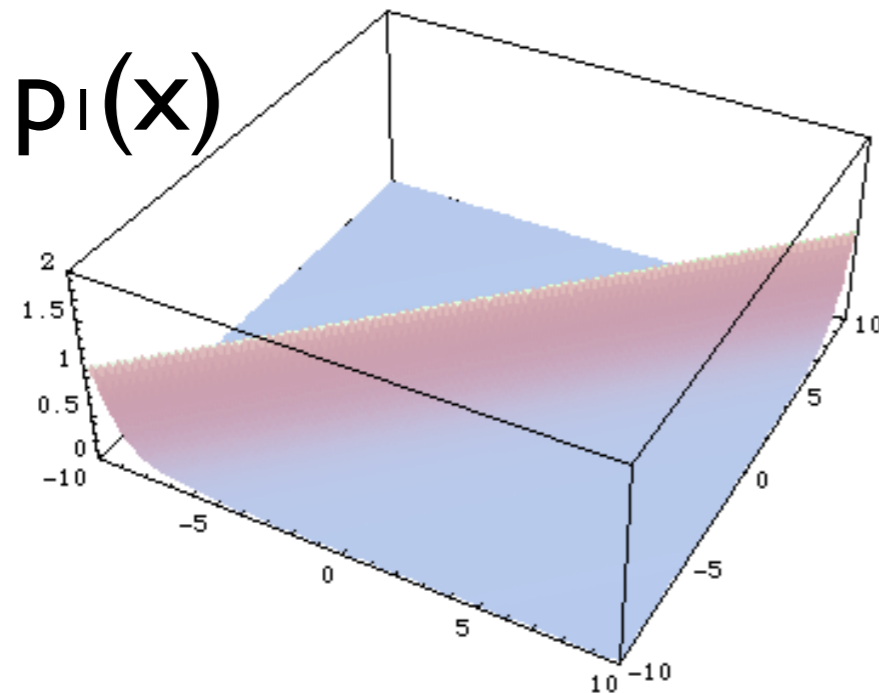
with each $p_i(x)$ taking care of one “peak” at the time

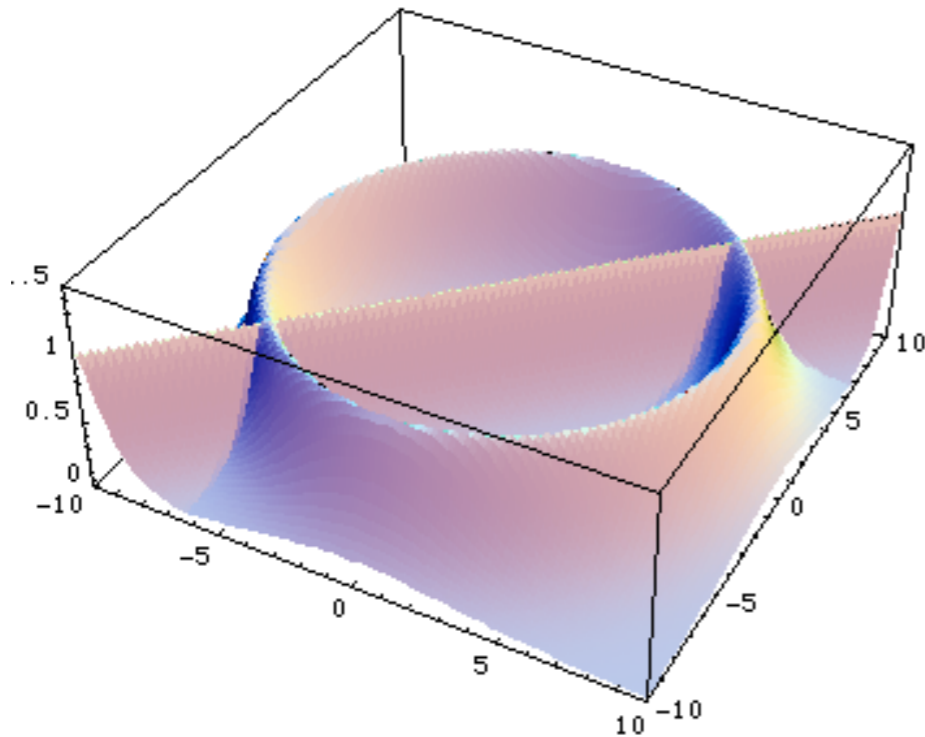


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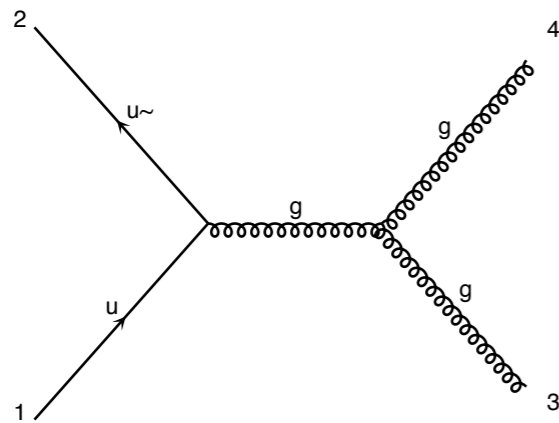
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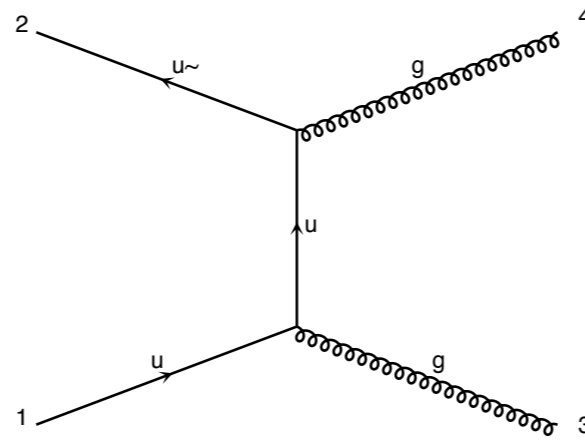
Then,

$$I = \int f(x) dx = \sum_{i=1}^n \alpha_i \int \frac{f(x)}{p(x)} p_i(x) dx$$

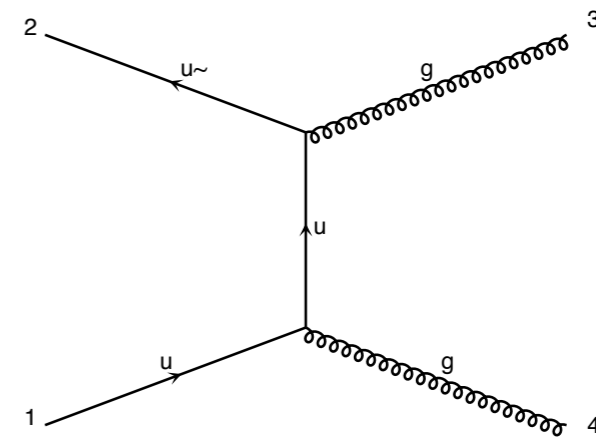
≈ 1



$$\propto \frac{1}{\hat{s}} = \frac{1}{(p_1 + p_2)^2}$$



$$\propto \frac{1}{\hat{t}} = \frac{1}{(p_1 - p_3)^2}$$



$$\propto \frac{1}{\hat{u}} = \frac{1}{(p_1 - p_4)^2}$$

Three very different pole structures contributing to the same matrix element.

Does a basis exist?

$$\int |M_{tot}|^2 = \int \frac{\sum_i |M_i|^2}{\sum_j |M_j|^2} |M_{tot}|^2 = \sum_i \int \frac{|M_i|^2}{\sum_j |M_j|^2} |M_{tot}|^2$$

Does a basis exist?

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Key Idea

- Any single diagram is “easy” to integrate (pole structures/ suitable integration variables known from the propagators)
- Divide integration into pieces, based on diagrams
- All other peaks taken care of by denominator sum

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N Integral

- Errors add in quadrature so no extra cost
- “Weight” functions already calculated during $|M|^2$ calculation
- Parallel in nature

- Phase-Space integration are difficult
- We need to know the function
 - ➔ Be careful with cut
- MadGraph split the integral in different contribution linked to the Feynman Diagram
 - ➔ Those are not the contribution of a given diagram

[P0_gg_hqq](#)

$s = 0.44288 \pm 0.00268$ (pb)

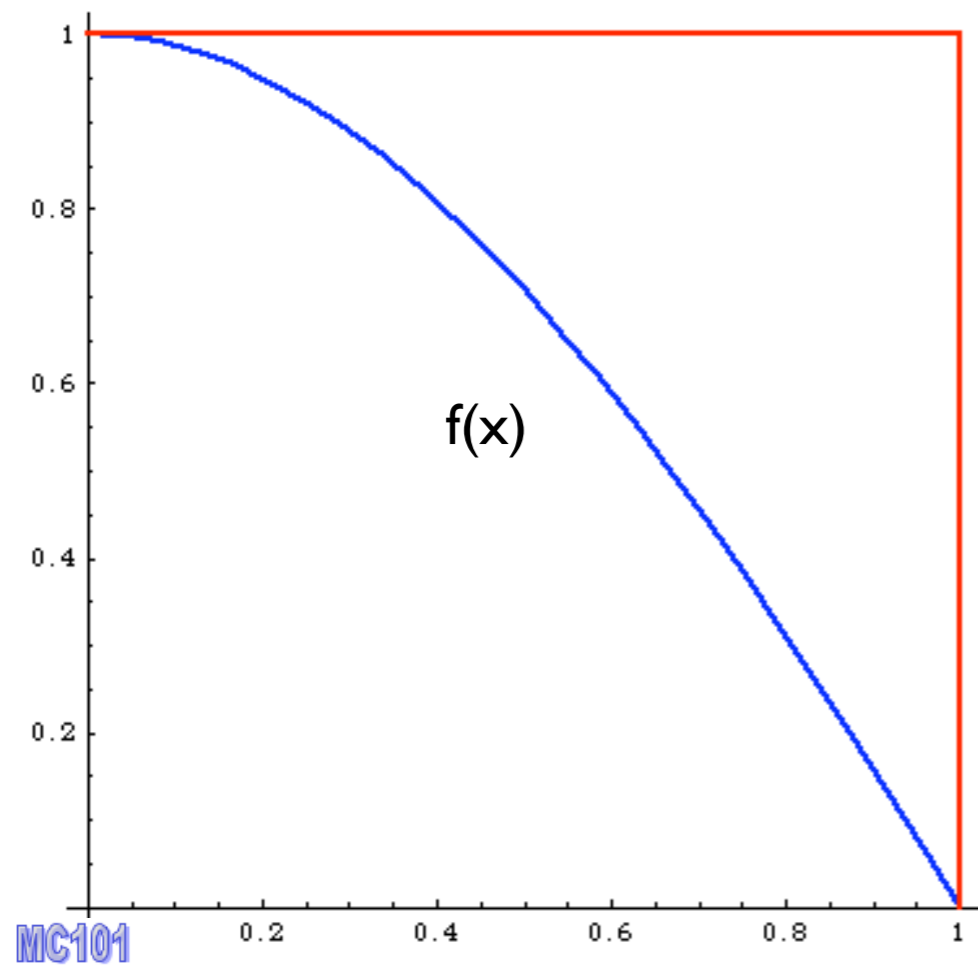
Graph	Cross-Section ↓	Error	Events (K)	Unwgt	Luminosity
G7	0.1263	0.00102	8.002	256.0	2.03e+03
G6	0.1225	0.00132	16.002	760.0	6.21e+03
G2	0.08464	0.0011	32.002	1931.0	2.28e+04
G4	0.08122	0.00169	32.002	101.0	1.25e+03
G1	0.02821	0.000563	8.002	144.0	5.1e+03

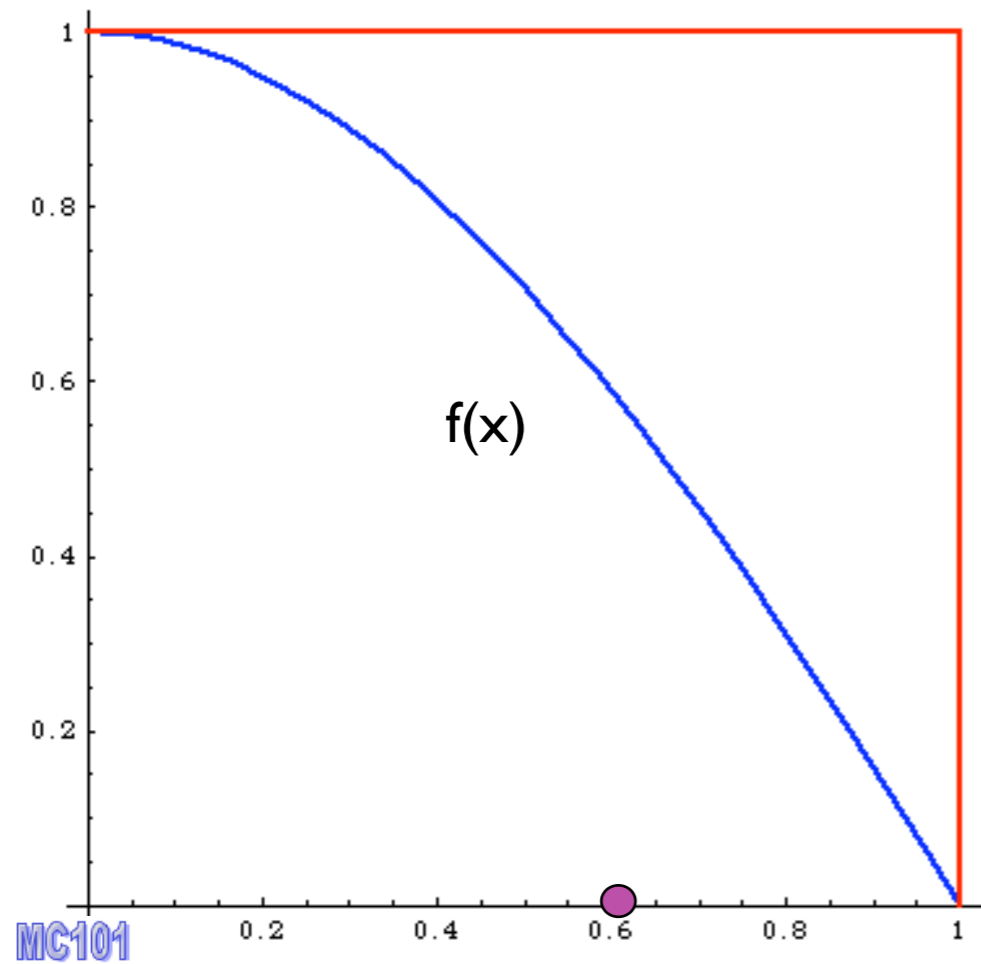
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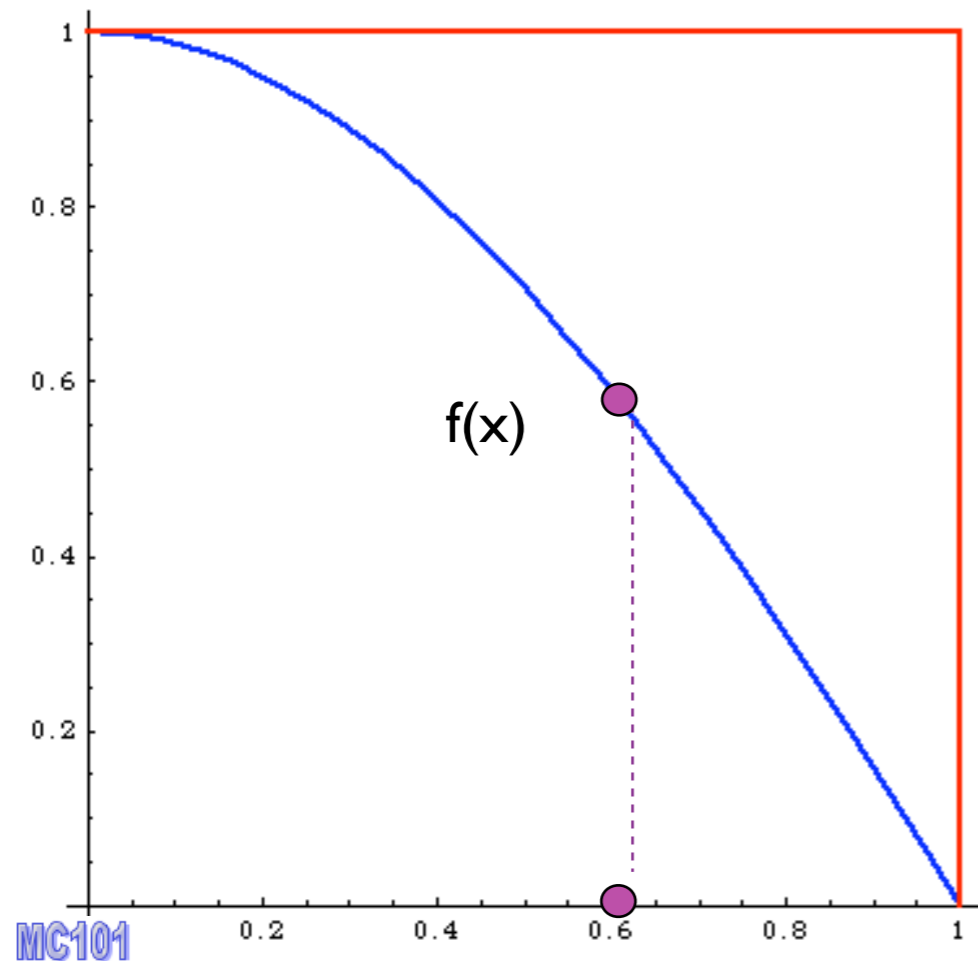
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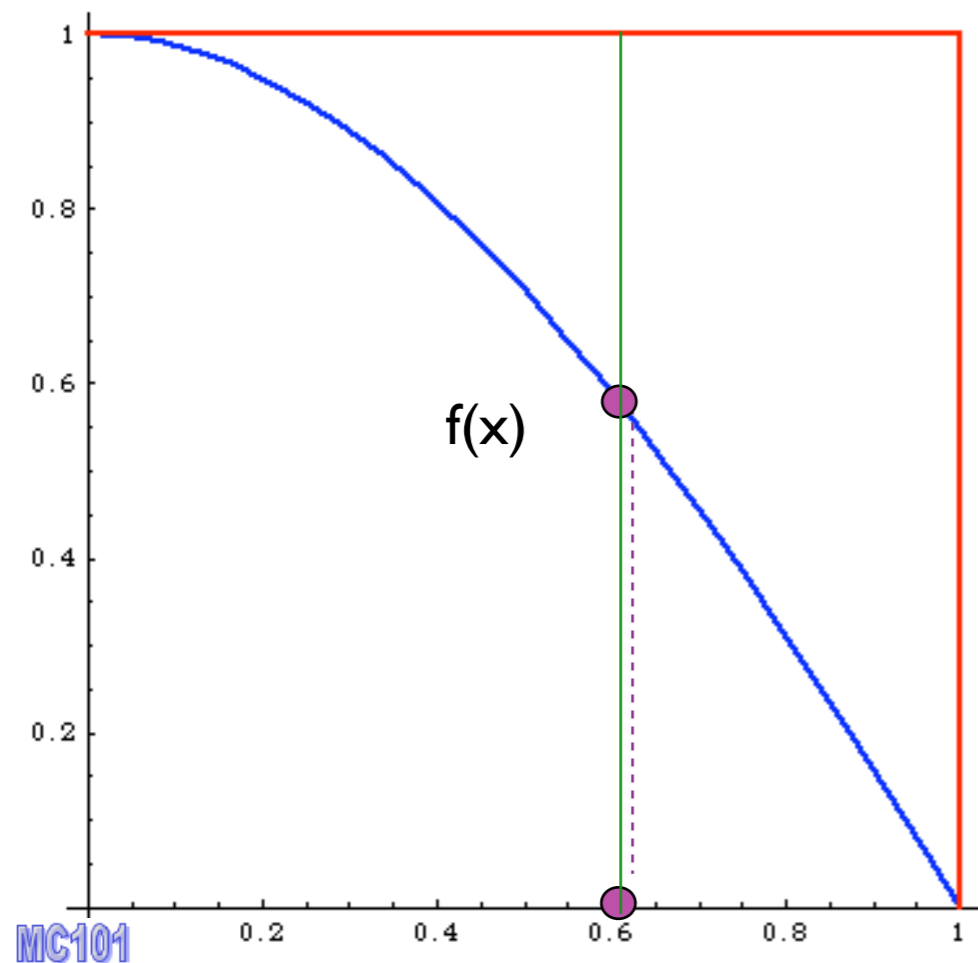




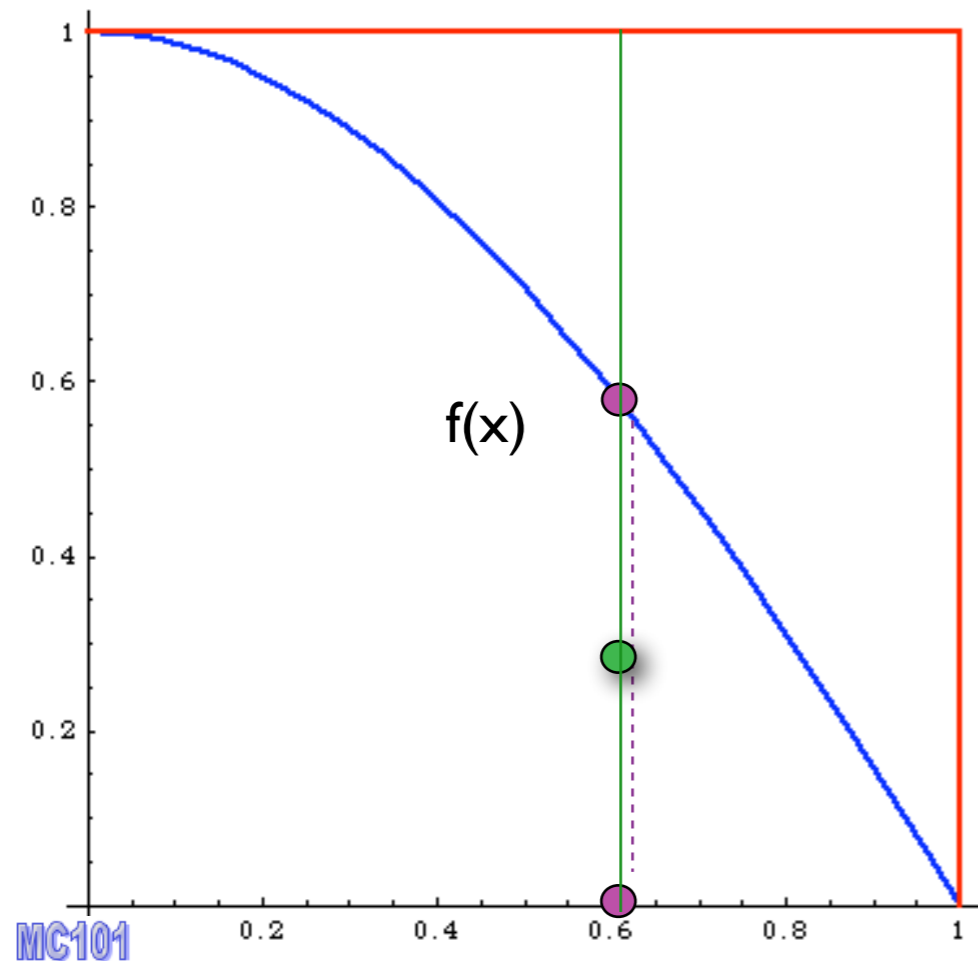
I. pick x



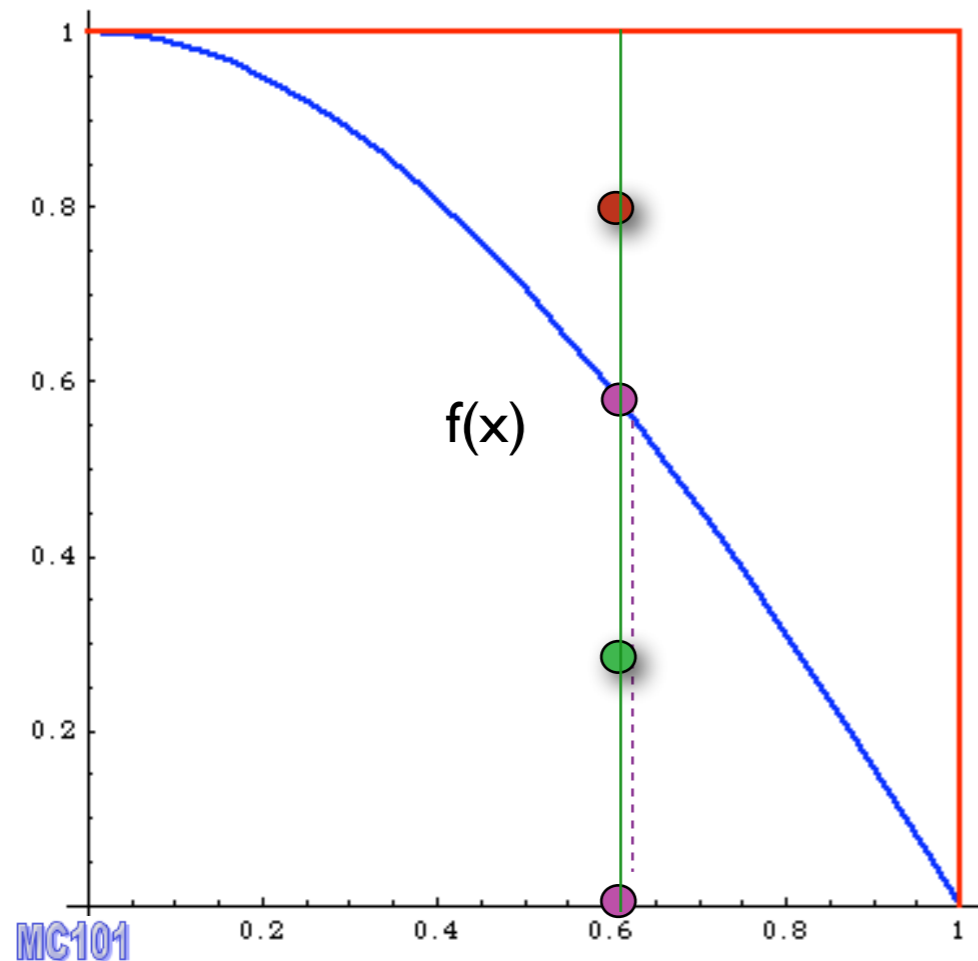
1. pick x
2. calculate $f(x)$



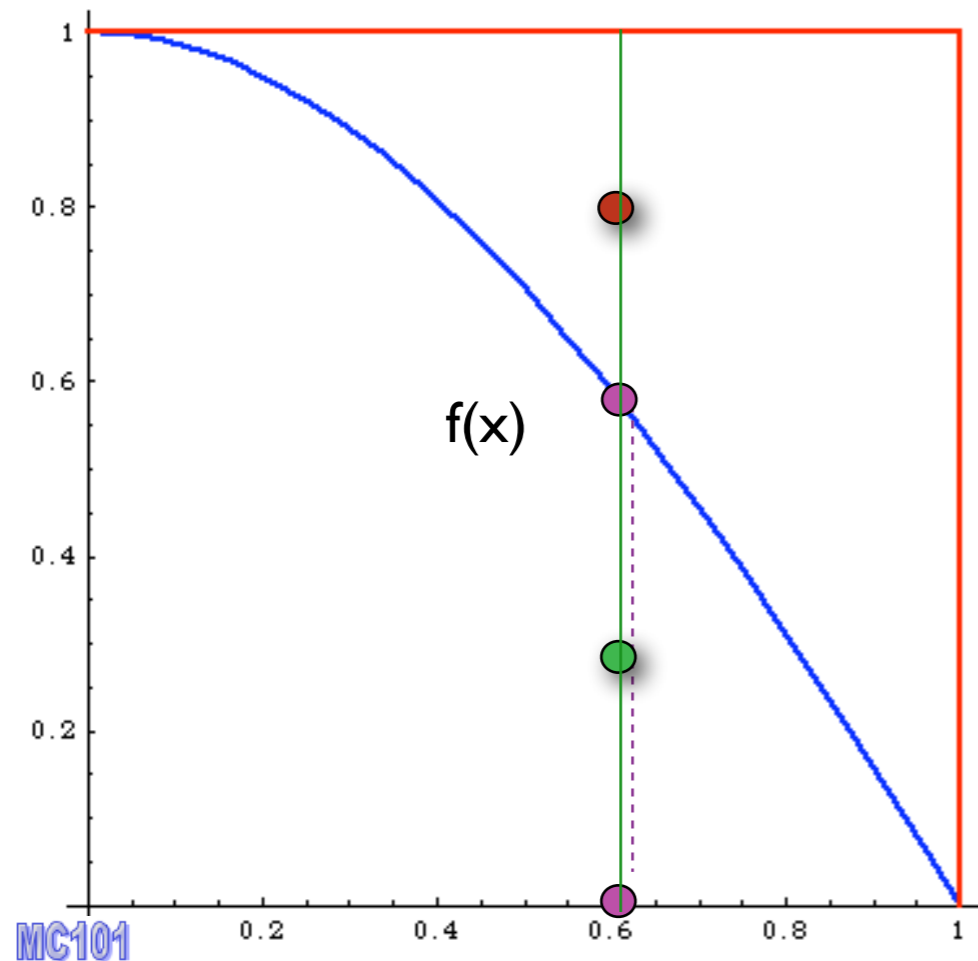
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if $f(x) > y$ accept event,



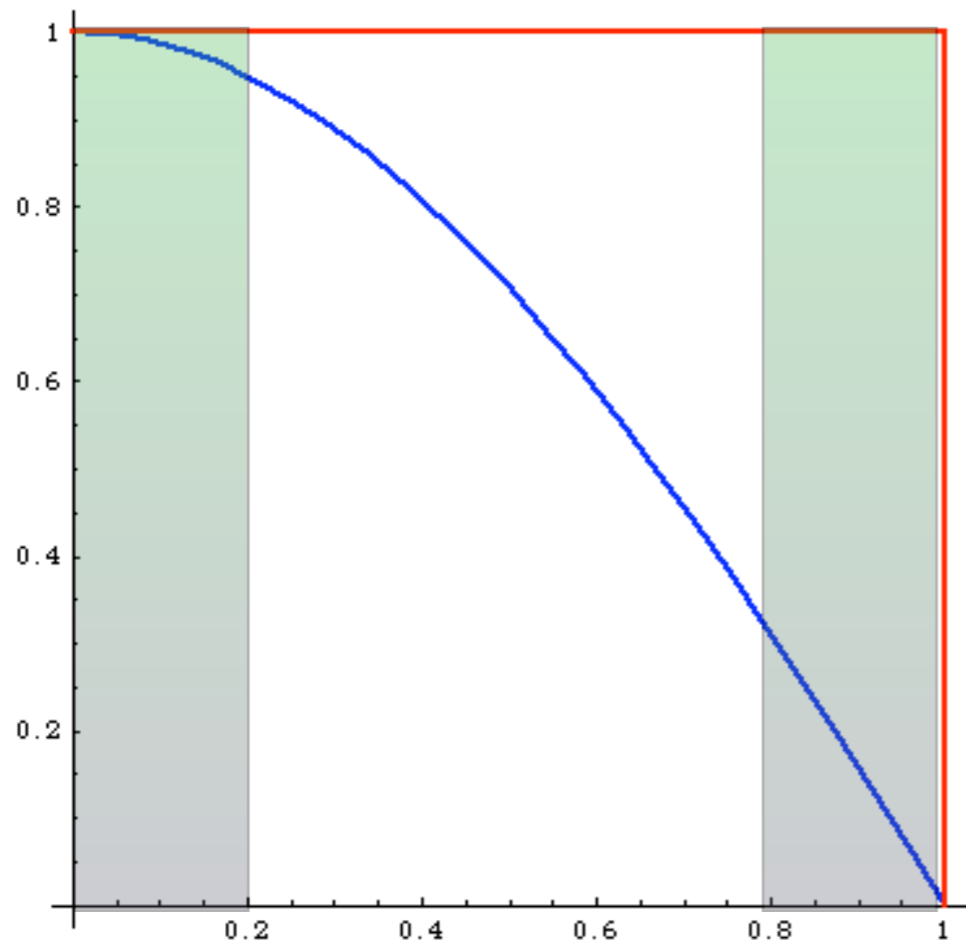
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if $f(x) > y$ accept event,
else reject it.

$$I = \frac{\text{accepted}}{\text{total tries}} = \text{efficiency}$$

Event generation

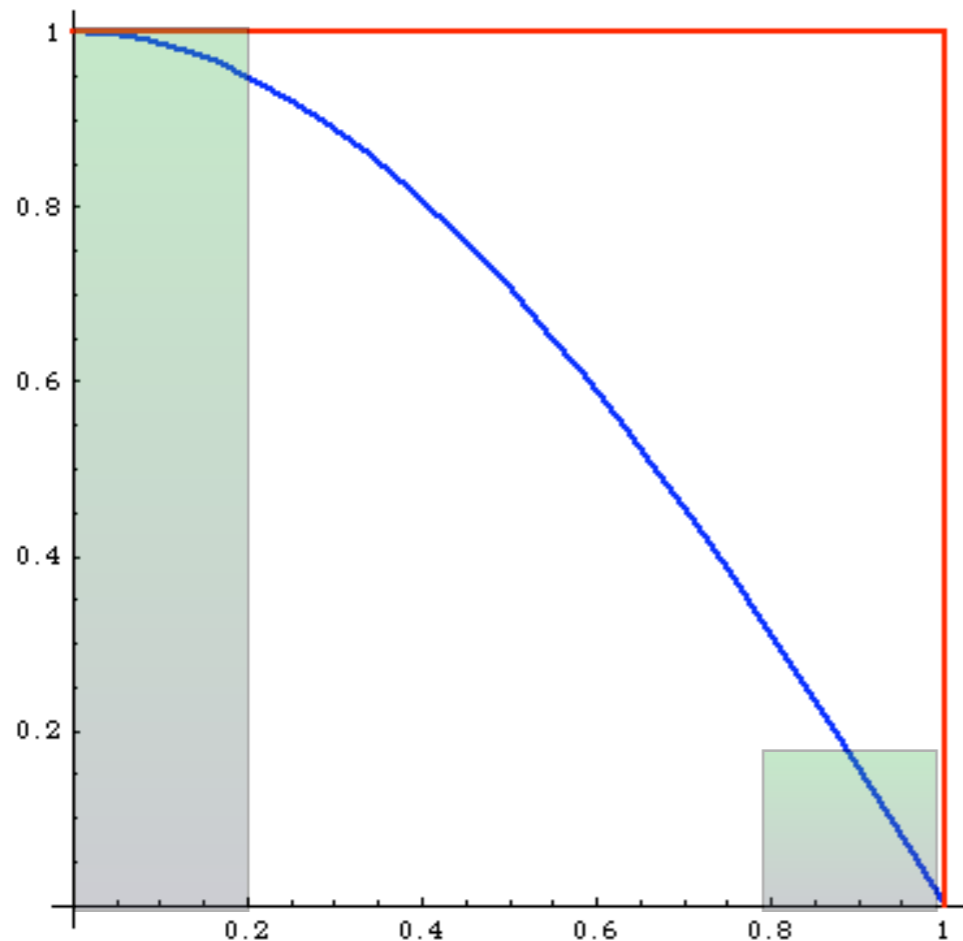


What's the difference between weighted and unweighted?

Weighted:

Same # of events in areas of phase space with very different probabilities: events must have different weights

Event generation



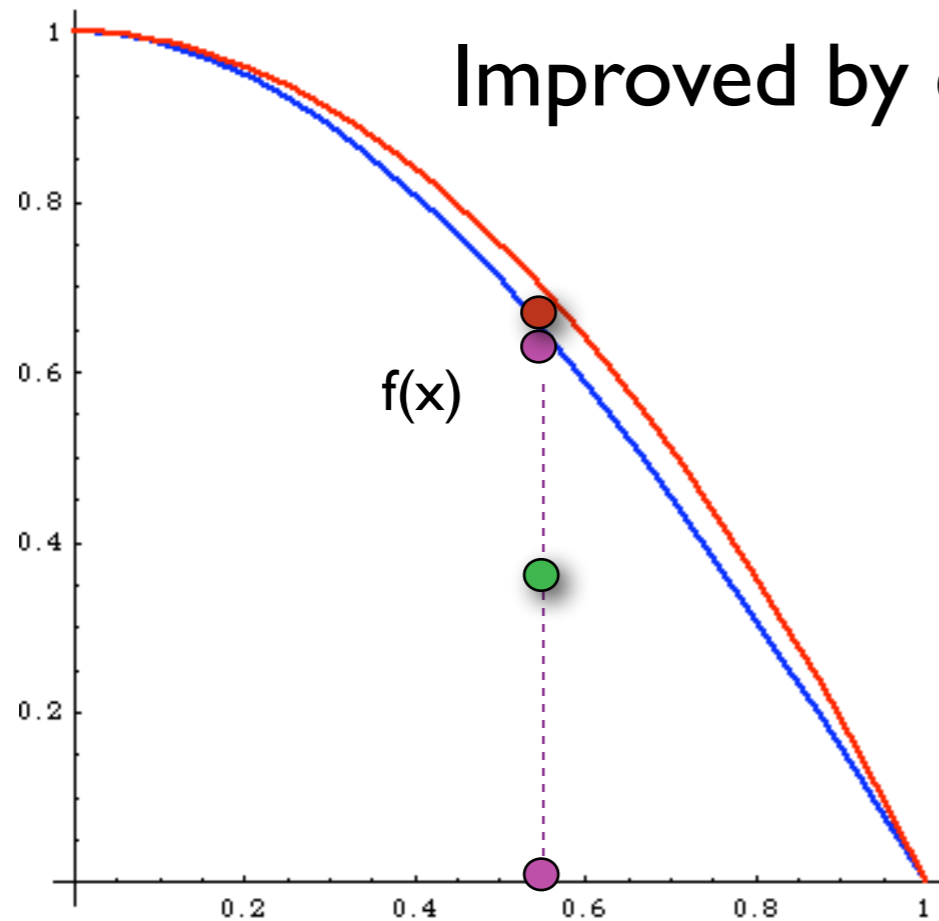
What's the difference between weighted and unweighted?

Unweighted:

events is proportional to the probability of areas of phase space:
events have all the same weight ("unweighted")

Events distributed as in nature

Event generation



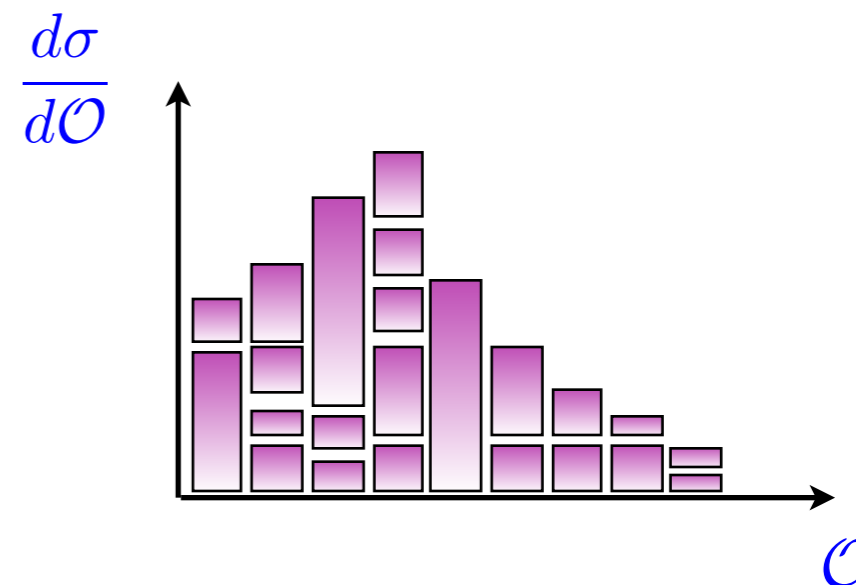
Improved by combining with importance sampling:

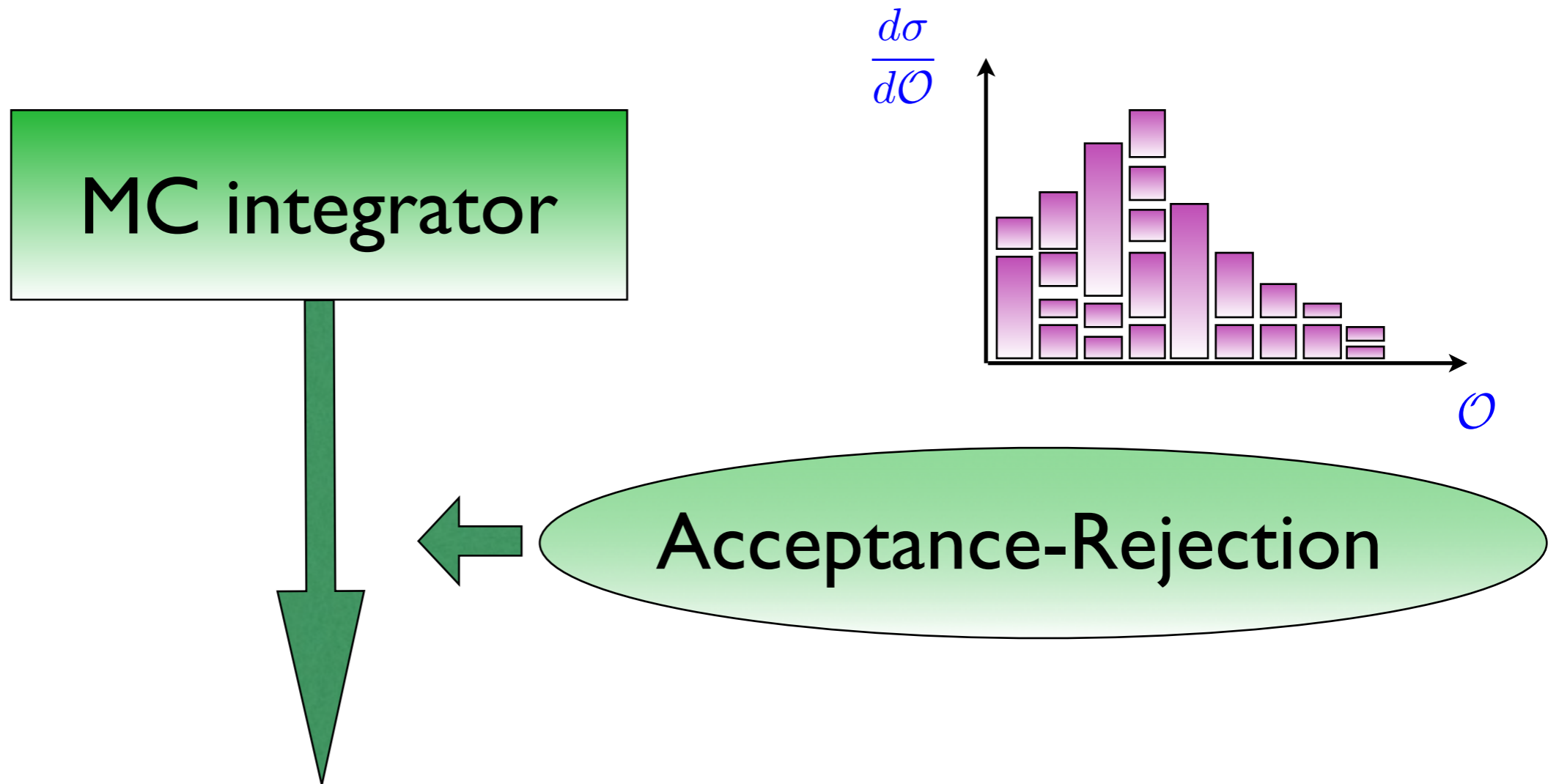
1. pick x distributed as $p(x)$
2. calculate $f(x)$ and $p(x)$
3. pick $0 < y < 1$
4. Compare:
if $f(x) > y p(x)$ accept event,
else reject it.

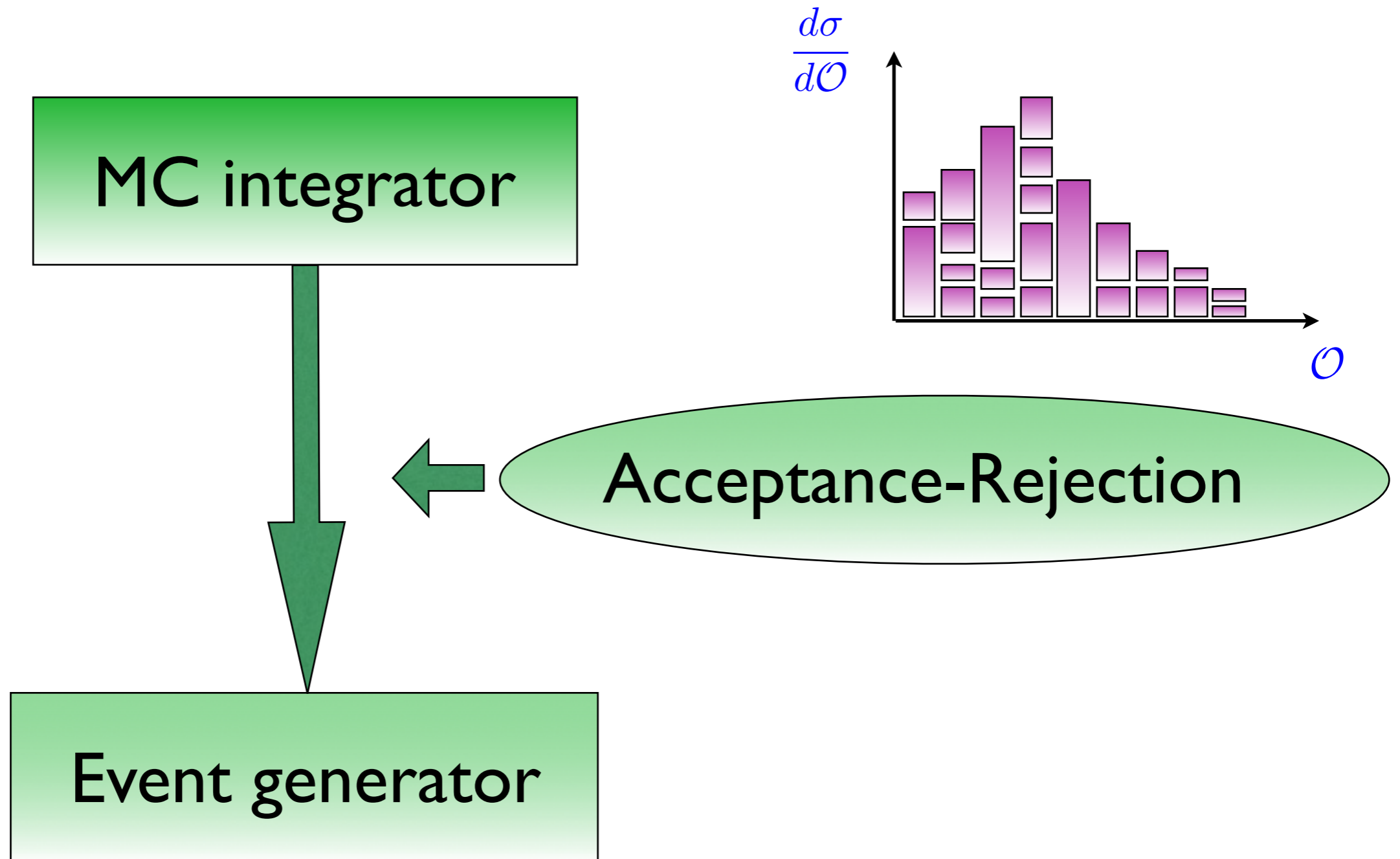
much better efficiency!!!

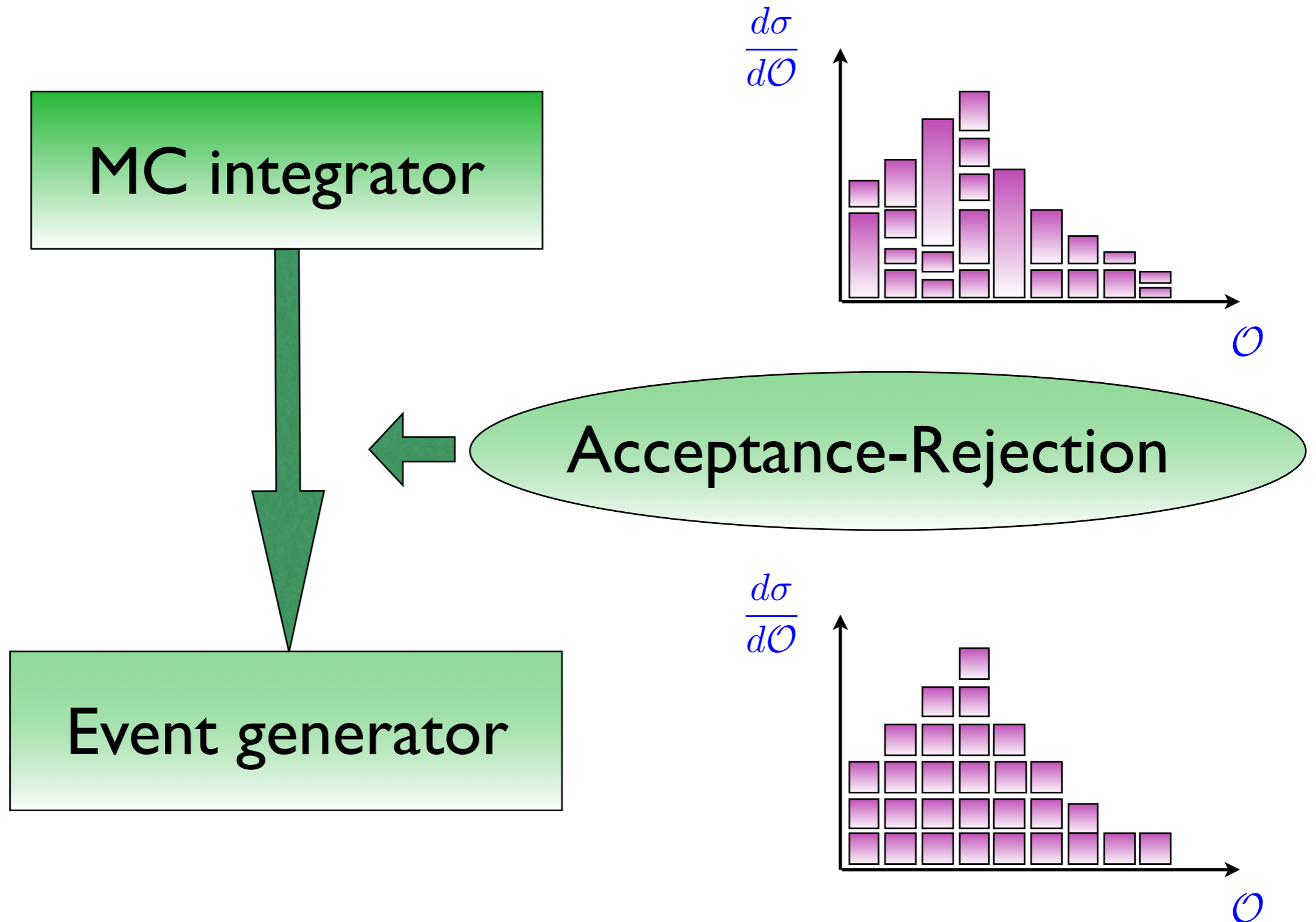
MC integrator

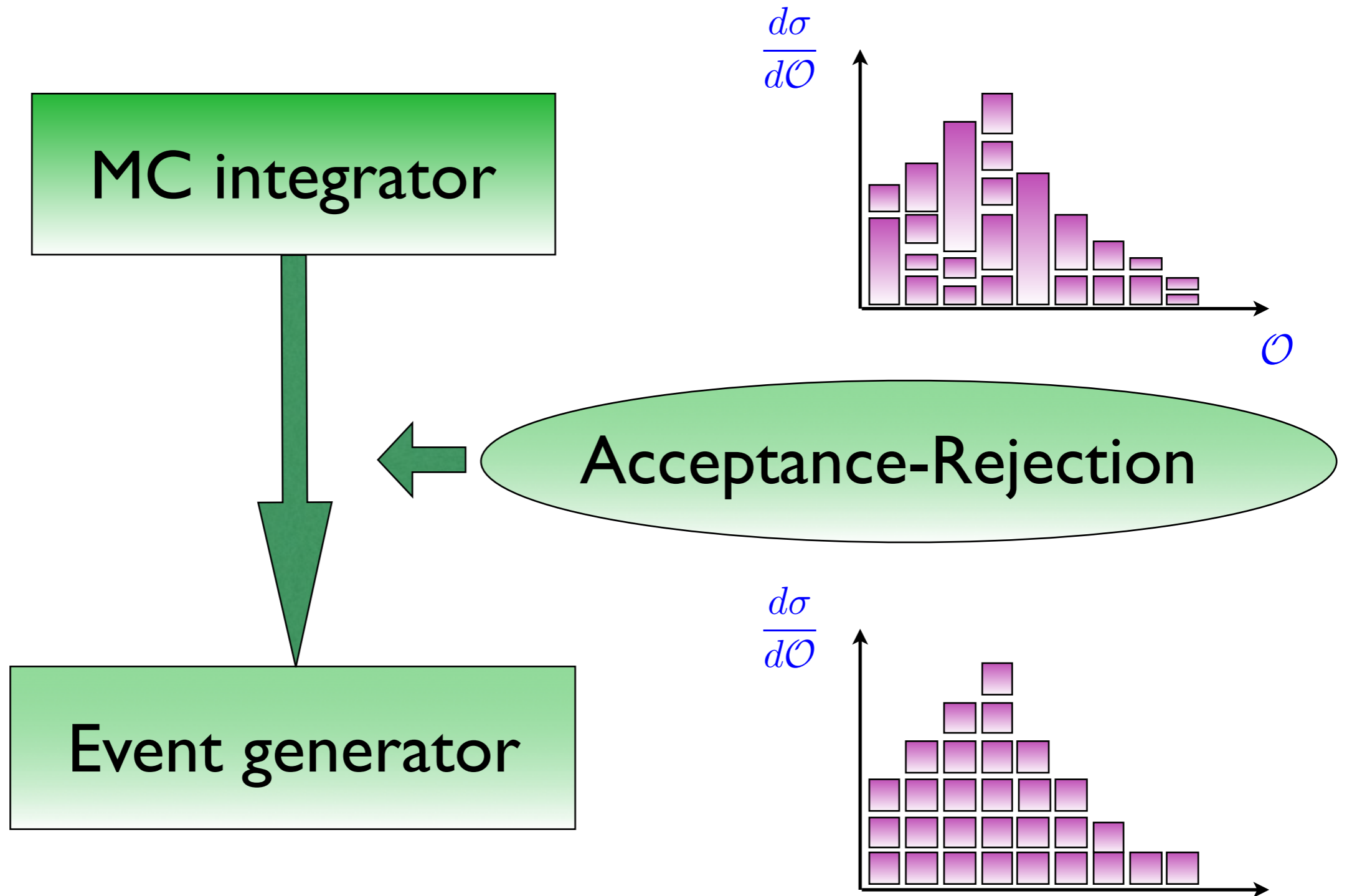
MC integrator











This is possible only if $f(x) < \infty$ AND has definite sign!

- Sample of unweighted events
 - ➔ Events distributed like nature
 - ➔ Need the function to be
 - Borne
 - Always positive
 - ➔ More efficient if the integration is more efficient
 - Same dependencies in the cut

Bad Point

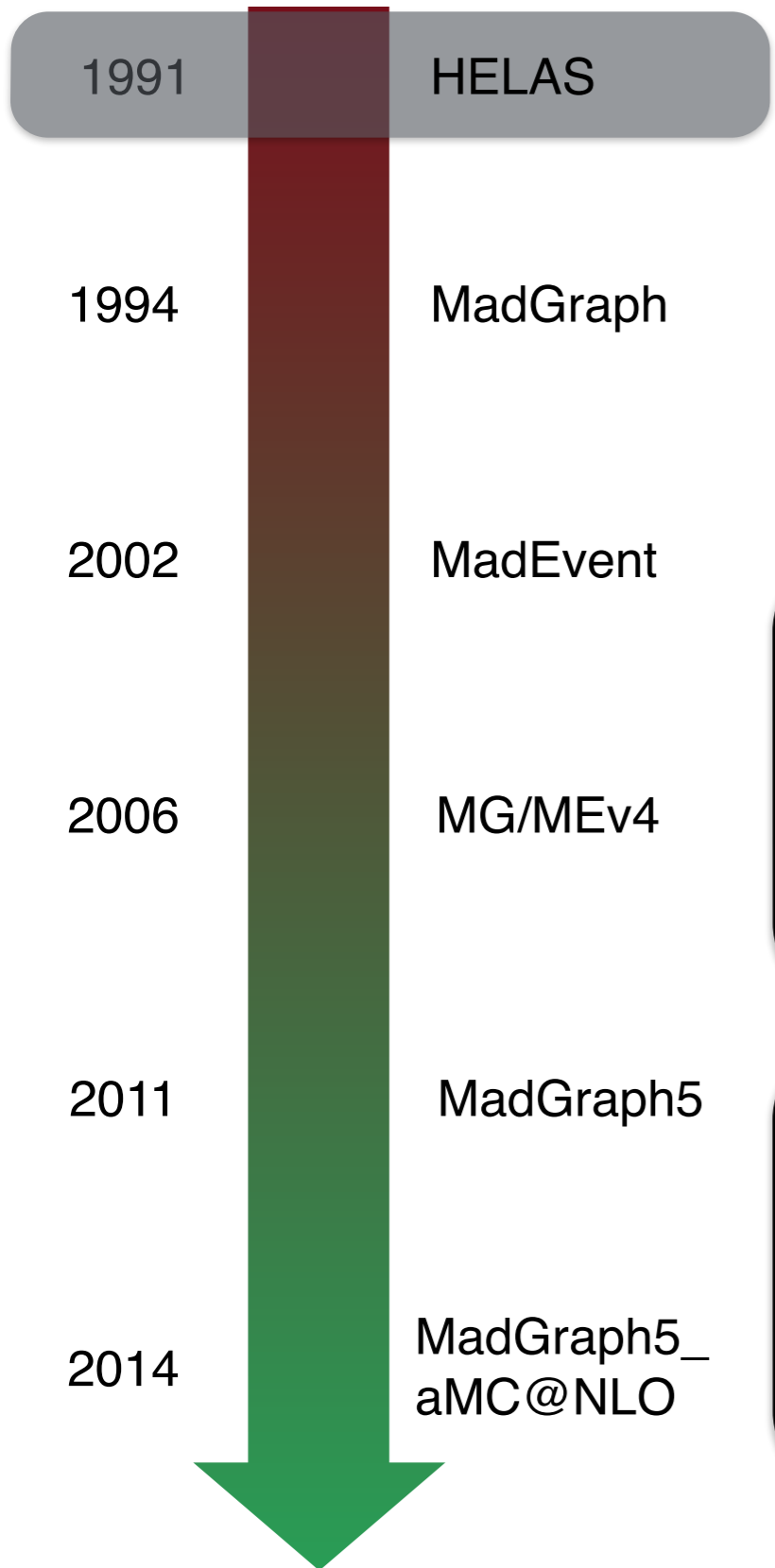
- Slow Convergence (especially in low number of Dimension)
- Need to know the function
 - Impact on cut

Bad Point

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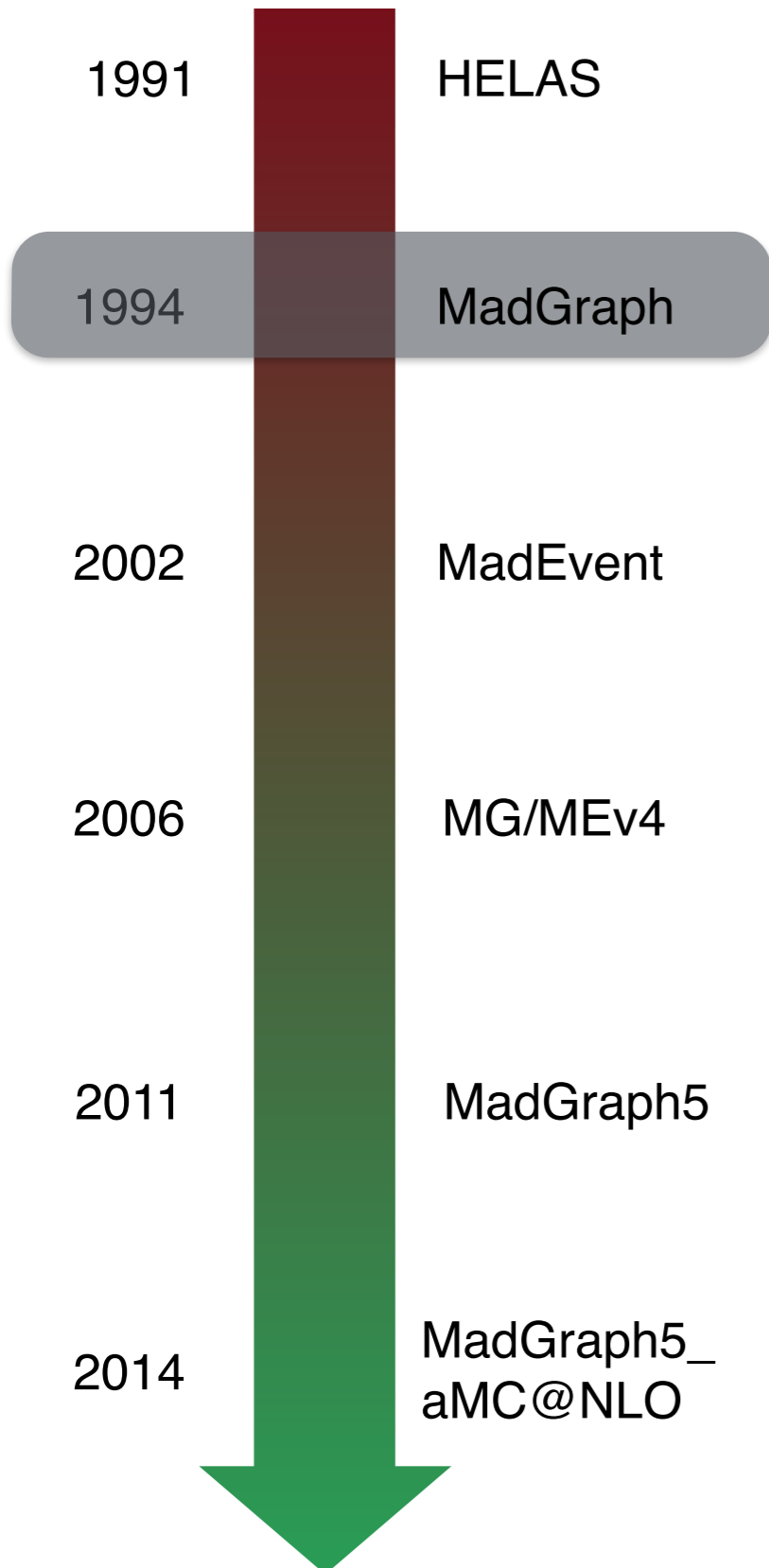
Good Point

- Complex area of Integration
- Easy Error estimate
- quick estimation of the integral
- Possibility to have **unweighted** events

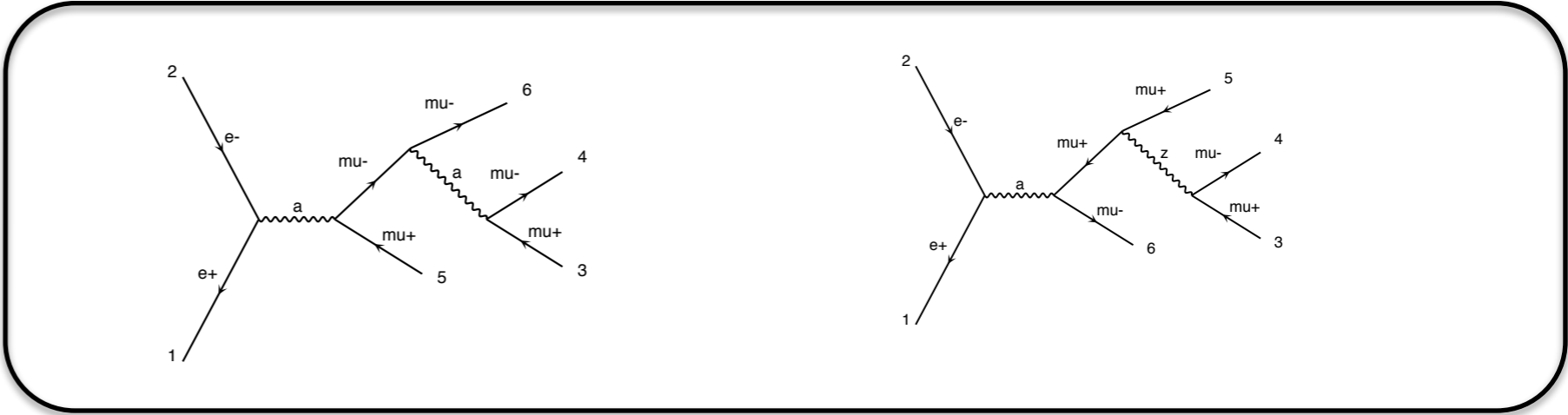


- Computing Matrix Element for a fixed Helicity and sum over the helicities.

- Suite of Routine, which allow to write the matrix element for any (SM) process

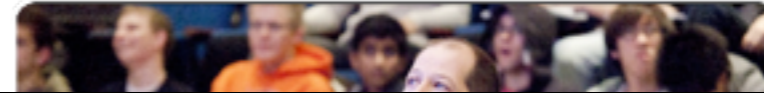


- Automate the creation of the diagram generation and the writing of the HELAS routine



1991

HELAS



MAD stands for Madison

1999



2000

2000

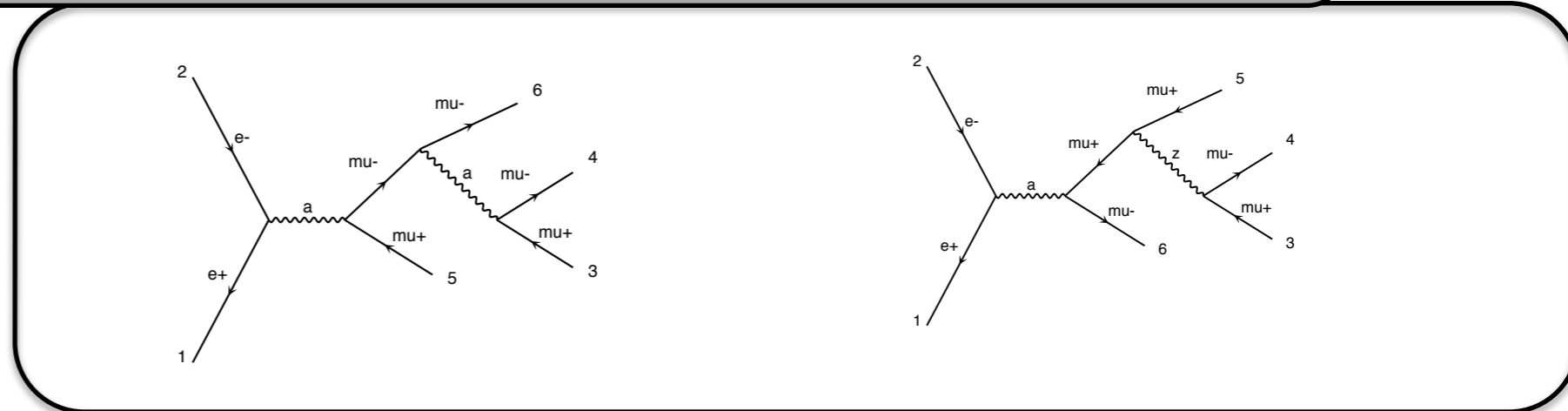
gram

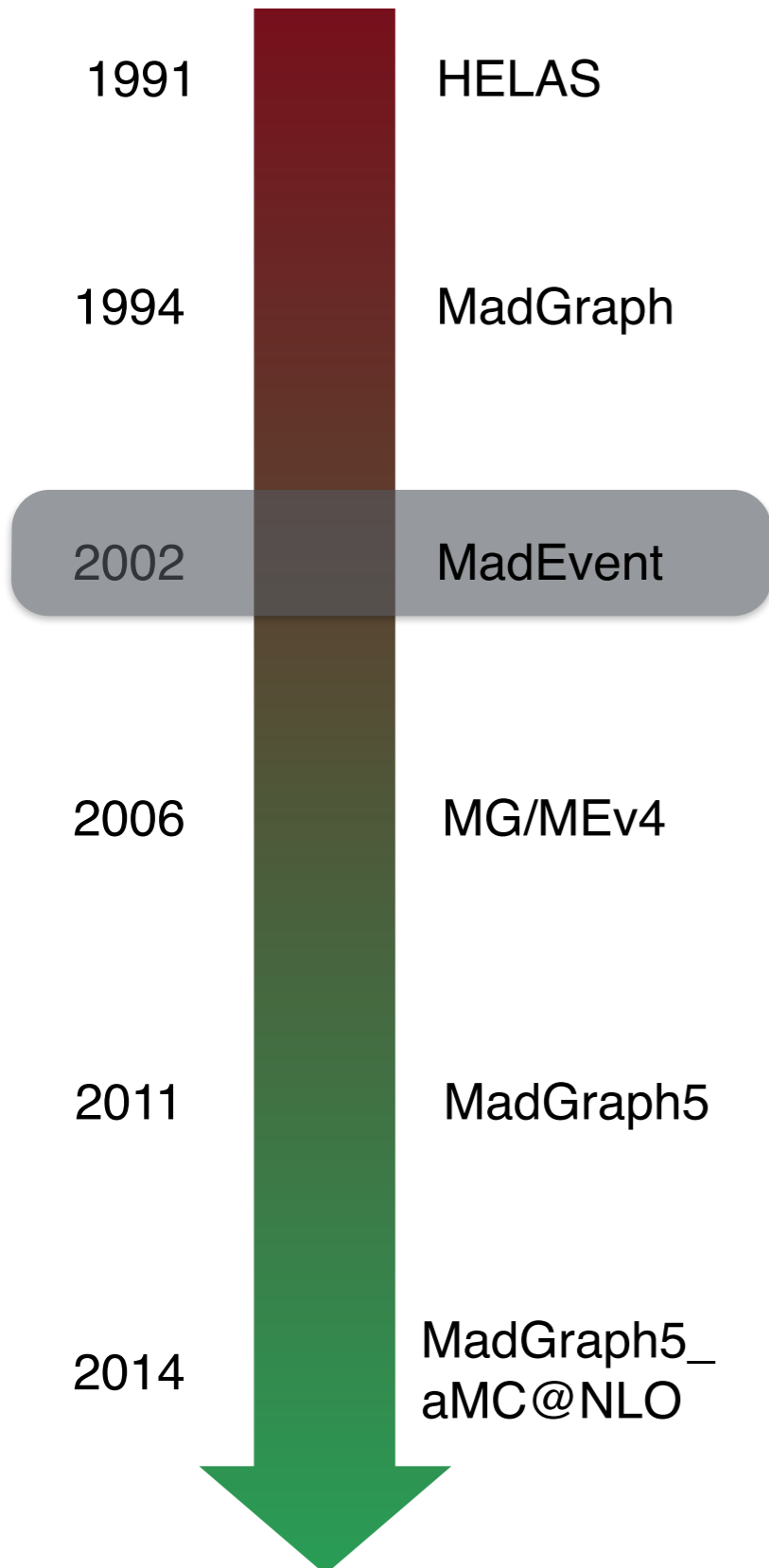
2011

MadGraph5

2014

MadGraph5_
aMC@NLO





- Multi-Channel Method!
- Automatic phase-space Integration
- Generation of Events

- Support for the MSSM (SMADGRAPH)





1991

HELAS

1994

MadGraph

2002

MadEvent

2006

MG/MEv4

2011

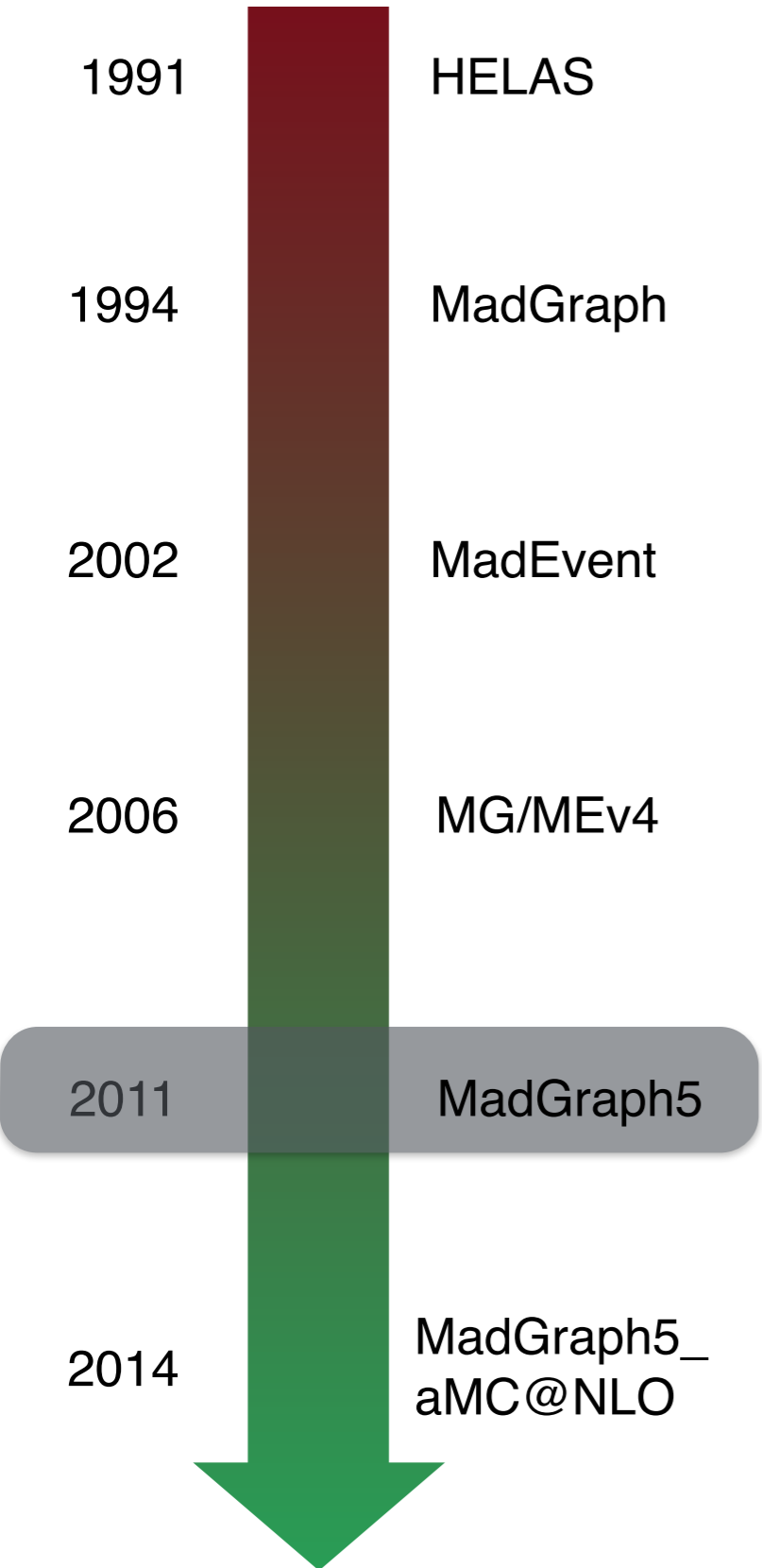
MadGraph5

2014

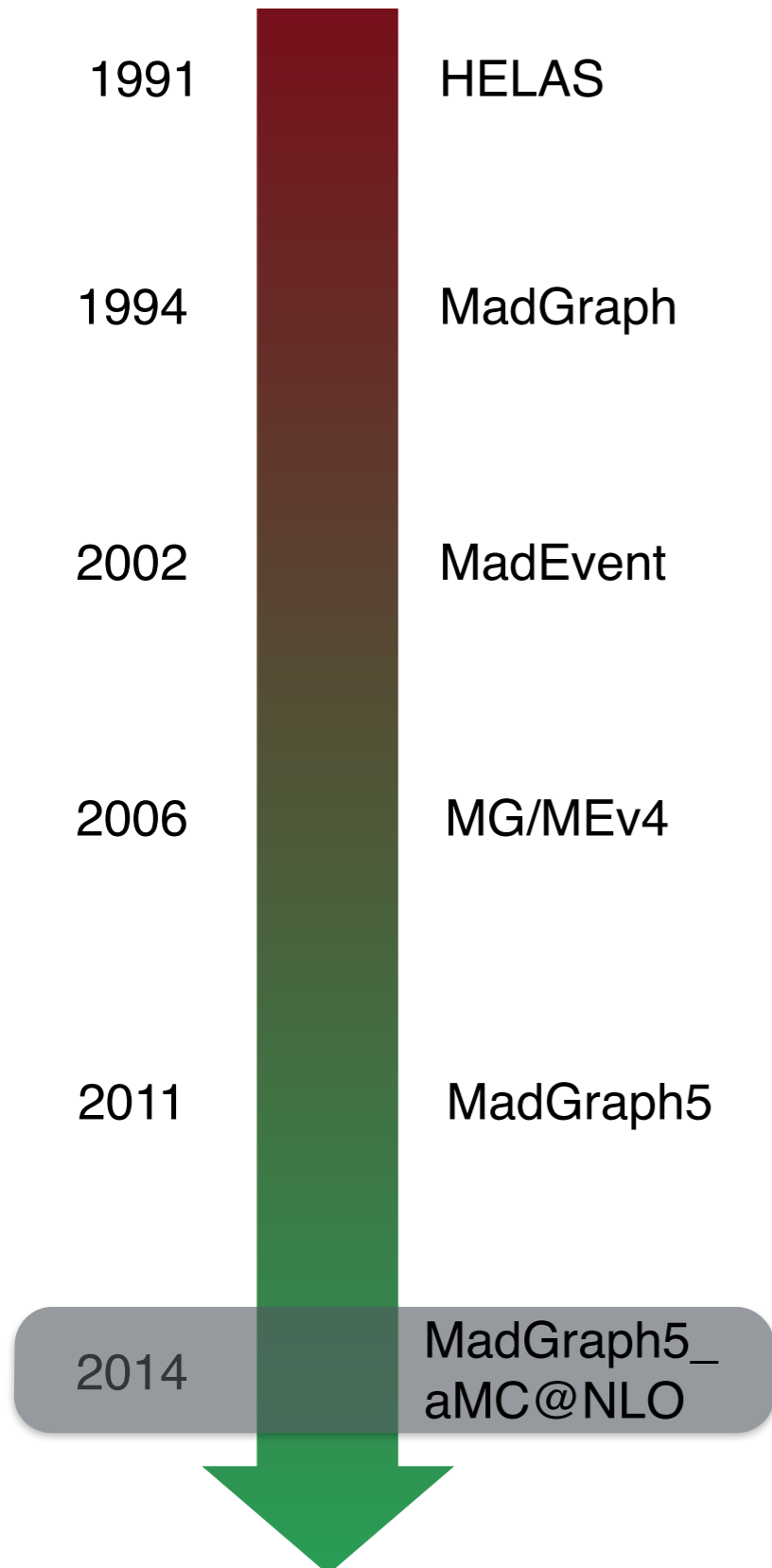
MadGraph5_
aMC@NLO

- Support for BSM
- Decay Chain
- Pass to a platform (MadOnia/MadWeight/...)
- Link to Pythia/PGS
- Matching/Merging

- Official/Main SM generator for CMS



- Full restart of the MadGraph part in Python
- Fully Automatic BSM
- Various Output Format
- Huge Improvement



- Fully Automatic computation at
 - NLO* (cross-section)
 - NLO* matched to PS

*NLO= NLO in QCD



- Leading Order Option
 - Support of BSM
 - Fermion Flow
- Computation of the Width
- Narrow width Approximation
 - Decay Chain
 - MadSpin
- Systematics
- NLO
 - SM with merging

$$\sum_{a,b} \int dx_1 dx_2 d\Phi_{\text{FS}} f_a(x_1, \mu_F) f_b(x_2, \mu_F) \hat{\sigma}_{ab \rightarrow X}(\hat{S}, \mu_F, \mu_R)$$

Phase-space integral
Parton density functions
Parton-level cross section

- The Importance of PDF
 - Defines the physics
- Evaluation of Matrix Element
 - Numerical method faster than analytical formula
 - cross-section prediction needs NLO
- Phase Space Integration
 - Need to know in advance what we integrate. Be careful with strong cuts!